# A CHARACTERIZATION OF WHITNEY FORMS 

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#### Abstract

We give a characterization of Whitney forms on an $n$-simplex $\sigma$ and prove that for every real valued simplicial $k$-cochain $c$ on $\sigma$, the form $W c$ is the unique differential $k$-form $\varphi$ on $\sigma$ with affine coefficients that pulls back to a constant form of degree $k$ on every $k$-face $\tau$ of $\sigma$, and satisfies $\int_{\tau} \varphi=\langle c, \tau\rangle$.


## 1. Introduction

Whitney forms have been extraordinarily useful in several areas of mathematics: algebraic topology [8, 6]; global analysis and spectral geometry [4, [3] numerical electromagnetism [1], 2]; vibrations of thin plates [7]. Their definition in Whitney's book [9, p. 140] appears somewhat mysterious. Attempts to gain a better insight into the definition have continued up to now. For example, the recent paper of Lohi and Kettunen [5] contains three different equivalent definitions. In this note we give a conceptual, easily stated characterization of Whitney forms.

We consider a triangulated differentiable manifold $M$ of $n$ dimensions with a triangulation $h: K \longrightarrow M$, cf. [9] p. 124]. We use Whitney's terminology exactly. Thus $h$ is a homeomorphism of a simplicial complex $K$ onto the manifold $M$ with the additional property that for every closed $n$-simplex $\sigma$ of $K$ there exists a coordinate system $\chi_{\sigma}$ defined in an open neighborgood $U_{\sigma}$ of the image $h(\sigma)$ so that the composition $\chi_{\sigma} \circ(h \mid \sigma)$ is an affine map of $\sigma$ into $\mathbb{R}^{n}$. One often identifies $K$ with $M$ via $h$ which usually does not lead to any confusion. We will do so here as well and regard a simplex $\sigma$ as a subset of $K, M$ or $\mathbb{R}^{n}$ without explicitly mentioning identifications given by $h$ or $\chi_{\sigma}$.

Now the Whitney form $W c$ corresponding to the cochain $c \in C^{k}(K)$ is an assignment of a smooth $k$-form $\omega_{\sigma}$ s to each closed $n$-simplex $\sigma$ that satisfies certain compatibility conditions. Namely, if $\tau$ is a common face of two top dimensional faces $\sigma_{1}$ and $\sigma_{2}$, then the pull-backs to $\tau$ of $\omega_{\sigma_{1}}$ and $\omega_{\sigma_{2}}$ coincide. Thus to describe the Whitney form $W c$ it suffices to give a description of $W c \mid \sigma=\omega_{\sigma}$ for every simplex $\sigma$ of top dimension. Note that the homeomorphism $h$ defines an affine structure on $\sigma$ and the induced affine structures on common faces of two $n$-simplexes agree. Thus the concept of an affine function on a simplex is well-defined and so is a notion of a "constant" form of degree $k$ on a $k$-simplex.

From now on we work on a fixed $n$-simplex $\sigma$. Our characterization of $W c$ is stated precisely in the Theorem below. It asserts that $W c$ restricted to $\sigma$ is the

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unique $k$-form on $\sigma$ with affine coefficients and constant pull-backs to $k$-faces whose integrals over $k$-faces $\tau$ are prescribed by the values $\langle c, \tau\rangle$ of $c$ on $\tau$.

## 2. Proof of the Theorem

A simplex $\tau=\left[p_{0}, p_{1}, \ldots, p_{k}\right]$ of $k$ dimensions is a convex hull of $k+1$ points in general position in $\mathbb{R}^{n}$. In particular, every simplex is closed. We will consider a fixed $n$-simplex $\sigma$ together with all its $k$-faces $\tau$ with $0 \leq k \leq n$. Thus a point $q \in \sigma$ is a convex linear combination

$$
\begin{gathered}
q=m_{0} p_{0}+m_{1} p_{1}+\ldots+m_{n} p_{n} \\
m_{i} \geq 0 \quad \text { for } \quad i=0,1 \ldots, n \\
m_{0}+m_{1} \ldots+m_{n}=1
\end{gathered}
$$

and the barycentric coordinate functions $\nu_{i}(q)$ are defined by

$$
\nu_{i}(q)=m_{i} .
$$

We observe that, if $q=\left(x^{1}, x^{2}, \ldots, x^{n}\right)$ the barycentric coordinates are affine functions of $x^{1}, x^{2}, \ldots, x^{n}$ i.e. are of the form $a_{1} x^{1}+a_{2} x^{2}+\ldots+a_{n} x^{n}+b$. We regard all simplices as oriented with the orientation determined by the order of vertices with the usual convention that $-\tau$ is $\tau$ with the opposite orientation and that under a permutation of vertices the orientation changes by the sign of the permutation. A cochain $c$ of degree $k$ is then defined as a formal linear combination with real coefficients of duals $\tau^{*}$ of the $k$-faces $\tau$ of $\sigma$ and we denote by $C^{k}(\sigma)=C^{k}$ the space of all such cochains. If $c=\sum_{\tau} a_{\tau} \tau^{*}$ we will write $a_{\tau}=\langle c, \tau\rangle$. Finally, we will denote by $\Lambda^{k}(\sigma)=\Lambda^{k}$ the space of all smooth exterior differential forms of degree $k$ on the simplex $\sigma$. With this notation, one defines the Whitney mapping

$$
W: C^{k} \longrightarrow \Lambda^{k}
$$

for all $k=0,1, \ldots n$, cf. [9] or [3] for a detailed discussion. We will call forms in the image of $W$ the Whitney forms. It follows immediately from the definition that the Whitney forms when expressed in terms of the coordinates of $\mathbb{R}^{n}$ have affine coefficients. We abuse the language and say that a form $\eta \in \Lambda^{k}(\tau)$ is constant if it is a constant multiple of the Euclidean volume element on $\tau$. For clarity, we emphsize that by the integral of a $k$-form over a submanifold of $k$ dimensions we always mean the integral of the pull-back of the form to the submanifold via the inclusion map. Thus, for example, in (3) below $\int_{\tau} \omega=\int_{\tau} \iota_{\tau}^{*} \omega$. After these preliminaries we state our theorem.

Theorem. Let $\sigma$ be a simplex of $n$ dimensions and c a cochain of degree $k$ on $\sigma$. $W c$ is the unique $k$-form $\omega$ on $\sigma$ satisfying the following conditions.
(1) $\omega$ has affine coefficients.
(2) The pull-back $\iota_{\tau}^{*} \omega$ is constant for every $k$-dimensional face $\tau$ of $\sigma$, where $\iota_{\tau}: \tau \hookrightarrow \sigma$ denotes the inclusion map.
(3) $\int_{\tau} \omega=\langle c, \tau\rangle$ for every $k$-face $\tau$ of $\sigma$.

Proof. We first observe that without any loss of generality we can assume that $\sigma$ is the standard simplex in $\mathbb{R}^{n}$ i.e. is given by

$$
\sigma=\left\{\left(x^{1}, x^{2}, \ldots, x^{n}\right) \in \mathbb{R}^{n} \mid x^{i} \geq 0 \quad \text { for } \quad i=1,2, \ldots n ; \quad \sum_{i=0}^{n} x^{i} \leq 1\right\}
$$

Thus $\sigma=\left[0, e_{1}, e_{2} \ldots, e_{n}\right]$ where $e_{i}$ is the point on the $i$-th coordinate axis with $x^{i}=1$. The barycentric coordinate functions restricted to $\sigma$ are then given by

$$
\begin{equation*}
\nu_{0}=1-\left(x^{1}+x^{2}+\ldots+x^{n}\right) \quad \text { and } \quad \nu_{i}=x^{i} \quad \text { for } \quad i=1,2, \ldots n . \tag{1}
\end{equation*}
$$

We first do a quick dimension count that makes the theorem plausible. The dimension of the space of $k$-forms with affine coefficients on $\sigma$ is $\binom{n}{k}(n+1)$. Requiring that $\iota_{\tau}^{*} \omega$ is constant on a $k$-simplex $\tau$ imposes $k$ conditions and the number of $k$-faces of an $n$-simplex is $\binom{n+1}{k+1}$. Thus, the dimension of the space of $k$-forms satisfying (1) and (2) above ought to be

$$
\binom{n}{k}(n+1)-\binom{n+1}{k+1} k=\binom{n+1}{k+1} .
$$

This last integer is the number of $k$-faces of $\sigma$, i.e. the dimension of the space $C^{k}(\sigma)$ of $k$-cochains.

It is instructive to consider the simplest cases $k=0$ and $k=n$ of the theorem. A 0 -cochain is a sum $c=\sum a_{i} p_{i}^{*}$ and

$$
\begin{aligned}
W c & =a_{0} \nu_{0}+a_{1} \nu_{1}+\ldots a_{n} \nu_{n} \\
& =a_{0}\left(1-\sum_{i=1}^{n} x^{i}\right)+\sum_{i=1}^{n} a_{i} x^{i} \\
& =a_{0}+\sum_{i=1}^{n}\left(a_{i}-a_{0}\right) x^{i}
\end{aligned}
$$

is the unique affine function $f$ taking prescribed values $f\left(p_{i}\right)=\int_{p_{i}} f=\left\langle c, p_{i}\right\rangle$, where the integration of a form of degree 0 over a vertex is just the evaluation.

If $k=n, \sigma$ is the only face of dimension $n$ so every cochain is a multiple of $\sigma^{*}$. For $c=\sigma^{*}$, we have

$$
\begin{aligned}
W c & =W \sigma^{*} \\
& =\left(n!\sum_{j=0}^{n}(-1)^{j} \nu_{j} d \nu_{0} \wedge \ldots \wedge \widehat{d \nu_{j}} \wedge \ldots \wedge d \nu_{n}\right) \\
& =n!d x^{1} \wedge \ldots \wedge d x^{n}
\end{aligned}
$$

where we used the explicit expressions of the barycentric coordinates (11) in terms of the coordinates $x^{1}, \ldots, x^{n}$ and the hat over a factor means that the factor is omitted. Since the volume of the standard $n$-simplex in $\mathbb{R}^{n}$ is equal to $1 / n!, \int_{\sigma} W\left(\sigma^{*}\right)=$ $\left\langle\sigma^{*}, \sigma\right\rangle=1, W \sigma^{*}$ is the unique constant form with prescribed integral equal to one.

We now consider the case when $1 \leq k \leq n-1$. We will write $\Lambda_{e}^{k}$ for the space of $k$-forms on $\sigma$ with affine coefficients and with constant pull-backs to $k$-faces of $\sigma$. It is obvious from the definition of $W c$ and from (1) that $W c$ has affine coefficients on $\sigma$ for every $c \in C^{k}(\sigma)$. Similarly, since $\iota_{\tau}^{*} W(c)$ is a form of maximal degree on $\tau$, the calculation above, with $k$ replacing $n$, shows that $\iota_{\tau}^{*} W(c)$ is constant on $\tau$ for every $k$-face $\tau$ of $\sigma$. It follows that $W C^{k} \subset \Lambda_{e}^{k}$. Now let $\varphi \in \Lambda_{e}^{k}$. We use the restriction of the de Rham map $R: \Lambda^{k}(\sigma) \longrightarrow C^{k}(\sigma)$,

$$
\langle R \omega, \tau\rangle=\int_{\tau} \omega
$$

to $\Lambda_{e}^{k}$ and consider the difference $\eta=\varphi-W R \varphi$. Clearly, $\eta \in \Lambda_{e}^{k}$. Moreover basic properties of the Whitney mapping (cf. [3, 9]) imply that $R \eta=R \varphi-R W R \varphi=$ $R \varphi-R \varphi=0$, i.e. $\eta$ integrates to zero on every $k$-face of $\sigma$. Since the pull-back $\iota^{*} \eta$ is constant on every such face $\tau, \iota_{\tau}^{*} \eta$ vanishes identically on every $k$-face $\tau$. Thus to show that $\varphi=W R \varphi$ (which would prove our theorem) it suffices to show that every form $\eta \in \Lambda_{e}^{k}$, whose pull-backs to all $k$-faces vanish, is itself identically zero on $\sigma$. Let $\eta$ be such a form. We express it in the standard coordinates of $\mathbb{R}^{n}$ as follows.

$$
\begin{equation*}
\eta=\sum_{I}\left(b_{I}+a_{I, 1} x^{1}+\ldots+a_{I, n} x^{n}\right) d x^{I} . \tag{2}
\end{equation*}
$$

Here $I$ is a multi-index $I=\left(i_{1}<i_{2}<\ldots<i_{k}\right), 1 \leq i_{j} \leq n$ for every $j$ and $d x^{I}=d x^{i_{1}} \wedge d x^{i_{2}} \wedge \ldots \wedge d x^{i_{k}}$. We will abuse the notation at times and think of $I$ as a set. Fix a multi-index $J$ and consider the coordinate plane of the variables $x^{j_{1}}, x^{j_{2}}, \ldots, x^{j_{k}}$.

Let $\tau_{J}$ denote the $k$-face of $\sigma$ contained in that plane. By assumption $\iota_{\tau_{J}}^{*} \eta$ is identically zero. The variables $x_{t}$ for $t \notin J$ vanish in this plane so that

$$
\begin{equation*}
\iota_{\tau_{J}}^{*} \eta=\sum_{t \in J}\left(a_{J, t} x^{t}+b_{J}\right) d x^{J} \equiv 0 . \tag{3}
\end{equation*}
$$

Since $J$ was arbitrary, $b_{J}=0$ and $a_{J, t}=0$ for all $J$ and all $t \in J$. It follows that we can rewrite (2) on $\sigma$ as follows.

$$
\begin{equation*}
\eta=\sum_{I} \sum_{j \notin I} a_{I, j} x^{j} d x^{I} . \tag{4}
\end{equation*}
$$

Again, fix the multi-index $L$, an integer $m \notin L, 1 \leq m \leq n-1$, and the simplex $\tau=\left[e_{m}, e_{l_{1}}, \ldots, e_{l_{k}}\right] . \tau$ is a $k$-simplex in the $(k+1)$-plane $P$ with coordinates $x^{m}, x^{l_{1}}, \ldots, x^{l_{k}}$ as in the figure below. Recall that on $\tau, x^{l_{1}}, \ldots, x^{l_{k}}$ can be taken as local coordinates since

$$
\begin{equation*}
x^{m}=1-\left(x^{l_{1}}+\ldots+x^{l_{k}}\right) . \tag{5}
\end{equation*}
$$

Moreover,

$$
\begin{equation*}
d x^{m}=-\left(d x^{l_{1}}+\ldots+d x^{l_{k}}\right) . \tag{6}
\end{equation*}
$$



We express the pull-back $\iota_{\tau}^{*} \eta$ in terms these coordinates using (5) and (6). Observe that if $I \cup\{j\} \neq L \cup\{m\}$ one of the indices in $I \cup\{j\}$ is not in $L \cup\{m\}$. The corresponding variable is identically zero on the plane $P$ so that the summand
$a_{I, j} x^{j} d x^{I}$ vanishes on $P$ and is therefore equal to zero when pulled back to $\tau$. Therefore

$$
\begin{equation*}
\iota_{\tau}^{*} \eta=\sum_{I \cup\{j\}=L \cup\{m\}} a_{I, j} x^{j} d x^{I} . \tag{7}
\end{equation*}
$$

Now consider the summand with $I=L$ and $j=m$. The coefficient of $d x^{L}$ in this term is

$$
a_{L, m} x^{m}+a_{L, l_{1}} x^{l_{1}}+\ldots+a_{L, l_{k}} x^{l_{k}}
$$

and we use (5) to eliminate $x^{m}$.
Thus, on $\tau$, the coefficient in question can be written as

$$
a_{L, m}-a_{L, m} \sum_{s=1}^{k} x^{l_{s}}+a_{L, l_{1}} x^{l_{1}}+\ldots+a_{L, l_{k}} x^{l_{k}}
$$

Remaining terms in the sum (7) have $j \neq m$. It follows that, for those terms, $x^{j}$ is one of $x^{l_{1}}, \ldots, x^{l_{k}}$ and $x^{m}$ enters only into the differential monomial $d x^{I}$ from which it can be eliminated using (6). It follows that

$$
\iota_{\tau}^{*} \eta=\left(a_{L, m}+\text { linear terms }\right) d x^{L} .
$$

Since $\iota_{\tau}^{*} \eta$ is assumed to be identically zero, $a_{L, m}=0 . L$ and $m$ were fixed but arbitrary so that $\eta \equiv 0$.

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## References

[1] Alain Bossavit, A uniform rationale for Whitney forms on various supporting shapes, Math. Comput. Simulation 80 (2010), no. 8, 1567-1577, DOI 10.1016/j.matcom.2008.11.005. MR 2647251
[2] Alain Bossavit, Computational electromagnetism, Electromagnetism, Academic Press, Inc., San Diego, CA, 1998. Variational formulations, complementarity, edge elements. MR1488417
[3] Józef Dodziuk, Finite-difference approach to the Hodge theory of harmonic forms, Amer. J. Math. 98 (1976), no. 1, 79-104, DOI 10.2307/2373615. MR407872
[4] Józef Dodziuk, de Rham-Hodge theory for $L^{2}$-cohomology of infinite coverings, Topology 16 (1977), no. 2, 157-165, DOI 10.1016/0040-9383(77)90013-1. MR445560
[5] Jonni Lohi and Lauri Kettunen, Whitney forms and their extensions, J. Comput. Appl. Math. 393 (2021), Paper No. 113520, 19, DOI 10.1016/j.cam.2021.113520. MR4229401
[6] Wolfgang Lück, $L^{2}$-invariants: Theory and applications to geometry and $K$-theory, Ergebnisse der Mathematik und ihrer Grenzgebiete. 3. Folge. A Series of Modern Surveys in Mathematics [Results in Mathematics and Related Areas. 3rd Series. A Series of Modern Surveys in Mathematics], vol. 44, Springer-Verlag, Berlin, 2002, DOI 10.1007/978-3-662-04687-6. MR1926649
[7] Santiago R. Simanca, The (small) vibrations of thin plates, Nonlinearity 32 (2019), no. 4, 1175-1205, DOI 10.1088/1361-6544/aaf3eb. MR3923164
[8] Dennis Sullivan, Cartan-de Rham homotopy theory, Colloque "Analyse et Topologie" en l'Honneur de Henri Cartan (Orsay, 1974), Astérisque, No. 32-33, Soc. Math. France, Paris, 1976, pp. 227-254. MR402729
[9] Hassler Whitney, Geometric integration theory, Princeton University Press, Princeton, NJ, 1957. MR 87148

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