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Some Current Topics on Nonlinear Conservation Laws

Lectures at the Morningside Center of Mathematics, 1

Ling Hsiao and Zhouping Xin, Editors

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Volume 15

Some Current Topics on Nonlinear Conservation Laws

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Series Preface

This is a subseries of AMS-IP publications dedicated to the research activities of the Morningside Center of Mathematics, Chinese Academy of sciences in Beijing, China. The Morningside Center supports research activities by inviting leading figures in different disciplines of mathematics to give lectures. These leaders direct programs in their disciplines. The participants are young Chinese mathematicians or some visitors from overseas. At the end of the programs, the lectures are collected and organized by the program directors. We believe these programs have been very successful and helpful in promoting mathematics as a whole.

We wish to thank the Morningside Foundation and the Chinese Academy of Sciences for providing the funding. We also thank the staff at our Morningside Center to help in our programs. We also like to take this opportunity to thank all the universities in China who have cooperated with us on these activities. Most of all, to the mathematicians all over the world who have provided unselfish help. We thank them for their efforts. We would also like to thank the International Press and American Mathematical Society for the help of publication.

S.T.Yau and YANG Lo

PREFACE

A year-long program on nonlinear conservation laws was organized by us at the Morningside Center of the Academic Sinica in Beijing during the academic year of 1997. The aim was to introduce young researchers in China to the frontiers of an important branch of applied mathematics that has been expanding rapidly. The program consisted of mini-courses on various current topics, short lecture series, invited talks, member seminars, and intensive working seminars. The emphasis was laid on theoretical analysis and applications. Important aspects, numerical methods for large scale discontinuous flows, were left for a future program. The program's success is partially reflected from the quality of this volume which consists of selected lecture notes for the mini-courses. The topics covered in this volume include: theory of L^1 -well-posedness for system of conservation laws by A. Bressan, theory of compartness methods by G. Chen, mathematical modelling of semi-conductor devices by P. Degond, kinetic formulation of conservation laws by B. Perthame, theory of viscous conservation laws by Zhouping Xin, and mathematical theory of ideal incompressible fluid by Yuxi Zheng. The volume also includes a survey article by Zhouping Xin.

The success of such an undertaking is due to the generous support of the Morningside Center of the Academic Sinica and the help of many individuals. We are particularly grateful to Professors S. T. Yau and L. Yang for their interests in this program and their support; to Professor Z. J. Lu for his constant practical support and help; and to Professor A. Bressan, G. Chen, P. Degond, B Perthame, and Y. Zheng for their lectures and contributions to this volume. We also wish to thank the junior participants in the program, Cheng He, Hailiang Li, R. H. Pan, Shaoqiang Tang, K. Zhang, Ping Zhang, P. Zhu, etc, especially Hai Liang Li, who have made various contributions to the success of the program and this volume.

Ling Hsiao and Zhouping Xin

Some Current Topics in Nonlinear Conservation Laws

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Many basic phenomena in continuium mechanics and various other branches of natural sciences are governed by nonlinear conservation laws of the form [33, 87, 178]:

$$\partial_x u + \nabla_x \cdot F(u) = \varepsilon A(u), \qquad u \in \mathbb{R}^n, x \in \mathbb{R}^m,$$

$$(0.1)$$

and its variants by taking into account other additional physical effects such as dispersion, relaxations, chemical reaction, and external sources, etc, besides the convection, compressions, and dissipation, where $\varepsilon \geq 0$, A is some elliptic operator, and $F(u) = (f_1(u), \dots, f_m(u))$ is a smooth nonlinear map. The convection is assumed to be hyperbolic in the sense that the $n \times n$ matrix $\xi \cdot \nabla_u F(u)$ has n real eigenvalues for all $\xi \in \mathbb{R}^m \setminus \{0\}$. The most important examples of (0.1) may be the following Navier-Stokes equations and many of their variants [33,87,174]:

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho u) = 0\\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla p = \operatorname{div}(T)\\ \partial_t(\rho E) + \operatorname{div}(\rho u E + up) = \operatorname{div}(uT) + k\Delta\theta \end{cases}$$
(0.2)

which express the fundamental physical laws in continuium mechanics: conservation of mass, momentum, and energy, where $x \in \mathbb{R}^d$, $t \in \mathbb{R}^1$; the unknown, ρ , $u \in \mathbb{R}^d$, p, eand θ denote the density, pressure, internal energy, and temperature respectively while E and T represent the total energy and stress tensor given by $E = \frac{1}{2}|u|^2 + e$

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and $T = \mu (\nabla u + (\nabla u)^t) + \mu'(\operatorname{div} u)I$ with μ and μ' being the first and second coefficients of viscosity respectively, and I is the $n \times n$ identity matrix; k is the coefficient of heat conduction. It is assumed that $\mu \ge 0$, $\mu' + \frac{2}{d}\mu \ge 0$, and $k \ge 0$. The relation between ρ , p, θ , and e is given by the equation of state for the fluid concerned and the second law of thermodynamics. For ideal inviscid fluids, $\mu' = \mu = k = 0$, then (0.2) becomes the well-known compressible Euler equations:

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho u) = 0\\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla p = 0\\ \partial_t(\rho E) + \operatorname{div}(\rho u E + u p) = 0 \end{cases}$$
(0.3)

which is one of the most important examples of systems of nonlinear hyperbolic conservation laws, which have been the main focus and the driving forces for mathematical theory of shock waves in the last centrary [33, 90, 178, 159]. Other examples of (0.1) include: the Euler and Navier-Stokes equations for incompressible fluids, the equations of elasticity, the equations of electro-magneto field dynamics for electrically conducting compression fluids, and equations for combustion, oil reservoir simulations, multiphase flows, etc. [87, 33, 178, 131, 129, 53, 135].

The most important feature of system (0.1) is that the speed of wave depends on the wave itself, which leads to great complexity and rich phenomena in the behavior of solutions to (0.1) [33, 178]. This can be revealed by understanding the dynamics and interactions of basic linear and nonlinear waves such as shock waves, rarefaction waves, solitons, diffusion waves, vortex sheets, and boundary layers, etc. [90, 129] The study on the properties of the solutions to (0.1) has important applications to many fields such as turbulence theory, geophysics, meteorology, material sciences, multi-phase flows, and chemical engineering, etc. This poses many problems which are both physically important and mathematically challenging.

Substantial progress has been made in solving these nonlinear conservation laws both theoretically and numerically, and in understanding and interpreting the behaviors of solutions to these systems by the important analytical works [90, 33, 50, 35, 40, 41, 55, 56, 97, 98, 102-104, 128, 141, 157, 159, 165-167, 13, 109, 57, 60, 162, 66, 2-3, 73, 16, 38, 86, 129, 136, 195, 32, 17], and many powerful modern high resolutions-numerical methods for calculating discontinuous flows (or waves with large gradients) have been developed [31, 63, 39, 147, 148, 70-72, 6, 7, 91, 93, 171]. This is particularly so in the one-space-dimension problems, and for suitably "small" solutions with some notable exceptions [4, 39, 98, 100, 85, 165]. Fundamental concepts such as basic linear and nonlinear waves, physical and Lax's geometrical entropy conditions, Riemann solutions, loss of regularity, dissipature mechanism of nonlinearity, compactness and irreversibility of the solution operator, etc, had been established for strictly hyperbolic systems (in one space dimension) almost three decades ago by the pioneering works in [50, 88-90, 159, 178]. A rather complete and satisfactory theory exists for scalar conservation laws in arbitrary

space dimensions in terms of well-posedness theory, regularity and compartness of solution operator, large time asymptotic behavior toward nonlinear waves and local structure of entropy weak solutions, and convergence of various approximate solutions generated by either physical perturbations (such as viscous, relaxational, and kinetic extensions, see the lectures by Chen, Perthame, and Xin in this volume and the references therein), or various numerical schemes [85, 93, 122,127, 142, 165, 168, 175, 76]. All the methods working in this case rely essentially on the availability of the maximum principle in this case. Thus the corresponding theory for systems of nonlinear conservation laws is extremely difficult. For some special systems where a weak maximum principle exists, the existence of weak solution in the space of bounded measurable functions had been proved by the theory of compensated-compactness (see the lectures by Chen and Perthame in this volume for details, also [77, 78, 133, 134, 41, 26, 27]). However, though this elegant approach requires no restrictions on the sizes of the initial (or boundary) data, it does demand that there exist sufficiently many entropy-entropy flux pairs, whose existence is only quaranteed for 2 x 2 systems [88, 155]. Furthermore, relatively little information can be obtained for such weak solutions obtained by the weak convergence methods compared with the ones obtained by the Glimm's method described below.

The most important breakthrough in the mathematical theory of shock waves is the celebrated random choice method pioneered by J. Glimm [55], which yields not only the existence of BV (bounded total variation) weak solutions for general I-D strictly hyperbolic systems of conservation laws with small BV data, but also precise asymptotic structures (both local and in large time) of such weak solutions [56, 36, 40, 104]. The Glimm theory is based on the fact that in one-space dimension, dialation invariant solutions, (Riemann solutions), which are superpositions of elemantary waves: shock waves, rarefaction waves, and contact discontinuities, dominate the asymptotic behavior of general flows, both locally and for large time, and thus form building blocks for general weak solutions. The deep understanding on the propogation and interactions of elemantary waves provides the insight of the construction of the now well-known Glimm's functional which is decreasing in time on the Glimm's approximate solutions consisting of piecewise Riemann solutions, as a consequence of the fact that the possible increase of the total wave strength is always balanced out by the decrease of the total future wave interaction potential. [55, 159]. Similar ideas play crucial roles in the later major refinements and improvments such as, the deterministic version of the Glimm scheme [103, 104], L^{∞} -norm estimate of the Glimm method [195, 167], and L^1 -stability theory of the Glimm's solutions [12, 13, 109]. Tremendous progress has been achieved on the structure, regularity, and large time asymptotic behavior of the Glimms solutions [12, 42, 56, 36, 104]. In particular, recently, Bressan successfully proved the uniqueness and L^1 -continuous dependence in the class of the viscosity solutions by a semigroup approach in a weighted space, this important recent development is discussed in detail in the lectures given by Bressan in this volume. More recently, the L^1 -continuous dependence of Glimm's solutions is shown by Liu and Yang by constructing a different interesting functional [109].

However, it seems that Glimm's method is limited to strictly hyperbolic systems in one-space-dimension since BV space is not a suitable solution space for hyperbolic system in more than one space dimensions [151]. Multi-dimensional theory for hyperbolic systems of conservation laws still has a long way to be developed with some notable exceptions: the short time structural stability of multi-dimensional gas dynamical shock fronts and rarefaction waves has been established [128, 131]; the theory of weakly nonlinear geometric optics has been systematically developed and applied to many interesting physical problems, such as reflection of shock wave by a wedge, and chemical reaction fluids, etc. (see [131, 68] and references therein); even rigorous justification of the asymptotic theory for resonant wave interactions has been achieved for some interesting cases [74, 131], the development of singularities in finite time from smooth data for Euler system (0,3) has been proved and the structure as shock wave has been verified in many important cases (such as irrotational flows) by the recent theory of geometrical blow-up for hyperbolic system in [2, 3, 156, 194]; and there have been many interesting studies on special multi-dimensional problems with various symmetries [26, 28, 59, 193].

In the last two decades or so, there have been interesting works by many people on the studies of non-strictly hyperbolic systems and equations that change types, which arise as models in Van der Waal gases, elastic-plastic materials, magnetohydro dynamics, multi-phase flows, and phase transitions [69, 82, 48, 107, 153, 157]. New types of shock waves are needed to solve Riemann problems [69, 153, 107]. These new shock waves behave drastically different from the classical shock waves, and their admissibility and nonlinear stability have been investivgated recently for some prototype models [48, 49, 106, 112], where new wave phenomena have been found. Even in the cases that the Riemann problems can be solved successfully, it seems that the great complexity in the interactions of elementary waves makes it almost impossible to apply Glimm's approach in general except in the case of small perturbation of strong waves [153]. For certain 2 x 2 systems, the theory of compensated compactness has been applied to obtain the general weak solutions [27, 78, 141]. A new approach seems to be desired to treat general non-strictly hyperbolic systems.

The inviscid Euler system (0.3) is expected to describe the large scale structures of general viscous flows which are governed by the more physical Navier-Stokes Equations. The studies on the general Navier-Stokes equations have been the center of the continuum mechanics. Two major issues on compressible Navier-Stokes equations have been investigated extensively in the past decades. One is the global (in time) well-posedness of smooth solutions and the large time asymptotic behavior for multi-dimensional compressible Navier-Stokes system with fixed viscosity and heat conduction which is of hyperbolic-parabolic type. This has been succesfully

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is shown to be governed by the linear diffusion waves [136, 159]. This theory has been generalized even for small discontinuous initial data in multi-dimensions in [64]. The other issue is to understand the effects of small scale dissipations on the behavior of large scale physical flows, i.e. to discribe the asymptotic relationship between the inviscid and viscous flows, in the presence of discontinuities and physical boundaries for either small (but non-zero) dissipation or large time (with fixed dissipation). This is a very important issue not only for the understanding of the asymptotic behavior of the solutions to the Navier-Stokes system, but also crucial for the theory of inviscid hyperbolic systems, which is due to the fact that it is the dissipation mechanism which selects the correct weak solutions to the inviscid system from other nonphysical solutions. Other considerations such as entropy production, entropy inequalities, or even the existence of visions shocks profile do not suffice, as examples in [48, 66] show. This is especially important for nonstrictly hyperbolic problems, problems in multi-dimensions and with boundaries. There has been a great success achieved in this direction in terms of nonlinear and linear waves in the past decade. Indeed, an almost complete theory on the nonlinear large time asymptotic stability toward basic viscous waves (viscous shock profiles, rarefaction waves, viscous contact waves, and diffusion waves) has been developed for quite general systems of viscous, strictly hyperbolic conservation laws [57, 80, 105, 108, 110-111, 137, 162, 163, 183, 186-187, 198-199], and even for some pro-type non-strictly hyperbolic models with viscosity [29, 49, 106, 112]. Some of these results have been generalized to multi-dimension for scalear equations and some special systems [58, 185]. Furthermore, many new interesting methods, such as rigorous mulit-scale matched asymptotic method, weighted-characteristic energy method, and approximate parametric method, etc., which build more hyperbolic elements into the conventional methods by taking into account more internal structure of the underlying flows, have been developed, and should be very useful for other problems. The insights gained in the nonlinear stability theory can be crucial in the studies of a related long standing problem - zero dissipation problem, which is to show that the regular solutions to the visious system converge (in a suitable topology) to entropy weak solutions of the corresponding inviscid system as the dissipation goes to zero, and to understand the structures and dynamics of the viscous solutions across discontinuities and boundaries for small but non zero viscosity. The study of this problem has been the main driving force for both theoretical analysis and numerical methods for scalar conservation laws [85, 93, 159]. However, major difficulties occur for general systems due to the lack of compactness of the viscous solutions except for some special 2 x 2 systems where some weak maximum principle is available so that the theory of the compensated compactness applies to show the convergence of viscous solutions to weak solutions to the corresponding inviscial systems, see the lectures by Chen and Perthame in this volume for details, also in [141]. In the past several years, a new approach, which is based on a multi-scale

matched asymptotic analysis and a-priori estimates motivated by the theory of nonlinear stability of viscous waves, has been developed to treat specific but physically interesting flows for general systems [60, 65, 181, 188, 190]. In particular, it is shown in [60] that piecewise smooth solutions with entropy-satisfying shock discontinuities to general strictly hyperbolic systems can be realized as limits of viscous solutions as viscosities tend to zero, and detailed local structures of the viscous solutions have been obtained, see also [196] for recent improvement. For more ideas on this approach, we refer the reader to the lectures by Xin in this volume. This approach has been applied to treat other singular perturbation problems [168]. The zero dissipation limit problem becomes more singular in the presence of physical boundaries due to the appearance of boundary layers. There have been extensive studies on the boundary layer theory due to its importance in the applications to high speed flows [154]. Most engineering literature deals with the formal asymptotic analysis of boundary layers for steady flows with some specific flow configurations [154, 47, 147]. In the case where the boundaries are uniformly non-characteristic, the multi-scale structure of the viscous solutions can be revealed by matched asymptotic analysis and the strong convergence of the viscous solutions to the inviscid flow away from the boundaries can be obtained by nonlinear stability analysis of the boundary layer before shock formation provided that the boundary layer is suitably weak [52, 62, 75, 5, 139, 181]. For the well-known non-slip boundary conditions for the Navier-Stokes equations, which is uniformly characteristic, the formal prandtl's boundary layer theory has existed since the early of this century, yet its validity has been proved rigorously for linearied flows in general Sobolev space in [182] and for nonlinear flows only in the analytical setting [152]. The study on the Prandtls boundary layer equations, which is a degenerate parabolic-elliptic system, is quite difficult in general Sobolev spaces. Only some partial results, such as finite time blow-up, long time existence of solutions for special class of monotonic initial data, and steady flows, exist [46, 47, 146].

Another major recent development in the theory of the compressible Navier-Stokes system is the global (in time) existence of large amplitude weak solutions to either the Cauchy problems or initial boundary value problems of the isentropic compressible Navier-Stokes equations with some special equations of state [98]. Although the uniqueness and regularities of such weak solutions remain open, some important physical properties, such as boundedness of the energy and the zero-mach number limit to the solutions of the incompressible Navier-Stokes equations, have been established for these weak solutions [98, 99]. The approach of this theory is a weak convergence method based on the standard energy estimate and an additional space-time higher estimate which is guaranteed by the special form of the equation of state [98]. It is remarkable that such a approach can work thus far since almost all the a -prior estimates required for this theory are well known. It should be noted that in the case where the initial data is not periodic, the existence of weak solutions [98] requires that the far fields must be vacuum, and thus the weak

solutions cannot be regular in general even for smooth initial data due to the fact that any non-trival smooth solution with compartly supported initial density must blow-up in finite time [180]. This somewhat surprising fact is a direct consequence of the observation that the pressure behaves dispersively for both Euler and Navier-Stokes systems in arbitrary space dimensions if the far fields are at vacuum [180, 132]. This instance shows the importance of understanding vacuum state in the theory of the compressible fluids besides the physical relevance, such as evolution of gaseous stars due to the high order degeneracy of both Euler and Navier-Stokes systems at vacuum. Some rich phenomena have been revealed in recent studies involving vacuum state, which include the dispersive behavior of the total pressure [180], the singular motions of the free surfaces separating the gases from vacuum, and the better regularity of the inviscid flows than the viscous ones, etc [113, 127].

The compressible Euler or Navier-Stokes systems either give macroscopic descriptions of interacting particles in a fluid motions as the mean free path tends to zero in the kinetic theory [21] or govern the equilibrium states of thermo-nonequilibrium process as the rate of relaxation goes to zero [22]. There have been increasing interests in studing the hydrodynamic limits for various kinetic models and the zero-relaxation limits for many thermo-non-equilibrium processes in the recent years [4, 18-21, 25, 61, 70-72, 81, 94, 101-102, 116-117, 123-124, 127, 133-134, 142-143, 145, 158-149, 157, 174, 176-177, 184, 188-189, 192, 199]. Besides the obvious physical significance of such limits problems, one would also hope that the complexities and many great difficulties such as loss of regularity, non-uniqueness of weak solutions, global entropy condition, and considerable difficulties in numerical calculations, encounted at the macroscopic level, can be illuminated and overcome by looking into the microscopic discriptions of the fluid motions which are not only physically sound but also mathematically tractable. Despite the substantial progress such as Chapman-Enskog theory [61, 21], initial layer theory [20, 18], and existence of shock profiles [19], etc have been achieved for the hydro-dynamic limit problems for various models of the nonlinear Boltzmann type equations, the global nonlinarity of the interaction operators for Boltzmann type transport equations and the stiffness of the hydro-dynamic limit make it extremely difficult (if not impossible) to study the limit rigorously and to gain deep insight about the macroscopic flows, in particular, in the presence of discontinuities and physical boundaries [19, 21, 116-117, 174, 176, 184, 192] One of recent efforts has been concentrated on designing better (in the sense that they are physically reasonable and amiable to mathematical analysis both analytically and numerically) kinetic models or relaxation approximations for a given system of a hyperbolic conservation laws [11, 70-72, 79, 101, 143, 148-149, 192]. One such approach is the so called kinetic formulation of conservation laws, which artificially construct a collision operator in a similar spirit as that for the BGK model of Boltzmann equation [21] so that the Maxwellian satisfies the given conservation laws [149, 101]. This approach was successfully applied to multi-dimension scalar conservation laws and some special

systems, for more detailed discussions on this approach, we refer the reader to the lectures by B. Perthame in this volume and the references therein. Another approach is the local relaxations approximation methods introduced recently by Jin and Xin [70]. The basic idea is that for any given general system of nonlinear conservation laws, one can construct a corresponding linear hyperbolic system with a nonlinear stiff source term that approximates the original system of conservation laws with a small dissipative correction under the so-called "subcharacteristic condition" which ensures the relaxation system contains all the information of the original system. The special structure of the relaxation system proposed in [70], such as linear characteristic fields with constant waves speeds; localized nonlinear interaction with relaxational structure; and the stiff source terms being not-fully ranked and linear in the induced variables which approximate the fluxes of the original conservation laws, makes it more advantagerous to solve the relaxation system numerically than the original system of conservation laws. Indeed, the resulting relaxation schemes are high order accurate, entropy preserving, and requiring neither the Riemann solvers nor nonlinear algebraic solvers [70]. The main features of the relaxation schemes in [70-72] are their simplicity and generality. General systems of higher space dimensions can be treated in the same way as in the one dimension. Furthermore, there is no requirement of hyperbolicity in the formulation, thus the relaxation schemes should be useful in many problems such as in MHD, multiple phase flows, and elastic-plastic materials, etc., in which Riemann solvers (or ever approximate Riemam solvers) are not easily available. Various extensions of the local relaxation idea exist [71, 72, 79, 143], and rigorous analysis of the asymptotic equivalence between and relaxation system and the limiting system of conservation laws has been obtained in many interesting cases [123-124, 127, 142, 168, 176, 192].

On the other hand, in many pratical applications, the macroscopic models, are considerably simpler both conceptually and numerically than their macroscopic kinetic models [21, 22, 135, 174], and multi-scale expansion methods, such as Hilbert and Chapman-Enskog expansions are very powerful to derive useful macroscopic models from the more fundamental kinetic descriptions. One particular case is the mathematical modelling of electron transport in semi-conductor materials [135]. A semiconductor is a solid-state material with an intrinsically low electron conductivity. On the quantum level, the state of a free electron is governed by the Schrödinger equation. In the semi-classical limit, the kinetic description of the electron and hole population in a semiconductor device is given by the nonlinear Boltzmann equations governing the distribution functions with the electrostatic potential satisfying a Poission equation. However, such a detailed description for carrier transport in semiconductors is too complicated for pratical purposes, and many numerical codes based on Monte-Carlo or particle methods, are very expensive and cannot be routinely used to design components. Various macroscopic models have been sought. The mostly widely used model in practice is the so called drift-diffusion model, which is a parabolic equation for electron density and has been extensively studied both

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mathematically and numerically since early fifties [51, 125, 133, 135]. The improvement of the drift-diffusion model is the Energy-Transport model, which consists of a parabolic system for the electron density and energy, and has been derived from the Boltzmann equation by multiscale analysis only recently [8]. The Energy-Transport model, which includes the drift-diffusion model as a particular case, resembles in many ways the compressible Navier-Stokes equations and has been studies extensively recently [8, 37]. These topics and other models are discussed in great detail in the lectures by P. Degond in this volume.

In the limit of small Mach number, one can derive from the compressible Euler and Navier-Stokes systems the incompressible Euler and Navier-Stokes equations respectively [83, 99], which govern various flows such as water in the oceans and air in the atmosphere. The mathematical theory for the incompressible fluid is more mature than that for the compressible one (see [32, 86, 97, 129, 179]). In particular, the well-posedness theory for smooth solutions exists in 2-D for both Euler and Navier-Stokes systems thanks to the fact that the vorticity of the flow satisfies a maximum principle [97, 129]. There have been some significant results concerning physically important inviscid 2-D singular flows: the motions of vortex patch and vortex sheets [129, 160, 140, 84, 150, 164, 17, 119-120]. For the vortex patch problem, here the vorticity is bounded, the well-posedness is proved long age [197], even the preservation of the regularity of the boundary of the vortex patch has been achieved recently [23, 10]. While the progress has been much slow and limited for the vortex sheets problem, where the vorticity is a finite Radom measure concentrated on a curve which is smooth initially, the global (in time) existence of classical weak solutions, and convergence of approximate solutions generated by either viscous regularization or vortex method have been achieved only recently for the case that the initial vorticity has single sign [38, 44, 115, 118, 130] except for special flows [121], though the short time existence in the analytic setting have been extensively studied [17, 129, 140]. Many important questions such as uniqueness of weak solutions, structure of the vortex sheet after roll-up, and behavior of approximate solutions remain open. For more discussion on this and related problems, we refer to the lectures by Y. Zheng in this volume. Much less progress has been made for 3-D incompressible flows, although the Leray-Hopf weak solutions to the 3-D Navier-Stokes equation have been shown to exist globally in time for quite general data and exterior forces [32, 86], the partial regularity and large time behavior of the Leray-Hopf weak solutions have been obtained [16, 95, 96, 169, 170], and the existence of large amplitude strong solutions under certain symmetries has been established [173, 86]. Yet, most central questions in the incompressible fluid mechanics, such as whether finite time singularities do develop from smooth initial data for both inviscid and viscous flows, the uniqueness and regularity of Leray-Hopf weak solutions with large amplitudes, long time dynamics of viscous flows, and instability of visious boundary layers, etc., remain completely open and challenge the field for many years to come, though some special types of blow-up for

Navier-Stokes, such as self similar blow-up, has been ruled out recently [144, 170, 172]. Even the existence of weak solution to the inviscid Euler equations has not been proved except the cases with special symmetries [129, 173], and the existence of measure-valued solutions [45].

Another most significant achievement in the past several decades is the development of many powerful modern high resolution numerical methods for calculating large scale flows governed by (0.1) [29, 31, 63, 147, 93, 70, 6, 7, 120, 84]. This is an important but difficult task due to the possible appearance of transitional layers, (such as shock layers, boundary layers, and shear layers etc.) in the flows governed by (0.1). Rich phenomena such as smearing, post-shock oscillations, discrete shock profile, and numerical cell-entropy conditions, etc, have been found. For inviscid flows, the finite difference schemes can be classified into either front-tracking method, or shock-capturing method, and hybridizations of these two [30, 63, 91, 171]. The convergence of many pratical scheme has been proved for scalar equations and some special systems, see [93] and the references therein. However, since most of the lectures in this volume deal mainly with theoretical issues, I will omit the detailed discussions on this important topic.

Despite the important progress achieved in the past on the study of (0.1) and related equations, many fundamental questions remain to be answered and continue to challenge the field for many years to come. Since I have discussed some open problems for (0.1) in great detail in [191], I will sketch some general problems here, and refer the readers to [191] for more discussion on the motivation, related work, and significance of these problems.

1. Well-posedness and local structure of BV weak solutions for 1-D strictly hyperbolic systems (in particular, for the full gas dynamic system) with Cauchy data of arbitrary amplitude. Is the solution unique? Can a weak solution to the full gas dynamics equations in BV space blow-up in finite time [138]? Can a BV solution behave locally as an appropriate peturbation of a Riemann solution as Glimm's solutions do [103]?

2. Local structure and uniqueness of limits of zero dissipation. If a weak solution to the 1-D inviscid strictly hyperbolic system is a limit of solutions to the corresponding viscous system as the viscosity tends to zero, then what is the local structure of such a solution? can it be a viscosity solution defined by Bressan in [12, 13]? Can one show that such a solution is unique in the class of weak solutions obtained as limits of zero dissipation?

3. Spatial periodic solutions and homogenization theory for general 1-D strictly hyperbolic systems. For a given strictly hyperbolic systems without a Riemann invariant coordinates (the 3×3 full gas dynamical system in particular), when do shocks form from a smooth spatial periodic flow [90, 92]? Is there a global weak

solution with periodic initial data? How do the oscillations propogate for such system [56]?

4. Theory of well-posedness of weak solutions to general nonstrictly hyperbolic systems [27, 48, 69, 78, 82].

5. Multi-dimension inviscid shock wave theory. Due to the complexity and lack of understanding, one should concentrate on the multi-dimension compressible Euler system. Despite the recent progress in geometrical blow-up theory [1, 2, 3, 156, 194], the singularity structure of solutions to Euler system has not be found unless the flow is assumed to be irrotational [2, 194]. In general, do shocks form before shell singularities? Can one study some of the special physically relevant wave patterns where a lot of experimental data, numerical simulations, and asymptotic results are available, such as: shock reflection phenomena [9, 54, transonic flows and standing shocks in a nozzle [59], self-similar flows, and other flows with various symmetries [28, 33, 193, 200].

6. Global (in time) well-posedness of the Cauchy problem or initial-boundary value problem for the compressible Navier-Stokes equation with large data in multidimension. Can smooth solution develop finite time sigularity [180]? Can the vacuum state be formed in finite for non-vacuum date? Can one obtain regularity for the weak solution constructed by P. L. Lions [98, 180]?

7. Nonlinear stability of planary viscous nonlinear waves for viscous systems is multi-dimension (in particular, for 2D or 3D compressible Navier-Stokes equations). When are basic nonlinear wave patterns: such as planar viscous shock profiles and viscous rarefaction waves, nonlinearly stable under generic multi-dimensional perturbations? What are the large time asymptotic behavior toward a contact wave for solutions to the compressible Navier-Stokes equations?

8.Shock and boundary layer theory for the compressible Navier-Stokes equations. Can one prove the asymptotic equivalence of the Euler and Navier-Stokes systems in the limit of small viscosity and heat conduction in the case that the inviscid flow contains finitely many shock discontinuities [60]? Can one justify the Prandtl's boundary layer theory [182, 154]?

9. 2-D vortex sheets motion. What are the structures of the approximate solutions generated by either Navier-Stokes solutions or vortex methods with general vortex sheets initial data [38, 130, 115, 118]? Doses concentration-cancellation always occur [44]? Does there exist dynamical energy-concentration for such approximate solutions with initial vorticity either one sign or both signs [44]? In the case of bounded domain, can vorticity concentrate near the physical boundary [121]?

10. Fluid-dynamic limits. Study the fluid-dynamic limit problems for general models of the nonlinear Boltzmann equations in the presence of discontinuities and physical boundaries [116-117, 158, 177, 184, 188-189]. What will be the fluid-dynamic limit for the renormalized weak solutions to the Boltzmann equation obtained by Diperna-Lions in [43]?

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