# A NOTE ON MATRIX SUMMABILITY OF A CLASS OF FOURIER SERIES 

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1. Recently several papers by Rajagopal [7] ,Varshney [11] and others have been written, on Nörlund summability of Fourier series, in order to unify some of the classical results. Though lower semitriangular matrix ( $\Lambda$ ) summability method has been known for quite some time no attempt has yet been made to apply it to Fourier series. The object here is to determine a necessary and sufficient condition for ( $\Lambda$ ) summability of Fourier series and to include a wider class of known results.

A Fourier series, of a Lebesgue-integrable function, is said to be summable at a point by triangular matrix method $(\Lambda)$, defined by Hardy [1], if $\Lambda_{n, k}=0$ for $k>n, \sum \Lambda_{n, k} \rightarrow 1$ as $n \rightarrow \infty$ and $\sum_{k=0}^{n}\left|\Lambda_{n, k}\right| \leqq M$, where $M$ is a constant, and the point is in a certain subset of the Lebesgue set.

The following main theorem has been proved here.
Theorem. Let a sequence $\left\{\Lambda_{n, k}\right\}$ be defined in terms of

$$
\begin{align*}
& \Lambda_{n}(u), \text { monotonic decreasing and strictly positive for all } u \geqq 0,  \tag{1.1}\\
& \Lambda_{n, u} \equiv \Lambda_{n}(u)
\end{align*}
$$

and if

$$
\begin{equation*}
\Phi(t) \equiv \int_{0}^{t}|\phi(u)| d u=o\left(\frac{t}{\psi(1 / t)}\right) \text { as } t \rightarrow+0 \tag{1.2}
\end{equation*}
$$

and $\psi(t)$ be positive, nondecreasing with $t$; then a necessary and sufficient condition for $(\Lambda)$ summability of Fourier series, to 0 or

$$
\begin{equation*}
t_{n} \equiv\left\{\sum_{k=0}^{n} \Lambda_{n, k} S_{k}\right\} \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{1.3}
\end{equation*}
$$

is

$$
\begin{equation*}
\int_{1}^{n} \frac{\bar{\Lambda}_{n}(u)}{u \psi(u)} d u=O(1) \tag{1.4}
\end{equation*}
$$

[^0]where $\bar{\Lambda}_{n, k} \equiv \sum_{r=1}^{k} \Lambda_{n, r} \equiv \bar{\Lambda}_{n, k}(k)$.
2. We shall need the following lemma.

Lemma. If the sequence $\left\{\Lambda_{n, k}\right\}$ is defined as in (1.1), then

$$
\begin{equation*}
\left|\frac{\sum_{k=0}^{n} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u}{\operatorname{Sin}(u / 2)}\right|<c \frac{\bar{\Lambda}_{n}(1 / u)}{u}, \tag{2.1}
\end{equation*}
$$

where $c$ is a constant, not the same at each occurrence.
Proof. If we choose $m=$ integral part of $1 / u$ and suppose that $1 / n \leqq u \leqq \delta$, we get $m \operatorname{Sin} \frac{1}{2} u>m u / \pi$. Now for $u>0$ and $m \leqq n$ we have

$$
\begin{aligned}
& \left\lvert\, \begin{array}{|l}
\left\lvert\, \sum_{k=0}^{n} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u\right. \\
\\
\quad<\frac{1}{\operatorname{Sin}(u / 2)}\left[\left|\sum_{k=0}^{m} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u\right|+\left|\sum_{m}^{n} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u\right|\right] \\
\quad<\frac{1}{\operatorname{Sin} \frac{u}{2}}\left[\sum_{k=0}^{m} \Lambda_{n, k}\left|\operatorname{Sin}\left(k+\frac{1}{2}\right) u\right|+\bar{\Lambda}_{n, m} \operatorname{Max}_{m \leqq k \leqq n} \sum_{m}^{n} \operatorname{Sin}\left(k+\frac{1}{2}\right) u\right] \\
\quad=\frac{\bar{\Lambda}_{n, m}}{\operatorname{Sin} \frac{u}{2}}+\Lambda_{n, m} \operatorname{Max}_{m \leqq k \leqq n} \frac{\operatorname{Cos}\left(k+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{u}{2}} \\
\quad<\frac{\bar{\Lambda}_{n}(m)}{\operatorname{Sin} \frac{1}{2} u}+\frac{c \bar{\Lambda}_{n}(m)}{m\left(\operatorname{Sin} \frac{1}{2} u\right)^{2}} \\
\quad=\frac{\bar{\Lambda}_{n}[1 / u]}{\operatorname{Sin} \frac{1}{2} u}+\frac{c \bar{\Lambda}_{n}[1 / u]}{m\left(\operatorname{Sin} \frac{1}{2} u\right)^{2}} \\
\quad<\frac{\bar{\Lambda}_{n}[1 / u]}{\operatorname{Sin} \frac{1}{2} u}+\frac{c \bar{\Lambda}_{n}[1 / u]}{\operatorname{Sin}\left(\frac{1}{2} u\right)} \\
\quad<\frac{c \bar{\Lambda}_{n}(1 / u)}{u},
\end{array}\right.
\end{aligned}
$$

which proves the lemma.
3. Proof of the theorem. To prove the sufficiency part, first, we see [2] that the $n$th partial sum of Fourier series is given by

$$
\begin{aligned}
S_{n} & =\frac{1}{\pi} \int_{0}^{\pi} \phi(u) \frac{\operatorname{Sin}\left(n+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{1}{2} u} d u \\
& =\frac{1}{\pi} \int_{0}^{\delta} \phi(u) \frac{\operatorname{Sin}\left(n+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{1}{2} u} d u+o(1) .
\end{aligned}
$$

Using (1.3), and the last expression, we get

$$
\begin{align*}
t_{n} & =\frac{1}{\pi} \int_{0}^{\delta} \phi(u) \sum_{k=0}^{n} \Lambda_{n, k} \frac{\operatorname{Sin}\left(k+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{1}{2} u} d u+o(1) \\
& =\frac{1}{\pi}\left[\int_{0}^{1 / n}+\int_{1 / n}^{\delta}\right] \phi(u) \sum_{k=0}^{n} \Lambda_{n, k} \frac{\operatorname{Sin}\left(k+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{1}{2} u} d u+o(1) \\
& =I_{1}+I_{2}+o(1), \text { say } \tag{3.1}
\end{align*}
$$

by virtue of (1.1).
Considering $I_{1}$, we get

$$
\begin{align*}
I_{1} & =\frac{1}{\pi} \int_{0}^{1 / n} \frac{\phi(u)}{\operatorname{Sin} \frac{1}{2} u} \sum_{k=0}^{n} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u d u \\
& =O(n) \int_{0}^{1 / n}|\phi(u)| d u  \tag{3.2}\\
& =o\left(\frac{1}{\Psi(n)}\right), \quad \text { by (1.2). }
\end{align*}
$$

Next, by the above lemma

$$
\begin{aligned}
I_{2}= & O(1) \int_{1 / n}^{\delta}|\phi(u)|\left|\frac{\sum_{k=0}^{n} \Lambda_{n, k} \operatorname{Sin}\left(k+\frac{1}{2}\right) u}{\operatorname{Sin} \frac{1}{2} u}\right| d u \\
= & O(1) \int_{1 / n}^{\delta}|\phi(u)| \frac{\bar{\Lambda}_{n}(1 / u)}{u} d u \\
= & O(1)\left\{\left[\Phi(u) \frac{\bar{\Lambda}_{n}(1 / u)}{u}\right]_{1 / n}^{\delta}+\int_{1 / n}^{\delta} \Phi(u) d\left[\frac{\bar{\Lambda}_{n}(1 / u)}{u \Psi(1 / u)} \cdot \Psi\left(\frac{1}{u}\right)\right]\right\} \\
= & o\left(\frac{\bar{\Lambda}_{n}(n)}{\Psi(n)}\right)+o(1) \\
& +o(1) \int_{1 / n}^{\delta} \frac{\bar{\Lambda}_{n}(1 / u)}{\{\Psi(1 / u)\}^{2}} d \Psi\left(\frac{1}{u}\right)+o(1) \int_{1 / n}^{\delta} u d\left\{\frac{\bar{\Lambda}_{n}(1 / u)}{u \Psi(1 / u)}\right\} \\
= & o(1)+o(1) \int_{1 / n}^{\delta} \frac{d \Psi(1 / u)}{\{\Psi(1 / u)\}^{2}}+o(1) \int_{1 / n}^{\delta} \frac{\Lambda_{n}(1 / u)}{u \Psi(1 / u)} d u \\
= & o(1), \quad \text { by virtue of }(1.4) .
\end{aligned}
$$

Now the first part of the proof is complete by virtue of (3.1), (3.2) and (3.3), when $n \rightarrow \infty$.

To prove the necessary part, a look at the proof of sufficiency part shows that, it is sufficient to show here that

$$
\begin{equation*}
\int_{1 / n}^{\delta} \frac{u}{\Psi(1 / n)} d\left\{\frac{\bar{\Lambda}_{n}(1 / u)}{u}\right\}=O(1) \tag{3.4}
\end{equation*}
$$

Considering the left-hand side and proceeding as above, we have

$$
\begin{aligned}
\int_{1 / n}^{\delta} \frac{u}{\Psi(1 / u)} & d\left\{\frac{\bar{\Lambda}_{n}(1 / u)}{u}\right\} \\
& =O(1) \int_{1 / n}^{\delta} u d\left\{\frac{\bar{\Lambda}(1 / u)}{u \Psi(1 / u)}\right\}+O(1) \int_{1 / n}^{\delta} \frac{\bar{\Lambda}_{n}(1 / u)}{\{\Psi(1 / u)\}^{2}} d \Psi(1 / u) \\
& =O\left(\frac{1}{\Psi(n)}\right)+O\left(\bar{\Lambda}_{n}(n)\right)+O(1) \int_{1 / n}^{\delta} \frac{\bar{\Lambda}_{n}(1 / u)}{u \Psi(1 / u)} d u
\end{aligned}
$$

which is bounded by virtue of (1.4) and this completes the proof.
4. The results due to Hille and Tamarkin [4], Iyengar [6] and Siddiqi [10] on harmonic summability form the particular cases of the theorem for $\Lambda_{n, k} \equiv(1 / \log (n+1)) \cdot(1 /(k+1))$ and $\Psi(1 / t) \equiv \log (1 / t)$. The case on ( $H, p$ )-summability due to Sahney [8] can be obtained by considering the case

$$
\Psi(u)=\prod_{q=0}^{p-1}(\log )^{q+1}(u)
$$

and

$$
\Lambda_{n, k}=\frac{1}{(\log )^{p}(n+1) \prod_{q=0}^{p-1}(\log )^{q}(k+1)} .
$$

Results on Cesàro summability, due to Fejer, Lebesgue and Hardy [2] can be obtained if we choose $\Psi(u) \equiv 1$ and

$$
p_{n}=\frac{\sqrt{n+\alpha}}{\sqrt{(n+1)} \sqrt{\alpha}}
$$

for $0<\alpha<1$ and $p_{n}$ is as defined below.
Lastly the different results on Nörlund summability, for $\Lambda_{n, k}$ $=p_{n-k} / P_{n}$, can be obtained by considering the cases $\Psi(u) \equiv 1$ and $\Psi(u) \equiv \log u$ which are due to Hille-Tamarkin [4], Rajagopal [7] and Varshney [11], respectively.

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