

## Determination of the $D^{1/2}$ -Norm of the SOR Iterative Matrix for the Unsymmetric Case

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**Abstract.** This paper is concerned with the determination of the Jordan canonical form and  $D^{1/2}$ -norm of the SOR iterative matrix derived from the coefficient matrix  $A$  having the form

$$A = \begin{pmatrix} D_1 & -H \\ H^T & D_2 \end{pmatrix}$$

with  $D_1$  and  $D_2$  symmetric and positive definite. The theoretical results show that the Jordan form is not diagonal, but has only  $q$  principal vectors of grade 2 and that the  $D^{1/2}$ -norm of  $\mathcal{L}_{\omega_b}$  ( $\omega_b$ , the optimum parameter) is less than unity if and only if  $\bar{\mu} = \rho(B)$ , the spectral radius of the associated Jacobi iterative matrix, is less than unity. Here  $q$  is the multiplicity of the eigenvalue  $i\bar{\mu}$  of  $B$ .

**1. Introduction.** For the iterative solution of the linear system of equations,

$$(1.1) \quad Ax = b,$$

the Jacobi and Gauss-Seidel methods are well known. They are very simple from a computational point of view since only matrix-vector multiplications and linear combinations of vectors are needed. This is also valid for the modification called "Successive Overrelaxation" or "SOR" method, where a relaxation factor is introduced for accelerating the convergence. Let

$$(1.2) \quad A = D - A_L - A_U,$$

where  $D$  is the block diagonal part of  $A$ ,  $-A_L$  and  $-A_U$  are the remaining strictly lower and upper triangular parts of  $A$ ; then, if  $D$  is nonsingular, the SOR method is given by

$$(1.3) \quad x_{k+1} = \mathcal{L}_\omega x_k + \omega(I - \omega L)^{-1} D^{-1} b, \quad k \geq 0.$$

Here,  $x_0$  is an initial vector,  $\mathcal{L}_\omega$  the iterative matrix given by

$$(1.4) \quad \mathcal{L}_\omega = (I - \omega L)^{-1} [(1 - \omega)I + \omega U],$$

and

$$(1.5) \quad L = D^{-1} A_L, \quad U = D^{-1} A_U.$$

Now (1.3) converges if and only if the spectral radius of  $\mathcal{L}_\omega$  is less than unity, and the asymptotic rate of convergence is given by

$$(1.6) \quad R_\infty(\mathcal{L}_\omega) = -\log(\rho(\mathcal{L}_\omega)).$$

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The SOR method has been extensively studied for a symmetric positive definite matrix  $A$  (see, e.g., Varga [5] and Young [7]). For a positive definite and consistently ordered matrix  $A$ , from [5] and [7] we have:

S1.  $\rho(\mathcal{L}_\omega) < 1 \Leftrightarrow \bar{\mu} < 1$  and  $0 < \omega < 2$ .

S2.

$$\rho(\mathcal{L}_\omega) = \begin{cases} \{\omega\bar{\mu} + [\omega^2\bar{\mu}^2 - 4(\omega - 1)^{1/2}]^{1/2}\}^2/4 & \text{if } 0 < \omega < \omega'_b, \\ \omega - 1 & \text{if } \omega'_b \leq \omega < 2. \end{cases}$$

S3.  $\rho(\mathcal{L}_{\omega'_b}) < \rho(\mathcal{L}_\omega)$  if  $\omega \neq \omega'_b$ .

Here,

$$\bar{\mu} = \rho(B) = \rho(L + U), \quad \omega'_b = 2/[1 + (1 - \bar{\mu}^2)^{1/2}].$$

Young [6] has shown that if  $A$  is consistently ordered and the eigenvalues of  $B$  are real and less than unity in modulus, then the Jordan canonical form of  $\mathcal{L}_{\omega'_b}$  is not diagonal. Therefore, in this case, the SOR method converges slower than expected based on the spectral radius  $\rho(\mathcal{L}_{\omega'_b})$ . When  $A$  is symmetric and positive definite, and when  $A$  has the form

$$(1.7) \quad A = \begin{pmatrix} D_1 & -H \\ -K & D_2 \end{pmatrix},$$

Young [7, Chapter 7] determined the  $D^{1/2}$ -norm and  $A^{1/2}$ -norm of  $\mathcal{L}_{\omega'_b}$  (the spectral norms of  $D^{1/2}\mathcal{L}_{\omega'_b}D^{-1/2}$  and  $A^{1/2}\mathcal{L}_{\omega'_b}A^{-1/2}$ , respectively) and pointed out that the  $D^{1/2}$ -norm of  $\mathcal{L}_{\omega'_b}$  is greater than unity in general. Moreover,  $\|\mathcal{L}_{\omega'_b}^m\|$  (the spectral norm of  $\mathcal{L}_{\omega'_b}^m$ ) behaves much like  $\|\mathcal{L}_{\omega'_b}^m\|_{D^{1/2}}$ . However, for large  $m$ ,  $\|\mathcal{L}_{\omega'_b}^m\|_{D^{1/2}} < 1$ , and eventually  $\|\mathcal{L}_{\omega'_b}^m\|_{D^{1/2}}$  tends to zero, though considerably more slowly than  $\rho(\mathcal{L}_{\omega'_b}^m)$ .

For the matrix  $A$  in (1.1) the unsymmetric case is by far not as common as the symmetric one, but nevertheless, unsymmetric matrices appear, e.g., in the numerical solution of the biharmonic equation [1] and the computation of cubic splines, [3] and [4, Chapter 3]. If the matrix  $A$  (1.2) is consistently ordered and  $B$ , given by

$$(1.8) \quad B = L + U,$$

is similar to a skew-symmetric matrix and has either zero eigenvalues or purely imaginary eigenvalues, then from [1], [3], and [4], or the theory of Young [7], we have:

US1.  $\rho(\mathcal{L}_\omega) < 1 \Leftrightarrow 0 < \omega < 2/(1 + \bar{\mu})$ .

US2.

$$\rho(\mathcal{L}_\omega) = \begin{cases} 1 - \omega & \text{if } 0 < \omega \leq \omega_b, \\ \left[ \frac{\bar{\mu}\omega + \sqrt{\omega^2\bar{\mu}^2 + 4(\omega - 1)}}{2} \right]^2 & \text{if } \omega_b < \omega < \frac{2}{1 + \bar{\mu}}. \end{cases}$$

US3.  $\rho(\mathcal{L}_{\omega_b}) < \rho(\mathcal{L}_\omega)$  if  $\omega \neq \omega_b$ .

Here,

$$(1.9) \quad \bar{\mu} = \rho(B), \quad \omega_b = 2/(1 + \sqrt{1 + \bar{\mu}^2}).$$

Notice that in this case we can always choose the relaxation factor  $\omega$  such that  $\rho(\mathcal{L}_\omega) < 1$ , no matter how large  $\bar{\mu}$  is. This is very different from the symmetric

case. Another difference between the two cases is that the optimum factor  $\omega_b$  for the unsymmetric case is less than unity and the optimum factor  $\omega'_b$  for the symmetric case is greater than unity. It is also important to note that overestimating  $\omega'_b$  is better than an underestimation, but for  $\omega_b$  an underestimate is better than overestimating.

However, to our knowledge, the Jordan canonical form and  $D^{1/2}$ -norm of  $\mathcal{L}_\omega$  for the unsymmetric case are not discussed in the literature.

In this paper we will investigate these problems under the assumption that in (1.1) the matrix  $A$  has the special form (1.7) with  $D_1$  and  $D_2$  symmetric and positive definite and  $K^T = -H$ . We will obtain some results similar to those for the symmetric case.

In the next section we review some properties for skew-symmetric matrices required for their application in the later sections. In Section 3 we construct the basis of eigenvectors of the associated Jacobi matrix  $B$  which is similar to a skew-symmetric matrix. In Section 4 we will show that the Jordan canonical form of  $\mathcal{L}_{\omega_b}$  is not a diagonal matrix, but has only  $q$  principal vectors of grade 2 associated with  $\omega_b - 1$ , the eigenvalues of  $\mathcal{L}_{\omega_b}$ . Here,  $q$  is the multiplicity of the eigenvalue  $i\bar{\mu}$  ( $= i\rho(B)$ ) of  $B$ . Hence,  $\|\mathcal{L}_{\omega_b}^m\|$ , the spectral norm, behaves like  $m \cdot \rho(\mathcal{L}_{\omega_b})^{m-1}$  rather than  $\rho(\mathcal{L}_{\omega_b})^m$ .

In Section 5 we will determine the  $D^{1/2}$ -norm of  $\mathcal{L}_\omega$  and point out that if  $\bar{\mu} = \rho(B) \geq 1$ , then  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} \geq 1$ . However, in Section 6, we will show that for any  $\bar{\mu} > 0$ , for  $m$  large enough,  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} < 1$ . Eventually,  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$  converges to zero, though considerably more slowly than  $\rho(\mathcal{L}_{\omega_b}^m)$ .

In this paper, almost all the notations used are the same as those adopted by Young [7], and all our work is based on the theory of Young [7].

**2. Some Properties of Skew-Symmetric Matrices.** Let  $A \in R^{n \times n}$  and

$$(2.1) \quad A^T = -A.$$

It is well known that  $A$  has the following properties:

- (a) All diagonal elements of  $A$  are zero.
- (b)  $A$  has either zero eigenvalues or purely imaginary eigenvalues, that is, any eigenvalue  $\mu$  of  $A$  has the form

$$(2.2) \quad \mu = i\xi.$$

Here  $\xi$  is real. Also,  $-\mu = -i\xi$  is an eigenvalue of  $A$ .

- (c)  $A$  is a normal matrix, that is,

$$(2.3) \quad A^T A = A A^T.$$

- (d)  $A$  is unitarily similar to a diagonal matrix.

All the above properties are easy to prove and can be found in any textbook of linear algebra, e.g., see [2].

**3. The Eigenvectors of the Associated Jacobi Iteration Matrix  $B$ .** Consider  $A$  in (1.1) to have the special form

$$(3.1) \quad A = \begin{pmatrix} D_1 & -H \\ -K & D_2 \end{pmatrix},$$

where  $D_1 (\in R^{r \times r})$  and  $D_2 (\in R^{s \times s})$  are symmetric positive definite and

$$(3.2) \quad H^T = -K;$$

the associated Jacobi iterative matrix  $B$  has the form

$$(3.3) \quad B = \begin{pmatrix} 0 & F \\ G & 0 \end{pmatrix}.$$

Here,

$$(3.4) \quad G = D_2^{-1}K, \quad F = D_1^{-1}H.$$

Because  $D_1$  and  $D_2$  are positive definite, we can choose symmetric and positive definite matrices  $D_1^{1/2}$  and  $D_2^{1/2}$  such that

$$(3.5) \quad D_1^{1/2}D_1^{1/2} = D_1, \quad D_2^{1/2}D_2^{1/2} = D_2.$$

If we write

$$(3.6) \quad D = \begin{pmatrix} D_1^{1/2} & 0 \\ 0 & D_2^{1/2} \end{pmatrix},$$

then we have

$$(3.7) \quad D^{1/2}BD^{-1/2} = \begin{pmatrix} 0 & D_1^{-1/2}HD_2^{-1/2} \\ D_2^{-1/2}KD_1^{-1/2} & 0 \end{pmatrix}.$$

Hence  $B$  is similar to a skew-symmetric matrix, and thus unitarily similar to a diagonal matrix.

In this section we will construct a basis of eigenvectors for  $B$ . From (3.3) we have

$$(3.8) \quad B^2 = \begin{pmatrix} FG & 0 \\ 0 & GF \end{pmatrix}.$$

Evidently,  $B^2$  is also similar to a diagonal matrix and, in fact, the  $(r \times r)$  matrix  $FG$  and the  $(s \times s)$  matrix  $GF$  are also similar to diagonal matrices, where  $r+s = n$ , the order of the matrix  $A$ . Also note that  $FG$  and  $GF$  have nonpositive eigenvalues. Let the  $p$  eigenvectors of  $FG$  associated with the nonzero eigenvalues  $\nu_1, \nu_2, \dots, \nu_p$  be  $\xi_1, \xi_2, \dots, \xi_p$ , i.e.,

$$(3.9) \quad FG\xi_j = \nu_j\xi_j, \quad j = 1, 2, \dots, p.$$

If we let

$$(3.10) \quad \eta_j = G\xi_j, \quad j = 1, 2, \dots, p,$$

then  $\eta_j \neq 0$ , and  $\eta_j$  is an eigenvector of  $GF$  associated with  $\nu_j$ , i.e.,

$$(3.11) \quad GF\eta_j = \nu_j\eta_j, \quad j = 1, 2, \dots, p.$$

Moreover, since the  $\xi_j, j = 1, 2, \dots, p$ , are linearly independent, then so are the  $\eta_j, j = 1, 2, \dots, p$ , since

$$\sum_{j=1}^p c_j\eta_j = 0$$

implies that

$$0 = F \left( \sum_{j=1}^p c_j\eta_j \right) = \sum_{j=1}^p \nu_j c_j \xi_j = 0,$$

and hence the  $c_j$ ,  $j = 1, 2, \dots, p$ , vanish because of the linear independence of the  $\xi_j$ ,  $j = 1, 2, \dots, p$ . Evidently, there can be no more than  $p$  eigenvectors of  $GF$  associated with nonzero eigenvalues; otherwise, there would be more than  $p$  linearly independent eigenvectors of  $FG$  associated with the nonzero eigenvalues. Thus, we have

$$(3.12) \quad p \leq \min\{r, s\}.$$

Since  $\nu_j < 0$ ,  $j = 1, 2, \dots, p$ , if we let

$$(3.13) \quad \mu_j = i|\nu_j|^{1/2}, \quad x_j = \mu_j \xi_j, \quad y_j = \eta_j, \quad v_j = \begin{pmatrix} x_j \\ y_j \end{pmatrix}, \quad j = 1, 2, \dots, p,$$

where  $i^2 = -1$ , then using (3.9), (3.11) and (3.13), we have

$$(3.14) \quad Bv_j = \begin{pmatrix} Fy_j \\ Gx_j \end{pmatrix} = \begin{pmatrix} \nu_j \xi_j \\ \mu_j \eta_j \end{pmatrix} = \begin{pmatrix} \mu_j \mu_j \xi_j \\ \mu_j \eta_j \end{pmatrix} = \mu_j v_j, \quad j = 1, 2, \dots, p.$$

Notice that  $\mu_j$ ,  $j = 1, 2, \dots, p$ , have positive imaginary parts.

Let us now define for  $j = p + 1, p + 2, \dots, 2p$

$$(3.15) \quad x_j = x_{j-p}, \quad y_j = -y_{j-p}, \quad v_j = \begin{pmatrix} x_j \\ y_j \end{pmatrix}, \quad \mu_j = -\mu_{j-p}.$$

Evidently, we have

$$(3.16) \quad Bv_j = \mu_j v_j, \quad j = p + 1, p + 2, \dots, 2p.$$

If we let  $FGx = 0$ , where  $x \neq 0$ , then by (3.4) we have

$$D_1^{-1}HD_2^{-1}Kx = 0, \quad \text{or} \quad HD_2^{-1}Kx = 0.$$

Thus, we have

$$(3.17) \quad -HD_2^{-1/2}D_2^{-1/2}H^T x = 0 \quad \text{or} \quad (D_2^{-1/2}H^T x)^*(D_2^{-1/2}H^T x) = 0.$$

Here  $*$  stands for the conjugate transpose of a matrix. Hence from (3.17) we have  $D_2^{-1/2}H^T x = 0$ , or  $D_2^{-1}H^T x = Gx = 0$ . Therefore, we have that if  $FGx = 0$ , where  $x \neq 0$ , then

$$(3.18) \quad B \begin{pmatrix} x \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ Gx \end{pmatrix} = 0.$$

Thus, if the eigenvectors of  $FG$  associated with the eigenvalue zero are  $x_{2p+1}, x_{2p+2}, \dots, x_{p+r}$ , then the vectors

$$(3.19) \quad v_j = \begin{pmatrix} x_j \\ 0 \end{pmatrix}, \quad j = 2p + 1, 2p + 2, \dots, p + r,$$

are eigenvectors of  $B$  associated with the eigenvalue zero. Similarly, if the eigenvectors of  $GF$  associated with the eigenvalue zero are  $y_{p+r+1}, y_{p+r+2}, \dots, y_{s+r}$ , then the vectors

$$(3.20) \quad v_j = \begin{pmatrix} 0 \\ y_j \end{pmatrix}, \quad j = p + r + 1, p + r + 2, \dots, n = r + s,$$

are eigenvectors of  $B$  associated with the eigenvalue zero.

We have thus constructed a basis of eigenvectors for  $B$ ,

$$(3.21) \quad v_j = \begin{pmatrix} x_j \\ y_j \end{pmatrix}, \quad j = 1, 2, \dots, n,$$

and moreover, we have by (3.14) and (3.16)

$$(3.22) \quad Gx_j = \mu_j y_j, \quad Fy_j = \mu_j x_j, \quad j = 1, 2, \dots, n.$$

We also have

$$\begin{aligned} \text{for } \mu_j &= i|\nu_j|^{1/2}, \quad v_j = \begin{pmatrix} x_j \\ y_j \end{pmatrix}, \quad j = 1, 2, \dots, p; \\ \text{for } \mu_j &= -i|\nu_j|^{1/2}, \quad v_j = \begin{pmatrix} x_{j-p} \\ -y_{j-p} \end{pmatrix}, \quad j = p+1, \dots, 2p; \\ \text{for } \mu_j &= 0, \quad v_j = \begin{pmatrix} x_j \\ 0 \end{pmatrix}, \quad j = 2p+1, \dots, r+p; \\ \text{for } \mu_j &= 0, \quad v_j = \begin{pmatrix} 0 \\ y_j \end{pmatrix}, \quad j = p+r+1, \dots, n. \end{aligned}$$

**4. The Principal Vectors of  $\mathcal{L}_\omega$ .** We now seek the eigenvectors and principal vectors of  $\mathcal{L}_\omega$  for  $\omega \neq 0$ . Because  $A$  has the form of (3.1), from (3.3) we have

$$(4.1) \quad L = \begin{bmatrix} 0 & 0 \\ G & 0 \end{bmatrix}, \quad U = \begin{bmatrix} 0 & F \\ 0 & 0 \end{bmatrix}.$$

Thus we have

$$\begin{aligned} \mathcal{L}_\omega &= (I - \omega L)^{-1}((1 - \omega)I + \omega U) \\ &= \begin{bmatrix} I_1 & 0 \\ -\omega G & I_2 \end{bmatrix}^{-1} \begin{bmatrix} (1 - \omega)I_1 & \omega F \\ 0 & (1 - \omega)I_2 \end{bmatrix} \\ (4.2) \quad &= \begin{bmatrix} I_1 & 0 \\ \omega G & I_2 \end{bmatrix} \begin{bmatrix} (1 - \omega)I_1 & \omega F \\ 0 & (1 - \omega)I_2 \end{bmatrix} \\ &= \begin{bmatrix} (1 - \omega)I_1 & \omega F \\ \omega(1 - \omega)G & \omega^2 GF + (1 - \omega)I_2 \end{bmatrix}, \end{aligned}$$

where  $I_1$  and  $I_2$  are identity matrices of the same sizes as  $D_1$  and  $D_2$ , respectively. For each nonzero eigenvalue  $\mu$  of  $B$ , let  $\lambda_+^{1/2}$  and  $\lambda_-^{1/2}$  be the roots of

$$(4.3) \quad \lambda + \omega - 1 = \omega\mu\lambda^{1/2}.$$

Since

$$B \begin{pmatrix} x \\ y \end{pmatrix} = \mu \begin{pmatrix} x \\ y \end{pmatrix},$$

the vectors

$$(4.4) \quad w = \begin{pmatrix} x \\ \lambda_+^{1/2} y \end{pmatrix}, \quad z = \begin{pmatrix} x \\ \lambda_-^{1/2} y \end{pmatrix}$$

are the eigenvectors of  $\mathcal{L}_\omega$ , since by (4.2), (3.22) and (4.3) we have

$$(4.5) \quad \begin{aligned} \mathcal{L}_\omega \begin{pmatrix} x \\ \lambda_+^{1/2} y \end{pmatrix} &= \begin{pmatrix} (1 - \omega + \omega\mu\lambda_+^{1/2})x \\ [\omega\mu(1 - \omega + \omega\mu\lambda_+^{1/2}) + (1 - \omega)\lambda_+^{1/2}]y \end{pmatrix} \\ &= \lambda_+ w \end{aligned}$$

and

$$(4.6) \quad \mathcal{L}_\omega z = \lambda_- z.$$

If we let

$$(4.7) \quad v = \begin{pmatrix} x \\ y \end{pmatrix}, \quad \hat{v} = \begin{pmatrix} x \\ -y \end{pmatrix},$$

then we have

$$(4.8) \quad \begin{aligned} w &= \frac{1}{2}(v + \hat{v}) + \frac{1}{2}\lambda_+^{1/2}(v - \hat{v}) = \frac{1}{2}(1 + \lambda_+^{1/2})v + \frac{1}{2}(1 - \lambda_+^{1/2})\hat{v}, \\ z &= \frac{1}{2}(v + \hat{v}) + \frac{1}{2}\lambda_-^{1/2}(v - \hat{v}) = \frac{1}{2}(1 + \lambda_-^{1/2})v + \frac{1}{2}(1 - \lambda_-^{1/2})\hat{v}. \end{aligned}$$

If  $\lambda_-^{1/2} \neq \lambda_+^{1/2}$ , then  $w$  and  $z$  are linearly independent. But for  $\omega\mu \neq 0$ , the discriminant  $\omega^2\mu^2 - 4(\omega - 1)$  of (4.3) does not vanish unless

$$(4.9) \quad \omega^2|\mu|^2 + 4(\omega - 1) = 0.$$

On the other hand, if (4.9) holds and if  $\omega\mu \neq 0$ , then  $\lambda_+^{1/2} = \lambda_-^{1/2} = \lambda^{1/2} = \omega\mu/2 \neq 0$ , and  $w$  and  $z$  are not linearly independent. Notice that in this case,

$$(4.10) \quad \lambda_+ = \lambda_- = \lambda = \omega^2\mu^2/4 = \omega - 1.$$

If we let

$$(4.11) \quad \hat{z} = \frac{1}{2} \cdot \frac{1}{\lambda^{1/2}} \begin{pmatrix} 0 \\ y \end{pmatrix},$$

then we have

$$(4.12) \quad \begin{aligned} \mathcal{L}_\omega \hat{z} &= \begin{pmatrix} (1-\omega)I_1 & \omega F \\ \omega(1-\omega)G & \omega^2 GF + (1-\omega)I_2 \end{pmatrix} \cdot \frac{1}{2\lambda^{1/2}} \begin{pmatrix} 0 \\ y \end{pmatrix} \\ &= \frac{1}{2\lambda^{1/2}} \begin{pmatrix} \omega F y \\ \omega^2 GF y + (1-\omega)y \end{pmatrix} = \frac{1}{2\lambda^{1/2}} \begin{pmatrix} \omega\mu x \\ \omega\mu^2 y + (1-\omega)y \end{pmatrix} \\ &= \frac{1}{2\lambda^{1/2}} \begin{pmatrix} \omega\mu x \\ [\omega^2\mu^2 + 1 - \omega - \lambda + \lambda]y \end{pmatrix} \\ &= \frac{1}{2\lambda^{1/2}} \begin{pmatrix} 2 \cdot \frac{\omega\mu}{2} x \\ [\omega^2\mu^2 + 1 - \omega - (\omega - 1)]y \end{pmatrix} + \lambda \begin{pmatrix} 0 \\ \frac{1}{2\lambda^{1/2}} y \end{pmatrix} \\ &= \frac{1}{2\lambda^{1/2}} \begin{pmatrix} 2 \cdot \lambda^{1/2} x \\ [4(\omega - 1) + 2(1 - \omega)]y \end{pmatrix} + \lambda \hat{z} = \frac{1}{2\lambda^{1/2}} \begin{pmatrix} 2\lambda^{1/2} x \\ 2(\omega - 1)y \end{pmatrix} + \lambda \hat{z} \\ &= \begin{pmatrix} x \\ \lambda^{1/2} y \end{pmatrix} + \lambda \hat{z} = w + \lambda \hat{z}. \end{aligned}$$

Hence,  $\hat{z}$  is a principal vector of grade 2. Moreover, we have

$$(4.13) \quad w = \frac{1}{2}(1 + \lambda^{1/2})v + \frac{1}{2}(1 - \lambda^{1/2})\hat{v}, \quad \hat{z} = \frac{1}{4} \cdot \frac{1}{\lambda^{1/2}}v - \frac{1}{4} \cdot \frac{1}{\lambda^{1/2}}\hat{v}.$$

Thus  $w$  and  $\hat{z}$  are linearly independent.

If we let, for  $\omega\mu \neq 0$ ,

$$(4.14) \quad w_j = \begin{pmatrix} x_j \\ (\lambda_j)_+^{1/2} y_j \end{pmatrix}, \quad j = 1, 2, \dots, p;$$

$$w_{j+p} = \begin{cases} \begin{pmatrix} x_j \\ (\lambda_j)_-^{1/2} y_j \end{pmatrix}, & j = 1, 2, \dots, p, \text{ if } \omega^2|\mu_j|^2 + 4(\omega - 1) \neq 0, \\ \frac{1}{2} \cdot \frac{1}{(\lambda_j)_+^{1/2}} \begin{pmatrix} 0 \\ y_j \end{pmatrix}, & j = 1, 2, \dots, p, \text{ if } \omega^2|\mu_j|^2 + 4(\omega - 1) = 0; \\ w_j = v_j, & j = 2p + 1, 2p + 2, \dots, n, \end{cases}$$

then we can easily prove that  $w_j$ ,  $j = 1, 2, \dots, n$ , are linearly independent and hence form a basis of the  $n$ -dimensional complex vector space  $\mathbf{C}^n$ . Therefore, the matrix whose columns are the  $w_j$  reduces  $\mathcal{L}_\omega$  to Jordan canonical form. Thus we have proved:

**THEOREM 1.** *If the matrix  $A$  has the form of (3.1) and  $i\bar{\mu} = i\rho(B)$  is an eigenvalue of multiplicity  $q$  of  $B$  of (3.3), then the Jordan canonical form of  $\mathcal{L}_{\omega_b}$  has  $n-2q$  ( $1 \times 1$ ) sub-Jordan blocks and  $q$  ( $2 \times 2$ ) sub-Jordan blocks which correspond to the eigenvalue  $\omega_b - 1$ .*

Notice that if  $q = 1$ , then the Jordan canonical form of  $\mathcal{L}_{\omega_b}$  has one nondiagonal element.

From Theorem 3.1 of [5, p. 65] and Theorem 3-7.1 of [7, p. 85] we have

$$(4.15) \quad \|\mathcal{L}_{\omega_b}^m\| \sim J(\mathcal{L}_{\omega_b}) \cdot m \cdot \rho(\mathcal{L}_{\omega_b})^{m-1}.$$

Here,  $\|A\|$  is the spectral norm of the matrix  $A$  and  $J(\mathcal{L}_{\omega_b})$  is the Jordan condition number of the matrix  $\mathcal{L}_{\omega_b}$ , defined by Young [7, p. 85] and given by

$$(4.16) \quad J(\mathcal{L}_{\omega_b}) = \inf_{V \in S_1} \kappa(V),$$

where  $\kappa(V)$  is the spectral condition number of the matrix  $V$  and  $S_1$  the set of all matrices such that

$$(4.17) \quad V^{-1} \mathcal{L}_{\omega_b} V = J.$$

Here,  $J$  is the Jordan canonical form of the matrix  $\mathcal{L}_{\omega_b}$ .

**5. Determination of  $\|\mathcal{L}_\omega\|_{D^{1/2}}$ .** Let

$$(5.1) \quad \hat{\mathcal{L}}_\omega = D^{1/2} \mathcal{L}_\omega D^{-1/2};$$

then from (4.2) and (3.6) we have

$$(5.2) \quad \hat{\mathcal{L}}_\omega = \begin{pmatrix} (1-\omega)I_1 & \omega D_1^{1/2} F D_2^{-1/2} \\ \omega(1-\omega)D_2^{1/2} G D_1^{-1/2} & \omega^2 D_2^{1/2} G F D_2^{-1/2} + (1-\omega)I_2 \end{pmatrix}.$$

If we let

$$(5.3) \quad \begin{aligned} \hat{F} &= D_1^{1/2} F D_2^{-1/2} = D_1^{-1/2} H D_2^{-1/2}, \\ \hat{G} &= D_2^{1/2} G D_1^{-1/2} = D_2^{-1/2} K D_1^{-1/2}, \end{aligned}$$

then

$$(5.4) \quad \hat{G}^T = -\hat{F}$$

and

$$(5.5) \quad \hat{\mathcal{L}}_\omega = \begin{pmatrix} (1-\omega)I_1 & \omega \hat{F} \\ \omega(1-\omega)\hat{G} & \omega^2 \hat{G} \hat{F} + (1-\omega)I_2 \end{pmatrix}.$$

Hence,  $\hat{\mathcal{L}}_\omega$  is the SOR iterative matrix corresponding to the matrix

$$(5.6) \quad \hat{A} = D^{1/2} A D^{-1/2} = \begin{pmatrix} I_1 & -\hat{F} \\ -\hat{G} & I_2 \end{pmatrix}$$

with the associated Jacobi iterative matrix

$$(5.7) \quad \hat{B} = \begin{pmatrix} 0 & \hat{F} \\ \hat{G} & 0 \end{pmatrix} = D^{1/2} B D^{-1/2}.$$

Therefore,  $\rho(\hat{B}) = \rho(B)$ , and  $\omega_b$  is the same for  $\hat{A}$  as for  $A$ . Moreover, if we let  $\mathcal{L}_\omega[A]$  stand for the SOR iterative matrix associated with the matrix  $A$  and  $D[A]$  for the diagonal block of the matrix  $A$ , then we have

$$(5.8) \quad \|\mathcal{L}_\omega^m[A]\|_{D[A]^{1/2}} = \|\hat{\mathcal{L}}_\omega^m[A]\| = \|\mathcal{L}_\omega^m[\hat{A}]\| = \|\mathcal{L}_\omega^m[\hat{A}]\|_{D[\hat{A}]^{1/2}}.$$

Thus, it is sufficient to assume  $A$  of (3.1) with  $D_1 = I_1$  and  $D_2 = I_2$ . Otherwise, we consider  $\hat{A}$  (5.6). Notice that when  $D_1 = I_1$  and  $D_2 = I_2$  then  $F = H$ ,  $G = K$ , and  $F^T = -G$ .

Since

$$(5.9) \quad \|\mathcal{L}_\omega^T\|_{D^{1/2}} = \|\mathcal{L}_\omega\|_{D^{1/2}} = [\rho(\mathcal{L}_\omega \mathcal{L}_\omega^*)]^{1/2},$$

according to the expression (4.2) for  $\mathcal{L}_\omega$  we first study the eigenvalues of products of matrices of the form

$$(5.10) \quad \begin{pmatrix} a_{11}(FG) & a_{12}(FG)F \\ a_{21}(GF)G & a_{22}(GF) \end{pmatrix},$$

where  $a_{11}$  and  $a_{12}$  are polynomials in  $FG$  and  $a_{21}$  and  $a_{22}$  polynomials in  $GF$ . By an analogy with Theorem 7-2.1 of Young [7, p. 239] we have:

**THEOREM 2.** *If  $B$  is a matrix of the form (3.3), then*

(a) *The matrix*

$$(5.11) \quad Q = \begin{pmatrix} a_{11}(FG) & a_{12}(FG)F \\ a_{21}(GF)G & a_{22}(GF) \end{pmatrix}$$

*is nonsingular if*

$$(5.12) \quad \tau(B^2) = a_{11}(B^2)a_{22}(B^2) - a_{21}(B^2)a_{12}(B^2)B^2$$

*in nonsingular. Moreover,  $\tau(B^2)$  is nonsingular if and only if for each eigenvalue  $\mu$  of  $B$  the matrix*

$$(5.13) \quad R(\mu) = \begin{pmatrix} a_{11}(\mu^2) & a_{12}(\mu^2)\mu \\ a_{21}(\mu^2)\mu & a_{22}(\mu^2) \end{pmatrix}$$

*is nonsingular.*

(b) *Let*

$$(5.14) \quad G_m = \prod_{k=m}^1 \begin{pmatrix} a_{11}^{(k)}(FG) & a_{12}^{(k)}(FG)F \\ a_{21}^{(k)}(GF)G & a_{22}^{(k)}(GF) \end{pmatrix}^{\nu_k},$$

*where for each  $k$ ,  $\nu_k = \pm 1$ . It is assumed that for any  $k$  the matrix*

$$(5.15) \quad \tau^{(k)}(B^2) = a_{11}^{(k)}(B^2)a_{22}^{(k)}(B^2) - a_{21}^{(k)}(B^2)a_{12}^{(k)}(B^2)B^2$$

*is nonsingular for  $\nu_k = -1$ . For each eigenvalue  $\mu$  of  $B$ , let*

$$(5.16) \quad M_m(\mu) = \prod_{k=m}^1 \begin{pmatrix} a_{11}^{(k)}(\mu^2) & a_{12}^{(k)}(\mu^2)\mu \\ a_{21}^{(k)}(\mu^2)\mu & a_{22}^{(k)}(\mu^2) \end{pmatrix}^{\nu_k}.$$

*If  $\mu$  is a nonzero eigenvalue of  $B$  and if  $\lambda$  is an eigenvalue of  $M_m(\mu)$ , then  $\lambda$  is an eigenvalue of  $G_m$ . If  $\mu = 0$  is an eigenvalue of  $B$ , then at least one of the eigenvalues of  $M_m(0)$  is an eigenvalue of  $G_m$ .*

(c) If  $\lambda$  is an eigenvalue of  $G_m$ , then there exists an eigenvalue  $\mu$  of  $B$  such that  $\lambda$  is an eigenvalue of  $M_m(\mu)$ .

Notice that although the matrix  $B$  considered here and the matrix  $B$  considered in Theorem 7-2.1 of Young [7, p. 239] are not the same type of matrices—the former is similar to a skew-symmetric matrix and the latter a symmetric matrix—the statement of these two theorems are the same and the proofs are also the same. Hence the proof of Theorem 2 is omitted.

From (4.2) we have

$$(5.17) \quad \begin{aligned} \mathcal{L}_\omega^* &= \mathcal{L}_\omega^T = \begin{pmatrix} (1-\omega)I_1 & 0 \\ -\omega G & (1-\omega)I_2 \end{pmatrix} \begin{pmatrix} I_1 & -\omega F \\ 0 & I_2 \end{pmatrix} \\ &= \begin{pmatrix} (1-\omega)I_1 & -\omega(1-\omega)F \\ -\omega G & \omega^2 GF + (1-\omega)I_2 \end{pmatrix}. \end{aligned}$$

Thus, from (4.2) and (5.7),  $\mathcal{L}_\omega$  and  $\mathcal{L}_\omega^*$  have the required form for the applicability of Theorem 2. From Theorem 2 we know that the eigenvalues of  $\mathcal{L}_\omega \mathcal{L}_\omega^*$  are the same as the eigenvalues of  $M(\omega, \mu)M^*(\omega, \mu)$ , where

$$(5.18) \quad \begin{aligned} M(\omega, \mu) &= \begin{pmatrix} 1 & 0 \\ \omega\mu & 1 \end{pmatrix} \begin{pmatrix} 1-\omega & \omega\mu \\ 0 & 1-\omega \end{pmatrix} \\ &= \begin{pmatrix} 1-\omega & \omega\mu \\ (1-\omega)\omega\mu & \omega^2\mu^2 + 1-\omega \end{pmatrix}. \end{aligned}$$

If we notice  $\bar{\mu} = -\mu$  (here  $\mu$  is purely imaginary), and if we let

$$(5.19) \quad M(\omega, \mu)M^*(\omega, \mu) = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix},$$

then

$$\begin{aligned} m_{11} &= (1-\omega)^2 - \omega^2\mu^2, \\ m_{12} &= \omega^3\mu^3 + \omega\mu(1-\omega) - (1-\omega)^2 \cdot \omega\mu = \mu\omega^2[1-\omega + \omega\mu^2], \\ m_{21} &= -m_{12} = -\mu\omega^2[1-\omega + \omega\mu^2], \\ m_{22} &= \omega^4\mu^4 + (1-\omega)^2 + 2\omega^2\mu^2(1-\omega) - \omega^2\mu^2(1-\omega)^2 \\ &= \omega^4\mu^4 + (1-\omega)^2 + \omega^2\mu^2(1-\omega)(1+\omega). \end{aligned}$$

Since

$$(5.20) \quad \begin{aligned} m_{11} + m_{22} &= 2(1-\omega)^2 + \omega^4\mu^4 - \omega^4\mu^2, \\ m_{12}m_{21} &= -\mu^2\omega^4[(1-\omega)^2 + 2\omega\mu^2(1-\omega) + \omega^2\mu^4], \\ m_{22}m_{11} &= -\omega^6\mu^6 - 2\omega^4\mu^4 \cdot \omega \cdot (1-\omega) - \omega^2\mu^2(1-\omega)^2 \cdot \omega^2 + (1-\omega)^4 \\ &= -\omega^4\mu^2[\omega^2\mu^4 + 2\omega\mu^2(1-\omega) + (1-\omega)^2] + (1-\omega)^4 \\ &= m_{21}m_{12} + (1-\omega)^4, \end{aligned}$$

we have

$$(5.21) \quad m_{22}m_{11} - m_{21}m_{12} = (1-\omega)^4.$$

Thus, if we let

$$(5.22) \quad \det(\lambda I - M(\omega, \mu)M^*(\omega, \mu)) = \lambda^2 - T(\mu^2)\lambda + c = 0,$$

then

$$(5.23) \quad T(\mu^2) = 2(1 - \omega)^2 + \omega^4 \mu^4 - \omega^4 \mu^2, \quad c = (1 - \omega)^4.$$

Notice that  $\mu^2 \leq 0$ , so that  $T(\mu^2)$  is an increasing function of  $|\mu|$ . Therefore, by Lemma 6-2.9 of Young [7, p. 186], it follows that for a given  $\omega$ , the largest value of the root radius of (5.22) is assumed for  $\mu = i\rho(B)$  (or  $\mu = -i\rho(B)$ ). From (5.22) and (5.23) we know that if  $\lambda$  satisfies (5.22), then  $t = \lambda^{1/2}$  satisfies

$$(5.24) \quad t^2 - (\omega - 1)^2 = \omega^2 |\mu| (1 + |\mu|^2)^{1/2} t.$$

Note  $|\mu| = \rho(B)$ , and if we let  $\bar{\mu} = \rho(B)$  and

$$(5.25) \quad d = \bar{\mu} (1 + \bar{\mu}^2)^{1/2},$$

then by Lemma 6-2.1 of Young [7, p. 171] the root radius of (5.24) is less than unity if and only if we have

$$|\omega - 1| < 1 \quad \text{and} \quad \omega^2 d < 1 - (\omega - 1)^2 = \omega(2 - \omega),$$

or, equivalently,

$$(5.26) \quad 0 < \omega < \min \left\{ 2, \frac{2}{1 + d} \right\} = \frac{2}{1 + d} \quad (\bar{\mu} > 0).$$

Thus we have proved

**THEOREM 3.** *If  $A$  has the form (3.1) with  $D_1$  and  $D_2$  symmetric and positive definite and  $H$  and  $K$  satisfying (3.2), then  $\|\mathcal{L}_\omega\|_{D^{1/2}} < 1$  if and only if  $\omega$  satisfies (5.26). Moreover, we have*

$$(5.27) \quad \|\mathcal{L}_\omega\|_{D^{1/2}} = \frac{\omega^2 d + \sqrt{\omega^4 d^2 + 4(1 - \omega)^2}}{2}.$$

We now determine the minimum value of  $\|\mathcal{L}_\omega\|_{D^{1/2}}$ . If we let

$$f(\omega) = \omega^2 d + \sqrt{\omega^4 d^2 + 4(1 - \omega)^2},$$

then the derivative of  $f(\omega)$  is given by

$$f'(\omega) = 2\omega d + [4\omega^3 d^2 + 8(\omega - 1)] / 2\sqrt{\omega^4 d^2 + 4(1 - \omega)^2}.$$

Assume  $f'(\omega) = 0$ ; then

$$(5.28) \quad -\omega d \sqrt{\omega^4 d^2 + 4(1 - \omega)^2} = \omega^3 d^2 + 2(\omega - 1).$$

Notice that (5.28) means

$$(5.29) \quad g(\omega) = \omega^3 d^2 + 2(\omega - 1) < 0.$$

By Descartes' rule we know that  $g(\omega)$  has only one positive root  $\omega_u$ . Thus, if  $\omega \in (0, \omega_u)$ , then (5.29) holds. Moreover, if  $\omega \geq \omega_u$ , we have

$$(5.30) \quad f'(\omega) > 0.$$

If  $0 < \omega < \omega_u$ , and from (5.28), we have

$$(5.31) \quad \omega^2 d^2 + \omega - 1 = 0.$$

Evidently, the positive root  $\omega_+$  of (5.31) is given by

$$(5.32) \quad \omega_+ = [-1 + \sqrt{1 + 4d^2}] / 2d^2 = \frac{2}{1 + \sqrt{1 + 4d^2}}.$$

One can examine

$$(5.33) \quad \omega_+ < \min \left\{ 2, \frac{2}{1+d} \right\} = \frac{2}{1+d}.$$

Thus, we obtain

$$f'(\omega) \begin{cases} < 0 & \text{if } 0 < \omega < \omega_+, \\ = 0 & \text{if } \omega = \omega_+, \\ > 0 & \text{if } \omega > \omega_+, \end{cases}$$

because we have

$$(5.34) \quad \omega_+ < \omega_u.$$

In fact, if  $\omega_+ \geq \omega_u$ , then from (5.31) and (5.32) we have  $\omega^2 d^2 + \omega - 1 < 0$  for  $0 < \omega < \omega_u$ . Thus we can prove  $f'(\omega) < 0$  for  $0 < \omega < \omega_u$ . Owing to the continuity property of  $f'(\omega)$ , we have  $f'(\omega_u) \leq 0$ , which contradicts (5.30). Hence (5.34) holds. We have now proved the following theorem.

**THEOREM 4.** *Under the assumptions of Theorem 3 we have*

$$\|\mathcal{L}_{\omega_+}\|_{D^{1/2}} < \|\mathcal{L}_\omega\|_{D^{1/2}} \quad \text{for } \omega \neq \omega_+.$$

Here,  $\omega_+$  is given by (5.32).

It is important to note that from (1.9), (5.25), and (5.26) we have

$$\begin{aligned} \omega_b &< \frac{2}{1+d} & \text{if } \bar{\mu} = \rho(B) < 1, \\ \omega_b &\geq \frac{2}{1+d} & \text{if } \bar{\mu} \geq 1. \end{aligned}$$

Thus, when  $\bar{\mu} < 1$ , we also have  $\|\mathcal{L}_{\omega_b}\|_{D^{1/2}} < 1$ .

In fact we have proved

**THEOREM 5.** *Under the assumptions of Theorem 3 we have*

$$\|\mathcal{L}_{\omega_b}\|_{D^{1/2}} < 1 \quad \text{if and only if } \bar{\mu} = \rho(B) < 1.$$

But when  $\bar{\mu} \geq 1$ , we have  $\|\mathcal{L}_{\omega_b}\|_{D^{1/2}} \geq 1$ . However, in the next section, we will prove that for any  $\bar{\mu} > 0$ ,  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} < 1$  if  $m$  is large enough, and that  $\lim_{m \rightarrow \infty} \|\mathcal{L}_{\omega_b}^m\| = 0$ .

**6. Determination of  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$ .** In this section we continue with the theory of Young [7, Chapter 7] to investigate  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$ . From the discussion of the last section it is sufficient to consider  $A$  (3.1) with  $D_1 = I_1$  and  $D_2 = I_2$ . Since the eigenvalues of  $\mathcal{L}_{\omega_b}^m (\mathcal{L}_{\omega_b}^m)^*$  are the same as those of  $M^m(\omega_b, \mu) [M^m(\omega_b, \mu)]^*$ , where  $M(\omega_b, \mu)$  is given by (5.18), we first develop an expression for  $M^m(\omega, \mu)$ . If we define the polynomials  $S_0(\mu), S_1(\mu), \dots$  by the recursion formula

$$(6.1) \quad \begin{aligned} S_k(\mu) &= \omega\mu S_{k-1}(\mu) + (1-\omega)S_{k-2}(\mu), & k \geq 2, \\ S_0(\mu) &= 1, & S_1(\mu) = \omega\mu, \end{aligned}$$

then by a result of Young [7, p. 248] we have

$$(6.2) \quad M(\omega, \mu)^m = \begin{pmatrix} (1-\omega)S_{2m-2} & S_{2m-1} \\ (1-\omega)S_{2m-1} & S_{2m} \end{pmatrix}.$$

Notice that  $\mu$  is purely imaginary. By (6.1) one can see that  $S_{2k}(\mu)$  are real and  $S_{2k+1}$  purely imaginary. Also from the result of Young [7, p. 249, Eq. (4.7)] we have

$$(6.3) \quad S_k(\mu) = \sum_{j=0}^k \alpha_1^{k-j} \alpha_2^j = \begin{cases} \frac{\alpha_1^{k+1} - \alpha_2^{k+1}}{\alpha_1 - \alpha_2} & \text{if } \alpha_1 \neq \alpha_2, \\ (k+1)\alpha^k & \text{if } \alpha_1 = \alpha_2. \end{cases}$$

Here,  $\alpha_1$  and  $\alpha_2$  are the solution of the quadratic equation

$$(6.4) \quad \alpha^2 - \omega\mu\alpha + \omega - 1 = 0.$$

Now we prove that if  $\omega = \omega_b$  and  $r = (1 - \omega_b)$  then

$$(6.5) \quad S_k(i\bar{\mu}) = S_k(i\rho(B)) = (i)^k (r^{1/2})^k \cdot (k+1),$$

$$(6.6) \quad \max_{\substack{\mu=i\beta \\ -\bar{\mu} \leq \beta \leq \bar{\mu}}} |S_k(\mu)| = |S_k(i\bar{\mu})| = (k+1)(r^{1/2})^k.$$

Let  $\mu = i\beta$ ; then the roots  $\alpha_1$  and  $\alpha_2$  of (6.4) are given by

$$\alpha_{1,2} = [i\beta\omega_b \pm \sqrt{-\beta^2\omega_b^2 - 4(\omega_b - 1)}]/2 = i[\beta\omega_b \pm \sqrt{\beta^2\omega_b^2 + 4(\omega_b - 1)}]/2.$$

By (1.9) we have

$$\bar{\mu}^2\omega_b^2 + 4(\omega_b - 1) = 0.$$

Thus, if  $\beta = \bar{\mu}$ , then  $\alpha_1 = \alpha_2 = i\bar{\mu}\omega_b/2 = ir^{1/2}$ . Hence (6.5) follows from (6.3). If  $|\beta| \leq \bar{\mu}$ , then  $\beta^2\omega_b^2 + 4(\omega_b - 1) \leq 0$ . Therefore,  $|\alpha_1| = |\alpha_2| = (1 - \omega_b)^{1/2} = r^{1/2}$ . Again by (6.3), (6.6) follows.

From (6.2) we have

$$(6.7) \quad \begin{aligned} &M^m(\omega, \mu)[M^m(\omega, \mu)]^* \\ &= \begin{pmatrix} (1-\omega)S_{2m-2} & S_{2m-1} \\ (1-\omega)S_{2m-1} & S_{2m} \end{pmatrix} \begin{pmatrix} (1-\omega)S_{2m-2} & -(1-\omega)S_{2m-1} \\ -S_{2m-1} & S_{2m} \end{pmatrix} \\ &= \begin{pmatrix} (1-\omega)^2 S_{2m-2}^2 - S_{2m-1}^2 & S_{2m}S_{2m-1} - (1-\omega)^2 S_{2m-1}S_{2m-2} \\ (1-\omega)^2 S_{2m-1}S_{2m-2} - S_{2m}S_{2m-1} & -(1-\omega)^2 S_{2m-1}^2 + S_{2m}^2 \end{pmatrix}. \end{aligned}$$

Evidently, the characteristic equation for  $M^m(\omega_b, \mu)[M^m(\omega_b, \mu)]^*$  is

$$(6.8) \quad \lambda^2 - T_m(\omega_b, \mu)\lambda + \Delta = 0,$$

where

$$(6.9) \quad T_m(\omega_b, \mu) = (1 - \omega_b)^2 S_{2m-2}^2 - S_{2m-1}^2 - (1 - \omega_b)^2 S_{2m-1}^2 + S_{2m}^2$$

and

$$(6.10) \quad \Delta = \det\{M^m(\omega_b, \mu)[M^m(\omega_b, \mu)]^*\} = r^{4m} = (1 - \omega_b)^{4m}$$

by (5.18). Since  $[T_m(\omega_b, \mu)]^2 - 4\Delta \geq 0$ , because the eigenvalues of the Hermitian matrix  $M^m(\omega_b, \mu)[M^m(\omega_b, \mu)]^*$  are real, it follows that for fixed  $\Delta$  the modulus of the root of (6.8) is maximized when  $|T_m(\omega_b, \mu)|$ , considered as a function of  $\mu$ , is maximized. But, by (6.5) and (6.6),  $|T_m(\omega_b, \mu)|$  is maximized when  $\mu = i\bar{\mu}$ , and we have

$$(6.11) \quad \begin{aligned} |T_m(\omega_b, i\bar{\mu})| &= r^2 \cdot (2m - 1)^2 r^{2m-2} + r^{2m-1} \cdot (2m)^2 \\ &\quad + r^2 (2m)^2 r^{2m-1} + (2m + 1)^2 r^{2m} \\ &= 2r^{2m} [1 + 2m^2 (\sqrt{r} + r^{-1/2})^2]. \end{aligned}$$

Thus, from (6.8), (6.10), and (6.11) we have

$$(6.12) \quad (\lambda - r^{2m})^2 = 4m^2(r^{-1/2} + r^{1/2})^2 r^{2m} \cdot \lambda$$

and

$$(6.13) \quad \lambda - r^{2m} = 2m(r^{-1/2} + r^{1/2})r^m \lambda^{1/2}.$$

Hence we have proved the following

**THEOREM 6.** *Under the assumptions of Theorem 3, we have*

$$(6.14) \quad \begin{aligned} \|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} &= r^m \{m(r^{-1/2} + r^{1/2}) + [m^2(r^{-1/2} + r^{1/2})^2 + 1]^{1/2}\} \\ &= F_1(m), \end{aligned}$$

where

$$(6.15) \quad r = 1 - \omega_b, \quad \omega_b = \frac{2}{1 + \sqrt{1 + \bar{\mu}^2}}, \quad \bar{\mu} = \rho(B).$$

From (6.14) we know that for any  $\bar{\mu} = \rho(B) > 0$

$$\lim_{m \rightarrow \infty} \|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} = \lim_{m \rightarrow \infty} F_1(m) = 0.$$

But, for values of  $r$  close to unity, the function  $F_1(m)$  increases initially before eventually decreasing. For  $r$  close to unity we have

$$(6.16) \quad \begin{aligned} \|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}} &\sim 2mr^m(r^{-1/2} + r^{1/2}) \\ &= 2mr^m(r^{-1}r^{1/2} + r^{-1}r \cdot r^{1/2}) \\ &\sim 4mr^{m-1}. \end{aligned}$$

On the other hand, we have

$$\|\mathcal{L}_{\omega_b}^m\| = \|M^m(\omega_b, i\bar{\mu})\| \sim mJ(M(\omega_b, i\bar{\mu}))r^{m-1}$$

by Theorem 3-7.1 [7, p. 85]. Here,  $J(M(\omega_b, i\bar{\mu}))$  is the Jordan condition number of  $M(\omega_b, i\bar{\mu})$ . But by [7, Theorem 3-8.1, p. 89] we have

$$J(M(\omega_b, i\bar{\mu})) = \omega_b \bar{\mu} + (1 - \omega_b)_{\omega_b} \bar{\mu} = \omega_b \bar{\mu}(1 + 1 - \omega_b) = 2r^{1/2}(1 + r) \sim 4.$$

Hence,

$$(6.17) \quad \|\mathcal{L}_{\omega_b}^m\| \sim 4mr^{m-1}.$$

Therefore,  $\|\mathcal{L}_{\omega_b}^m\|$  behaves like  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$ .

Young [7, p. 255, Eq. (4.50)] has given  $m_0$ , the estimated number of iterations needed to reduce the  $D^{1/2}$ -norm of the error vector to a specified fraction  $\varepsilon$  of the  $D^{1/2}$ -norm of the initial error vector as follows:

$$(6.18) \quad \begin{aligned} m_0 &= \log((2\nu/\varepsilon) \cdot \log(2\nu/\varepsilon)) / \log(1/r), \\ \nu &= \frac{r^{1/2} + r^{-1/2}}{\log(1/r)}. \end{aligned}$$

*Final Remarks.* (a) Since  $\|\mathcal{L}_{\omega_b}\|_{D^{1/2}} \geq 1$  if  $\bar{\mu} \geq 1$ , one should expect that it may be better to use  $\omega = \omega_+$  rather than  $\omega = \omega_b$  in the initial steps. In this direction, an investigation is under way.

(b) By noting Theorem 6 and Theorem 7-4.1 of Young [7] one can find out that  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$  for the nonsymmetric case and  $\|\mathcal{L}_{\omega_b}^m\|_{D^{1/2}}$  for the symmetric and

positive definite case have the same expression in  $m$  and  $r$ . The only difference is that for the former,

$$(6.19) \quad r = 1 - \omega_b = \frac{\rho^2(B)}{(1 + \sqrt{1 + \rho^2(B)})},$$

and for the latter,

$$(6.20) \quad r = \omega'_b - 1 = \frac{\rho^2(B)}{(1 + \sqrt{1 - \rho^2(B)})}.$$

Especially for  $m = 1$ , we have

$$(6.21) \quad \begin{aligned} \|\mathcal{L}_{\omega_b}\|_{D^{1/2}} &= \|\mathcal{L}_{\omega'_b}\|_{D^{1/2}} = r\{(r^{-1/2} + r^{1/2}) + [1 + (r^{-1/2} + r^{1/2})^2]^{1/2}\} \\ &= r^{1/2}\{(1 + r) + [r + (1 + r)^2]^{1/2}\} = F(r). \end{aligned}$$

It is clear that  $F(r)$  is an increasing function of  $r$ . In fact one can prove

LEMMA. *Let  $F(r)$  be given by (6.21). Then*

$$F(r) < 1 \Leftrightarrow 0 \leq r < r_0 = 1/(1 + \sqrt{2})^2.$$

By means of the above lemma we can give another proof of Theorem 5. In fact, it follows from (6.19), (6.21) and the above lemma that  $\|\mathcal{L}_{\omega_b}\|_{D^{1/2}} < 1$  if and only if  $\rho^2(B)/(1 + \sqrt{1 + \rho^2(B)}) < 1/(1 + \sqrt{2})^2$ , or equivalently,  $\rho(B) < 1$ . Thus, Theorem 5 follows. However, we can give a similar result for the symmetric case. By noting (6.20), (6.21) and the above lemma, we have  $\|\mathcal{L}_{\omega'_b}\|_{D^{1/2}} < 1$  if and only if  $\rho^2(B)/(1 + \sqrt{1 - \rho^2(B)}) < 1/(1 + \sqrt{2})^2$ , or equivalently,  $\rho(B) < 1/\sqrt{2}$ . Thus we have proved

COROLLARY. *If  $A$  has the form (1.7) and is symmetric positive definite, then the  $D^{1/2}$ -norm of the corresponding optimum SOR iterative matrix  $\mathcal{L}_{\omega'_b}$  is less than unity if and only if  $\rho(B) < 1/\sqrt{2}$ .*

To our knowledge, the result of the above corollary is new. However, it should be noted that the result can be deduced from Theorem 7-3.1 of Young [7].

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