

ASYMPTOTICS OF SOBOLEV EMBEDDINGS AND SINGULAR PERTURBATIONS FOR THE p -LAPLACIAN

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To the memory of Carlos Cid

ABSTRACT. We consider the best constant $S(\Omega_\lambda)$ for the embedding of $W^{1,p}(\Omega_\lambda)$ into $L^q(\Omega_\lambda)$ where $1 < p < 2$, $p < q < \frac{Np}{N-p}$. Here $\Omega_\lambda = \lambda\Omega$ with Ω a smooth, bounded domain in \mathbb{R}^n and λ a large positive number. It is proven by the validity of the expansion

$$S(\Omega_\lambda) = S(\mathbb{R}_+^n) - \lambda^{-1}\gamma \max_{x \in \partial\Omega} H(x) + o(\lambda^{-1}),$$

as $\lambda \rightarrow \infty$, where γ is a positive constant depending on p, q and N . The behavior of associated extremals, which satisfy an equation involving the p -Laplacian operator, is also analyzed.

1. INTRODUCTION AND STATEMENT OF MAIN RESULTS

Let Ω be a bounded domain in \mathbb{R}^N , $N \geq 2$, with $\partial\Omega$ smooth. Let $p > 1$ and denote $p^* = \frac{Np}{N-p}$ if $p < N$, $p^* = +\infty$ otherwise. It is well known that for any $1 < q < p^*$ the Sobolev embedding of $W^{1,p}(\Omega)$ into $L^q(\Omega)$ holds, namely there exists a positive constant $S = S(p, q, \Omega)$ such that

$$S\|u\|_{L^q(\Omega)}^p \leq \|u\|_{W^{1,p}(\Omega)}^p$$

for all $u \in W^{1,p}(\Omega)$. The *best constant* for this embedding is the largest S for which the above relation holds, namely the number $S(\Omega)$ defined as

$$(1.1) \quad S(\Omega) = \inf_{u \in W^{1,p}(\Omega) \setminus \{0\}} \frac{\int_\Omega |\nabla u|^p + |u|^p}{\left(\int_\Omega |u|^q\right)^{\frac{p}{q}}}.$$

This embedding is compact, which implies the existence of *extremals* for it, namely, functions u at which this infimum is achieved.

Let us fix p and q as above, and a bounded smooth domain Ω . For a large positive number λ we consider the family of expanding domains

$$\Omega_\lambda = \lambda\Omega = \{\lambda x \mid x \in \Omega\}.$$

Our purpose in this paper is to describe the asymptotic behavior as $\lambda \rightarrow +\infty$ of the best constants $S(\Omega_\lambda)$ as well as that of the associated family of extremals u_λ .

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In what follows we shall denote by u_λ an extremal normalized so that the relation

$$\int_{\Omega_\lambda} |\nabla u_\lambda|^p + |u_\lambda|^p = \int_{\Omega_\lambda} |u_\lambda|^q$$

holds. Then the Euler-Lagrange equation satisfied by u_λ becomes

$$(1.2) \quad \begin{cases} \Delta_p u_\lambda - |u_\lambda|^{p-2} u_\lambda + |u_\lambda|^{q-2} u_\lambda = 0 & \text{in } \Omega_\lambda, \\ u > 0 \text{ in } \Omega_\lambda, \frac{\partial u}{\partial \eta} = 0 & \text{on } \partial\Omega_\lambda. \end{cases}$$

Here Δ_p stands for the p -Laplacian operator, $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2} \nabla u)$. Note that u_λ is one-signed in Ω . Indeed, $|u_\lambda|$ is also a minimizer of the Raleigh quotient above, hence also a solution to (1.2). Regularity theory for the p -Laplacian (see [5, 11]) applies to yield that $|u_\lambda|$ is actually of class $C^{1,\alpha}(\Omega)$. Then the strong maximum principle proved in [12] implies that $|u_\lambda|$ does not vanish in Ω , and therefore u_λ is one-signed. Henceforth we will assume $u_\lambda > 0$ in Ω .

Since Ω_λ expands toward entire space or to a half-space depending on the choice of origin, it is natural to relate the behavior of $S(\Omega_\lambda)$ and u_λ with best constant and extremals of the Sobolev embedding in \mathbb{R}^N ,

$$S(\mathbb{R}^N) = \inf_{u \in W^{1,p}(\mathbb{R}^N) \setminus \{0\}} \frac{\int_{\mathbb{R}^N} |\nabla u|^p + |u|^p}{\left(\int_{\mathbb{R}^N} |u|^q\right)^{\frac{2}{q}}}.$$

A concentration-compactness argument along the lines of [1], [7] shows that this infimum is achieved. Modulo normalization, extremals are positive solutions of the equation

$$(1.3) \quad \begin{aligned} \Delta_p w - |w|^{p-2} w + |w|^{q-2} w &= 0 && \text{in } \mathbb{R}^N, \\ w(x) &\rightarrow 0 && \text{as } |x| \rightarrow +\infty. \end{aligned}$$

It has been established in [2] that for $1 < p < 2$, positive solutions of (1.3) are radially symmetric around some point. Moreover, from a recent result in [10], the radial solution around the origin is unique. We shall denote it by $w_* = w_*(|x|)$ in what follows. Let us consider the half-space $\mathbb{R}_+^N = \{(x', x_N) \mid x_N > 0\}$. The best constants of \mathbb{R}_+^N and \mathbb{R}^N relate as

$$S(\mathbb{R}_+^N) = 2^{-\frac{q-p}{q}} S(\mathbb{R}^N).$$

Corresponding extremals are positive solutions of the problem

$$(1.4) \quad \begin{aligned} \Delta_p w - |w|^{p-2} w + |w|^{q-2} w &= 0 && \text{in } \mathbb{R}_+^N, \\ w(x) &\rightarrow 0 \text{ as } |x| \rightarrow +\infty, \quad \frac{\partial w}{\partial x_N} = 0 && \text{on } \partial\mathbb{R}_+^N. \end{aligned}$$

Even extension of such a solution to entire space corresponds to a solution of problem (1.3). It is natural to suspect that $S(\Omega_\lambda)$ converges to the corresponding quantity for the half-space, and u_λ to an associated extremal. Our principal result states that when $1 < p < 2$, $S(\Omega_\lambda)$ indeed approaches $S(\mathbb{R}_+^N)$, corrected by a negative factor of the maximum mean curvature of $\partial\Omega$.

Theorem 1.1. *Assume that $1 < p < 2$ and that $p < q < \frac{Np}{N-p}$. There is a constant $\gamma = \gamma(p, q, N) > 0$ such that the following expansion holds:*

$$(1.5) \quad S(\Omega_\lambda) = S(\mathbb{R}_+^N) - \lambda^{-1} \gamma \max_{x \in \partial\Omega} H(x) + o(\lambda^{-1}),$$

as $\lambda \rightarrow +\infty$. Here $H(x)$ denotes the mean curvature of the boundary at the point x . Moreover, there exist points $x^\lambda \in \partial\Omega$ such that

$$\sup_{y \in \Omega_\lambda} |u_\lambda(y) - w_*(|y - \lambda x^\lambda|)| \rightarrow 0$$

and

$$H(x^\lambda) \rightarrow \max_{x \in \partial\Omega} H(x)$$

as $\lambda \rightarrow \infty$.

The radial symmetry of the extremals of $S(\mathbb{R}^N)$ plays a crucial role in the proof of the theorem. For $p > 2$ the same result would hold if such a fact was true. The constant γ above is given explicitly as follows:

$$\gamma = \frac{2p}{q} \frac{N-1}{N+1} \omega_N \omega_{N-1} \int_0^\infty |w'_*(r)|^p r^N dr \int_0^\infty w_*(r)^q r^{N-1} dr.$$

For $p = 2$, these facts have been known since the works by Lin Ni and Takagi, and Ni and Takagi [6, 8, 9]. The proof devised in those works does not apply in the current situation. Strong use of linearity of the differential operator, as well as certain nondegeneracy properties of the linearized equation around w_* only known for $p = 2$, is used. A different proof of those results was found in [4]. We borrow ideas from that work in the proof of Theorem 1.1. See also [3] for a related result involving trace embeddings and $p = 2$.

2. PRELIMINARIES

Let us consider the best Sobolev constant $S(\Omega_\lambda)$ given by (1.1) for the embedding of $W^{1,p}(\Omega_\lambda)$ into $L^q(\Omega_\lambda)$. It is convenient for our purposes to obtain a further characterization of this value and its extremals in terms of the energy functional

$$(2.1) \quad J_\lambda(u) = \frac{1}{p} \int_{\Omega_\lambda} (|\nabla u|^p + u^p) dx - \frac{1}{q} \int_{\Omega_\lambda} u_+^q dx.$$

It is standard to check that nontrivial critical points of J_λ in $W^{1,p}(\Omega_\lambda)$ correspond precisely to the positive solutions of problem (1.2). Let us consider the number

$$(2.2) \quad c_\lambda \equiv \inf_{\substack{u \in W^{1,p}(\Omega_\lambda) \\ u \neq 0}} \sup_{t > 0} J_\lambda(tu).$$

It is easy to see that if $u_+ \neq 0$, the function $t \mapsto J_\lambda(tu)$ has a maximum $t = \bar{t} > 0$ which is its unique critical point. Then $\bar{t}u \in M_\lambda$, where

$$(2.3) \quad M_\lambda = \left\{ u \in W^{1,p}(\Omega_\lambda) / u \neq 0, \int_{\Omega_\lambda} |\nabla u|^p + u^p = \int_{\Omega_\lambda} u_+^q \right\},$$

is the so-called *Nehari's manifold* of J_λ . It follows from this fact that

$$c_\lambda = \inf_{u \in M_\lambda} J_\lambda(u).$$

Since all nontrivial solutions of (1.2) lie in M_λ , the above number is called the *least energy* value for J_λ and a solution u of (1.2) with $J_\lambda(u) = c_\lambda$, a *least energy solution*. These solutions and extremals of $S(\Omega_\lambda)$ are related in the following way: if u is a least energy solution, then it is an extremal of $S(\Omega_\lambda)$. Reciprocally, if $\bar{u} \geq 0$

minimizes the Raleigh quotient (1.1), then $u = t\bar{u}$ is a least energy solution of (1.2) where

$$t^{q-p} = \frac{\int_{\Omega_\lambda} |\nabla \bar{u}|^p + \bar{u}^p}{\int_{\Omega_\lambda} \bar{u}^q}.$$

In fact we always have the exact relation

$$(2.4) \quad c_\lambda = \left(\frac{1}{p} - \frac{1}{q}\right) S(\Omega_\lambda)^{\frac{q}{q-p}}.$$

As we have mentioned, for Ω_λ bounded the compactness of the associated embedding yields the existence of extremals for $S(\Omega_\lambda)$ and correspondingly of critical points of J_λ at level c_λ .

Now, we establish an L^∞ estimate for solutions of (1.2) in terms of their energy values.

Lemma 2.1. *Let u be a solution of (1.2). Then there are constants $B = B(\Omega, p, N)$ and $\theta = \theta(\Omega, p, N)$, independent of $1 \leq \lambda < \infty$, such that*

$$(2.5) \quad \|u\|_\infty \leq B J_\lambda(u)^\theta.$$

Proof. We consider a positive solution u of (1.2). Let us multiply (1.2) by $u^{\alpha p+1-p}$, where $\alpha \geq 1$. Integrating over Ω_λ , we find that

$$(2.6) \quad \frac{\alpha p + 1 - p}{\alpha^p} \int_{\Omega_\lambda} |\nabla u^\alpha|^p + \int_{\Omega_\lambda} u^{\alpha p} = \int_{\Omega_\lambda} u^{\alpha p+q-p}.$$

Noticing that $\frac{1}{2\alpha^{p-1}} \leq \min\{1, \frac{\alpha p+1-p}{2\alpha^p}\}$, we get

$$\frac{1}{2\alpha^{p-1}} \int_{\Omega_\lambda} |\nabla u^\alpha|^p + u^{\alpha p} \leq \int_{\Omega_\lambda} u^{\alpha p+q-p}.$$

Sobolev’s inequality applied for $v = u^\alpha$ yields

$$(2.7) \quad \left(\int_{\Omega_\lambda} u^{\alpha p^*}\right)^{\frac{p}{p^*}} \leq C \alpha^{p-1} \int_{\Omega_\lambda} u^{q+\alpha p-p}$$

where, we recall, $p^* = \frac{Np}{N-p}$ and the constant C is independent of λ . Next, we consider the sequence of positive numbers $\alpha_j, j = 0, 1, 2, \dots$, defined inductively as

$$(2.8) \quad q + \alpha_0 p - p = p^*, \quad q + \alpha_{j+1} p - p = \alpha_j p^*, \quad \forall j = 0, 1, \dots,$$

or, explicitly,

$$(2.9) \quad \alpha_j = \frac{(p^*/p)^{j+1}(p^* - q) + q - p}{p^* - p}.$$

Note that, by (2.7) with $\alpha = 1$,

$$(2.10) \quad \left(\int_{\Omega_\lambda} u^{p^*}\right)^{p/p^*} \leq C \left(\frac{1}{p} - \frac{1}{q}\right)^{-1} J_\lambda(u),$$

where J_λ is the energy functional given by (2.1). Now we will construct a suitable sequence of positive numbers M_j such that

$$(2.11) \quad \int_{\Omega_\lambda} u^{q+\alpha_j p-p} \leq M_j, \quad \forall j.$$

Inequality (2.10) gives us M_0 . Assuming that (2.11) holds, we have by (2.7) and (2.9) that

$$(2.12) \quad \int_{\Omega_\lambda} u^{q+\alpha_{j+1}p-p} \leq \left(C\alpha_j^{p-1} \int_{\Omega_\lambda} u^{q+\alpha_j p-p} \right)^{p^*/p} \leq (C\alpha_j^{p-1} M_j)^{p^*/p}.$$

(2.11) then holds for M_j defined as

$$(2.13) \quad M_0 = \left(C \left(\frac{1}{p} - \frac{1}{q} \right)^{-1} J_\lambda(v) \right)^{p^*/q}, \quad M_{j+1} = \left(C\alpha_j^{p-1} M_j \right)^{p^*/q} \quad \forall j \geq 0.$$

From these relations, the explicit form of the α_j 's given by (2.9), and a straightforward computation, we find constants C_1, C_2 depending only on C, p, q, N such that

$$(2.14) \quad M_j \leq \exp(C_1 \log M_0) \exp(C_2 \alpha_{j-1} (1 + \log M_0))$$

and using relations (2.7)-(2.9),

$$(2.15) \quad \|v\|_{L^{\alpha_j p^*}(\Omega_\lambda)} \leq (\exp(C_1 \log M_0))^{\frac{p}{\alpha_j p^*}} \exp(C_2/p^*(1 + \log M_0)).$$

Letting $j \rightarrow \infty$ we obtain the result. □

Now we claim that $J_\lambda(u_\lambda)$ is uniformly bounded independently of λ . In fact, let us consider $w_\lambda^z(x) = w_*(\lambda(x - z))$, where $z \in \partial\Omega$ and w_* is the radial least energy solution of equation (1.3). Using the definition of c_λ and a direct computation we then find that

$$(2.16) \quad c_\lambda \leq \sup_{t>0} J_\lambda(tw_\lambda^z) \leq \frac{c_*}{2} + o(1),$$

where c_* is the corresponding energy of w_* in the whole space \mathbb{R}^N , namely $c_*/2$ is that in the half-space. As a consequence, the above lemma yields a uniform L^∞ estimate for the solutions u_λ . Moreover, from the $C^{1,\alpha}$ estimates found in [5] and [11], we have the validity of the following

Lemma 2.2. *There is a constant C independent of λ such that for any least energy solution u_λ of (1.2)*

$$\|u_\lambda\|_{C^{1,\alpha}(\Omega_\lambda)} \leq C.$$

We end this section by establishing uniform exponential decay on the least energy solutions u_λ to (1.2).

Lemma 2.3. *There exist positive constants α, β such that*

$$u_\lambda(y) \leq \alpha \exp(-\beta|y - y^\lambda|)$$

for all $y \in \Omega_\lambda$ and λ sufficiently large. Here y^λ denotes any maximum point of u_λ .

Proof. First, we will see that the functions u_λ decay uniformly at infinity, namely that given $\varepsilon > 0$ there is an $R > 0$ such that $u_\lambda(y) < \varepsilon$ whenever $|y - y^\lambda| > R$. By contradiction, let us assume that for some $\varepsilon > 0$ there are sequences $\lambda_n \rightarrow \infty$ and $y^n \in \Omega_{\lambda_n}$ such that $|y^{\lambda_n} - y^n| \rightarrow \infty$ and $u_{\lambda_n}(y_n) \geq \varepsilon$. We claim that under these conditions,

$$(2.17) \quad \liminf_{n \rightarrow \infty} J_{\lambda_n}(u_n) \geq c_*$$

which is a contradiction with relation (2.16). Since u_{λ_n} is uniformly bounded in $C^{1,\alpha}$, we may assume passing to a subsequence that $u_{\lambda_n}(y^{\lambda_n} + y) \rightarrow u(y)$ uniformly

over compacts, where, thanks to our contradiction assumption, u is a positive solution of (1.3) or, after a rotation and translation, a positive solution of (1.4), depending on whether the distance $\text{dist}(y_{\lambda_n}, \partial\Omega_{\lambda_n}) \rightarrow +\infty$ or remained uniformly bounded. Let us introduce the notation

$$(2.18) \quad J_\Lambda(v) = \frac{1}{p} \int_\Lambda |\nabla v|^p + v^p - \frac{1}{q} \int_\Lambda v^q.$$

Then, given $\delta > 0$ we have that for all sufficiently large R ,

$$(2.19) \quad \lim_{n \rightarrow \infty} J_{B_R(y^{\lambda_n}) \cap \Omega_{\lambda_n}}(u_{\lambda_n}) \geq \frac{c_*}{2} - \delta.$$

Similarly, for all large R ,

$$(2.20) \quad \lim_{n \rightarrow \infty} J_{B_R(y^n) \cap \Omega_{\lambda_n}}(u_{\lambda_n}) \geq \frac{c_*}{2} - \delta.$$

Let us consider $R > 0$ and a smooth cut-off function η_R^n such that $\eta_R^n \equiv 0$ on $B_{R-1}(y^{\lambda_n}) \cup B_{R-1}(y^n)$, $0 \leq \eta_R^n \leq 1$, $\eta_R^n \equiv 1$ on $\mathbb{R}^N \setminus (B_R(y^{\lambda_n}) \cup B_R(y^n))$, and $|\nabla \eta_R^n| \leq C$, C independent of R and n .

We use $w_n = \eta_R^n u_{\lambda_n}$ as a test function for $J'_{\lambda_n}(u_{\lambda_n}) = 0$ to obtain

$$(2.21) \quad 0 = J'_{\lambda_n}(u_{\lambda_n})w_n = E_n + pJ_{\mathbb{R}^N \setminus B_R(y^{\lambda_n}) \cup B_R(y^n)}(u_{\lambda_n}) + \int_{\Omega_{\lambda_n}} g_n$$

where $g_n = (\frac{p}{q} - 1)u_{\lambda_n}^q \eta_R^n \leq 0$ and E_n is given by

$$E_n = \int_{A_R^n} |\nabla u_{\lambda_n}|^{p-2} \nabla u_{\lambda_n} \nabla (\eta_R^n u_{\lambda_n}) + \eta_R^n u_{\lambda_n}^p$$

where $A_R^n = \{y \in \Omega_{\lambda_n} / R - 1 < |y - y^{\lambda_n}| < R \text{ or } R - 1 < |y - y^n| < R\}$. The convergence of u_{λ_n} in the C^1 -sense over compacts around y_n and y^{λ_n} to functions in $W^{1,p}(\mathbb{R}^N)$ implies that for $R > 0$ sufficiently large $\lim_{n \rightarrow \infty} |E_n| \leq \delta$. It follows that for large enough R , $J_{\mathbb{R}^N \setminus B_{R+1}(y^{\lambda_n}) \cup B_{R+1}(y^n)}(u_{\lambda_n}) \geq -\delta$. Using this together with relations (2.19) and (2.20), (2.17) follows.

The desired exponential decay will be a consequence of the following

Claim. There exists $R_0 > 0$ and $\nu_0 > 0$ such that for all $R > R_0$

$$\sup_{|y-y^\lambda| \geq R} u_\lambda(y) \geq 2 \sup_{|y-y^\lambda| \geq R+\nu_0} u_\lambda(y)$$

for all λ sufficiently large.

By contradiction, let us assume that there exist sequences $R_n \rightarrow \infty$, $\bar{\lambda}_n \rightarrow \infty$, $\nu_n \rightarrow \infty$ and $\bar{y}^n \in \Omega_{\bar{\lambda}_n}$ with $|\bar{y}^n - y^{\bar{\lambda}_n}| \geq R_n + \nu_n$ such that $u_{\bar{\lambda}_n}(\bar{y}^n) = \mu_n > \frac{1}{2}M_n$, where

$$M_n = \sup_{|y-y^{\bar{\lambda}_n}| > R_n} u_{\bar{\lambda}_n}(y).$$

From the uniform decay established above, we see that $M_n, \mu_n \rightarrow 0$. Let us set $v_n(y) = \mu_n^{-1} u_{\bar{\lambda}_n}(y + \bar{y}^n)$. Then v_n is bounded, $v_n(0) = 1$ and satisfies

$$\Delta_p v_n - (1 - |u_{\bar{\lambda}_n}(y + \bar{y}^n)|^{q-p})v_n^{p-1} = 0 \quad \text{in } \Omega_{\bar{\lambda}_n} - \bar{y}^n,$$

with $\frac{\partial v_n}{\partial \nu} = 0$ on the boundary. Letting $n \rightarrow \infty$ we obtain a contradiction since v_n converges locally uniformly to a positive bounded solution v of the limiting problem $\Delta_p v - v^{p-1} = 0$, $v(0) = 1$, in the entire space \mathbb{R}^N or in the half-space \mathbb{R}_+^N , with $\frac{\partial v}{\partial \nu} = 0$ on \mathbb{R}^{N-1} . This object does not exist. Indeed, let us consider a sequence

x_n with $v(x_n) \rightarrow \sup_{x \in \mathbb{R}^N} v(x) \geq 1$, and define $w_n(x) = v(x_n + x)$. Since w_n is uniformly bounded, it follows that w_n also has a uniform $C^{1,\alpha}$ bound. Thus passing to a subsequence, we may assume that $w_n \rightarrow w$ uniformly over compacts, where again w solves $\Delta_p w - w^{p-1} = 0$, but now $w(0) = \max_{x \in \mathbb{R}^N} w(x)$. The strong maximum principle in [12] yields a contradiction. Thus, the claim holds and we have the uniform exponential decay of the functions u_λ . By standard arguments involving local elliptic estimates, we obtain the validity of the same property for the derivatives of u_λ . \square

Remark 2.1. Let y^λ be a point where u_λ reaches its maximum value. Then the above proof also shows that $\text{dist}(y^\lambda, \partial\Omega_\lambda)$ must remain bounded, for otherwise we would end up in the limit with an energy of a level at least c^* .

3. PROOF OF THEOREM 1.1

Let u_λ be a positive least energy solution of (1.2), that is a solution with $\lambda(u_\lambda) = c_\lambda$. Let y^λ be a point in $\bar{\Omega}_\lambda$ where u_λ maximizes, and $x^\lambda \in \partial\Omega$ is the closest point in $\partial\Omega$ to $\lambda^{-1}y^\lambda$. From Remark 2.1, we then see that $|\lambda x^\lambda - y^\lambda|$ remains uniformly bounded. Passing to a subsequence, we assume that $x^\lambda \rightarrow \bar{x} \in \partial\Omega$.

Let us set $v_\lambda(y) = u_\lambda(\lambda x_\lambda + y)$. Then v_λ converges in the $W^{1,p}$ -sense to w_* . Moreover, for certain positive constants a and b we have $v_\lambda(y) \leq ae^{-b|y|}$.

After a rotation and a translation λ -dependent we may also assume that $x^\lambda = 0$ and that Ω can be described in a fixed neighborhood V of \bar{x} as the set $\{(x', x_N) \mid x_N > G_\lambda(x')\}$ where G_λ is smooth, $G_\lambda(0) = 0$ and $G'_\lambda(0) = 0$. Further, we may also assume that G_λ converges locally in a C^2 -sense to G , a corresponding parametrization at \bar{x} .

Let us also set $\tilde{\Omega}_\lambda = \Omega_\lambda - \lambda x_\lambda$. From the variational characterization of c_λ , we have that

$$J_\lambda(u_\lambda) \geq J_\lambda(tu_\lambda) = J_{\tilde{\Omega}_\lambda}(tv_\lambda)$$

for all $t > 0$. Let us define the function \tilde{v}_λ on $\lambda(\mathbb{R}_+^N \cap V)$ as $\tilde{v}_\lambda(y', y_N) = v_\lambda(y', y_N)$ if $G_\lambda(y'/\lambda) > 0$ and $\tilde{v}_\lambda(y', y_N) = v_\lambda(y', \lambda G_\lambda(y'/\lambda))$ if $G_\lambda(y'/\lambda) \leq 0$. Then

$$J_{\tilde{\Omega}_\lambda}(tv_\lambda) \geq J_{\mathbb{R}_+^N \cap V_\lambda}(t\tilde{v}_\lambda) + J_{(\tilde{\Omega}_\lambda \cap V_\lambda) \setminus \mathbb{R}_+^N}(tv_\lambda) - J_{(\mathbb{R}_+^N \cap V_\lambda) \setminus \tilde{\Omega}_\lambda}(t\tilde{v}_\lambda).$$

Let us choose $t = t_\lambda$ so that $J_{\mathbb{R}_+^N \cap V_\lambda}(t\tilde{v}_\lambda)$ maximizes in t . Then, by definition of the number $c(\mathbb{R}_+^N) = c_*/2$ and the exponential decay of v_λ , one gets that

$$(3.1) \quad J_{\mathbb{R}_+^N \cap V_\lambda}(t_\lambda \tilde{v}_\lambda) \geq \frac{c_*}{2} + O(e^{-2\lambda\alpha})$$

for some $\alpha > 0$. Again using the exponential decay of v_λ we obtain

$$(3.2) \quad \begin{aligned} I_1 &= -J_{(\tilde{\Omega}_\lambda \cap V_\lambda) \setminus \mathbb{R}_+^N}(t_\lambda v_\lambda) \\ &= - \int_{B_\lambda} dy' \int_{\lambda G_\lambda(y'/\lambda)_-}^0 \left(\frac{1}{p} |\nabla v_\lambda|^p + \frac{1}{p} v_\lambda^p - \frac{1}{q} v_\lambda^q \right) (y', y_N) dy_N + O(e^{-2\lambda\alpha/}), \end{aligned}$$

where $B_\lambda = \{|y'| < \lambda\delta\}$. Similarly, we find that

$$(3.3) \quad \begin{aligned} I_2 &= J_{(\mathbb{R}_+^N \cap V_\lambda) \setminus \tilde{\Omega}_\lambda}(t_\lambda \tilde{v}_\lambda) \\ &= \int_{B_\lambda} dy' \int_0^{\lambda G_\lambda(y'/\lambda)_+} \left(\frac{1}{p} |\nabla v_\lambda|^p + \frac{1}{p} v_\lambda^p - \frac{1}{q} v_\lambda^q \right) (y', G_\lambda(y'/\lambda)) dy_N + O(e^{-2\lambda\alpha}). \end{aligned}$$

Here we have denoted $a_+ = \max\{a, 0\}$, $a_- = \min\{a, 0\}$. Now we note that $v_\lambda \rightarrow w$ C^1 -locally with uniform exponential decay. Then since $G_\lambda(0) = 0$ and $G'_\lambda(0) = 0$ and G_λ converges in a C^2 local sense to G , an application of dominated convergence yields

$$\begin{aligned} \lim_{\lambda \rightarrow \infty} \lambda(I_1 + I_2) &= \sum_{i,j=1}^{N-1} \int_{\mathbb{R}^{N-1}} G''_{ij}(0) y'_i y'_j \left(\frac{1}{p} |\nabla w_*|^p + \frac{1}{p} w_*^p - \frac{1}{q} w_*^q \right) (y', 0) dy' \\ (3.4) \qquad \qquad \qquad &= \tilde{\gamma} \sum_{i=1}^{N-1} G''_{ii}(0) = \tilde{\gamma} H(\bar{x}). \end{aligned}$$

Here

$$\tilde{\gamma} = (N-1)\omega_{N-1} \int_0^\infty \left\{ \frac{1}{p} (|w'_*(r)|^p + w_*(r)^p) - \frac{1}{q} w_*(r)^q \right\} dr.$$

Using $r^2 w'_*(r)$ as a test function in the equation satisfied by w_* , one obtains

$$\tilde{\gamma} = 2 \frac{N-1}{N+1} \omega_{N-1} \int_0^\infty r^N |w'_*(r)|^p dr.$$

We conclude that

$$c_\lambda \geq \frac{c_*}{2} - \lambda^{-1} \gamma H(\bar{x}) + o(\lambda^{-1}).$$

On the other hand, using a computation along the same lines as above, refining estimate (2.16) yields

$$c_\lambda \leq \sup_{t>0} J_\lambda(tw_\lambda^z) = \frac{c_*}{2} - \lambda^{-1} \gamma H(z) + o(\lambda^{-1}),$$

for any $z \in \partial\Omega$. Here $w_\lambda^z = w_*(\lambda(x - x_\lambda))$.

Combining these two estimates directly provides assertions (1) and (2) of the theorem, since in particular we conclude $H(\bar{x}) \geq H(z)$ for all $z \in \partial\Omega$. Finally, relation (2.4) yields the desired expansion (1.5) for $S(\Omega_\lambda)$, and the proof of Theorem 1.1 is concluded.

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