
Take an $n$-dimensional manifold $M$. Endow it with a distribution, by which I mean a smooth linear subbundle $D \subset TM$ of its tangent bundle $TM$. So, for $x \in M$, we have a $k$-plane $D_x \subset T_x M$, and by letting $x$ vary we obtain a smoothly varying family of $k$-planes on $M$. Put a smoothly varying family $g$ of inner products on each $k$-plane. The data $(M, D, g)$ is, by definition, a sub-Riemannian geometry.

Take the viewpoint that we can explore $M$ only by traveling along paths tangent to $D$. Call such paths horizontal. Since $g$ measures lengths of horizontal ($D$) vectors, we can measure lengths of horizontal paths and formulate the sub-Riemannian geodesic problem: to find the shortest horizontal path connecting two given points.

Gromov and his school use the term Carnot-Carathéodory geometry for what we call sub-Riemannian geometry. The sub-Riemannian analogues of Euclidean space are a class of Lie groups endowed with left-invariant sub-Riemannian metrics, christened Carnot groups by the Gromov school, homogeneous groups by Stein and his school, and their Lie algebras named symbol algebras by Tanaka’s school.

The first nontrivial example is the 3-dimensional Heisenberg group and is the simplest non-Euclidean Carnot group. Its sub-Riemannian geodesic problem is essentially the isoperimetric problem. Take $M$ to be standard 3-space $\mathbb{R}^3$. Define $D$ by the vanishing of the 1-form $\theta = dz - (1/2)(xdy - ydx)$; in other words $D_{(x,y,z)}$ is the 2-plane $\{(v_1, v_2, v_3) : v_3 - (1/2)(xv_2 - yv_1) = 0\}$. $D$ can be visualized as a kind of continuous spiral staircase. See the accompanying Figure 1, inspired by the description on p. 40 of [1]. Along the $z$-axis, the 2-planes of $D$ are parallel to the $xy$-plane. As we move out radially from the axis along any orthogonal ray, the 2-plane spins about the axis of the ray, rotating monotonically so that by the time we “reach” infinity, they have just rotated by 90 degrees, turning vertical. We have just described the standard contact structure on 3-space. The projection along the $z$-axis maps each plane linearly onto the $xy$-plane. Declaring these restricted projections to be isometries defines the family $g$ of inner products on $D$. Consider a horizontal path $\gamma$ connecting the origin to a point $A$ units up along the $z$-axis. The projection $c$ of $\gamma$ to the $xy$-plane is closed, and by definition of $g$ the sub-Riemannian length of $\gamma$ equals the standard Euclidean length of $c$. An application of Stokes theorem to the equation $A = (1/2) \int_c (xdy - ydx)$ shows that the (signed) area inside $c$ equals the height $A$. Thus, the sub-Riemannian geodesic problem with endpoint constraints $\gamma(0) = (0, 0, 0)$, $\gamma(1) = (0, 0, A)$ is equivalent to the dual of the isoperimetric problem: among all closed curves $c$ in the plane enclosing a fixed area $A$, find the one of minimum length.

Gauge theory and falling cats. My introduction to sub-Riemannian geometry started with a science fiction story [12], or at least what sounded like science fiction at the time. Jair Koiller (now working in Rio de Janiero) had told me how

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2000 Mathematics Subject Classification. Primary 53C17; Secondary 37J99, 53C60, 58E10, 70H05, 35H10, 53C20, 22E40.

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physicists Shapere and Wilczek were using gauge theory to understand the physical mechanism behind how some microorganisms swim. The story was true. I learned the basics of their paper and became engrossed with a simplified model of their idea, the problem of how a cat, dropped upside down, rights itself in midair so as to land on its feet.

A frozen cat will drop like a stone and lands on its back. If the living cat is to land on its feet, it must change its shape in midair. To understand how this works, first ask, What is a shape? A configuration of the cat is to be viewed as an actual cat, floating in space. When we choose a model of the cat, such as $N$-point masses

![Figure 1. The Heisenberg group and a few of its geodesics.](image1)

![Figure 2. A cat consisting of three mass points rotates 60 degrees, its angular momentum zero at every instant.](image2)
or a continuum model, the set of all configurations becomes a manifold $M$. Two configurations of the cat represent “the same shape” if there exists an orientation-preserving isometry of space taking one to the other. Thus, the space of shapes is the quotient space $M/G$, where $G$ is the group of orientation-preserving isometries of space. To rotate and land on its feet, the cat would love to just move purely in the $G$-direction. But it cannot, as the frozen cat illustrates. The reason such motions are impossible is conservation of its angular momentum $J$. The cat’s initial angular momentum is zero, and it remains constant throughout freefall. $J$ is a vector-valued function on $TM$ which is linear in the fiber variables—the velocities. Thus the assertion $J = 0$ defines a distribution $D$ on $M$. The angular momentum zero constraint forces the cat to move horizontally. The cat’s mass distribution induces a physically natural Riemannian metric on $M$, the kinetic energy metric, and relative to this metric, $D$ is orthogonal to the $G$-orbits: to rotate, the cat must move orthogonally to rotations. Figure 2, modified from Figure 8 of Chenciner [4], is a cartoon of how a cat consisting of three equal mass points might rotate by 60 degrees. If we restrict this metric (or any other) to $D$ we get a $g$ as above, and the cat’s problem has become a case of the sub-Riemannian geodesic problem.

I will divide sub-Riemannian geometry into three streams of research, according to the methods and motivations used. The first stream arises out of optimal control and surrounds the sub-Riemannian geodesic problem. I am most familiar with this stream. The second stream is partial differential equations (PDEs), exemplified by a theorem of Hormander. The third stream concerns sub-Riemannian geometries arising out of limits of Riemannian or other geometries.

**Optimal control.** The deepest results on the sub-Riemannian geodesic problem have been obtained by mathematicians trained in optimal control. The basic problem in nonlinear control theory is to construct paths $\gamma(t)$ connecting two points $q_0, q_1$ of a manifold $M$, subject to path constraints of the form $\dot{\gamma}(t) \in V(\gamma(t))$, where $V(x) \subset T_xM$ are subsets varying with $x$. The sets $V(x)$ are represented as images of a map $u \mapsto f(x; u)$ where the “control” $u$ varies over some auxiliary set, usually a subset of a Euclidean space. In applications $u$ is fuel, steering angle, temperature, etc. To relate this setup to sub-Riemannian geometry, set $V(x) = D_x$. If we throw in a function $L(x, u)$ and try to minimize $\int L(\gamma(t), u(t)) dt$ over the space of all controls $u(t)$ which “steer” from $q_0$ to $q_1$ (thus $\dot{\gamma} = f(\gamma, u(t))$, $\gamma(0) = q_0$, $\gamma(1) = q_1$), we have formulated the basic problem in optimal control. By taking the $V(x) = D_x$ and expressing $L$ using $g$, we have represented the sub-Riemannian geodesic problem as a problem in optimal control.

The first question that arises when facing the sub-Riemannian geodesic problem is accessibility. Is there any horizontal path at all which connects the two given points? If $D$ is involutive—meaning Lie brackets of horizontal vector fields remain horizontal—then, by the Frobenius integrability theorem, the answer is no for generic pairs of points. The Frobenius theorem asserts that the set of points connected to $q_0$ by a horizontal path form an immersed submanifold of dimension $k = \text{rank}(D)$ called the leaf of $D$ through $q_0$. A typical $q_1$ will not lie on this leaf and so cannot be horizontally connected to $q_0$. The first big theorem in the subject, the Chow-Rashevskii theorem [5], [11], asserts that for $D$ at the opposite extreme of involutive, the answer is yes, locally. The technical term for at the opposite extreme of involutive is bracket generating. A distribution $D$ is called bracket generating if
its horizontal vector fields, \( X_1, X_2, \ldots \), together with their iterated Lie brackets, 
\([X_i, X_j], [X_i, [X_j, X_k]], \ldots \), span the tangent bundle to \( M \).

**Subelliptic PDEs.** Chow-Rashevskii’s bracket generating condition reappeared
40 odd years later in Hörmander [8]. Form a “sub-Laplacian” \( \Delta = \sum_{i=1}^{k} X_i^2 \), where
the \( X_i \) are a (local) orthonormal frame for \( D \). Does \( \Delta \) “act” like the usual Laplacian,
meaning are there close analogues of the usual elliptic regularity theorems?
Hörmander’s theorem asserts that the answer is yes, provided that \( D \) is bracket generating.
In honor of his theorem, the bracket generating condition is often referred to as
Hörmander’s condition in the PDE literature.

Most problems in Riemannian geometric analysis involve, in one way or another,
the Riemannian Laplacian. So it is not a surprise that sub-Riemannian Laplacians
are central to sub-Riemannian geometric analysis. Over the last ten years there
has been enormous work on the sub-Riemannian versions of the minimal surface
equations, constant mean curvature equations, Yamabe’s problem, etc. Most of the
analysis has been on Carnot groups.

**Carnot and Carathéodory.** Carnot, working on heat engines, realized that certain
thermodynamic states are inaccessible from other states. At the urging of his friend
the physicist Max Born, Carathéodory used a local version of inaccessibility to
derive the second law of thermodynamics, and of entropy. What he established was
a kind of converse to Chow-Rashevskii for co-rank 1 (\( k = n - 1 \)) distributions \( D \).
Carathéodory showed that if every point in such an \( (M, D) \) has a nearby inaccessible
point, then \( D \) is involutive. See pp. 38–42 of [1] for a beautiful rendition of this
story.

**Nilpotentization and Carnot groups.** The infinitesimal model of a Riemannian
manifold is Euclidean space. The infinitesimal model of a sub-Riemannian geometry
\((M, D, g)\) is a Carnot group: a simply connected Lie group \( G \) whose underlying Lie
algebra \( \mathfrak{g} \) is graded nilpotent, i.e., \( \mathfrak{g} = \mathfrak{g}_1 \oplus \mathfrak{g}_2 \oplus \cdots \oplus \mathfrak{g}_r \) where \([\mathfrak{g}_i, \mathfrak{g}_j] \subset \mathfrak{g}_{i+j} \) and
\( \mathfrak{g}_s = 0 \) for \( s > r \). To construct such a \( \mathfrak{g} \) from the infinitesimal neighborhood of a
point \( p \) of a manifold with distribution \( D \), set \( \mathfrak{g}_1(p) = D_p \). Take \( \mathfrak{g}_k(p) \) to be the
vector space of \( T_pM \) spanned by all \( k \)-fold iterated Lie brackets of horizontal vector
fields, evaluated at \( p \), modulo the subspace of \((k-1)\)-fold iterated brackets of such
vector fields. We call \( p \) a regular point if the dimensions of the \( \mathfrak{g}_q(q) \) are constant for
\( q \) near \( p \). In this case, the Lie bracket of vector fields induces a Lie algebra structure
on \( \mathfrak{g}(p) = \mathfrak{g}_1(p) \oplus \cdots \oplus \mathfrak{g}_r(p) \). (An open dense subset of points of \( M \) are regular.)
The bracket generating condition on \( D \) implies that \( \mathfrak{g}_1(p) \) Lie generates \( \mathfrak{g}(p) \). We
call \( \mathfrak{g} \) the nilpotentization of the distribution at \( p \). Tanaka’s school [14] calls this
Lie algebra the symbol algebra of \( D \). The nilpotentization of the Heisenberg group
discussed in the beginning is the standard 3-dimensional Heisenberg algebra.

Since \( \mathfrak{g} \) is nilpotent, the exponential map yields a diffeomorphism from \( \mathfrak{g} \) to \( G \).
Translate \( \mathfrak{g}_1 \) about \( G \) to obtain a left-invariant distribution on \( G \). Now throw in
the metric \( g_p \) on \( D_p = \mathfrak{g}_1 \). We get a sub-Riemannian structure on \( G \) for which left
translation acts by isometry and whose nilpotenization is again \( \mathfrak{g} \). Various theorems
assert quantitatively how the Carnot group \( G \) approximates the sub-Riemannian
manifold \( M \) near \( p \). Unlike the Riemannian case, this approximation is generally
not Lipshitz.
There is a striking alternative definition of a Carnot group, which we will have
cause to use later. Scalar multiplication on a Euclidean space $E$ defines a 1-
parameter family of metric dilations, $x \mapsto tx = \delta_t(x)$, which are also group au-
tomorphisms when $E$ is viewed as an Abelian group. A Carnot group is a Lie
group endowed with a left-invariant metric which admits a 1-parameter family $\delta_t$
of metric dilations which are also group automorphisms. On the infinitesimal level,
$\delta_t(v) = (tv_1 \oplus t^2v_2 \oplus \cdots \oplus t^rv_r)$ for $v = (v_1 \oplus v_2 \oplus \cdots \oplus v_r) \in g = g_1 \oplus g_2 \oplus \cdots \oplus g_r$.

**Metric limits.** The most surprising applications of sub-Riemannian ideas can be
found in Gromov’s strange and beautiful book-length paper “Groups of Polynomial
Growth and Expanding Maps” [7]. Given a discrete group $\Gamma$ (preferably infinite)
with a finite set of generators, one forms its Cayley graph by declaring two elements
of the group to be connected by an edge if multiplying one element (on the right,
say) by a generator yields the other. Setting all edge lengths to 1 yields a metric on
$\Gamma$. The group is said to have polynomial growth if the volume (number of elements)
of balls of radius $R$ is bounded by a polynomial in $R$, as $R \to \infty$. A group is called
nilpotent if all sufficiently long iterated commutators are the identity. Bass, Milnor,
and others had proved that discrete nilpotent groups have polynomial growth, as
did their close cousins, the virtually nilpotent. Gromov proved the converse: any
group $\Gamma$ of polynomial growth is virtually nilpotent. Here is how he did it. He
shrunk all the edges of the graph to length $\epsilon$, thus getting scaled metric spaces $\Gamma_\epsilon$
on the same underlying group. He constructed a metric—now known as the Gromov-
Hausdorff metric—on the space of pointed metric spaces and showed that as $\epsilon \to 0$
is $\Gamma_\epsilon$ converged to something in this space of all pointed metric spaces. Then he
proved that this something was a Carnot group $G$ endowed with a sub-Riemannian geometry (more accurately, sub-Finsler). The distribution on $G$ corresponds to the
choice of generators. The original group $\Gamma$ (or a finite index subgroup thereof) is
embedded in $G$ as a lattice. Lattices in Carnot groups are nilpotent, finishing the
proof. (For the experts, apologies for playing a bit fast and loose with the full
truth.) More on Carnot groups momentarily.

Gromov was inspired by the appearance of sub-Riemannian geometries in proofs
of Mostow rigidity theorems [10]. Mostow’s original rigidity theorem asserts that
in dimensions $n \geq 3$ if two compact hyperbolic $n$-manifolds have isomorphic funda-
damental groups, then they are diffeomorphic. By Hadamard’s theorem, such
groups form lattices within the standard simply connected hyperbolic $n$-space, and
Mostow’s proof involves studying the asymptotics of these lattices. Replacing hy-
perbolic space by one of the other noncompact symmetric spaces, one can formulate
analogous rigidity theorems. The simplest noncompact symmetric space after hy-
perbolic space is complex hyperbolic space, a complete Riemannian metric space
whose curvatures are bounded between $-1/4$ and $-1$. The boundary at infinity
for complex hyperbolic space is an odd-dimensional sphere which inherits, by a
Riemannian limiting procedure, a contact distribution endowed with a conformal
sub-Riemannian structure. Isometries of complex hyperbolic space are in one-to-one
correspondence with conformal sub-Riemannian isometries of the boundary. Rigid-
ity problems for lattices in complex hyperbolic spaces can be rephrased as delicate
regularity problems concerning quasi-conformal maps of the boundary. A number
of deep works in sub-Riemannian analysis and geometry have been inspired by these
connections. Most of this work, as with most of the analysis in sub-Riemannian
PDEs, has concentrated on the case of Carnot groups.
In each of the research streams just described, there arise phenomena with no Riemannian counterparts. Associated with these phenomena are open problems. I discuss two.

**Geodesics that do not satisfy the geodesic equations.** All the geodesics for a Riemannian metric on $M$ arise out of a single Hamiltonian system: the geodesics are the projections to $M$ of solutions to a system of ODEs on the cotangent bundle $T^*M$. In sub-Riemannian geometry, there is a similar Hamiltonian system on $T^*M$ whose solutions, projected to $M$, are sub-Riemannian geodesics. For some time it had been believed that all sub-Riemannian geodesics arise in this way. Alleged theorems, published in the literature, claimed that “all sub-Riemannian geodesics satisfy the geodesic problem”.

I came up with a counterexample to these alleged theorems in the case $(k, n) = (2, 3)$. The distribution $D$ for my example is a Martinet distribution—a generic degeneration of a rank 2 contact distribution. Liu and Sussmann generalized my counterexample to the case of $(2, n)$, for general $n$. The resulting singular minimizers—geodesics that do not satisfy the geodesic equations—are stable: they cannot be perturbed away by perturbing the sub-Riemannian structure.

The key error in purported proofs of these alleged theorems was a misunderstanding of the fundamental result of optimal control theory, the maximum principle of Boltyanski-Gramkrelidze-Pontrjagin. (See the book by L.C. Young [15] for clear expositions around this principle.) The maximum principle is a fancy version of the method of Lagrange multipliers. Recall this method for the problem of minimizing a function $F$ subject to a constraint $G = \text{const}$. Write $x$ for the vector variable of $F$ and $G$. Introduce the multiplier $\lambda$. Solve the pair of equations $dF(x) + \lambda dG(x) = 0$, $G(x) = \text{const}$ for $x$. Sort through the resulting solutions by evaluating $F$ to find the minimizer. If the minimizer is not a singular point of the constraint surface $G = \text{const}$, then the method will find it. But if the constraint surface has singular points, one of these may be the minimizer, and the method will most likely miss it.

In the sub-Riemannian application, the variable $x$ is a horizontal path, varying over the space (a Banach manifold) of paths issuing from the initial point $q_0 = \gamma(0)$. $F$ is the length of the path. The constraint map $G$ is the endpoint map, $G(\gamma) = \gamma(1)$, so that $G = \text{const}$ imposes the final endpoint condition $\gamma(1) = q_1$. The minimizers caught by the method are the solutions to the sub-Riemannian geodesic equations as described by Hamilton’s equations. The ones missed—minimizers that are singular points of the endpoint map—are the singular or abnormal extremals appearing in the maximum principle. These curves embody phenomena not present in the Riemannian case where the corresponding endpoint map is always regular.

How rare are singular extremals, the singular curves for the endpoint map? Sard’s theorem, fundamental in differential topology, asserts that the set of critical values of a smooth mapping between finite-dimensional manifolds has measure zero. The endpoint mapping $G$ for a distribution is a map between an infinite-dimensional path space (a Banach manifold) and a finite-dimensional manifold, so we cannot just quote Sard. Does an analogous Sard theorem hold for the endpoint map? In other words, how “rare” are endpoints of singular extremals? This basic question in the theory remains open.

The dimension is bigger than the dimension. The notions of Hausdorff dimension and measure make sense for a general metric space. The Hausdorff dimension of a Riemannian $n$-manifold is $n$ and its Hausdorff measure is the Riemannian volume form, up to a constant scale factor.
Recall that a point $p$ of a sub-Riemannian manifold is called regular if the dimensions of the vector spaces $g_i$ used in constructing the nilpotentization are constant near $p$. John Mitchell proved that in a neighborhood of a regular point the Hausdorff dimension $N$ of the sub-Riemannian manifold is given by $N = \Sigma \dim(g_i)$. Since $n = \Sigma \dim(g_i)$, we have $N > n$: the metric dimension is larger than the manifold dimension. For example, 4 is the Hausdorff dimension of the 3-dimensional Heisenberg group.

For the case of a Carnot group we have a simple explanation for this dimensional discrepancy. The Lebesgue measure $d\text{vol}$ on an $n$-dimensional Euclidean vector space scales to $\lambda^n d\text{vol}$ under metric dilation by a factor $\lambda$. A Carnot group of topological dimension $n$ is diffeomorphic to $n$-dimensional Euclidean space, and under the diffeomorphism $d\text{vol}$ becomes the Haar measure on $G$ which is also its Hausdorff measure (up to scale). But under the Carnot metric dilation $\delta_\lambda$ by a factor $\lambda$, as described above, we compute that $d\text{vol} \mapsto \lambda^{N} d\text{vol}$ with $N$ as above. Thus, the Hausdorff dimension $N$ of $G$ is the usual scaling dimension of $G$, viewed as a metric space, as one finds it discussed in fractal textbooks.

In analysis and geometry, a basic and very useful fact is the Euclidean isoperimetric inequality: there is a universal constant $C$ such that for any bounded domain $\Omega$ in the $n$-dimensional Euclidean space with rectifiable boundary $\partial \Omega$, we have $\text{area}(\partial \Omega)^n \leq C \text{vol}(\Omega)^{n-1}$. Moreover, the optimal (smallest) constant is found by taking $\Omega$ to be the round ball and replacing the inequality by equality. The best constant is directly related to best constants in Sobolev inequalities.

Pansu uncovered the analogous assertion for the Heisenberg group. To understand it, recall that the central fact about the distribution on the Heisenberg group is its nonintegrability: there are no surfaces everywhere tangent to the distribution. Related to this fact is the fact that the Hausdorff dimension of any surface in the Heisenberg group is 3 rather than 2. One has a natural notion of the sub-Riemannian volume of such a surface $S$, which I will still write as $\text{area}(S)$, although it is not a constant multiple of the Euclidean area. Pansu’s Heisenberg isoperimetric inequality asserts that for domains in the Heisenberg group with sub-Riemannian rectifiable boundaries, we have $\text{area}(\partial \Omega)^4 \leq C \text{vol}(\Omega)^3$. It is an open problem to find the best constant $C$ in Pansu’s inequality, which is to say, the Heisenberg isoperimetric replacement for the ball. There is a candidate for this replacement—the bubble sets as described in [2].

The zoo of distributions. Forget the fiber inner product $g$ on $D$. The local study of distributions $D$ is extremely rich, perhaps too rich. This fact came as a huge surprise when I started working in the field.

I was brought up only knowing about contact distributions and involutive distributions. A contact distribution is one with $k = n - 1$, $n$ even, and whose nilpotentization is the Heisenberg algebra on $k$ generators. (We continue using $k$ for rank($D$) and $n = \dim(M)$.) The Darboux theorem asserts that every contact distribution in dimension $n$ is locally diffeomorphic to its nilpotentization. All contact distributions are locally isomorphic. There is no local theory. This fact is essential for the burgeoning field of contact topology. Similarly, the Frobenius theorem asserts that all involutive distributions (with fixed $(k, n)$) are locally diffeomorphic, allowing for the existence of the field of the topology of foliations on manifolds. For both contact and involutive distributions, the local symmetry algebra of the distribution $D$, meaning the Lie algebra of vector fields whose (local) flows map $D$ to itself, is infinite dimensional.
Now the contact condition is stable: if we perturb a contact distribution by a small amount, then the resulting distribution remains contact, and hence by Darboux, locally diffeomorphic to the original. On the other hand, if \((k,n)\) are in the range \(1 < k < n - 1\), \((k,n) \neq (2,4)\), then we can perturb the distribution by an arbitrarily small amount so as to obtain a new distribution not locally diffeomorphic to the original. The first occurring case is \((k,n) = (2,5)\) and is the subject of a famously difficult paper by Cartan [3], his “5 variables paper”. A generic \((2,5)\) distribution has nilpotentization (at a generic point) equal to the free graded three-step Lie algebra \(g\) on two generators \(X,Y\). Thus \(g\) has a basis \(X,Y,Z,U,V\) for which the nonzero bracket relations are \([X,Y] = Z, [X,Z] = U, [Y,Z] = V\). Cartan established the existence of curvature-type tensor invariants for such \(D\)'s, invariants which establish the existence of moduli of inequivalent distributions and provide quantitative means for establishing when two such distributions can be locally diffeomorphic. These Cartan curvatures vanish if and only if \(D\) is locally diffeomorphic to its nilpotentization. The local symmetry algebra of such a flat \(D\) is the first occurring exceptional simple Lie algebra: the 14-dimensional Lie algebra \(g_2\) (in its split real form). If the distribution’s curvatures are nonzero, then the symmetry algebra has dimension less than 14, so in particular the symmetry algebra of such a distribution is never infinite dimensional. (Generically, the symmetry algebra is trivial: there are no symmetries.)

We have just described the intricacies of the first occurring nonstable cases of \((k,n)\). As we increase \(k, n\) and \(n - k\), the situation gets much more complicated. Quickly moduli of inequivalent nilpotentizations arise. Even if we fix the nilpotentization type, Cartan’s curvature invariants become prohibitive even to write down. The situation appears beyond hope of systematic classification. The huge zoo of distributions in the range \(1 \ll k < n - 1\) has stymied the subject. Researchers tend to hide out in the contact case or the Carnot case. Only rarely do we go into situations in which we have to face the diversity of the zoo of inequivalent distributions. Exceptions include detailed work on the \((2,n)\) case by Bryant and Hsu, Liu and Sussmann, Zelenko, Agrachev, and others.

The book. I’ll begin with the good news. The book contains a wealth of worked examples of sub-Riemannian metrics, geodesics, solutions to the Hamilton-Jacobi equations and to the sub-Laplacian. Some appear potentially useful to students. They describe and compute Martinet geodesics in terms of elliptic functions, first uncovered by Kupka, in detail. I learned in Chapter 3 a new proof of the Chow-Rashevskii theorem, based on a theorem they attribute to one “Teleman” whom I did not know. Their section on the motivational thermodynamical work of Carathéodory is also worth spending some time on, although I prefer Born’s treatment [1], referenced earlier. Toward the end there is a chapter not found in other books on sub-Riemannian geometry. This chapter concerns Grushin metrics on \(\mathbb{R}^n\). Grushin metrics are not sub-Riemannian geometries in our sense, because the distribution drops rank at the origin, but they are an interesting class of geometric structures sharing many properties of sub-Riemannian geometries which are susceptible to explicit computation. There is also an interesting appendix on non-solvable linear subelliptic PDEs. I enjoyed puzzling my way through the example in Chapter 3 which is that of the rank 2 distribution defined by the Pfaffian system \(dx - xydz = 0\) on \(\mathbb{R}^3\). Any distribution decomposes its space up into accessible sets—the horizontally connected components. In this example those sets are two
half-spaces and the plane $x = 0$ dividing them. The distribution necessarily fails to be bracket generating along this plane.

Now for the bad news. This text contains at least two serious errors, errors that would confound probably anyone trying to learn the subject. Theorem 4.2.2, on p. 88, the book’s first theorem on sub-Riemannian geodesics, falsely asserts that if we are handed a solution $S$ to the Hamilton-Jacobi equation, then a curve is a geodesic $\gamma$ if and only if it satisfies $d\gamma/dt = \nabla S(\gamma)$. (The $\nabla$ is the horizontal gradient, defined by $dS_x(v) = g(\nabla S(x), v)$ for any horizontal vector $v$ at $x$. The Hamilton-Jacobi equation is the equation $\partial S/\partial t = H(dS_x)$, where $H$ is the sub-Riemannian kinetic energy, above.) It is difficult to list the number of ways this theorem is wrong. To give one, since $S$ is fixed in the theorem, if the theorem were true it would imply that all geodesics passing through a point at a fixed time must point in the same direction: $\nabla S$. (I have posted some other ways the theorem is wrong, and a version of how the theorem can be salvaged, on my web page.)

The other serious error I stumbled across is equation 7.1.2 on p. 138. That equation asserts that a certain differential expression which depends upon the choice of an orthonormal frame for $D$, is in fact independent of that choice. If true, the expression would yield a well-defined connection associated to any sub-Riemannian geometry, a fact unknown to previous researchers, and one which would be central to the theory: a great result if only it were true! The error occurs in the first line of the proof, also the first line on p. 139. It is an error most students of differential geometry have made when learning index notation, especially of the Einstein variety: there is a double sum, but the authors use the same dummy index, here “$\ell$”, for both, inadvertently setting the two summation indices equal. The error is tantamount to insisting that the $(dg)g^{-1}$ term arising in the transformation formula for connections under changes of gauge (local section) is never present, i.e., always zero. For the authors, $g$ is a function with values in the orthogonal group $O(k)$. They do successfully prove just before the offending line, the well-known fact that the diagonal entries of $(dg)g^{-1}$ are zero, but they miss the essential presence of the off-diagonal entries in their alleged derivation. This error destroys much of the sense of their Chapter 7.

By this point, I lost patience. I suspect that the authors have worked out carefully and explicitly many interesting and wonderful low-dimensional examples. As a reference text, it is infuriatingly inconsistent. I urge prospective readers to go carefully through results before using them. The book also lacks from leaving out all post-Carathéodory motivation for bothering with the subject. Gromov’s and Mostow’s application of sub-Riemannian ideas are not touched upon. Cartan’s “zoo” is not mentioned.

Other books. I know of three books in the field. Gromov wrote [6], a book contained within a book. Acres of research can probably still be mined out of his ideas. I wrote [9] which focuses on sub-Riemannian geodesics and outgrowths of the falling cat problem. Finally, the multiauthored book [2] focuses on geometric analysis on the Heisenberg group, with a grand finale of various proofs of Pansu’s Heisenberg isoperimetric inequality. It is well written in a detailed yet friendly manner. As an added bonus it explains a proposed application of Heisenberg geometry toward understanding the workings of a certain part of the visual cortex. I recommend any of these books over the one being reviewed. For the reader who wants just a
taste more of the subject without diving into a book, I recommend [13], with its beautiful exposition and unique way of putting the subject in context.

ACKNOWLEDGMENTS

I would like to thank Alain Chenciner and Rick Moeckel for useful comments on a draft of this book review. I would also like to thank Gregoire Vion for redrawing Figure 1.

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