

BOUNDS FOR LINEAR FUNCTIONALS¹

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We shall obtain upper and lower bounds for certain functionals associated with linear equations involving positive operators. Attention is focused on these functionals because of their considerable physical significance in applications. A bound from one side is furnished by the usual variational principle. For boundary value problems the reciprocal variational principle introduced by Friedrichs, and later modified by Diaz, provides a complementary bound. In the present article we extend these ideas to an integral equation over a domain E . Our procedure requires information (which is often available) for the same integral equation over some larger domain E' . This approach bears resemblance to the one used by Weinstein and Aronszajn in a series of papers dealing with eigenvalue problems.

Suppose then that we wish to estimate

$$I = \int_E f(x)u(x)dx$$

where

$$(1) \quad u(x) + \int_E k(x, y)u(y)dy = f(x), \quad x \in E.$$

We assume that we know how to solve the integral equation

$$(2) \quad Az = z(x) + \int_{E'} k(x, y)z(y)dy = h(x), \quad x \in E',$$

for some domain $E' \supset E$.

The situation described above occurs frequently in applications. For instance, if the domains are one-dimensional and the kernel is a difference kernel $k(x-y)$, then the integral equation (2) is easily solved if (a) $k(x)$ has period T and E' is an interval of length T , or (b) $k(x)$ is Fourier transformable and E' is the whole real axis.

Since the method we employ is not restricted to integral equations, we describe it in a slightly more abstract setting.

Let A be a real, self-adjoint, positive operator on the space of real L_2 functions over E' . The usual inner product of two functions $v(x)$

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and $w(x)$ is written $\langle v, w \rangle$. We denote by P the projection operator defined by

$$\begin{aligned} Ph &= h(x), & x \in E, \\ &= 0, & x \in (E' - E). \end{aligned}$$

The integral equation (1) can then be rewritten

$$(3) \quad PAu = f, \quad \text{with } Pu = u, \quad Pf = f.$$

If u has been found, then Au can be calculated for all $x \in E'$ and we have

$$(4) \quad Au = f + g \quad \text{with } Pg = 0.$$

The equation (2) takes the form $Az = h$, where A^{-1} is regarded as known.

We wish to estimate $I = \langle f, u \rangle = \langle PAu, u \rangle = \langle Au, u \rangle$. It is convenient to introduce a new inner product $[v, w] = \langle v, Aw \rangle$, in terms of which $I = [u, u]$.

In what follows we need the Schwarz inequality

$$[u, u] \geq [v, u]^2 / [v, v] \quad \text{for all } v \ni [v, v] \neq 0,$$

and Bessel's inequality in the simple form

$$[u, u] \leq [v, v], \quad \text{for all } v \ni [v - u, u] = 0.$$

From the Schwarz inequality we find the well-known lower bound

$$(5) \quad I \geq \langle f, v \rangle^2 / \langle v, Av \rangle \quad \text{for all } v \ni v = Pv.$$

It is clear that the maximum is actually attained for $v = cu$, where c is any nonzero constant.

To apply the Bessel inequality, we must first characterize functions v for which $[v - u, u] = \langle A(v - u), Pu \rangle = 0$. This condition is certainly satisfied if $PA(u - v) = 0$ or $PAv = f$. We should observe that this last equation is not identical with (3) since we do *not* require $Pv = v$. The Bessel inequality then yields the following upper bound:

$$(6) \quad I \leq \langle v, Av \rangle \quad \text{for all } v \ni PAv = f.$$

By choosing $v = u$, the minimum in (6) is obviously attained.

We now rewrite (6) with a view toward the application of the Rayleigh-Ritz method. Let v_0 be an arbitrary fixed function such that $PAv_0 = f$, that is,

$$Av_0 = f + g_0, \quad \text{with } Pg_0 = 0.$$

Since A^{-1} is known, this equation can be solved for any g_0 , but convenience or physical considerations will usually dictate the choice of g_0 . We then define

$$I_0 = \langle v_0, Av_0 \rangle = \langle v_0, f + g_0 \rangle.$$

Any function v for which $PAv=f$ can be written $v=v_0+w$, with $Aw=q$ and $Pq=0$. Substituting in (6), we find

$$(7) \quad I \leq I_0 + 2\langle v_0, q \rangle + \langle q, A^{-1}q \rangle; \quad Pq = 0.$$

We note that the right side of this inequality reduces to I when q is chosen equal to $g-g_0$, where g is defined from (4).

To apply the Rayleigh-Ritz procedure to (7), we introduce an independent set of functions ψ_1, \dots, ψ_n with the property $P\psi_k=0$, $k=1, \dots, n$. Then, for any choice of c_1, \dots, c_n ,

$$I \leq I_0 + 2\langle v_0, \sum_{k=1}^n c_k \psi_k \rangle + \langle \sum_{k=1}^n c_k \psi_k, A^{-1} \left(\sum_{j=1}^n c_j \psi_j \right) \rangle.$$

The values of the coefficients which minimize the right side of the above inequality are obtained from the Galerkin equations

$$(8) \quad \langle v_0, \psi_k \rangle = - \sum_{j=1}^n c_j \langle A^{-1} \psi_j, \psi_k \rangle; \quad k = 1, \dots, n.$$

The corresponding approximation, call it q^* , to $g-g_0$ is then

$$(9) \quad q^* = \sum_{k=1}^n c_k \psi_k,$$

where the $\{c_k\}$ are calculated from (8).

We observe that q^* satisfies the reciprocity principle

$$\langle v_0, q^* \rangle = - \langle A^{-1} q^*, q^* \rangle$$

so that

$$I \leq I_0 - \langle q^*, A^{-1} q^* \rangle = I_0 + \langle v_0, q^* \rangle.$$

If we use a one term approximation $q^* = c\psi$, we find

$$I \leq I_0 - \langle v_0, \psi \rangle^2 / \langle \psi, A^{-1} \psi \rangle; \quad P\psi = 0.$$

In conjunction with (5), we have

$$\langle \phi, f \rangle^2 / \langle \phi, A\phi \rangle \leq I \leq I_0 - \langle v_0, \psi \rangle^2 / \langle \psi, A^{-1} \psi \rangle,$$

where $P\phi = \phi$ and $P\psi = 0$. In practice, the trial functions ϕ and ψ

should be chosen to be reasonable approximations to u and $g-g_0$, respectively.

Applications to physical problems will be described elsewhere.

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