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Centre de Recherches Mathématiques Montréal

# Function Theory on Symplectic Manifolds 

Leonid Polterovich Daniel Rosen

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## Preface

The symplectic revolution of the 1980s gave rise to the discovery of surprising rigidity phenomena involving symplectic manifolds, their subsets, and their diffeomorphisms. These phenomena have been detected with the help of a variety of novel powerful methods, including Floer theory, a version of Morse theory on the loop spaces of symplectic manifolds. A number of recent advances show that there is yet another manifestation of symplectic rigidity, taking place in function spaces associated to a symplectic manifold. These spaces exhibit unexpected properties and interesting structures, giving rise to an alternative intuition and new tools in symplectic topology, and providing a motivation to study the function theory on symplectic manifolds, which forms the subject of the present book.

Recall that a symplectic structure on a $2 n$-dimensional manifold $M$ is given by a closed differential 2 -form $\omega$ which in appropriate local coordinates is given by $\omega=\sum_{j=1}^{n} \mathrm{~d} p_{j} \wedge \mathrm{~d} q_{j}$. The Poisson bracket of a pair of smooth compactly supported functions $F, G$ on $M$ is a canonical operation given by

$$
\{F, G\}=\sum_{j}\left(\frac{\partial F}{\partial q_{j}} \frac{\partial G}{\partial p_{j}}-\frac{\partial F}{\partial p_{j}} \frac{\partial G}{\partial q_{j}}\right)
$$

The Poisson bracket, which is one of our main characters, plays a fundamental role in symplectic geometry and its applications. For instance, it governs Hamiltonian mechanics. The symplectic manifold $M$ serves as the phase space of a mechanical system. The evolution (or Hamiltonian flow) $h_{t}: M \rightarrow M$ of the system is determined by its time-dependent energy $H_{t} \in C^{\infty}(M)$. Hamilton's famous equation describing the motion of the system is given, in the Heisenberg picture, by $\dot{F}_{t}=\left\{F_{t}, H_{t}\right\}$, where $F_{t}=F \circ h_{t}$ stands for the time evolution of an observable function $F$ on $M$ under the Hamiltonian flow $h_{t}$. The diffeomorphisms $h_{t}$ coming from all possible energies $H_{t}$ form a group $\operatorname{Ham}(M, \omega)$, called the group of Hamiltonian diffeomorphisms. For closed simply connected manifolds this group is just the identity component of the symplectomorphism group. The group Ham can be considered as an infinite-dimensional Lie group. The function space $C^{\infty}(M)$ is, roughly speaking, the Lie algebra of this group, and the Poisson bracket is its Lie bracket.

The structure of the function theory we are going to develop can be illustrated with the help of the following picture. Fix your favorite $t>0$


Figure 0.1. Two opposite regimes
and consider the natural mapping $C^{\infty}(M) \rightarrow \operatorname{Ham}(M)$ which takes a (timeindependent) function $H$ to the time- $t$ map $h_{t}$ of the corresponding Hamiltonian flow. In principle, this mapping enables one to translate information about Hamiltonian diffeomorphisms (which nowadays is quite a developed subject, see Chapter 4) into the language of function spaces. This naive plan works successfully in two opposite regimes, infinitesimal (when $t \rightarrow 0$ ) and asymptotic (when $t \rightarrow \infty$ ) (see Figure 0.1).

Working in the infinitesimal regime, one arrives at a surprising phenomenon of $C^{0}$-robustness of the Poisson bracket. Observe that the expression for the Poisson bracket involves the first derivatives of the functions $F$ and $G$. Nevertheless, the functional $\Phi(F, G):=\|\{F, G\}\|$, where $\|\cdot\|$ stands for the uniform norm of a function, exhibits robustness with respect to $C^{0}{ }^{-}$ perturbations. In particular, as we shall show in Chapter 2, $\Phi$ is lower semicontinuous in the uniform norm. Even though this result sounds analytical in nature, it turns out to be closely related to a remarkable bi-invariant geometry on the group $\operatorname{Ham}(M, \omega)$ discovered by Hofer in 1990. We shall discuss various facets of $C^{0}$-robustness of the Poisson bracket. One of them is the Poisson bracket invariant of a quadruple of subsets of a symplectic manifold discussed in Chapter 7. Its definition is based on an elementary looking variational problem involving the functional $\Phi$, while its study involves a variety of methods of "hard" symplectic topology. Another facet is symplectic approximation theory, discussed in Chapter 8 . Its basic objective is to find an optimal uniform approximation of a given pair of functions by a pair of (almost) Poisson commuting functions.

The asymptotic regime gives rise to the theory of symplectic quasi-states presented in Chapter 5. A symplectic quasi-state is a monotone functional $\zeta: C^{\infty}(M) \rightarrow \mathbb{R}$ with $\zeta(1)=1$ which is linear on every Poisson-commutative subalgebra, but not necessarily on the whole function space. The origins of this notion go back to foundations of quantum mechanics and Aarnes' theory of topological quasi-states, an interesting branch of abstract functional analysis. In our context, nonlinear quasi-states on higher-dimensional manifolds are provided by Floer theory, the cornerstone of modern symplectic topology. Interestingly enough, symplectic quasi-states are closely related to quasi-morphisms on the group of Hamiltonian diffeomorphisms $\operatorname{Ham}(M, \omega)$.

Roughly speaking, a quasi-morphism on a group is "a homomorphism up to a bounded error." This group-theoretical notion coming from bounded cohomology has been intensively studied in the past decade due to its various applications to geometry and dynamics. We discuss it in Chapter 3, A recent survey of quasi-states and quasi-morphisms in symplectic topology can be found in Entov's ICM-2014 talk [57].

Quasi-states serve as a useful tool for a number of problems in symplectic topology such as symplectic intersections, Hofer's geometry on groups of Hamiltonian diffeomorphisms, and Lagrangian knots. These applications are presented in Chapter 6. In addition, quasi-states provide yet another insight into robustness of the Poisson brackets, see Section 4.6 ,

Besides applications to some mainstream problems in symplectic topology, function theory on symplectic manifolds opens up a prospect of using "hard" symplectic methods in quantum mechanics. Mathematical quantization and, mostnotably, the quantum-classical correspondence principle provide a tool which enables one to translate basic notions of classical mechanics into quantum language. In general, a meaningful translation of symplectic rigidity phenomena involving subsets and diffeomorphisms faces serious analytical and conceptual difficulties. However, such a translation becomes possible if one shifts the focus from subsets and morphisms of manifolds to function spaces. We present some first steps in this direction in Chapter 9 ,

The book is a fusion of a research monograph on function theory on symplectic manifolds and an introductory survey of symplectic topology. On the introductory side, the first chapter discusses some basic symplectic constructions and fundamental phenomena, including the Eliashberg-Gromov $C^{0}$-rigidity theorem, Arnold's symplectic fixed point conjecture, and Hofer's metric, while in the last three chapters the reader will find an informal crash course on Floer theory. Even though our intention was to make the book as self-contained as possible, the reader is encouraged to consult earlier symplectic literature, such as the classical monographs [107, 108] by McDuff and Salamon. We also refer the reader to the manuscript by Oh [121] on Floer theory. The reader is assumed to have familiarity with basic differential and algebraic topology.

Most of the results presented in the book are based on a number of joint papers by L.P. with Michael Entov. L.P. expresses his gratitude to Michael for long years of pleasant collaboration. Furthermore, some central results of the book are joint with Lev Buhovsky (Poisson bracket invariants and symplectic approximation), Yakov Eliashberg (Lagrangian knots), and Frol Zapolsky (Poisson bracket inequality and rigidity of partitions of unity). L.P. cordially thanks all of them.

Parts of the material have been taught by L.P. in graduate courses at University of Chicago and Tel Aviv University, in a lecture series at UCLA, and (with the assistance of D.R.) in a mini-course at University of Melbourne. We thank these institutions for such an invaluable opportunity. We
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Figure 0.2. Subject road map. Numbers next to arrows indicate relevant chapters.

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## Notation Index

| $a \cap b$ | intersection product in homology, 156 |
| :---: | :---: |
| $A(x)$ | operator valued expectation, 134 |
| $A * B$ | quantum product, 175 |
| $\mathcal{A}_{F}$ | action functional, 51 |
| $\hat{A}$ | POVM associated to a von Neumann observable, 131 |
| $\mathcal{B}(F)$ | singly generated subalgebra of $C(X), 70$ |
| $B \rightarrow A$ | smearing of POVMs, 139 |
| $c(a, F)$ | spectral numbers of smooth functions, 147 spectral invariants in Floer theory, 179 |
| $c(a, \phi)$ | spectral invariants in Floer theory, 180 |
| $c_{1}$ | first Chern class, 16 |
| Cal | Calabi homomorphism, 49 |
| $\mathrm{Cal}_{U}$ | Calabi homomorphism on $\mathcal{G}(U), 57$ |
| $\mathrm{CF}(F)$ | Floer chain complex, 169 |
| $\mathrm{CF}_{i}(F)$ | $\mathbb{Z}_{2}$-graded Floer chain complex, 169 |
| $C_{c}^{\infty}(M)$ | smooth compactly supported functions, 21 |
| $\mathrm{CM}_{*}(F)$ | Morse chain complex, 153 |
| $C(M)$ | space of continuous functions, 65 |
| $\mathrm{CN}_{*}(\alpha)$ | Morse-Novikov chain complex, 159 |
| $\mathrm{CO}(X)$ | subsets of $X$ that are open or closed, 72 |
| CovDim | covering dimension, 74 |
| Crit(F) | set of critical points of $F, 151$ |
| $\operatorname{Crit}_{k}(F)$ | set of critical points of Morse index $k, 151$ |
| CZ | normalized Conley-Zehnder index, 44 |
| $\overline{C Z}$ | Conley-Zehnder index, 44 |
| $\Delta_{A}$ | noise operator, 134 |
| Diff( $M$ ) | group of compactly supported diffeomorphisms, 1 |
| $\mathbb{E}(\hat{A}, \xi)$ | expectation of a von Neumann observable at a pure state, 135 |
| $e_{c}$ | spectral displacement energy, 54 |
| $e_{H}$ | (Hofer's) displacement energy, 13 |
| End ( $H$ ) | space of bounded linear operators on $H, 134$ |
| $\mathcal{F}$ | product function space, 119 |
| $\mathcal{F}_{4}\left(X_{0}, X_{1}, Y_{0}, Y_{1}\right)$ | auxiliary function space, 103 |
| $\mathcal{F}_{4}^{\prime}\left(X_{0}, X_{1}, Y_{0}, Y_{1}\right)$ | auxiliary function space, 104 |
| Fix $(f)$ | fixed point set of $f, 5$ |
| $\mathcal{G}(U)$ | image of $\widetilde{\operatorname{Ham}}(U) \rightarrow \widetilde{\operatorname{Ham}}(M), 55$ |


| $\mathrm{GW}_{j}$ | Gromov-Witten invariants, 173 |
| :---: | :---: |
| $\mathbb{H}$ | hyperbolic half-plane, 37 |
| $\mathcal{H}$ | space of normalized Hamiltonians, 6 as the Lie algebra of Ham, 9 |
| $\operatorname{Ham}(M, \omega)$ | group of Hamiltonian diffeomorphisms, 6 as a Lie group, 9 |
| $\widetilde{\operatorname{Ham}}(M, \omega)$ | universal cover of $\operatorname{Ham}(M, \omega), 13$ |
| $\mathrm{HF}_{*}(F)$ | Floer homology of $F, 170$ |
| $\mathrm{HF}_{*}(M)$ | Floer homology of $M, 170$ |
| $\mathrm{HM}_{*}(F, \rho)$ | Morse homology of a generic pair, 154 |
| $\mathrm{HM}_{*}(M)$ | Morse homology of M, 155 |
| ind | Morse index, 151 |
| $I_{\omega}$ | linear isomorphism $E \rightarrow E^{*}$ induced by $\omega, 2$ bundle isomorphism $T M \rightarrow T^{*} M$ induced by $\omega, 3$ |
| $\mathcal{K}$ | field of Laurent series, 159 |
| $\kappa_{\text {cl }}$ | classical measure of noncommutativity, 128 |
| $\kappa_{\text {q }}$ | quantum measure of noncommutativity, 130 |
| $\mathcal{L}(H)$ | space of Hermitian operators on a Hilbert space, 68 |
| $\Lambda\left(\mathbb{R}^{2 n}\right)$ or $\Lambda$ | Lagrangian Grassmannian, 40 |
| $\Lambda M$ | space of contractible loops in $M, 52$ |
| $\mathcal{L}_{X}$ | Lie derivative, 4 |
| $\mathcal{M}\left(\tilde{z}_{-}, \tilde{z}_{+}\right)$ | space of Floer connecting trajectories, 164 |
| $\overline{\mathcal{M}}\left(\tilde{z}_{-}, \tilde{z}_{+}\right)$ | moduli space of unparametrized Floer trajectories, 165 |
| $\widetilde{\mathcal{M}}\left(\tilde{z}_{-}, \tilde{z}_{+}\right)$ | compactified moduli space, 169 |
| $\operatorname{Maslov}(\gamma)$ | Maslov index of a loop in $\operatorname{Sp}(2 n), 44$ |
| $\mathcal{M}_{J}(\alpha, \beta, \gamma ; j)$ | space of holomorphic spheres passing through 3 cycles, 173 |
| $\mu$ | Maslov quasi-morphism on $\widetilde{\mathrm{Sp}}(2 n), 43$ |
|  | Maslov class of a Lagrangian submanifold, 95 |
| $\mu_{L}$ | (non-homogeneous) Maslov quasi-morphism, 43 |
| $\mathcal{N}(A)$ | inherent noise of a POVM, 140 |
| $\mathcal{N}(A, \xi)$ | inherent noise of a POVM at a pure state, 139 |
| $N_{J}$ | Nijenhuis tensor, 15 |
| $\\|F\\|$ | uniform norm, 11 |
| $\nu_{c}$ | spectral pseudo-norm, 53 |
| $([\omega], A)$ | symplectic area of $A \in \pi_{2}(M), 161$ |
| [ $\omega$ ] | symplectic area class of a Lagrangian, 94 symplectic area functional $\pi_{2}(M) \rightarrow \mathbb{R}, 161$ |
| $\omega^{*}$ | symplectic form on the dual space, 30 |
| $\omega_{0}$ | standard symplectic form on $\mathbb{R}^{2 n}, 3$ |
| $\Omega_{N}$ | set of integers $1, \ldots, N, 131$ |
| ${ }_{\sim}^{\mathcal{P}}(F)$ | set of contractible periodic orbits, 164 |
| $\widetilde{\mathcal{P}}(F)$ | all lifts of contractible periodic orbits, 169 |
| $\mathrm{pb}_{4}$ | Poisson bracket invariant of a quadruple, 103 |
| $\mathrm{pb}(\mathcal{U})$ | Poisson bracket invariant of a cover, 128 |
| $\Phi_{\text {PSS }}$ | PSS map, 178 |
| $\bar{\pi}_{2}(M)$ | $\pi_{2}(M) / \operatorname{ker}[\omega], 162$ |
| $\Pi(F, G)$ | measure of noncommutativity, 59 |


| $\operatorname{PSL}(2, \mathbb{R})$ | isometry group of the hyperbolic half-plane, 37 |
| :---: | :---: |
| QH( $M$ ) | quantum homology, 173 |
| $Q_{N}$ | unit cube $[0,1]^{N}, 138$ |
| $Q(\Omega)$ | measurable functions $\Omega \rightarrow[0,1], 138$ |
| $r(F)$ | growth of one-parameter subgroup in Hofer's norm, 97 |
| $r_{F, G}$ | profile function, 119 |
| $r_{L}$ | Robbin-Salamon index, 41 |
| rot | Poincare rotation number, 39 |
| $S$ | positive generator of $\bar{\pi}_{2}(\mathrm{M}), 162$ |
| sgrad | symplectic gradient, 3 |
| $\Sigma_{L}$ | Maslov-Arnold cycle, 41 |
| sign | signature of a quadratic form, 41 |
| $\mathfrak{s p}(2 n)$ | Lie algebra of $\operatorname{Sp}(2 n)$, 44 |
| Sp( $2 n$ ) | symplectic linear group, 30 |
| $\operatorname{Sp}(E, \omega)$ | group of linear symplectomorphisms, 30 |
| spec | action spectrum, 52 |
| $\operatorname{Symp}(M, \omega)$ | group of compactly supported symplectomorphisms, 1 |
| $\mathbb{V}(\hat{A}, \xi)$ | variance of a von Neumann observable at a pure state, 135 |
| $w(U)$ | spectral width, 57 |
| $W^{\text {s }}(x)$ | stable manifolds of a critical point, 151 |
| $W^{\mathrm{u}}(x)$ | unstable manifolds of a critical point, 151 |
| $[x, y]$ | commutator of group elements, 23 |

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This is a book on symplectic topology, a rapidly developing field of mathematics which originated as a geometric tool for problems of classical mechanics. Since the 1980s, powerful methods such as Gromov's pseudo-holomorphic curves and Morse-Floer theory on loop spaces gave rise to the discovery of unexpected symplectic phenomena. The present book focuses on function spaces associated with a symplectic manifold. A number of recent advances show that these spaces exhibit intriguing properties and structures, giving rise to an alternative intuition and new tools in symplectic topology. The book provides an essentially self-contained introduction into these developments along with applications to symplectic topology, algebra and geometry of symplectomorphism groups, Hamiltonian dynamics and quantum mechanics. It will appeal to researchers and students from the graduate level onwards.
I like the spirit of this book. It formulates concepts clearly and explains the relationship between them. The subject matter is important and interesting.
—Dusa McDuff, Barnard College, Columbia University
This is a very important book, coming at the right moment. The book is a remarkable mix of introductory chapters and research topics at the very forefront of actual research. It is full of cross fertilizations of different theories, and will be useful to Ph.D. students and researchers in symplectic geometry as well as to many researchers in other fields (geometric group theory, functional analysis, mathematical quantum mechanics). It is also perfectly suited for a Ph.D.-students seminar.
-Felix Schlenk, Université de Neuchâtel


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