and curvature types of triply infinite families of curves. The reresults are summarized in the following table; the two types compared are named in the left-hand column; their intersection is identified in the center; and the number, that is, the infinitude, of (projectively different) common families is given at the right.
$\begin{array}{lcc}\text { Dynamical Sectional: } & \begin{array}{c}\text { Special central fields or } \\ \text { General cones }\end{array} & \infty^{f(1)} \\ \text { Dynamical Curvature: } & \text { Any central field } \\ \text { Sectional Curvature: } & \begin{array}{c}\text { General cones and } \\ \text { Quadric surfaces }\end{array} & \infty^{f(2)} \\ & \infty^{f(1)+2}\end{array}$
The 2 in the exponent of $\infty$ refers of course to two arbitrary constants, while (according to a notation which I proposed in this Bulletin in 1912, in a review of Riquier's treatise on partial differential equations) $f(1)$ means an arbitrary function of one independent variable, and $f(2)$ an arbitrary function of two independent variables.

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## ON NEVANLINNA'S WEAK SUMMATION METHOD $\dagger$

## BY A. F. MOURSUND

1. Introduction. Our principal object in this note is to discuss the function

$$
\begin{array}{r}
\left.\rho_{n}(\beta)=\frac{2}{\pi} \int_{0}^{\pi / 2} \right\rvert\, \int_{0}^{1} \beta(\log C)^{\beta}(1-t)^{-1}(\log C /(1-t))^{-\beta-1} \\
 \tag{1}\\
\left.\times \frac{\sin (2 n t+1) s}{\sin s} d t \right\rvert\, d s
\end{array}
$$

which, for $\beta>0$ and the "dummy" constant $C \geqq e^{\beta+1}$, plays a role in the theory of summation of Fourier series by Nevanlinna’s weak method $\ddagger$ analogous to the role the Lebesgue constants

[^0]\[

$$
\begin{equation*}
\rho_{n}=\frac{2}{\pi} \int_{0}^{\pi / 2}\left|\frac{\sin (2 n+1) s}{\sin s}\right| d s, \quad(n=0,1, \cdots) \tag{2}
\end{equation*}
$$

\]

play in the theory of convergence of such series. $\dagger$ Nevanlinna's weak method is essentially the same as the Bosanquet-Linfoot method of zero order. $\ddagger$

Our principal results concerning the function $\rho_{n}(\beta)$ are given by the following theorems.

Theorem 1. For each $n \geqq 0$, the function $\rho_{n}(\beta) \rightarrow \rho_{n}$ as $\beta \rightarrow 0$.
Theorem 2. When $\beta>1$, the function $\rho_{n}(\beta)$ is uniformly bounded with respect to $n$ for all $n \geqq 0$.

Theorem 3. For $0<\beta<1$,

$$
\rho_{n}(\beta)=\frac{4}{\pi^{2}} \frac{(\operatorname{llog} C)^{\beta}}{1-\beta}(\log n)^{1-\beta}+O(1),
$$

and

$$
\rho_{n}(1)=\frac{4}{\pi^{2}} \log C \log \log n+O(1) .
$$

2. Nevanlinna's Weak Summation Method. Applied to the Fourier series generated by a Lebesgue integrable function $f(x)$, Nevanlinna's weak method consists in forming from the well known expression for the sum of $n$ terms of the series the $N_{\beta}$ transform

$$
\begin{equation*}
N_{\beta} S_{n}(x)=\frac{1}{2 \pi} \int_{0}^{1} N_{\beta}(t) d t \int_{-\pi}^{\pi} f(s) \frac{\sin (2 n t+1)(x-s) / 2}{\sin (x-s) / 2} d s \tag{3}
\end{equation*}
$$

where

[^1]$$
N_{\beta}(t)=\beta(\log C)^{\beta}(1-t)^{-1}(\log C /(1-t))^{-\beta-1}
$$
with $\beta>0$ and $C \geqq e^{\beta+1}$, and in taking the limit
\[

$$
\begin{equation*}
\lim _{n \rightarrow \infty} N_{\beta} S_{n}(x) \tag{4}
\end{equation*}
$$

\]

Bosanquet and Linfoot have given an example, which will serve equally well for Nevanlinna's weak method, of a continuous function $f(x)$ whose Fourier series diverges at $x=0$ when summed by their zero order method with $0<\beta \leqq 1 . \dagger$ For $\beta>1$, Nevanlinna's method, and consequently the Bosanquet-Linfoot zero order method, is Lebesgue effective. $\ddagger$
3. The Function $\rho_{n}(\beta)$. Upon changing the order of integration in (3), we see that, for the values of $\beta$ for which the function

$$
\begin{equation*}
\rho_{n}(\beta, x)=\frac{1}{2 \pi} \int_{-\pi}^{\pi}\left|\int_{0}^{1} N_{\beta}(t) \frac{\sin (2 n t+1)(x-s) / 2}{\sin (x-s) / 2} d t\right| d s \tag{5}
\end{equation*}
$$

becomes infinite with $n$, functions continuous on $(-\pi, \pi)$ can be constructed for which (i) the limit (4) does not exist at the point $x$, (ii) the limit (4) exists at $x$ but does not exist uniformly in any neighborhood of that point. § Setting $x=0$ in (5) we obtain, after slight simplification, the function $\rho_{n}(\beta)$ defined in the introduction.
4. Lemmas. In the statements and proofs of our lemmas and theorems, $\beta>0, C \geqq e^{\beta+1}$, and $n \geqq 0$ unless otherwise stated. Proofs which the reader can readily supply are merely indicated or omitted entirely.

Lemma 1. For $2 n s \geqq \pi$

$$
\left.\begin{array}{rl}
\mid \int_{0}^{1} N_{\beta}(t) & \sin \\
\cos
\end{array}\right\} 2 n s t d t\left|\left\lvert\, \begin{array}{l}
1-\pi /(2 n s) \\
\\
\end{array}\right.\right.
$$

[^2]Lemma 2. $\rho_{n}(\beta)=\rho_{n}{ }^{*}(\beta)+O(1)$, where

$$
\begin{equation*}
\rho_{n}^{*}(\beta)=\frac{2}{\pi} \int_{0}^{1 / 2} \frac{1}{s}\left|\int_{0}^{1} N_{\beta}(t) \sin 2 n s t d t\right| d s \tag{6}
\end{equation*}
$$

and the $O(1)$ terms are uniformly bounded with respect to $\beta, C$, and $n . \dagger$

Lemma 3. For $v>0$,

$$
\int_{0}^{\infty}\left(\frac{\log C}{v}+1+t\right)^{-\beta-1} \cos e^{-t v} d t=\frac{1}{\beta}+O\left(\frac{1}{v}\right)
$$

Proof. We write

$$
\begin{aligned}
\frac{1}{\beta}-\int_{0}^{\infty}= & \int_{0}^{\infty}\left[(1+t)^{-\beta-1}-\left(\frac{\log C}{v}+1+t\right)^{-\beta-1} \cos e^{-t v} d t\right] \\
\leqq & \int_{0}^{\infty}\left\{\left[(1+t)^{-\beta-1}-\left(\frac{\log C}{v}+1+t\right)^{-\beta-1}\right]\right. \\
& \left.+\left(\frac{\log C}{v}+1+t\right)^{-\beta-1} e^{-2 t v} / 2\right\} d t \\
\leqq & (\beta+1) \int_{0}^{\infty} d t \int_{0}^{\log C / v}(1+t+u)^{-\beta-2} d u+\int_{0}^{\infty} e^{-2 t v} / 2 d t \\
\leqq & (\beta+1) \frac{\log C}{v} \int_{0}^{\infty}(1+t)^{-\beta-2} d t+\frac{1}{4 v}=\frac{\log C}{v}+\frac{1}{4 v}
\end{aligned}
$$

Lemma 4. For r sufficiently large,
$\int_{0}^{1} N_{\beta}(1-t) \cos r t d t=\beta(\log C)^{\beta}\left\{(\log r)^{-\beta} / \beta+O(\log r)^{-\beta-1}\right\}>0$,
$\int_{0}^{1} N_{\beta}(1-t) \sin r t d t=O\left[(\log r)^{-\beta-1}\right]$.
Proof. The lemma, except for the term $O\left[(\log r)^{-\beta-1}\right]$ in the first part, follows immediately from a theorem concerning Fourier coefficients due to U. S. Haslam-Jones. $\ddagger$ Upon inspection of

[^3]the proof given by Haslam-Jones, the reader will see how, with the aid of Lemma 3, our result can be obtained.

Lemma 5.

$$
|\sin r|=\frac{2}{\pi}-\frac{4}{\pi} \sum_{m=1}^{\infty} \frac{\cos 2 m r}{4 m^{2}-1} \dagger
$$

Lemma 6. For $M \geqq e$ and $n \geqq M$,

$$
\sum_{m=1}^{\infty} \frac{1}{4 m^{2}-1} \int_{M}^{n} \frac{1}{r}(\log r)^{-\beta} \cos 2 m r d r \leqq \frac{1}{2 M}
$$

Proof. It can be shown by means of the second mean value theorem that each of the integrals $\leqq 1 / M$. The lemma follows, because

$$
\sum \frac{1}{4 m^{2}-1}=\frac{1}{2}
$$

5. Proof of Theorems. In this paragraph we prove our theorems.

Proof of Theorem 1. Integrating by parts, we have

$$
\begin{aligned}
\rho_{n}(\beta)=\frac{2}{\pi} \int_{0}^{\pi / 2} \left\lvert\, 1+(\log C)^{\beta} \frac{2 n s}{\sin s} \int_{0}^{1}\right. & (\log C /(1-t))^{-\beta} \\
& \times \cos (2 n t+1) s d t \mid d s
\end{aligned}
$$

and the theorem follows by elementary theorems concerning limits.

Proof of Theorem 2. When $n \leqq \pi$, we have at once $\rho_{n}{ }^{*}(\beta) \leqq 2$; and when $n>\pi$, we have for $\beta>1$, using Lemma 1,

$$
\rho_{n}^{*}(\beta) \leqq O(1)+\frac{2}{\pi}(\log C)^{\beta} \int_{\pi /(2 n)}^{1 / 2} \frac{1}{s}(\log 2 C n s / \pi)^{-\beta} d s=O(1) .
$$

The theorem follows by Lemma 2.
Proof of Theorem 3. By Lemma 2,

$$
\rho_{n}(\beta)=\rho_{n}^{*}(\beta)+O(1)
$$

[^4]Using Lemmas 4, 5, and 6, we have, for a fixed sufficiently large $M$ and $n>M$,

$$
\begin{aligned}
\rho_{n}^{*}(\beta)= & \frac{2}{\pi}\left[\int_{0}^{M /(2 n)}+\int_{M /(2 n)}^{1 / 2}\right] \frac{1}{s}\left|\int_{0}^{1} N_{\beta}(t) \sin 2 n s t d t\right| d s \\
= & O(1)+\frac{2}{\pi} \int_{M}^{n} \frac{1}{r}\left|\int_{0}^{1} N_{\beta}(t) \sin r t d t\right| d r \\
= & \left.\frac{2}{\pi} \int_{M}^{n} \frac{1}{r} \right\rvert\, \sin r \int_{0}^{1} N_{\beta}(1-t) \cos r t d t \\
& -\cos r \int_{0}^{1} N_{\beta}(1-t) \sin r t d t \mid d r+O(1) \\
= & \frac{2}{\pi} \int_{M}^{n} \frac{|\sin r|}{r} d r \int_{0}^{1} N_{\beta}(1-t) \cos r t d t+O(1) \\
= & \frac{2 \beta(\log C)^{\beta}}{\pi}\left\{\frac{2}{\pi} \int_{M}^{n} \frac{1}{r} \cdot \frac{(\log r)^{-\beta}}{\beta} d r\right. \\
& -\frac{4}{\pi} \sum_{m=1}^{\infty} \frac{1}{4 m^{2}-1} \int_{M}^{n} \frac{(\log r)^{-\beta}}{\beta} \cos 2 m r d r \\
& \left.+\int_{M}^{n} \frac{|\sin r|}{r} O\left[(\log r)^{-\beta-1}\right] d r\right\}+O(1) \\
= & \frac{4}{\pi^{2}}(\log C)^{\beta} \int_{M}^{n} \frac{1}{r} \cdot(\log r)^{-\beta} d r+O(1) .
\end{aligned}
$$

The lemma follows when we carry out the integration.
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[^0]:    $\dagger$ Presented to the Society, June 20, 1934.
    $\ddagger$ F. Nevanlinna, Über die Summation der Fourier'schen Reihen und Integrale, Översikt av Finska Vetenskaps-Societetens Förhandlingar, vol. 64 (1921-22), A, No. 3, 14 pp. A. F. Moursund, On the Nevanlinna and BosanquetLinfoot summation methods, Annals of Mathematics, (2), to appear.

[^1]:    $\dagger$ L. Fejér, Lebesguesche Konstanten und Divergente Fourierreihen, Journal für Mathematik, vol. 138 (1910), pp. 22-53. Fejér shows in that paper that $\rho_{n} \sim\left(4 / \pi^{2}\right) \log n+O(1)$. T. H. Gronwall, Über des Lebesgueschen Konstanten bei den Fourierschen Reihen, Mathematische Annalen, vol. 72 (1912), pp. 244-261. G. Szegö, Über die Lebesgueschen Konstanten bei den Fourierschen Reihen, Mathematische Zeitschrift, vol. 9 (1921), pp. 163-166.
    $\ddagger$ L. S. Bosanquet and E. H. Linfoot, On the zero order summability of Fourrer series, Journal of the London Mathematical Society, vol. 6 (1931), pp. 117-126. L. S. Bosanquet and E. H. Linfoot, Generalized means and the summability of Fourier series, Quarterly Journal of Mathematics (Oxford Series), vol. 2 (1931), pp. 207-229. Moursund, loc. cit.

[^2]:    $\dagger$ Loc. cit., first paper.
    $\ddagger$ A. F. Moursund, On a method of summation of Fourier series, Annals of Mathematics, (2), vol. 33 (1932), pp. 773-784.
    § E. W. Hobson, The Theory of Functions of a Real Variable, 2d ed., vol. 2 Chapter 7.
    || See Moursund, second loc. cit., pp. 779-780. Lemma 5.1 holds for $N_{\beta}(t)$ is non-negative and monotone increasing on $(0,1)$.

[^3]:    $\dagger$ It can be shown by using Lemma 1 that the $O(1)$ terms are $o(1)$ as $n \rightarrow \infty$.
    $\ddagger$ U. S. Haslam-Jones, A note on the Fourier coefficients of unbounded functions, Journal of the London Mathematical Society, vol. 2 (1927), pp. 151154 (Theorem 2).

[^4]:    $\dagger$ Szegö, loc. cit., uses this Fourier series expansion in obtaining his formula for $\rho_{n}$.

