## CHANGING SURFACE TO VOLUME INTEGRALS.

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The note of Dr. Poor on "Transformation theorems in the theory of the linear vector function" in this Bulletin, January, 1916, page 174, raises the question: Why not make the work short by using other methods?

The equation* $\int d S()=-\int d \tau \nabla()$ is an obvious identity because

$$
\iint \boldsymbol{i} d y d z(\quad)=-\iiint \boldsymbol{i} d y d z d x \frac{\partial}{\partial x}(\quad)
$$

is merely a partial integration.
If $\Phi$ be a linear vector function,

$$
\nabla(\Phi \cdot \boldsymbol{u})=\nabla_{\Phi}(\Phi \cdot \mathbf{u})+\nabla_{\mathbf{u}}(\Phi \cdot \boldsymbol{u})=-\nabla_{M}(\Phi \cdot \mathbf{u})+\nabla \boldsymbol{u} \cdot \Phi_{C}
$$

where the subscripts $\Phi$ and $\boldsymbol{u}$ mean that the differentiation affects only the function indicated and the subscript $M$ means that the differentiation is with respect to the point $M$ of which $\boldsymbol{u}$ is independent (other differentiations are with respect to $P$ ). Hence, integrating with no sign, with dot, and with cross,

$$
\begin{array}{rlr}
\int d \boldsymbol{S} \Phi \cdot \boldsymbol{u} & =\int d \tau \nabla_{M}(\Phi \cdot \boldsymbol{u})-\int d \tau \nabla \boldsymbol{u} \cdot \Phi_{C}, & \text { Theorem 3, } \\
\int d \boldsymbol{S} \cdot \Phi \cdot \boldsymbol{u} & =\int d \tau \nabla_{M} \cdot(\Phi \cdot \boldsymbol{u})-\int d \tau \nabla \boldsymbol{u}: \Phi, \quad \text { Theorem } 2, \\
\int d \boldsymbol{S} \times \Phi \cdot \boldsymbol{u} & =\int d \tau \nabla_{M} \times(\Phi \cdot \boldsymbol{u})-\int d \tau\left(\nabla u \cdot \Phi_{C}\right)_{\times}
\end{array}
$$

Theorem 1.
Next if $\Phi \cdot d \Psi=d \Psi \cdot \Phi$, then $d(\Phi \cdot \Psi)=d \Phi \cdot \Psi+d \Psi \cdot \Phi$ and $\nabla(\Phi \cdot \Psi)=\nabla \Phi \cdot \Psi+\nabla \Psi \cdot \Phi$. Hence on integrating, we have

$$
\int d \boldsymbol{S} \Phi \cdot \Psi=-\int d \tau \nabla \Phi \cdot \Psi-\int d \tau \nabla \Psi \cdot \Phi, \quad \text { not given }
$$

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$$
\begin{aligned}
& \int d S \cdot \Phi \cdot \Psi=-\int d \tau \nabla \cdot \Phi \cdot \Psi-\int d \tau \nabla \cdot \Psi \cdot \Phi \\
& \int d S \times \Phi \cdot \Psi=-\int d \tau \nabla \times \Phi \cdot \Psi-\int d \tau \nabla \times \Psi \cdot \Phi, \\
& \text { Theorem 6, } \\
& \text { Theorem } 7
\end{aligned}
$$
\]

Finally we may write the identities

$$
\begin{aligned}
\nabla \cdot(\nabla \Phi \cdot \mathbf{u}) & =\nabla \cdot \nabla \Phi \cdot \boldsymbol{u}+\nabla_{\mathbf{u}} \cdot \nabla \Phi \cdot \boldsymbol{u}, \\
\nabla \cdot \nabla_{\mathbf{u}}(\Phi \cdot \boldsymbol{u}) & =\nabla_{\Phi} \cdot \nabla_{\mathbf{u}} \Phi \cdot \mathbf{u}+\Phi \cdot(\nabla \cdot \nabla \boldsymbol{u}) .
\end{aligned}
$$

The terms $\nabla_{\mathbf{u}} \cdot \nabla \Phi \cdot \boldsymbol{u}$ and $\nabla_{\Phi} \cdot \nabla_{\mathbf{u}} \Phi \cdot \boldsymbol{u}$ are the same, since the order in a scalar product is immaterial. Hence, by subtraction and integration,

$$
\begin{aligned}
& \nabla \cdot(\nabla \Phi \cdot \boldsymbol{u})-\nabla \cdot\left(\nabla \boldsymbol{u} \cdot \Phi_{C}\right)= \nabla \cdot \nabla \Phi \cdot \boldsymbol{u}-\Phi \cdot(\nabla \cdot \nabla \boldsymbol{u}) \\
&-\int d \boldsymbol{S} \cdot\left(\nabla \Phi \cdot \boldsymbol{u}-\nabla \boldsymbol{u} \cdot \Phi_{C}\right)=\int d \tau \Delta_{M}(\Phi \cdot \boldsymbol{u}) \\
& \quad-\int d \tau \Phi \cdot \Delta \boldsymbol{u}, \text { Theorem } 5 .
\end{aligned}
$$

In these theorems I have used the notation of Gibbs and the results are therefore in some cases conjugates of Poor's; for he uses $d=(\quad) \nabla \cdot d \boldsymbol{r}$ instead of $d=d \boldsymbol{r} \cdot \nabla(\quad)$. Tait and McAulay have employed a notation with subscripts so that the symbol $\nabla$ and its operand may occur in any positions; for instance $\boldsymbol{u}_{1} \nabla_{1}$ means the conjugate of $\nabla \boldsymbol{u}$. Poor's Theorem 4, without the integral sign, is in this notation $\Phi_{2} \cdot \boldsymbol{u}_{1} \nabla_{1} \nabla_{2} \cdot \boldsymbol{x}$ $=\left(\Phi_{2} \nabla_{2} \cdot \boldsymbol{x}\right) \cdot \boldsymbol{u}_{1} \nabla_{1}$ and represents an identity just as $(a b) c=a(b c)$ represents an identity in ordinary algebra.

The formal side of multiple algebra has been highly developed by Grassmann, Hamilton, Tait, Gibbs, McAulay, Clebsch, Aronhold, Study, and Shaw, not to mention a host of others. Why deny ourselves the advantages of the formal methods? The use of words like grad, div, rot is hampering: we no longer write Cubus $\bar{m}$ Census $\bar{p} 16$ rebus æquatur 40 for $x^{3}-8 x^{2}+16=40$.

Massachusetts Institute of Technology, January 18, 1916.


[^0]:    * Reference may be made to my review, "The unification of vectorial notations," this Bulletin, vol. 16, May, 1910, p. 428, where I use $d \mathrm{~S}$ as an exterior normal instead of an interior normal as here.

