

## A FORMULA FOR DISTRIBUTION TRACE CHARACTERS ON NILPOTENT LIE GROUPS

L. CORWIN AND F. P. GREENLEAF

**ABSTRACT.** The distribution trace character  $\theta_\pi$  of an irreducible representation  $\pi$  on a simply connected nilpotent Lie group  $N$  is described as a superposition of unitary characters on certain subgroups, in terms of the canonical objects introduced by R. Penney [8]. For  $l \in \mathfrak{n}^*$ , define  $\mathfrak{h}_1(l)$  = smallest ideal containing the radical  $\mathfrak{r}(l)$ , and  $\mathfrak{h}_{k+1}(l) = \mathfrak{h}_k(l) \cap \mathfrak{h}_k(l)$ . These subalgebras terminate in a subordinate subalgebra  $\mathfrak{h}_\infty(l)$  after finitely many steps. If  $H_\infty = \exp(\mathfrak{h}_\infty)$ ,  $\chi_\infty = (e^{2\pi i l}) \circ \log|H_\infty$ , and  $(\chi_\infty, H_\infty) \cdot n = (\chi_\infty \cdot n, H_\infty \cdot n)$ , where  $H_\infty \cdot n = n^{-1}H_\infty n$ ,  $\chi_\infty \cdot n(h') = \chi_\infty(nh'n^{-1})$  on  $H_\infty \cdot n$ , then  $\langle \theta_\pi, \phi \rangle = \int_{K_\infty \backslash N} \langle (\chi_\infty, H_\infty) \cdot n, \phi \rangle dn$ , where the pair  $(\chi_\infty, H_\infty) \cdot n$  is regarded as the tempered distribution  $\langle (\chi_\infty, H_\infty) \cdot n, \phi \rangle = \int_{H_\infty} \chi_\infty(h) \phi(n^{-1}hn) dh$ , and where  $\mathfrak{k}_\infty = \{X \in \mathfrak{n}; l[h_\infty, X] = 0\}$  gives the stabilizer of the pair  $(\chi_\infty, H_\infty)$ . The integral over  $K_\infty \backslash N$  is absolutely convergent for any Schwartz function  $\phi$  on  $N$ .

If  $N$  is a simply connected nilpotent Lie group and  $\pi \in \hat{N}$  an irreducible representation, then  $\pi_\phi$  is trace class for any  $\phi \in C_c^\infty(N)$  and the trace character  $\langle \theta_\pi, \phi \rangle = \text{Tr } \pi_\phi$  is a tempered distribution on  $N$ . Its Fourier transform has been described in terms of coadjoint orbits in  $\mathfrak{n}^*$  by Kirillov [6]. Various attempts have been made to describe the  $\theta_\pi$  directly as distributions on  $N$ , but this seems quite difficult. Dixmier [3] computed a number of examples, and some asymptotic properties were discussed in [1] for special  $N$ .

This note shows how  $\theta_\pi$  can be represented as a superposition of ordinary characters on certain canonical subgroups on  $N$  related to the canonical objects introduced by Penney in [8], and exploited in [2]. The formula we get is clearly suggested by [2], but a very direct calculation can be made for the  $\theta_\pi$ . This formula can be used to estimate the order of  $\theta_\pi$  as a Schwartz distribution, and may prove to have uses in the harmonic analysis of smooth functions on  $N$ .

Let  $\mathcal{S}(N)$  be the Schwartz functions on  $N$  (i.e.  $f \circ \exp$  is Schwartz on the Lie algebra  $\mathfrak{n}$ ). If  $\pi \in \hat{N}$ , let  $l$  be an element in the coadjoint orbit  $\mathcal{O} = \mathcal{O}_\pi \subseteq \mathfrak{n}^*$ ,  $\mathfrak{m}$  a maximal subordinate subalgebra for  $l$ , and  $\chi = (e^{2\pi i l}) \circ \log|M$  the resulting character on  $M = \exp(\mathfrak{m})$ . Let  $\{X_1, \dots, X_m, \dots, X_n\}$  be a Jordan-Hölder basis for  $\mathfrak{n}$  that passes through  $\mathfrak{m}$ ;  $\mathfrak{n}_j = \mathbf{R} - \text{span}\{X_1, \dots, X_j\}$  a subalgebra with  $\mathfrak{n}_j \triangleleft \mathfrak{n}_{j+1}$  for all  $j$ , and  $\mathfrak{n}_m = \mathfrak{m}$ . The map  $\gamma: \mathbf{R}^k \rightarrow N$ ,

$$\gamma(t) = \exp(t_1 X_{m+1}) \cdots \exp(t_k X_{m+k}), \quad k = n - m,$$

---

Received by the editors August 20, 1982 and, in revised form, March 21, 1983.  
 1980 *Mathematics Subject Classification*. Primary 22E27.

cross sections the  $M \setminus N$  cosets, carries Lebesgue measure  $dt$  to right invariant measure  $dn$  on  $M \setminus N$ , and allows us to realize the Hilbert space of the induced representation  $\pi = \text{Ind}(M \uparrow N, \chi)$  isometrically as  $L^2(\mathbf{R}^k)$ , as follows: the function  $\tilde{f}$  on  $G$  corresponds to  $f \in L^2(\mathbf{R}^k)$  via restriction to the cross-section  $\gamma(\mathbf{R}^k)$  and pullback to  $\mathbf{R}^k$ ,

$$\tilde{f}(n) = \tilde{f}(m\gamma(t)) = \chi(m)f(t).$$

For  $\phi \in \mathfrak{S}(N)$  we may explicitly compute the kernel  $K_\phi(s, t)$  of the operator  $\pi_\phi$ :

$$\begin{aligned} \pi_\phi f(t) &= \pi_\phi \tilde{f}(\gamma(t)) = \int_N \phi(n) \pi_n \tilde{f}(\gamma(t)) \, dn = \int_N \phi(\gamma(t)^{-1}n) \tilde{f}(n) \, dn \\ &= \int_{\mathbf{R}^k} \int_M \phi(\gamma(t)^{-1}m\gamma(s)) \tilde{f}(m\gamma(s)) \, dm \, ds \\ &= \int_{\mathbf{R}^k} \left[ \int_M \chi(m) \phi(\gamma(t)^{-1}m\gamma(s)) \, dm \right] f(s) \, ds \\ &= \int_{\mathbf{R}^k} K_\phi(t, s) f(s) \, ds. \end{aligned}$$

Since  $\pi_\phi$  is trace class [9, p. 108 ff.] and  $K_\phi$  is continuous, the trace is just the integral of the diagonal values

$$(1) \quad \langle \theta_\pi, \phi \rangle = \text{Tr } \pi_\phi = \int_{\mathbf{R}^k} K_\phi(t, t) \, dt$$

and the integral is absolutely convergent (see [4, III.10] for example).

Let  $(\chi, M) \cdot n = (\chi \cdot n, M \cdot n)$ , where  $M \cdot n = n^{-1}Mn$  and  $\chi \cdot n(m') = \chi(nm'n^{-1})$  for  $m' \in M \cdot n$ , to obtain a right action of  $N$  on pairs  $(\chi, M)$  where  $\chi$  is a character on  $M$ . Fixing a Haar measure  $dm$  on  $M$ , we may regard each pair in an  $N$  orbit as a tempered distribution

$$\langle (\chi, M) \cdot n, \phi \rangle = \int_M \chi(m) \phi(n^{-1}mn) \, dm, \quad \text{all } \phi \in \mathfrak{S}(N).$$

From (1) and the fact that  $(\chi, M) = (\chi, M) \cdot n \Leftrightarrow n \in M$ , we see that  $\theta_\pi$  is a superposition of all pairs in the orbit  $(\chi, M) \cdot N$ :

$$\begin{aligned} (2) \quad \langle \theta_\pi, \phi \rangle &= \int_{\mathbf{R}^k} K_\phi(t, t) \, dt = \int_{\mathbf{R}^k} \left[ \int_M \chi(m) \phi(\gamma(t)^{-1}m\gamma(t)) \, dm \right] dt \\ &= \int_{M \setminus N} \langle (\chi, M) \cdot n, \phi \rangle \, d\dot{x}, \end{aligned}$$

the formula being absolutely convergent for all  $\phi \in \mathfrak{S}(N)$ . Formula (2) is not canonical since it involves a choice of  $l$  and a polarization  $m$ . It must involve a great deal of hidden cancellation, as the canonical version below indicates.

For  $l \in \mathfrak{n}^*$  we define, following Penney [8], the canonical objects:  $\mathfrak{h}_1(l)$  = smallest ideal containing the radical  $\mathfrak{r}(l)$ , and  $\mathfrak{h}_{k+1}(l) = \mathfrak{h}_1(l | \mathfrak{h}_k(l))$ . These subalgebras stabilize eventually,  $\mathfrak{h}_1 \supseteq \mathfrak{h}_2 \supseteq \dots \supseteq \mathfrak{h}_k = \mathfrak{h}_{k+1} = \dots = \mathfrak{h}_\infty$ ; all contain the radical

$r(l)$  and are transformed covariantly:

$$\mathfrak{h}_k(l \cdot n) = \mathfrak{h}_k(l) \cdot n \quad \text{where } l \cdot n = \text{Ad}^*(n^{-1})l \text{ and } X \cdot n = \text{Ad}(n^{-1})X,$$

and  $\mathfrak{h}_\infty$  is subordinate to  $l$ . Let  $\mathfrak{f}_\infty(l) = \{X \in \mathfrak{n} : l[\mathfrak{h}_\infty, X] = 0\}$ ; this turns out to be a subalgebra, and if we set  $l_\infty = l|_{\mathfrak{f}_\infty}$ , then  $r(l_\infty) = \mathfrak{h}_\infty(l) \triangleleft \mathfrak{f}_\infty(l)$  so that  $l_\infty \in \mathfrak{f}_\infty^*$  has flat orbit (see [7]). Let  $\chi_\infty = (e^{2\pi i l}) \circ \log|H_\infty$ ; by fixing a Haar measure  $dh$  on  $H_\infty$  we may regard  $(\chi_\infty, H_\infty)$  and all its conjugates  $(\chi_\infty, H_\infty) \cdot n$  as tempered distributions. The significance of the following formula is that  $\mathfrak{h}_\infty$  can be much smaller than any polarization  $\mathfrak{m}$  for  $l$ .

**THEOREM 1.** *Let  $l \in \mathfrak{n}^*$  and define  $H_\infty, K_\infty, \chi_\infty$  as above. Fixing  $dh$  on  $H_\infty$  and  $d\mathfrak{n}$  on  $K_\infty \setminus N$ , and regarding each  $(\chi_\infty, H_\infty) \cdot n$  as a tempered distribution*

$$\langle (\chi_\infty, H_\infty) \cdot n, \phi \rangle = \int_{H_\infty} \chi(h) \phi(n^{-1}hn) dh, \quad \phi \in \mathfrak{S}(N),$$

we get

$$(3) \quad \langle \theta_\pi, \phi \rangle = \int_{K_\infty \setminus N} \langle (\chi_\infty, H_\infty) \cdot n, \phi \rangle d\mathfrak{n}, \quad \text{all } \phi \in \mathfrak{S}(N),$$

where  $d\mathfrak{n}$  is suitably normalized. The integral is absolutely convergent.

**PROOF.** Recall that  $K_\infty = \{x \in N : (\chi_\infty, H_\infty) \cdot n = (\chi_\infty, H_\infty)\}$  (see [2]). Let  $M$  be any maximal subordinate subgroup for  $l_\infty \in \mathfrak{f}_\infty^*$ ; it must also be maximal subordinate for  $l \in \mathfrak{n}^*$ , by dimension counting. Let  $\chi = (e^{2\pi i l}) \circ \log|M$  and  $\pi = \text{Ind}(M \uparrow N, \chi)$ . We start with a remark on the special case:  $\pi$  square integrable modulo  $Z(N)$ . This is precisely the case in which  $H_\infty(l) = R(l) = Z(N)$ . As is well known [7], we have flat orbit  $\mathcal{O}(l) = l + \mathfrak{z}^\perp$ . From a simple version of Kirillov's orbital integral formula and Poisson summation, we get

$$\begin{aligned} \theta_\pi &= \text{Euclidean Fourier transform of Lebesgue measure on } l + \mathfrak{z}^\perp \\ &= \text{scalar multiple of } \chi_\infty \text{ weighting Haar measure on } Z(N), \end{aligned}$$

or  $\langle \theta_\pi, \varphi \rangle = \langle (\chi_\infty, H_\infty), \varphi \rangle$  where  $dh$  on  $H_\infty$  is suitably chosen. Since  $\mathfrak{f}_\infty = \mathfrak{n}$  here, we have proved (3) in this special case.

We now use the noncanonical formula (2) to complete the proof of (3).

Take an arbitrary  $l \in \mathfrak{n}^*$ , consider a system  $H_\infty \subseteq M \subseteq K_\infty$ , and let

$$\pi_\infty = \text{Ind}(M \uparrow K_\infty, \chi) \in \widehat{K_\infty}.$$

Applying (2) to this situation, we get

$$(4) \quad \langle \theta_{\pi_\infty}, \phi \rangle = \int_{M \setminus K_\infty} \langle (\chi, M) \cdot k, \phi \rangle dk, \quad \text{all } \phi \in \mathfrak{S}(K_\infty)$$

(or  $\phi \in \mathfrak{S}(N)$ , for that matter). We claim that

$$(5) \quad \langle \theta_{\pi_\infty}, \phi \rangle = \langle (\chi_\infty, H_\infty), \phi \rangle, \quad \text{all } \phi,$$

so the superposed conjugates of  $(\chi, M)$  by elements in  $K_\infty$  collapse to the single character  $(\chi_\infty, H_\infty)$  through cancellations. Once (5) is proved, we finish proving (3) as follows: by Fubini,

$$\begin{aligned} \langle \theta_\pi, \phi \rangle &= \int_{M \setminus N} \langle (\chi, M) \cdot n, \phi \rangle \, d\dot{n} \\ &= \int_{K_\infty \setminus N} \left[ \int_{M \setminus K_\infty} \langle (\chi, M) \cdot k, \phi \circ \text{Ad}(x) \rangle \, d\dot{k} \right] \, d\dot{x} \\ &= \int_{K_\infty \setminus N} \langle (\chi_\infty, H_\infty), \phi \circ \text{Ad}(x) \rangle \, d\dot{x} = \int_{K_\infty \setminus N} \langle (\chi_\infty, H_\infty) \cdot x, \phi \rangle \, d\dot{x} \end{aligned}$$

with all integrals absolutely convergent.

To prove (5) it suffices to prove the following general statement about functionals  $l \in \mathfrak{n}^*$  with flat orbit (i.e.  $\mathfrak{r}(l) \triangleleft \mathfrak{n}$ , so that  $H_\infty = R(l)$ ,  $K_\infty = N$  and  $\pi_\infty = \pi = \text{Ind}(M \uparrow N, \chi)$ ).

(6) If  $l \in \mathfrak{n}^*$  has  $\mathfrak{r}(l) \triangleleft \mathfrak{n}$ ,  $\mathfrak{m}$  is any maximal subordinate subalgebra,  $\chi = (e^{2\pi i l} \circ \log) |M$ , and  $\pi = \text{Ind}(M \uparrow N, \chi)$ , then  $\theta_\pi = (\chi_\infty, H_\infty)$ .

Here the orbit is flat, and  $\Theta_\pi = l + \mathfrak{H}_\infty^\perp$ ; the  $N$ -invariant measure on  $\Theta_\pi$  is just Euclidean measure on  $\mathfrak{H}_\infty^\perp$  lifted to this coset. By Poisson summation and Kirillov's description of  $\theta$ , we get (6). This completes the proof of Theorem 1. Q.E.D.

The canonical formula (3) leads to the following conclusion about flat parts in coadjoint orbits. If  $\mathfrak{m}$  is maximal subordinate for  $l \in \mathfrak{n}^*$ , it is well known that  $\Theta$  is a disjoint union of the flat fibers  $l \cdot x + \mathfrak{m}^\perp \cdot x = l \cdot Mx$ ,  $x \in M \setminus N$ . In general  $\mathfrak{h}_\infty \subseteq \mathfrak{m}$ , so  $\mathfrak{h}_\infty^\perp$  may be larger than  $\mathfrak{m}^\perp$ . From (3) we see that

$$\Theta = \cup \{ (l + \mathfrak{h}_\infty^\perp) \cdot x : x \in K_\infty \setminus N \}.$$

This is a disjoint union: if  $x, y \in N$  give non-disjoint fibers, then  $(l + \mathfrak{h}_\infty^\perp) \cdot xy^{-1}$  meets  $l + \mathfrak{h}_\infty^\perp$ . But by dimension counting we can show that  $l + \mathfrak{h}_\infty^\perp = l \cdot K_\infty$ , so there exist  $k_1, k_2 \in K_\infty$  such that  $l \cdot k_1(xy^{-1})k_2 = l$ . Thus  $k_1(xy^{-1})k_2 \in R(l) \subseteq H_\infty \subseteq K_\infty$ , and  $xK_\infty = yK_\infty$  as required. Formula (3) also shows that  $\text{supp}(\theta_\pi) \subseteq$  the smallest connected (normal) subgroup containing  $\cup_{n \in N} n^{-1}H_\infty n$ . Since  $\mathfrak{h}_1 =$  ideal generated by  $\mathfrak{r}(l) \supseteq \mathfrak{h}_\infty \supseteq \mathfrak{r}(l)$ , this amounts to saying that  $\text{supp}(\theta_\pi) \subseteq H_1$ . By other means, Pukanszky [10] has recently shown that a normal Lie subgroup  $H$  contains  $\text{supp}(\theta_\pi) \Leftrightarrow H \supseteq H_1$ , so further reduction of the support through cancellation cannot occur.

REFERENCES

1. L. Corwin and F. P. Greenleaf, *Singular Fourier integral operators and representations of nilpotent Lie groups*, Comm. Pure Appl. Math. **31** (1978), 681-705
2. L. Corwin, F. P. Greenleaf and R. Penney, *A canonical formula for the distribution kernels of primary projections in  $L^2$  of a nilmanifold*, Comm. Pure Appl. Math. **30** (1977), 355-372.
3. J. Dixmier, *Sur les représentations unitaires des groupes de Lie nilpotents*. IV, Canad. J. Math. **11** (1959), 321-344.
4. I. C. Gohberg and M. G. Krein, *Introduction to the theory of linear nonselfadjoint operators*, Transl. Math. Mono., vol. 18, Amer. Math. Soc., Providence, R. I., 1969.

5. R. Howe, *On a connection between nilpotent groups and oscillatory integrals associated to singularities*, Pacific J. Math. **73** (1977), 329–364.
6. A. A. Kirillov, *Unitary representations of nilpotent Lie groups*, Uspehi Mat. Nauk **17** (1962), 57–110.
7. C. C. Moore and J. Wolf, *Square integrable representations of nilpotent groups*, Trans. Amer. Math. Soc. **185** (1973), 445–462.
8. R. Penney, *Canonical objects in the Kirillov theory of nilpotent Lie groups*, Proc. Amer. Math. Soc. **66** (1977), 175–178.
9. L. Pukanszky, *Leçons sur les représentations des groupes*, Dunod, Paris, 1967.
10. \_\_\_\_\_, *On Kirillov's character formula*, J. Reine Angew. Math. **311/312** (1979), 408–440.

DEPARTMENT OF MATHEMATICS, RUTGERS UNIVERSITY, NEW BRUNSWICK, NEW JERSEY 08903

COURANT INSTITUTE OF MATHEMATICAL SCIENCES, NEW YORK UNIVERSITY, NEW YORK, NEW YORK 10012