

Numerical Evaluation of an Isoperimetric Constant

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1. Let C be a rectifiable curve in $E^n (n > 2)$, let $L(C)$ be its length and $V[H(C)]$ the volume of its convex hull $H(C)$, and consider the problem of maximizing $V[H(C)]$ subject to the condition $L(C) = \text{constant}$. Depending on the parity of n and on whether C is an open arc, or a closed curve, we have four cases to consider: (A) C is closed and n is even, (B) C is open and n is even, (C) C is open and n is odd, (D) C is closed and n is odd.

Under certain restrictive assumptions on C the case (A) (which includes, as a particular example, the isoperimetric problem of the circle) was solved by Schoenberg [1] who proved that then

$$(1) \quad V[H(C)]/L^n(C) \leq [(\pi n)^{n/2} n!(n/2)!]^{-1}$$

and the inequality is strict except when C is similar to the hypercircle given parametrically by

$$(2) \quad x_{2j-1}(t) = (\sin jt)/j, \quad x_{2j}(t) = (\cos jt)/j \quad (j = 1, \dots, n/2)$$

where $0 \leq t < 2\pi$. In the case (B) one can use the reflection principle to show that here the maximizing curve is also given by (2) but with $0 \leq t \leq \pi$ (a semihypercircle), and the isoperimetric inequality is

$$(3) \quad V[H(C)]/L^n(C) \leq 2^{n/2-1} [(\pi n/2)^{n/2} (n/2)! n!]^{-1}.$$

The case (C) with $n = 3$ was treated by Egerváry [2]: the isoperimetric inequality is

$$(4) \quad V[H(C)]/L^3(C) \leq (18\pi \cdot 3^{1/2})^{-1}$$

and it is strict except when C is one turn of a circular helix of pitch $2^{-1/2}$. Similar treatment applies in general: the maximizing curve is one turn of a hypercircular helix.

The case (D) with $n = 3$ was considered by the author in [3]. Under certain restrictive conditions on C it is possible to express $V[H(C)]$ as an integral:

$$(5) \quad V[H(C)] = 4 \int_0^{s/4} xyz' ds,$$

here accent denotes differentiation with respect to the independent variable s which is the arc-length. The integral in (5) is to be maximized subject to the arc-length condition $x'^2 + y'^2 + z'^2 - 1 = 0$, this leads to the Euler-Lagrange equations

$$(6) \quad 4\lambda^2 x'' = xy^2, \quad 4\lambda^2 y'' = -yx^2, \quad 2\lambda z' = xy,$$

where λ is a constant (the Lagrange multiplier), and the initial conditions are

$$(7) \quad x(0) = c_1, \quad x'(0) = 0, \quad y(0) = 0, \quad y'(0) = c_2, \quad z(0) = 0.$$

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Hence, in the isoperimetric inequality for the case (D),

$$(8) \quad V[H(C)]/L^3(C) \leq B$$

the equality occurs if and only if C is a periodic solution of (6) and (7). It is shown in [3] that for all c_1 and c_2 there is a real-analytic solution (x, y, z) of (6), valid for all s , and for each c , there is $c_2 = f(c_1)$, such that the solution is periodic. Unlike in the cases (A), (B), and (C), the maximizing curve in (D) does not appear to be expressible by elementary, or standard transcendental, functions. Our purpose in this note is to compute numerically the corresponding isoperimetric constant B of (8) (which we propose to call the *Baggins constant*).

2. With suitable scaling the system (6) may be written as

$$(8a) \quad x'' = -xy^2, \quad y'' = -yx^2, \quad z' = xy$$

and the initial conditions are

$$(8b) \quad x(0) = 1, \quad x'(0) = 0, \quad y(0) = 0, \quad y'(0) = b, \quad z(0) = 0.$$

The independent variable t is now proportional to the arc-length s : $s = kt$. The constants k and b are very simply related: by (8a) there exists a first integral $x'^2 + y'^2 + x^2y^2 = b^2$ and since $z' = xy$, we have $k = b$. To determine b we consider the differential system $x'' = -xy^2, y'' = -yx^2, x(0) = 1, x'(0) = y(0) = 0, y'(0) = b$; as shown in [3], the periodic solution of it is symmetric with respect to the lines $x = 0, y = 0, x \pm y = 0$. Writing $x = x(t, b), y = y(t, b)$ and letting t_1 be the smallest positive value of t , for which $x(t, b) = y(t, b)$, we find that t_1 and b are determined by the equations

$$(9) \quad x(t_1, b) - y(t_1, b) = 0, \quad x'(t_1, b) + y'(t_1, b) = 0.$$

For the numerical work we take advantage of the analyticity of x and y in t . Let

$$(10) \quad x(t) = \sum_0^\infty a_n t^n, \quad y(t) = \sum_0^\infty b_n t^n,$$

then by (8a) and (8b) x turns out to be an even function and y an odd one. Substituting the series (10) into (8a) we get

$$(11) \quad \begin{aligned} a_{n+2} &= -[(n+1)(n+2)]^{-1} \sum_{i+j+k=n} a_i b_j b_k, \\ b_{n+2} &= -[(n+1)(n+2)]^{-1} \sum_{i+j+k=n} b_i a_j a_k. \end{aligned}$$

Suppose next that for $m = 0, 1, \dots, n$

$$(12) \quad |a_m| \leq \alpha \lambda^m, \quad |b_m| \leq \alpha \lambda^m.$$

Then, by (11) and by the parities of x and y , we have the estimates

$$|a_{n+2}| \leq [(n+1)(n+2)]^{-1} \alpha^3 \lambda^n N(n), \quad |b_{n+2}| \leq [(n+1)(n+2)]^{-1} \alpha^3 \lambda^n M(n),$$

where $N(n)$ (resp. $M(n)$) is the number of representations of an even (resp. odd) integer n in the form $i + j + k$, with $0 \leq i, j, k$, and i is even (resp. odd) while j and k are odd (resp. even). Therefore, for $n \geq 2$

$$(13) \quad |a_{n+2}| \leq \alpha^3 \lambda^n / 4, \quad |b_{n+2}| \leq \alpha^3 \lambda^n / 4.$$

Hence (12) will have been proved inductively to hold for all m if $1 \leq \alpha$, $b \leq \alpha\lambda$, (initialization, by (8b) and (12)), $\alpha \leq 2\lambda$, (induction, by (12) and (13)).

To optimize the estimates (12) we make λ possibly small by letting $\alpha = 2\lambda$ and $b = \alpha\lambda$, to get $\lambda = (b/2)^{1/2}$, $\alpha = (2b)^{1/2}$ so that (12) becomes

$$(14) \quad |a_m| \leq (2b)^{1/2}(b/2)^{m/2}, \quad |b_m| \leq (2b)^{1/2}(b/2)^{m/2}.$$

Now we determine the constants b and t_1 from (9), by taking a square grid of values for b and t_1 (11 by 11), solving the differential system numerically, and refining then the grid over the square where (9) appears to have its roots. After this is repeated two or three times, we may evaluate x and y numerically for a grid of three points only and fit planes through these triples of points. In this way b and t_1 are determined to be

$$(15) \quad b = 0.92114882, \quad t_1 = 1.22036757.$$

Observe that by (14) the power-series (10) converge for $t < (2/b)^{1/2} = 1.473 \dots$ which exceeds the calculated value t_1 so that the solution procedure by using power series is justified on the interval $0 \leq t \leq t_1$. Moreover, the estimates (14) lead to simple error estimates for truncating the power series for x and y and their derivatives.

3. The evaluation of B is now simple. First, we notice the four-fold inversion symmetry of the curve C given by (8a) and (8b): if it is rotated by $\pm 90^\circ$ about the z -axis the effect is the same as reflecting C in the plane $z = z_{\max}/2$. Therefore, the arc of C given by $0 \leq t \leq t_1$ (where t_1 is given by (15)) has the length $bt_1 = L(C)/8$, and by (5) the volume of $H(C)$ is

$$V[H(C)] = 8 \int_0^{bt_1} xyz' dt.$$

The integral above is easily computed by integrating the product of truncated power series for x , y , and z' , and we get finally $B = 0.0031816877$.

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