where \( S_r(a), S_r(b), S_r'(a) \) and \( S_r'(b) \) stand for the sum of all possible combinations of \( r \) factors taken at a time from the sequence of numbers \( q, q^a_1, q^a_2, \ldots, q^a_r; q^b_1, q^b_2, \ldots, q^b_r; q, q^{1-b_1}, q^{1-b_2}, \ldots, q^{1-b}, \) and \( q^{1-a_1}, q^{1-a_2}, \ldots, q^{1-a_r} \) respectively, we can sum up both the series on the right-hand side of (4.6) and obtain

\[
\psi_r \left[ \frac{q^{(a)}_r}{q^{1+(b)_r}} ; z \right]_N = \left[ 1 - q^{N+1}(a)_r ; N + 1 \right] [z - 1][q^{1+(b)_r} ; N] \\
+ \left[ 1 - q^{M+1}(b)_r ; M + 2 \right] \frac{[q^{(2)}_r - \Sigma(b)_r - \Sigma(a)_r]}{[q^{1-(a)_r} ; M + 1]z^{M+1}},
\]

under the set of conditions (4.7).

On the other hand, if in the set of conditions (7a) and (7b) we assume all the conditions except those corresponding to \( k = 1 \) and \( t = 1 \) we can still find the sum of \( M + N + 1 \) terms of a bilateral hypergeometric function by using the result (5).

Department of Mathematics
Lucknow University
Lucknow, India


The Fast Fourier Transform Recursive Equations for Arbitrary Length Records

By G. D. Bergland

1. Introduction. In Cooley and Tukey’s paper on the Fast Fourier Transform [1], a procedure was suggested which allows operation on a time series of length \( N \), where \( N \) is the product of an arbitrary number of integers greater than 1. An alternate procedure is reported in this paper which has the advantage of leading directly to a set of recursive equations for this general case. The notation and methods used have been patterned as much as possible to those of Cooley and Tukey so that the reader can go directly from their paper to this one. These results were noted and used by the author in his Ph.D. dissertation [2].

2. The Recursive Equations when \( N = r_1r_2 \cdots r_m \). Consider the problem of evaluating a complex Fourier series of the form

\[
X(j) = \sum_{k=0}^{N-1} A(k)W^j
\]

where \( W = e^{2\pi i/N}, j = 0, 1, 2, \ldots, N - 1 \).

First \( j \) and \( k \) must be expressed in the following variable radix representation.

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\[ j = j_{m-1}(r_1 r_2 \cdots r_{m-1}) + j_{m-2}(r_1 r_2 \cdots r_{m-2}) + \cdots + j_1 r_1 + j_0, \]
\[ k = k_{m-1}(r_2 r_3 \cdots r_m) + k_{m-2}(r_2 r_3 \cdots r_m) + \cdots + k_1 r_m + k_0. \]
The ranges on the variables are chosen to give a unique representation of each decimal integer (e.g. \( j_0 = 0, 1, \cdots, r_1 - 1; j_1 = 0, 1, \cdots, r_2 - 1; k_0 = 0, 1, \cdots, r_m - 1; k_1 = 0, 1, \cdots, r_{m-1} - 1; \) etc.). This allows (1) to be written as
\[ X(j_{m-1}, j_{m-2}, \cdots, j_1, j_0) = \sum_{k_0} \sum_{k_1} \cdots \sum_{k_{m-1}} A(k_{m-1}, k_{m-2}, \cdots, k_0) W^{jk}. \]
Note that
\[ W^{jk} = W^{j[k_{m-1}(r_2 r_3 \cdots r_m) + \cdots + k_0]} \]
but
\[ W^{jk_{m-1}(r_2 r_3 \cdots r_m)} = W^{j[k_{m-1}(r_2 r_3 \cdots r_m)]} W^{j_0 k_{m-1}(r_2 r_3 \cdots r_m)}. \]
When the product in the exponent is formed, the term may be expressed in the following form
\[ W^{jk_{m-1}(r_2 r_3 \cdots r_m)} = [W^{j_1 r_2 r_3 \cdots r_m}] [k_{m-1}(r_2 r_3 \cdots r_m)]^{j_1} k_{m-1} W^{j_0 k_{m-1}(r_2 r_3 \cdots r_m)}. \]
Note that \( r_1 r_2 r_3 \cdots r_m = N \) and
\[ W^N = \left(e^{2\pi i/N}ight)^N = 1, \]
therefore the bracketed term of (6) taken to any power is still equal to 1 and we have
\[ W^{jk_{m-1}(r_2 r_3 \cdots r_m)} = W^{jk_{m-1}(r_2 r_3 \cdots r_m)} \]
therefore
\[ W^{jk} = W^{j_0 k_{m-1}(r_2 r_3 \cdots r_m)} W^{j[k_{m-2}(r_3 \cdots r_m)]} W^{j_1 r_2 r_3 \cdots r_m} \]
This allows (3) to be written in the form
\[ X(j_{m-1}, \cdots, j_0) = \sum_{k_0} \sum_{k_1} \cdots \sum_{k_{m-2}} \left[ \sum_{k_{m-1}} A(k_{m-1}, \cdots, k_0) W^{j_0 k_{m-1}(r_2 r_3 \cdots r_m)} \right] W^{j[k_{m-2}(r_3 \cdots r_m)]} W^{j_1 r_2 r_3 \cdots r_m}. \]
If the expression in brackets is written as
\[ A_1(j_0, k_{m-2}, \cdots, k_0) = \sum_{k_{m-1}} A(k_{m-1}, \cdots, k_0) W^{j_0 k_{m-1}(r_2 r_3 \cdots r_m)}, \]
(10) may be expressed as
\[ X(j_{m-1}, \cdots, j_0) = \sum_{k_0} \sum_{k_1} \cdots \sum_{k_{m-2}} A_1(j_0, k_{m-2}, \cdots, k_0) W^{j[k_{m-2}(r_3 \cdots r_m)]}. \]
By applying (7) again we see that
\[ W^{jk_{m-2}(r_3 r_4 \cdots r_m)} = W^{j(j_1 r_1 + j_0) k_{m-2}(r_3 r_4 \cdots r_m)}. \]
This allows the innermost sum to be written as
\[ A_2(j_0, j_1, k_{m-3}, \cdots, k_0) = \sum_{k_{m-2}} A_1(j_0, k_{m-2}, \cdots, k_0) W^{j(j_1 r_1 + j_0) k_{m-2} r_3 r_4 \cdots r_m}, \]
leaving (12) in the form

\[ X(j_{m-1}, \ldots, j_0) = \sum_{k_0} \sum_{k_1} \cdots \sum_{k_{m-2}} A_2(j_0, j_1, k_{m-3}, \ldots, k_0) W^{j_0[k_{m-3}r_1 + \cdots + r_m] + \cdots + k_0}. \]

Proceeding in a similar fashion, a set of recursive equations are obtained of the form

\[ A_p(j_0, j_1, \ldots, j_{p-1}, k_{m-p-1}, \ldots, k_0) = \sum_{k_{m-p}} W^{j_0[k_{m-p}r_{p} + \cdots + r_m] k_{m-p}r_{p+1} + \cdots + k_0} W^{j_0[k_{m-p}r_{p} + \cdots + r_m]} \]

\[ (16) \quad p = 1, 2, \ldots, m. \]

Note that the last array calculated gives the Fourier sums as

\[ X(j_{m-1}, \ldots, j_0) = A_m(j_0, \ldots, j_{m-1}). \]

It should be noted that an alternate set of recursive equations is obtained if the \( jk \) product is combined by expanding \( j \) instead of \( k \). If we represent

\[ (17) \quad W^{j_k} = W^{j_k} W^{j_0[r_1 + j_2(r_1 + \cdots + j_{m-1}(r_1 + \cdots + r_m)]}, \]

a set of \( A_1 \) summations can be defined as

\[ A_1(j_0, k_{m-2}, \ldots, k_0) = \sum_{k_{m-1}} A(k_{m-1}, \ldots, k_0) W^{j_0[k_{m-1}r_2 + \cdots + r_m] + \cdots + k_0}. \]

After making use of (7) the \( A_2 \) summations can be defined as

\[ A_2(j_0, j_1, k_{m-3}, \ldots, k_0) = \sum_{k_{m-2}} A_1(j_0, k_{m-2}, \ldots, k_0) W^{j_1[j_{m-2}r_2 + \cdots + r_m] + \cdots + k_0}. \]

The general form of these recursive equations is

\[ A_p(j_0, j_1, \ldots, j_{p-1}, k_{m-p-1}, \ldots, k_0) = \sum_{k_{m-p}} A_{p-1}(j_0, j_1, \ldots, j_{p-2}, k_{m-p}, \ldots, k_0) W^{j_0[k_{m-p}r_{p} + \cdots + r_m] + \cdots + k_0}, \]

\[ (21) \quad p = 1, 2, \ldots, m, \]

where again the last array calculated gives the Fourier sums as

\[ X(j_{m-1}, \ldots, j_0) = A_m(j_0, \ldots, j_{m-1}). \]

3. Discussion of Results. Note that (16) can be easily specialized to the recursive equations for \( N = 2^m \) which Cooley and Tukey reported. Also, their upper bound of \( N(r_1 + r_2 + \cdots + r_m) \) operations to compute \( N \) spectral estimates can be applied to either (16) or (21). Of course this bound can be reduced in cases where the symmetries of the complex exponential weights can be exploited.

Bell Telephone Laboratories, Incorporated
Whippany, New Jersey
