

## MOTIONS IN LINEARLY CONNECTED TWO-DIMENSIONAL SPACES

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1. **Introduction.** A linearly connected space  $L_n$  is characterized by its coefficients of (asymmetric) connection  $L_{jk}^i$  which transform according to the law [2, p. 3]<sup>1</sup>

$$(1.1) \quad L_{jk}^i(x) \frac{\partial \bar{x}^a}{\partial x^i} = \bar{L}_{bc}^a(\bar{x}) \frac{\partial \bar{x}^b}{\partial x^j} \frac{\partial \bar{x}^c}{\partial x^k} + \frac{\partial^2 \bar{x}^a}{\partial x^j \partial x^k}.$$

An  $L_n$  will admit an infinitesimal motion defined by the infinitesimal transformation

$$(1.2) \quad \bar{x}^i = x^i + \xi^i(x) \delta t$$

if

$$(1.3a) \quad \frac{\partial^2 \xi^i}{\partial x^j \partial x^k} + \frac{\partial \Gamma_{jk}^i}{\partial x^h} \xi^h + \Gamma_{hk}^i \frac{\partial \xi^h}{\partial x^j} + \Gamma_{jh}^i \frac{\partial \xi^h}{\partial x^k} - \Gamma_{jk}^h \frac{\partial \xi^i}{\partial x^h} = 0,$$

$$(1.3b) \quad \frac{\partial \Omega_{jk}^i}{\partial x^h} \xi^h + \Omega_{hk}^i \frac{\partial \xi^h}{\partial x^j} + \Omega_{jh}^i \frac{\partial \xi^h}{\partial x^k} - \Omega_{jk}^h \frac{\partial \xi^i}{\partial x^h} = 0,$$

where

$$(1.4) \quad L_{jk}^i = \Gamma_{jk}^i + \Omega_{jk}^i,$$

and  $\Gamma_{jk}^i, \Omega_{jk}^i$  are the symmetric and skew-symmetric parts respectively of  $L_{jk}^i$  [1, p. 231].

In case  $\Omega_{jk}^i = 0$ , the  $L_n$  reduces to a symmetrically connected space  $A_n$  in which (1.3a) defines the infinitesimal affine collineations [2, p. 125], so that motions in an  $L_n$  can be considered as generalizations of such collineations of an  $A_n$  [3].

In a previous paper [3] we obtained all (symmetric) two-dimensional spaces  $A_2$  admitting real continuous groups of affine collineations, these being obtained by solving (1.3a) for the  $\Gamma_{jk}^i$ , the  $\xi^i$  being known and obtained from Lie's classification [4, pp. 71-73, 379-380] giving all real continuous groups  $G_r$  in two variables. In §2 we carry through a like procedure to obtain all linearly connected  $L_2$  admitting complete groups of motions. This will involve the solution of the

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<sup>1</sup> Numbers in brackets refer to the bibliography at the end of the paper.

combined system (1.3a-b) for the  $\Gamma_{jk}^i$  and  $\Omega_{jk}^i$ . The results of [3] will of course eliminate much of the analysis which would otherwise be necessary. We shall exclude all solutions of the system (1.3) which lead to  $A_2$  spaces ( $\Omega_{jk}^i = 0$ ).

In §3 we consider the special case of spaces of absolute parallelism  $T_n$ , and determine all  $T_2$  admitting motion groups. The conditions on an  $L_n$  to reduce to a  $T_n$  are [1, p. 234]

$$(1.5) \quad L_{jkm}^i \equiv \frac{\partial L_{jm}^i}{\partial x^k} - \frac{\partial L_{jk}^i}{\partial x^m} + L_{jm}^h L_{hk}^i - L_{jk}^h L_{hm}^i = 0,$$

that is, the  $L_n$  is of zero curvature. In this case there exists an ennuple of (absolutely) parallel vector fields  $h_{\alpha|}^i$ , that is,

$$(1.6) \quad h_{\alpha|i}^i \equiv \frac{\partial h_{\alpha|}^i}{\partial x^j} + h_{\alpha|}^m L_{mj}^i = 0 \quad (\alpha = 1, \dots, n).$$

These equations (1.6) are solved for the  $h_{\alpha|}^i$  using for the  $L_{jk}^i$  the values found from the solutions of (1.3). (The integrability conditions of (1.6) are (1.5).)

The problem of motions in a  $T_n$  has been treated by Robertson [5] who obtained the conditions for infinitesimal motions of a  $T_n$  in the form

$$(1.7) \quad \frac{\partial h_{\alpha|}^i}{\partial x^j} \xi^j - h_{\alpha|}^j \frac{\partial \xi^i}{\partial x^j} = \epsilon_{\alpha}^{\beta} h_{\beta|}^i,$$

where  $\epsilon_{\alpha}^{\beta}$  are constants. (See also [6].)

It is shown in §2 that there are 8 possible types of  $L_2$  admitting groups of motions, the maximum number of parameters being four. In §3 it is shown there are 7 types of motions for a  $T_2$ .

**2. Motions in an  $L_2$ .** We proceed to the solution of (1.3) for the case  $n = 2$ . It is easily shown that (1.3b) reduces to

$$(2.1a) \quad \frac{\partial A}{\partial x} \xi^1 + \frac{\partial A}{\partial y} \xi^2 + A \frac{\partial \xi^2}{\partial y} - B \frac{\partial \xi^1}{\partial y} = 0,$$

$$(2.1b) \quad \frac{\partial B}{\partial x} \xi^1 + \frac{\partial B}{\partial y} \xi^2 + B \frac{\partial \xi^1}{\partial x} - A \frac{\partial \xi^2}{\partial x} = 0,$$

where

$$(2.2) \quad A = \Omega_{12}^1, \quad B = \Omega_{12}^2 \quad (x, y = x^1, x^2).$$

Corresponding to each  $G_r$  of Lie's classification we solve (2.1) for  $A$  and  $B$ , excluding all solutions for which  $A = B = 0$ . (The solutions of (1.3a) for the  $\Gamma_{jk}^i$  have been obtained in [3].) For each of the nonexcluded solutions the  $G_r$  is tested for completeness as explained in [3].

As an illustration we consider the  $G_3[p, q, xq]$  where  $p = \partial f / \partial x$ ,  $q = \partial f / \partial y$ . Using  $X_1 f = p$ ,  $X_2 f = q$ ,  $X_3 f = xq$ , with  $\xi_{11}^i = \delta_1^i$ ,  $\xi_{21}^i = \delta_2^i$ ,  $\xi_{31}^1 = 0$ ,  $\xi_{31}^2 = x$  in (2.1) we obtain  $A = 0$ ,  $B = B_0$  an arbitrary constant.

To obtain the complete group, put  $A = 0$ ,  $B = B_0$  in (2.1), and solve for  $\xi^i$ . This gives  $\xi^1 = k$ ,  $\xi^2 = \xi^2(x, y)$ , with  $k$  an arbitrary constant and  $\xi^2$  an arbitrary function. Now as determined in [3] the  $\Gamma$ 's obtained as solutions of (1.3a) corresponding to the above  $G_3$  will admit as their complete group the  $G_4[p, q, xq, yq]$ . It is thus seen that by taking  $k = 0$ ,  $\xi^2 = y$ , we obtain the fourth generator  $yq$  of the  $G_4$ , and hence this  $G_4$  will be the complete group of motions of the  $L_2$  given by

$$\begin{aligned} L_{11}^1 &= 2a, & L_{12}^1 &= L_{21}^1 = L_{22}^1 = L_{11}^2 = L_{22}^2 = 0, \\ L_{12}^2 &= a + B_0, & L_{21}^2 &= a - B_0 \end{aligned} \quad (B_0 \neq 0).$$

The  $\Gamma$ 's used to give these  $L$ 's are obtained from group [4.1] of [3].

By proceeding as above with each  $G_r$  and making use of the results of [3] we find eight possible types of  $L_2$  spaces admitting real groups of motions. These with their respective groups are given below.

*Asymmetrically connected space  $L_2$  and their complete groups of motions*

[A1.1]  $[p]$

$$L_{jk}^i = L_{jk}^i(y) \quad (\text{arbitrary functions of } y).$$

[A2.1]  $[p, q]$

$$L_{jk}^i = \text{arbitrary constants.}$$

[A2.2]  $[p, xp + yq]$

$$L_{jk}^i = (1/y)A_{jk}^i \quad (A_{jk}^i \text{ arbitrary constants}).$$

[A3.1]  $[p, 2xp + yq, x^2p + xyq]$

$$L_{11}^1 = \frac{2b}{y^2}, \quad L_{12}^1 = -\frac{1}{y}, \quad L_{21}^1 = -\frac{1}{y}, \quad L_{22}^1 = 0,$$

$$L_{11}^2 = \frac{a}{y^3}, \quad L_{12}^2 = \frac{b+c}{y^2}, \quad L_{21}^2 = \frac{b-c}{y^2}, \quad L_{22}^2 = -\frac{2}{y} \quad (c \neq 0).$$

$$[A3.2] \quad [q, xq, yq]$$

$$\begin{aligned} L_{11}^1 &= 2f(x), \quad L_{12}^1 = L_{21}^1 = L_{22}^1 = L_{11}^2 = L_{22}^2 = 0, \\ L_{12}^2 &= f(x) + B(x), \quad L_{21}^2 = f(x) - B(x) \end{aligned} \quad (B \neq 0).$$

$$[A4.1] \quad [p, q, xq, yq]$$

$$\begin{aligned} L_{11}^1 &= 2a, \quad L_{12}^1 = L_{21}^1 = L_{22}^1 = L_{11}^2 = L_{22}^2 = 0, \\ L_{12}^2 &= a + b, \quad L_{21}^2 = a - b \end{aligned} \quad (b \neq 0).$$

$$[A4.2] \quad [p, q, yq, e^x q]$$

$$\begin{aligned} L_{11}^1 &= 1/3 + 2a, \quad L_{12}^1 = L_{21}^1 = L_{22}^1 = L_{11}^2 = L_{22}^2 = 0, \\ L_{12}^2 &= a + c - 1/3, \quad L_{21}^2 = a - c - 1/3 \end{aligned} \quad (c \neq 0).$$

$$[A4.3] \quad [p, yq, (e^{ax} \cos x)q, (e^{ax} \sin x)q]$$

$$\begin{aligned} L_{11}^1 &= 2a/3 + 2b, \quad L_{12}^1 = L_{21}^1 = L_{22}^1 = L_{22}^2 = 0, \\ L_{11}^2 &= (1 + a^2)y, \quad L_{12}^2 = -2a/3 + b + c, \\ L_{21}^2 &= -2a/3 + b - c \end{aligned} \quad (c \neq 0).$$

In the above  $a, b, c$  are arbitrary constants,  $f(x), B(x)$  are arbitrary functions subject to indicated restrictions.

**3. Motions in spaces  $T_2$  of absolute parallelism.** For each of the 8 types of  $L_2$  obtained in the previous section, we impose conditions (1.5), and then solve (1.6) for the  $h_{\alpha 1}^i$ .

The case [A3.2] will be used as an illustration. Conditions (1.5) reduce to  $(f+B)' - (f+B)^2 = 0$ , so that  $f(x) + B(x) = -1/(x+c)$  with  $c$  an arbitrary constant. Equations (1.6) are

$$(3.1a) \quad \frac{\partial h^1}{\partial x} + 2h^1 f(x) = 0, \quad (3.1c) \quad \frac{\partial h^2}{\partial x} + h^2(f(x) - B(x)) = 0,$$

$$(3.1b) \quad \frac{\partial h^1}{\partial y} = 0, \quad (3.1d) \quad \frac{\partial h^2}{\partial y} - \frac{h^1}{x+c} = 0,$$

which can be solved to give  $h^1 = ag(x)$ ,  $h^2 = (ay+b)g(x)/(x+c)$ , where  $g(x) = \exp \{ -2 \int f(x) dx \}$ .

The complete solution is then

$$(3.2) \quad h_{\alpha 1}^1 = a_{\alpha} g(x), \quad h_{\alpha 1}^2 = \frac{a_{\alpha} y + b_{\alpha}}{x+c} g(x) \quad (\alpha = 1, 2),$$

with  $a_{\alpha}, b_{\alpha}, c$  arbitrary constants such that  $a_1 b_2 - a_2 b_1 \neq 0$ . (The special case  $L_{12}^2 = f+B=0$  gives the solution  $h_{\alpha 1}^1 = a_{\alpha} g(x)$ ,  $h_{\alpha 1}^2 = b_{\alpha} g(x)$ .)

It is of interest to verify that the  $T_2$  will admit the finite group generated by the infinitesimal group. To show this we first obtain the finite equations of the group, these being

$$(3.3) \quad \bar{x} = x, \quad \bar{y} = \lambda x + \mu y + \nu$$

with the 3 parameters  $\lambda, \mu, \nu$ .

Now we must have [5]

$$(3.4) \quad h_{\alpha 1}^i(\bar{x}, \bar{y}) = A_{\alpha}^{\beta} h_{\beta 1}^j(x, y) \frac{\partial \bar{x}^i}{\partial x^j},$$

where the  $A_{\alpha}^{\beta}$  must be functions  $A_{\alpha}^{\beta}(\lambda, \mu, \nu)$  of the parameters only. From (3.2), (3.3), and (3.4) we obtain the relations

$$(3.5) \quad A_{\alpha}^{\beta} a_{\beta} = a_{\alpha}, \quad A_{\alpha}^{\beta} b_{\beta} = c_{\alpha},$$

where  $c_{\alpha} = [a_{\alpha}(\nu - \lambda c) + b_{\alpha}]/\mu$ . It follows that the  $A_{\alpha}^{\beta}$  exist and are of the required form.

The results from the solutions of (1.6) are given below.

*Spaces  $T_2$  of absolute parallelism admitting complete groups of motions*

[T1.1] [p]

$$\begin{aligned} \text{(a)} \quad h_{\alpha 1}^1 &= a_{\alpha}^j \exp \left\{ \int (\lambda_j L_{2i}^1 - L_{1i}^1) dx^i \right\}, \\ h_{\alpha 1}^2 &= - \sum_j a_{\alpha}^j \lambda_j \exp \left\{ \int (\lambda_j L_{2i}^1 - L_{1i}^1) dx^i \right\}; \\ \text{(b)} \quad h_{\alpha 1}^1 &= \exp \left\{ \int (\lambda L_{2i}^1 - L_{1i}^1) dx^i \right\} [a_{\alpha} \cos \rho + b_{\alpha} \sin \rho], \\ h_{\alpha 1}^2 &= - \exp \left\{ \int (\lambda L_{2i}^1 - L_{1i}^1) dx^i \right\} [a_{\alpha}(\lambda \cos \rho - \mu \sin \rho) \\ &\quad + b_{\alpha}(\lambda \sin \rho + \mu \cos \rho)]; \\ \text{(c)} \quad h_{\alpha 1}^1 &= \exp \left\{ \int (\lambda L_{2i}^1 - L_{1i}^1) dx^i \right\} \left[ a_{\alpha} \int \frac{L_{2i}^1 dx^i}{L_{21}^1} + b_{\alpha} \right], \\ h_{\alpha 1}^2 &= - \exp \left\{ \int (\lambda L_{2i}^1 - L_{1i}^1) dx^i \right\} \\ &\quad \cdot \left[ \frac{a_{\alpha}}{L_{21}^1} + \lambda \left( a_{\alpha} \int \frac{L_{2i}^1 dx^i}{L_{21}^1} + b_{\alpha} \right) \right]. \end{aligned}$$

In (a),  $\lambda_1 \neq \lambda_2$  are the real roots of  $f(\lambda) \equiv \lambda^2 L_{21}^1 + \lambda(L_{21}^1 - L_{11}^1) - L_{11}^1 = 0$ ;

in (b),  $\lambda \pm i\mu$  are its roots; and in (c),  $f(\lambda)$  is assumed to have equal roots,  $\lambda, \lambda$ . In (b),  $\rho = \int \mu L_{21}^1 dx^1$ . The  $L_{jk}^1(y)$  are otherwise arbitrary subject to (1.5), which have as a consequence that  $L_{11} + L_{21}$  and  $L_{11}L_{21} - L_{21}L_{11}$  must be constant. The  $a_\alpha, b_\alpha, a'_\alpha$  are arbitrary constants such that  $h_{\alpha 1}^1$  are linearly independent. For special values of the  $L_{jk}^1$ , for example,  $L_{21}^1 = 0$ , the solutions for the  $h_{\alpha 1}^1$  can be easily obtained.

[T2.1]  $[p, q]$

The  $h_{\alpha 1}^1$  are obtained as in the 3 cases of [T1.1] except the  $L_{jk}^1$  are now considered constants.

[T2.2]  $[p, xp + yq]$

$$h_{\alpha 1}^1 = y^m (a_\alpha u + b_\alpha), \quad h_{\alpha 1}^2 = y^m (c_\alpha v + d_\alpha),$$

with  $c_\alpha A_{11}^2 + a_\alpha A_{11}^1 = 0$ ,  $d_\alpha A_{21}^1 + b_\alpha A_{11}^1 = 0$ ,  $m = -(1 + A_{12}^1 + A_{22}^2)/2$ , and  $u = xA_{21}^1 + yA_{22}^1$ ,  $v = xA_{11}^2 + yA_{12}^2$ . The constants  $A_{jk}^1$  are arbitrary subject to (1.5). As in [T1.1] the solutions for  $h_{\alpha 1}^1$  corresponding to special values of the  $A_{jk}^1$  can be obtained without difficulty.

[T3.1]  $[p, 2xp + yq, x^2p + xyq]$

$$h_{\alpha 1}^1 = y(a_\alpha x + b_\alpha), \quad h_{\alpha 1}^2 = (y^2 + 2bx)a_\alpha + 2bb_\alpha.$$

The  $a, b, c$  of [A3.1] must satisfy  $c = 3b$ ,  $a = 4b^2$ ,  $b \neq 0$ .

[T3.2]  $[q, xq, yq]$

$$h_{\alpha 1}^1 = a_\alpha g(x), \quad h_{\alpha 1}^2 = (a_\alpha y + b_\alpha)g(x)/(x + c).$$

Here the  $L_{jk}^1$  of [A3.2] satisfy  $g(x) = \exp \{-\int 2f(x)dx\}$ , and  $f(x) + B(x) = -1/(x+c)$ .

[T4.1]  $[p, q, xq, yq]$

$$h_{\alpha 1}^1 = a_\alpha e^{-2ax}, \quad h_{\alpha 1}^2 = b_\alpha e^{-2ax},$$

with  $a+b=0$ , ( $b \neq 0$ ) in [A4.1].

[T4.2]  $[p, q, yq, e^x q]$

$$\begin{aligned} h_{\alpha 1}^1 &= a_\alpha e^{(2c-1)x}, & h_{\alpha 1}^2 &= b_\alpha e^{2cx} & (a+c = 1/3 \text{ in [A4.2]}), \\ h_{\alpha 1}^1 &= a_\alpha e^{(2c+1)x}, & h_{\alpha 1}^2 &= (a_\alpha y + b_\alpha) e^{(2c+1)x} & (a+c = -2/3 \text{ in [A4.2]}). \end{aligned}$$

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