THE SECOND VARIATIONAL FORMULA
FOR THE FUNCTIONAL $\int u^{(6)}(g)dV_g$

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Abstract. In this paper, we compute the second variational formula for the functional $\int u^{(6)}(g)dV_g$, which was introduced by Graham-Juhl; the first variational formula was obtained by Chang-Fang. We also prove that an Einstein manifold (with dimension $\geq 7$) is a strict local maximum within its conformal class unless the manifold is isometric to a round sphere with the standard metric up to a multiple of a constant.

1. Introduction

In the following, we let $(M^n, g)$ denote a compact, connected, smooth Riemannian manifold without boundary. We denote the Ricci curvature and scalar curvature by $\text{Ric}$ and $R$, respectively. Recall that the Schouten tensor $P_{ij}$ is defined by

$$P_{ij} = \frac{1}{n-2} \left( R_{ij} - \frac{R}{2(n-1)} g_{ij} \right),$$

and the Riemann curvature tensor can be written by

$$\text{Riem} = W + P \odot g,$$

where $W$ is the Weyl curvature and $\odot$ is the Kulkarni-Nomizu product, which is defined by

$$(\alpha \odot \beta)_{ijkl} = \alpha_{ik} \beta_{jl} + \alpha_{jl} \beta_{ik} - \alpha_{il} \beta_{jk} - \alpha_{jk} \beta_{il}, \quad \forall \text{ symmetric 2-tensors } \alpha, \beta.$$

The $\sigma_k(g)$ curvature is defined to be the $k$-th elementary symmetric polynomial of the eigenvalue of the Schouten tensor $P$. In [7], Viaclovsky began a study of the variational problems of the functional $\int_M \sigma_k(g)dv_g$. He proved that the first variation of the functional $\int_M \sigma_k(g)dv_g$ ($k = 1, 2$) within a conformal class subject to the constraint $\text{Vol}(M, g) = 1$ is a metric satisfying $\sigma_k(g) \equiv \text{const}$, and if $k \geq 3$ and the Riemannian manifold is locally conformally flat, the same result follows. However, for $k \geq 3$ and for a manifold that is not locally conformally flat, $\sigma_k(g) \equiv \text{const}$ is not the Euler-Lagrange equation of the functional $\int_M \sigma_k dv$ within a conformal class subject to the constraint $\text{Vol}(M, g) = 1$.

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The renormalized volume coefficients of $g$, denoted here by $v^{(2k)}(g)$, arose in the late 1990s in the physics literature. They are defined in terms of the expansion of the ambient or Poincaré metric associated to $g$. If the Riemannian manifold is locally conformally flat, these quantities coincide with the $\sigma_k(g)$ up to a constant. More precisely, it is known that (see [3], [1], or [4])

$$v^{(2)}(g) = -\frac{1}{2}\sigma_1(g), \quad v^{(4)}(g) = \frac{1}{4}\sigma_2(g),$$

$$v^{(6)}(g) = -\frac{1}{8}[\sigma_3(g) + \frac{1}{3(n-4)}(P_g)^{ij}(B_g)_{ij}],$$

where

$$(B_g)_{ij} := \frac{1}{n-3}\nabla^k
abla^l W_{ilkj} + \frac{1}{n-2}R^{kl}W_{ilkj}$$

is the Bach tensor of the metric. Just as $\int_M \sigma_k(g^{-1}\circ A_g)\,dv_g$ is conformally invariant when $2k = n$ and $(M, g)$ is locally conformally flat, Graham showed in [2] that $\int_M v^{(2k)}(g)\,dv_g$ is also conformally invariant on a general manifold when $2k = n$. Chang and Fang showed in [1] that, for $n \neq 2k$, the Euler-Lagrange equations for the functional $\int_M v^{(2k)}(g)\,dv_g$ under conformal variations subject to the constraint $Vol_g(M) = 1$ satisfy $v^{(2k)}(g) = \text{const.}$, which is a generalized characterization for the curvatures $\sigma_k(g^{-1}\circ A_g)$ when $(M, g)$ is locally conformally flat, as given by Viaclovsky [7].

We note that Graham [2] also gives an explicit expression of $v^{(8)}(g)$, but the explicit expression of $v^{(2k)}(g)$ for general $k$ is not known because they are algebraically complicated (see page 1958 of [2]). Thus the study of the $v^{(2k)}(g)$ curvatures involves significant challenges not shared by that of $\sigma_k(g)$: firstly, for $k \geq 3$, $v^{(2k)}(g)$ depends on derivatives of the curvature of $g$; in fact, for $k \geq 3$, $v^{(2k)}(g)$ depends on derivatives of curvatures of order up to $2k - 4$. Secondly, the $v^{(2k)}(g)$ are defined via an indirect highly nonlinear inductive algorithm (see [2]). We aim to study the stability of the critical metric of the functional

$$\mathcal{F}_3[g] = \frac{\int_M v^{(6)}(g)dv_g}{(\int_M dv_g)^{(n-6)/n}}$$

within a conformal class. First we recall the theorem of Chang-Fang [1] (also see Graham [2]).

**Theorem 1.1** ([1]). Let $(M^n, g)$ be an $n$-dimensional ($n \geq 7$) compact Riemannian manifold. Then the functional $\mathcal{F}_3(g)$ is variational within the conformal class; i.e. the critical metric in $[g]$ satisfies the equation

$$v^{(6)} \equiv \text{const.}$$

If $n = 6$, $\mathcal{F}_3[g]$ is a constant in the conformal class $[g]$.

In this paper, we compute the second variational formula of $\mathcal{F}_3[g]$ within its conformal class $[g]$. Our results are

**Theorem 1.2.** Let $(M^n, g)$ be an $n$-dimensional ($n \geq 7$) compact Riemannian manifold with $v^{(6)}(g) = \text{const}$. Then the second variational formula of the functional
The second variational formula for $\int u^{(6)}(g) dv_g$

$F_3[g]$ within its conformal class at $g$ is

$$\frac{d^2}{dt^2} F_3[g_t] = \left[(n-6)V^{-(n-6)/n}\sum_{i} g_{ij}(\nabla_i \phi)^2 + \frac{1}{8} T_{2ij} \phi_{ij} - \frac{1}{12} P_{ij} C_{ijk} \phi_k \right] dv_g,$$

where $g_t = e^{2u_t} g$, $\frac{\partial}{\partial t} |_{t=0} u_t = \phi$, and $\phi = \phi - \int_M \phi^2 dv_g / \int_M dv_g$, $T_{2ij}$ and $C_{ijk}$ are defined in section 2.

Theorem 1.3. Let $(M^n, g)$ be an $n$-dimensional ($n \geq 7$) compact Einstein manifold. Then it is a strict local maximum within its conformal class $[g]$ unless $(M^n, g)$ is isometric to $S^n$ with the standard metric up to a multiple of a constant.

Remark 1.4. When $(M^n, g)$ is locally conformally flat, $u^{(6)}(g) = -\frac{1}{8} \sigma_3(g)$, for the functional $\int_M \sigma_3(g) dv_g$, J. Viaclovsky [7] proved that a positive constant sectional curvature metric is a strict local minimum, unless the manifold is isometric to $S^n$ with the standard metric. Our result coincides with his at the locally conformally flat Einstein metrics; however, ours does not need the locally conformally flat assumption.

2. Preliminaries

Let $(M^n, g)$ be an $n$-dimensional compact Riemannian manifold. Throughout this paper, we use the convention that a repeated index means summation over 1 to $n$. First we recall the transformation law of various curvatures under a conformal change of metrics. Let $\tilde{g} = e^{2u} g$, $u \in C^\infty(M)$. Then the Riemannian curvature tensors satisfy

$$Riem(\tilde{g}) = e^{2u}(Riem(g) - \alpha \odot g),$$

where $\alpha_{ij} = u_{ij} - u_i u_j + \frac{1}{2} u^2 g_{ij}$ (note that $u_{ij}$ means the covariant derivative with respect to the fixed metric $g$). By contracting, we see that the Ricci curvature and scalar curvature satisfy

$$R_{ij}(\tilde{g}) = R_{ij} - (n-2) \alpha_{ij} - \left( \sum_k \alpha_{kk} \right) g_{ij}, \quad R(\tilde{g}) = e^{-2u} R - 2(n-1) e^{-2u} \sum_k \alpha_{kk}.$$

From (2.1) and the definition of the Schouten tensor, we see that

$$\tilde{P}_{ij} = P_{ij} - \alpha_{ij},$$

where we denote $P(\tilde{g})$ by $\tilde{P}$ for convenience of notation.

Lemma 2.1. We have the following formulae (see e.g. [1]):

1. $\nabla^i W_{ijkl} = -(n-3)C_{ijkl}$, where $C_{ijkl}$ is the Cotton tensor defined by $P_{ij,k} - P_{ik,j}$;
2. $\nabla^i B_{ij} = (n-4) \sum_{k,l} P_{kl} C_{kl};$
3. $B_{ij} = B_{ji}$, where $B_{ij}$ is defined by (1.1).

The proof of Lemma 2.1 is a direct calculation, and one can find it in [1].
Let $V$ be a vector space, $A : V \to V$ a linear map. Define the Newton transformation $T_k(A) : V \to V$ by

$$T_k(A) := \sigma_k(A)I - \sigma_{k-1}(A)A + \cdots + (-1)^kA^k = \sum_{i=0}^{k} \sigma_{k-i}(A)(-1)^iA^i,$$

where $I$ is the identity map and $\sigma_k(A)$ is the $k$-th elementary symmetric polynomial of the eigenvalues of $A$. Under an orthonormal basis of $V$, $T_k$ can be written as

$$T_{kij} = \frac{1}{k!} \delta^{j_1 \ldots j_k}_{i_1 \ldots i_k} A_{i_1 j_1} \cdots A_{i_k j_k},$$

where $\delta^{j_1 \ldots j_k}_{i_1 \ldots i_k}$ is the generalized Kronecker notation. We recall some well-known results in this respect, which we will need in our later arguments.

**Lemma 2.2.** The Newton transformations $T_k$ satisfy (5), (4):

1. Newton’s formula: $(k + 1)\sigma_{k+1}(A) = \text{tr}(T_kA)$;
2. $\frac{d}{dt} \sigma_k(A_t) = \text{tr}(T_{k-1}\frac{d}{dt}A_t)$, for any family of transformations $A_t : V \to V$;
3. $\text{tr}(T_k) = (n - k)\sigma_k(A)$.

In the following we denote $T_k(g^{-1} \circ P)$ simply by $T_k$. We have the following formula, which is a direct calculation (see [4]):

$$\sum_{i} T_{2ij,i} = -\sum_{k,l} P_{kl}C_{klj}.$$ 

### 3. The First Variational Formula and Proof of Theorem 1.1

In this section, we will compute the Euler-Lagrange equation for the functional $\mathcal{F}_3(g)$ within the conformal class. For convenience we denote the numerator of $\mathcal{F}_3(g)$ by

$$F(g) = \int_M v^{(6)}(g)dv_g.$$

Under the conformal change of metrics $g_t = e^{2u(t)}g$, by use of (2.2), we see that in local coordinates (see [4])

$$\delta \widehat{P}_i^j = e^{-2u(t)}(P_i^j - \alpha_i^j),$$

$$\delta \widehat{B}_i^j = e^{-4u(t)}\left(B_i^j + (n - 4)u^k(g^{jl}C_{ilk} + g^{jl}C_{ljk}) + (n - 4)u^kg^pjW_{ikpl}\right),$$

where we write $\alpha_{ij} = u_{ij}(t) - u_i(t)u_j(t) + \frac{1}{2}|\nabla g u(t)|^2g_{ij}$, and we use the convention that $\widehat{P}_i^j = P_{ij}(g_t)$, $\widehat{B}_i^j = B_{ij}(g_t)$, etc., for convenience of notation.

For convenience of notation we denote $\frac{d}{dt}$ by $\delta$. Denote $\frac{d}{dt}\bigg|_{t=0} u = \phi$ and $\frac{d^2}{dt^2}\bigg|_{t=0} u = \psi$. With the above preparations, we have

$$\delta \widehat{P}_i^j = -2(\delta u)\widehat{P}_i^j - e^{-2u}\delta \alpha_i^j,$$

$$\delta \widehat{B}_i^j = -4(\delta u)\widehat{B}_i^j + (n - 4)e^{-4u}\left((\delta u)^k(g^{jl}C_{ilk} + g^{jl}C_{ljk}) + (\delta u)^kg^pjW_{ikpl}\right) + (\delta u)^kp^jg^{pj}W_{ikpl} + u^k(\delta u)^lg^{pj}W_{ikpl},$$

$$\delta (dv_g) = n(\delta u)dv_g.$$
Now we derive the first variation formula for \( F \). First we have

\[
-8\delta F = \int \delta(\sigma_3(g_\ast)dv_{g_\ast}) + \frac{1}{3(n-4)} \delta(\tilde{P}_i^j \tilde{B}_j^i dv_{g_\ast}).
\]

By use of Lemma 2.2, we have

\[
\delta(\sigma_3(g_\ast)) = T_{2i}^j \delta \tilde{P}_i^j = T_{2i}^j (-2(\delta u) \tilde{P}_i^j - e^{-2u} \delta \alpha_i^j)
\]

\[
= -6(\delta u)\sigma_3(g_\ast) - e^{-2u} T_{2i}^j \delta \alpha_i^j.
\]

Hence

\[
\int \delta(\sigma_3(g_\ast)dv_{g_\ast}) = \int \left[ -6(\delta u)\sigma_3(g_\ast) - e^{-2u} T_{2i}^j \delta \alpha_i^j + n(\delta u)\sigma_3(g_\ast) \right] dv_{g_\ast}
\]

\[
= \int \left[ (n-6)(\delta u)\sigma_3(g_\ast) - e^{-2u} T_{2i}^j \delta \alpha_i^j \right] dv_{g_\ast}.
\]

On the other hand, the second term of (3.1) is

\[
\int \frac{1}{3(n-4)} \delta(\tilde{P}_i^j \tilde{B}_j^i dv_{g_\ast})
\]

\[
= \int \frac{1}{3(n-4)} \left[ \delta \tilde{P}_i^j \tilde{B}_j^i + \tilde{P}_i^j \delta \tilde{B}_j^i + n(\delta u)\tilde{P}_i^j \tilde{B}_j^i \right] dv_{g_\ast}
\]

\[
= \int \frac{1}{3(n-4)} \left[ \tilde{B}_j^i \left( -2(\delta u) \tilde{P}_i^j - e^{-2u} \delta \alpha_i^j \right) \right.
\]

\[
+ \tilde{P}_i^j \left( -4(\delta u) \tilde{B}_j^i + (n-4)e^{-4u} ((\delta u)^k (g^{il}C_{ijkl} + C_{ijkl})
\]

\[
+ 2(\delta u)^k u^l g^{plW_{pkjl}}) \right) + n(\delta u)\tilde{P}_i^j \tilde{B}_j^i \right] dv_{g_\ast}
\]

\[
= \int \frac{1}{3(n-4)} \left[ (n-6)(\delta u)\tilde{P}_i^j \tilde{B}_j^i - e^{-2u} \tilde{B}_j^i (\delta \alpha_i^j) \right.
\]

\[
+ 2(n-4)e^{-4u} \tilde{P}_i^j \left( (\delta u)^k g^{il}C_{ijkl} + (\delta u)^k u^l g^{plW_{pkjl}} \right) \right] dv_{g_\ast}.
\]

From calculations in (3.2) and (3.3), we have the following formula, which will be used in section 4:

\[
\delta \left( v^{(6)}(g_\ast)dv_{g_\ast} \right)
\]

\[
= -\frac{1}{8} \left[ (n-6)(\delta u)\sigma_3(g_\ast) - e^{-2u} T_{2i}^j (\delta \alpha_i^j) \right] dv_{g_\ast} - \frac{1}{24(n-4)} \left[ (n-6)(\delta u)\tilde{P}_i^j \tilde{B}_j^i \right.
\]

\[
- e^{-2u} \tilde{B}_j^i (\delta \alpha_i^j) + 2(n-4)e^{-4u} \tilde{P}_i^j \left( (\delta u)^k g^{il}C_{ijkl} + (\delta u)^k u^l g^{plW_{pkjl}} \right) \right] dv_{g_\ast}.
\]
Thus we have

\[
\delta F = \int \delta(v^6) g_t dv_{g_t}
\]

(3.5)

\[
= -\frac{1}{8} \int \left\{ (n-6)(\delta u)\sigma_3(g_t) - e^{-2u} \tilde{T}^i_{2i} \delta \alpha_i^j + \frac{1}{3(n-4)} [(n-6)(\delta u) \tilde{B}^i_{j} \tilde{B}^j_i - e^{-2u} \tilde{B}^i_{j} (\delta \alpha_i^j)^2 + 2(n-4) e^{-4u} \tilde{P}^i_{j} \left( (\delta u)^k g^i j C_{ijk} + (\delta u)^k u^l g^j p W_{p k j l} \right)] \right\} dv_{g_t}
\]

(3.6)

Proof of Theorem 1.1. Noting that \( u(0) = 0 \), we conclude that the first variational formula of \( F_3[g_t] \) within the conformal class \([g]\) is (see [1] or [2])

\[
\frac{d}{dt} \left|_{t=0} \right. F_3(g_t) = \frac{d}{dt} \left|_{t=0} \right. F \cdot V^{-\frac{n-6}{n}} - \frac{n-6}{n} V^{-\frac{n-6}{n}} \left( \int v^6 dv \right) \int n \phi dv
\]

\[
= V^{-\frac{n-6}{n}} \left\{ (n-6) \int \phi v^6(g) + \frac{1}{8} \int T_{2ij} \phi_{ij} + \frac{1}{8} \int \frac{B_{ij} \phi_{ij}}{3(n-4)} - \frac{1}{12} \int P_{ij} \phi_{ij} C_{ijk} - (n-6)V^{-1} \left( \int v^6(g) \right) \int \phi dv \right\}
\]

\[
= (n-6) V^{-\frac{n-6}{n}} \left\{ \int \phi \left( v^6 - V^{-1} \int v^6 dv \right) dv \right\},
\]

where we have used (2.3) and (2) of Lemma 2.1 and integration by parts. Here \( V = \int dv_g \). Hence, we see that the Euler-Lagrange equation of the functional \( F_3(g) \) within the conformal class \([g]\) is

\[
v^6(g) = V^{-1} \int v^6(g) dv_g \equiv \text{const},
\]

and we get Theorem 1.1. \( \square \)

4. The second variational formula and proofs of Theorems 1.2 and 1.3

In this section, we will calculate the second variational formula for the functional \( F_3 \) within the conformal class \([g]\). For convenience, we separate each term in the first variational equation (3.5) and compute them respectively.
For the derivative of the first term in (3.4), by use of (3.4), we have

\begin{align*}
(4.1) \quad (n-6) \frac{d}{dt} \bigg|_{t=0} (\delta u) v^{(6)} (g_t) dv_{g_t},
= (n-6) \int \left\{ \psi v^{(6)}(g) - \frac{1}{8} (\delta u) \left[ (n-6)(\delta u) g_3 (g_t) + e^{-2u} \hat{T}^j_i \delta \alpha^i_j \right] dv_{g_t} \bigg|_{t=0} \\
- \frac{(\delta u)}{24(n-4)} \left[ (n-6)(\delta u) \hat{P}^i_j \hat{B}^j_i - e^{-2u} \hat{B}^j_i \delta \alpha^i_j \right] + 2(n-4) e^{-4u} \hat{P}^i_j \left( (\delta u) k^l g^i_{ljk} + (\delta u) k^l u^i g^j_{ljk} \right) \bigg|_{t=0} \right. \\
= (n-6) \int \left\{ \psi v^{(6)}(g) + (n-6) \phi^2 v^{(6)}(g) - \frac{1}{8} \left[ - \phi T_{2ij} \phi^i_j \\
+ \frac{\phi B_{ij} \phi_{ij}}{3(n-4)} + \frac{2}{3} \phi \phi_{ij} P_{ij} C_{ij} \right] \right\} dv.
\end{align*}

For the derivative of the second term in (3.5), we need the following formula of the variation of the Newton transformation:

\begin{align*}
(4.2) \quad \frac{d}{dt} \bigg|_{t=0} (\hat{T}^j_i) = \frac{1}{2} \sum_{i} \frac{\delta^{j}_{i12} \hat{P}^i_{1j} \hat{P}^i_{2j}}{d \xi_{i12}^j} = \delta^{j}_{i12} \hat{P}^i_{1j} \hat{P}^i_{2j} \\
= \delta^{j}_{i12} \hat{P}^i_{1j} \hat{P}^i_{2j} (-2 \phi P_{1ij} - \phi_{12j}) = -4 \phi T_{2ij} - \delta^{j}_{i12} \hat{P}^i_{1j} \phi_{12j}.
\end{align*}

Therefore, the variation of the second term of (3.5) is given by

\begin{align*}
(4.3) \quad \frac{1}{8} \frac{d}{dt} \bigg|_{t=0} \int e^{-2u} \hat{T}^j_i \delta \alpha^i_j dv_{g_t} = \frac{1}{8} \int \left\{ -2 \phi T_{2ij} \phi^i_j + \frac{d}{dt} \bigg|_{t=0} \hat{T}^j_i \phi^i_j \\
+ T_{2ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} + n \phi T_{2ij} \phi_{ij} \right\} dv \\
= \frac{1}{8} \int \left\{ (n-2) \phi T_{2ij} \phi_{ij} + T_{2ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} + \phi_{ij} \left( -4 \phi T_{2ij} - \delta^{j}_{i12} \hat{P}^i_{1j} \phi_{12j} \right) \right\} dv \\
= \int \left\{ \frac{(n-6)}{8} \phi T_{2ij} \phi_{ij} + \frac{1}{8} T_{2ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} - \frac{1}{8} \delta^{j}_{i12} \hat{P}^i_{1j} \phi_{12j} \phi_{ij} \right\} dv.
\end{align*}

The variation of the third term of (3.5) is

\begin{align*}
(4.4) \quad \frac{1}{24(n-4)} \frac{d}{dt} \bigg|_{t=0} \int e^{-2u} \hat{B}^j_i \delta \alpha^i_j dv_{g_t},
= \frac{1}{24(n-4)} \int \left\{ -2 \phi B_{ij} \phi_{ij} + \frac{d}{dt} \bigg|_{t=0} \hat{B}^j_i \phi_{ij} + B_{ij} \frac{d}{dt} \bigg|_{t=0} \delta \alpha_{ij} + n \phi B_{ij} \phi_{ij} \right\} dv
\end{align*}
\[
\frac{1}{24(n-4)} \int \left\{ (n-2)\phi B_{ij}\phi_{ij} + B_{ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} \\
+ \phi_{ij} \left( -4\phi B_{ij} + 2(n-4)\phi_k C_{ijk} \right) \right\} dv
\]

\[
= \int \left\{ \frac{n-6}{24(n-4)} \phi B_{ij}\phi_{ij} + \frac{B_{ij}}{24(n-4)} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} + \frac{1}{12} \phi_k \phi_{ij} C_{ijk} \right\} dv.
\]

The variation of the fourth term of (3.5) is

\[
\frac{1}{12} \frac{d}{dt} \bigg|_{t=0} \int \left[ e^{-4u} \tilde{P}_{ij}^k \left( (\delta u)^k g^{jl} C_{ilk} + (\delta u)^k d^j W_{ikpl} \right) \right] dv.
\]

\[
= -\frac{1}{12} \int \left\{ -4\phi P_{ij} \phi_k C_{ijk} + \frac{d}{dt} \bigg|_{t=0} \tilde{P}_{ij}^k \phi^l g^{jl} C_{ijk} \\
+ P_{ij} \psi_k C_{ijk} + P_{ij} \phi_k \phi_l W_{ikjl} + n\phi P_{ij} \phi_k C_{ijk} \right\} dv
\]

\[
= -\frac{1}{12} \int \left\{ -(n-6)\phi P_{ij} \phi_k C_{ijk} - \frac{1}{12} P_{ij} \phi_k \phi_l W_{ikjl} + \frac{1}{12} \phi_k \phi_{ij} C_{ijk} - \frac{1}{12} P_{ij} \psi_k C_{ijk} \right\} dv.
\]

Combining (4.1), (4.3), (4.4) and (4.5), we have

\[
\frac{d^2}{dt^2} \bigg|_{t=0} F(g_t)
\]

\[
= \int \left\{ (n-6)\psi v^{(6)}(g) + (n-6)^2 \phi^2 v^{(6)}(g) + \frac{n-6}{4} \phi T_{2ij}\phi_{ij} - \frac{1}{8} \left[ -\phi B_{ij} \phi_{ij} \right] + \frac{4(n-6)}{3} \phi_k \phi_{ij} C_{ijk} - T_{2ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} - \frac{B_{ij}}{3(n-4)} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} \\
+ \delta_{ij}^k \tilde{P}_{ij}^l \phi_{ijl} \phi_{kij} - \frac{1}{3} \phi_k \phi_{ij} C_{ijk} + \frac{2}{3} P_{ij} \phi_k \phi_l W_{ikjl} + \frac{2}{3} P_{ij} \psi_k C_{ijk} \right\} dv.
\]
Since \( \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} = \psi_{ij} - 2\phi_i\phi_j + |\nabla \phi|^2 g_{ij} \), by use of the divergence theorem we obtain

\[
(4.7) \quad - \int T_{2ij} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij} - \int \frac{B_{ij}}{3(n-4)} \frac{d^2}{dt^2} \bigg|_{t=0} \alpha_{ij}
\]

\[
= - \int T_{2ij}(\psi_{ij} - 2\phi_i\phi_j + |\nabla \phi|^2 g_{ij}) - \int \frac{B_{ij}(\psi_{ij} - 2\phi_i\phi_j + |\nabla \phi|^2 g_{ij})}{3(n-4)}
\]

\[
= \int -\frac{2}{3} P_{kl} C_{kli} \psi_i + 2T_{2ij} \phi_i \phi_j + \frac{2}{3(n-4)} B_{ij} \phi_i \phi_j - |\nabla \phi|^2 T_{2kk}
\]

\[
= \int -\frac{2}{3} P_{kl} C_{kli} \psi_i - |\nabla \phi|^2 T_{2kk} - 2T_{2ij,j} \phi_i \phi_j - 2T_{2ij} \phi_i \phi_j
\]

where we have used the following identity in the second equality:

\[
\int T_{2ij} \psi_{ij} dv + \frac{1}{3(n-4)} \int B_{ij} \psi_{ij} dv = \int \frac{2}{3} P_{kl} C_{kli} \psi_i,
\]

which can be checked by use of (2.3), (2) of Lemma 2.1 and integration by parts.

Substituting (4.7) into (4.6) and making some cancellations, we conclude that

\[
(4.8) \quad \frac{d^2}{dt^2} \bigg|_{t=0} F(g_t) = \int \left\{ (n-6)\psi(v^{(6)}) + (n-6)^2 \phi^2 v(v^{(6)}) (g) - \frac{1}{8} \left[ -2(n-5) \phi T_{2ij} \phi_{ij} \right.ight.
\]

\[
- \frac{2}{3(n-4)} \phi B_{ij} \phi_{ij} + \frac{4}{3(n-5)} \phi P_{ij} \phi_k C_{ijk} - \frac{4}{3} \phi_k \phi_{ij} C_{ijk} \right.
\]

\[
\left. + \delta_{kli}^{mnj} P_{km} \phi_{ln} \phi_{ij} + \frac{2}{3} P_{ij} \phi_k \phi_l W_{ikjl} - (n-2)|\nabla \phi|^2 \sigma_2(g) \right\} dv,
\]

where we have used the identity that \( T_{2kk} = (n-2)\sigma_2(g) \) (see Lemma 2.2). It remains to study the last four terms on the right hand side of (4.8). By definition,

\[
\delta_{kli}^{mnj} = \det \begin{pmatrix} \delta_{km} & \delta_{kn} & \delta_{kj} \\ \delta_{lm} & \delta_{ln} & \delta_{lj} \\ \delta_{im} & \delta_{in} & \delta_{ij} \end{pmatrix}
\]

\[
= \delta_{km} \delta_{ln} \delta_{ij} - \delta_{km} \delta_{lj} \delta_{in} - \delta_{lm} \delta_{kn} \delta_{ij} + \delta_{lm} \delta_{in} \delta_{kj} + \delta_{im} \delta_{kn} \delta_{lj} - \delta_{im} \delta_{ln} \delta_{kj}.
\]

We compute by use of the divergence theorem

\[
(4.9) \quad \int \delta_{kli}^{mnj} P_{km} \phi_{ln} \phi_{ij} = \int \phi \left( \delta_{kli}^{mnj} (P_{mk} \phi_{nl})_{ij} \right)
\]

\[
= \int \phi \delta_{kli}^{mnj} \left( P_{km} \phi_{nl} + 2P_{km} \phi_{nl,j} + P_{km} \phi_{nl,i,j} \right).
\]
Now we compute the integrands of the right hand side of (4.9) respectively. The first term is

\[
(4.10)
\]

\[
\phi \delta_{kli}^{mnj} P_{km,ij} \phi_{nl} = \phi P_{km,ij} \phi_{nl} (\delta_{km} \delta_{in} \delta_{ij} - \delta_{km} \delta_{ijn} \delta_{in} - \delta_{lm} \delta_{kn} \delta_{in} \\
+ \delta_{lm} \delta_{kn} \delta_{kj})
\]

\[
= \phi P_{kk,ij} \phi_{nn} - \phi P_{kk,ij} \phi_{ij} - \phi P_{kl,ij} \phi_{kl} + \phi P_{kl,ik} \phi_{il} + \phi P_{kl,ik} \phi_{il}
\]

\[
= \phi \phi_{nn} C_{iik,k} + \phi \phi_{kl} C_{ilk,i} + \phi \phi_{il} C_{iki,l} = \phi \phi_{kl} C_{ilk,i}.
\]

The second term is

\[
(4.11)
2 \phi \delta_{kli}^{mnj} P_{km,i} \phi_{nlj} = 2 \phi P_{kk,i} \phi_{nni} - 2 \phi P_{kk,i} \phi_{ijj} - 2 \phi P_{km,i} \phi_{kmi} \\
+ 2 \phi P_{km,i} \phi_{imk} + 2 \phi P_{km,m} \phi_{kl} - 2 \phi P_{km,m} \phi_{lk}
\]

\[
= 2 \phi P_{km,i} (\phi_{mik} - \phi_{mki}) = 2 \phi P_{km,i} \phi_{j} R_{jmik}.
\]

The third term is

\[
(4.12)
\phi \delta_{kli}^{mnj} P_{km} \phi_{nl,i} = \phi \phi_{kk,ii} P_{nn} - \phi \phi_{kkli} P_{kl} - \phi \phi_{kl} P_{kl} + \phi \phi_{kll} P_{kl} - \phi \phi_{klik} P_{ll}
\]

\[
= \phi P_{nn} (\phi_{kk} - \phi_{kkii}) + \phi P_{kl} (\phi_{kl} - \phi_{kl}) + \phi P_{kl} (\phi_{kl} - \phi_{kl}) \\
= \phi P_{nn} (\phi_{mk} R_{mk} - \phi_{m} R_{mk,k}) + \phi P_{kl} (\phi_{ml} R_{mlk} + \phi_{m} R_{mlk,l}) \\
+ \phi P_{kl} (\phi_{mk} R_{mk} + \phi_{m} R_{mk,l}),
\]

where we have used the Ricci identity in the last equality.

Substituting the identities

\[
R_{ij} = (n - 2) P_{ij} + \frac{R}{2(n - 1)} g_{ij}, \quad R_{ij,i} = \frac{R_{ij}}{2}, \quad P_{kk} = \frac{R}{2(n - 1)};
\]

\[
R_{ijkl} = W_{ijkl} + P_{ik} g_{jl} + P_{jk} g_{il} - P_{ij} g_{kl}, \quad P_{kk,i} = P_{ik,k};
\]

\[
R_{ikml,i} = R_{kl,m} - R_{km,l} = (n - 2) C_{kml} + \frac{\nabla_{m} R_{lkm} - \nabla_{l} R_{gkm}}{2(n - 1)}
\]

into (4.11) and (4.12) and after making some cancellations, we see that the left hand side of (4.9) becomes

\[
(4.13)
\int \phi \delta_{kli}^{mnj} P_{km} \phi_{nl} \phi_{ij}
\]

\[
= \int \left[ - \phi \phi_{kl} C_{ik,i} + \frac{4 - n}{2(n - 1)} \phi R_{mk} P_{mk} - \frac{R^{2}}{4(n - 1)^{2}} \phi \Delta \phi - \frac{n - 2}{4(n - 1)^{2}} \phi R_{m} R_{m} \\
+ \phi P_{kl} \phi_{ml} W_{mlk} + \phi |P_{kl}|^{2} \Delta \phi + (n - 4) \phi P_{kl} \phi_{ml} P_{mk} + n \phi P_{kl} \phi_{ml} P_{kl,m} \\
+ 2 \phi P_{km,i} \phi_{j} W_{jmik} - 2 \phi P_{km,i} \phi_{k} P_{mi} \right].
\]
On the other hand, by the divergence theorem, we see that one of the remaining terms on the last line of (4.8) is

$$
\frac{2}{3} \int P_{ij} \phi_k \phi_l W_{ikjl} = \frac{2}{3} \int -\phi P_{ij,k} \phi_l W_{ikjl} - \phi P_{ij} \phi_{kl} W_{ikjl} - \phi P_{ij} \phi_l W_{ikjl,k} \\
= \int -\frac{2}{3} \phi P_{ij,k} \phi_l W_{ikjl} - \frac{2}{3} \phi P_{ij} \phi_{kl} W_{ikjl} - \frac{2(n - 3)}{3} \phi P_{ij} \phi_l C_{ijkl},
$$

$$
(4.14)
$$

$$
- \frac{4}{3} \int \phi_k \phi_{ij} C_{ijk} = \int \frac{4}{3} \phi \phi_{ij} C_{ijk} + \frac{4}{3} \phi \phi_{ij} C_{ijk,k},
$$

$$
(4.15)
$$

$$
- \int (n - 2) |\nabla \phi|^2 \sigma_2(g) \\
= \int (n - 2) \phi \Delta \sigma_2(g) + (n - 2) \phi \sigma_2(g),
$$

$$
(4.16)
$$

By combining equations (4.13), (4.14), (4.15) and (4.16) and doing some cancellations, we conclude that the last four terms on the right hand side of (4.8) are equal to

$$
\frac{1}{8} \int \left[ \frac{1}{3} \phi_{kl} B_{kl} + (n - 4) \phi T_{2ij} \phi_{ij} - \frac{2(n - 6)}{3} \phi P_{ij} \phi_k C_{ijk} \\
+ \frac{4}{3} \phi C_{km} \phi_j W_{jkml} + \frac{4}{3} \phi \phi_{ij} C_{ijkl} \right] dv,
$$

where we have used $T_{2ij} = \sigma_2(g) \delta_{ij} - \sigma_1(g) P_{ij} + P_{ij} P_{kj}$ and $B_{ij} = C_{ij,k,k} + P_{kl} W_{ikjl}$. Moreover,

$$
\frac{4}{3} \phi C_{ijk} \phi_{jk} = \frac{4}{3} \phi C_{ijk} (\phi_{jki} + \phi_m R_{mjk}) \\
= \frac{4}{3} \phi C_{ijk} \phi_m (W_{mjk} + P_{mi} g_{jk} + P_{jk} g_{mi} - P_{mk} g_{ij} - P_{ij} g_{mk}) \\
= \frac{4}{3} \phi C_{ijk} \phi_m W_{mjk} + \frac{4}{3} \phi C_{ij} \phi_m P_{mi} + \frac{4}{3} \phi C_{ijk} \phi_k P_{jk} \\
- \frac{4}{3} \phi C_{ijk} \phi_m P_{mk} - \frac{4}{3} \phi C_{ijk} \phi_j P_{ik} \\
= -\frac{4}{3} \phi \phi_m C_{ijk} W_{mjk} - \frac{4}{3} \phi C_{ijk} P_{ij} \phi_k,
$$

where we used $\sum_i C_{ijk} = 0$ and $C_{ijk} = -C_{ikj}$. 


Thus it follows that (4.17) is equal to

\[(4.18)\]
\[-\frac{1}{8} \int \left[ \frac{1}{3} \phi \phi_{kl} B_{kl} + (n - 4) \phi \phi_{ij} T_{2ij} - \frac{2(n - 6)}{3} \phi \phi_k P_{ij} C_{ijk} - \frac{4}{3} \phi \phi_k C_{ijk} P_{ij} \right] dv.\]

Substituting (4.18) into (4.8), we conclude that

\[(4.19)\]
\[\frac{d^2}{dt^2} \bigg|_{t=0} F(g_t) = \int \left\{ (n - 6) \psi v^{(6)} + (n - 6)^2 \phi^2 v^{(6)}(g) - \frac{1}{8} \left[ -2(n - 5) \phi \phi_{ij} T_{2ij} - \frac{2(n - 5)}{3(n - 4)} \phi \phi_{ij} B_{ij} + \frac{4(n - 5)}{3} \phi \phi_k P_{ij} C_{ijk} - \frac{4}{3} \phi \phi_k C_{ijk} P_{ij} \right] \right\} \]

\[= \int \left\{ (n - 6) \psi v^{(6)} + (n - 6)^2 \phi^2 v^{(6)} + \frac{1}{8} \left[ (n - 6) \phi \phi_{ij} T_{2ij} + \frac{n - 6}{3(n - 4)} \phi \phi_{ij} B_{ij} - \frac{2(n - 6)}{3} \phi \phi_k P_{ij} C_{ijk} - \frac{4}{3} \phi \phi_k C_{ijk} P_{ij} \right] \right\}.\]

**Proof of Theorem 1.2.** By Theorem 1.1 at the critical metric of the functional \(F_3(g),\) it follows that \(v^{(6)}(g)\) should be constant, and it follows that \(F(g) = V v^{(6)}(g)\). By our notation, \(F_3[g_t] = \frac{F(g_t)}{\left( \int dv_v \right)^{(n-6)/n}}.\) By use of (4.19) and (3.6), at the critical metric \(g,\) we have

\[(4.20)\]
\[\frac{d^2}{dt^2} \bigg|_{t=0} F_3[g_t] = \frac{d^2}{dt^2} \bigg|_{t=0} F(g_t) \cdot V^{\frac{n-6}{n}} + 2 \frac{d}{dt} \bigg|_{t=0} F(g_t) \cdot \frac{d}{dt} \bigg|_{t=0} V(g_t) - \frac{n-6}{n} \int n \phi + F \left\{ (n - 6)(2n - 6) V^{\frac{(3n-6)}{n}} \left( \int \phi \right)^2 - n(n - 6) V^{\frac{(2n-6)}{n}} \int \phi^2 \right\} - (n - 6) V^{\frac{(2n-6)}{n}} \int \psi \right\} = V^{\frac{n-6}{n}} \left\{ \frac{d^2}{dt^2} \bigg|_{t=0} F(g_t) = 2(n - 6)^2 v^{(6)}(g) V^{-1} \left( \int \phi \right)^2 + (n - 6)(2n - 6) v^{(6)}(g) V^{-1} \left( \int \phi \right)^2 - n(n - 6) v^{(6)}(g) \int \phi^2 \right\} - (n - 6) v^{(6)}(g) \int \psi \right\}\]
\[ V^{-\frac{n-6}{n}} \left\{ \frac{d^2}{dt^2} \bigg|_{t=0} F(g) + 6(n-6)v^{(6)}(g)V^{-1} \left( \int \phi \right)^2 - n(n-6)v^{(6)}(g) \int \phi^2 - (n-6)v^{(6)}(g) \int \psi \right\} \]

\[ = V^{-\frac{n-6}{n}} \left\{ \int \left[ -6v^{(6)}(g) \left( \phi - V^{-1} \int \phi \right)^2 + \frac{\phi \phi_{kl}}{24(n-4)} B_{kl} + \frac{1}{8} \phi \phi_{mk} T_{2mk} - \frac{1}{12} \phi C_{ijk} P_{ij} \phi_k \right] dv \right\}. \]

If we define an operator \( L \) by

\[ L(f) := \frac{B_{ij} f_{ij}}{24(n-4)} + \frac{1}{8} T_{2ij} f_{ij} - \frac{1}{12} P_{ij} C_{ijk} f_k, \]

for \( f \in C^\infty(M) \), it is easy to see that \( L \) is self-adjoint with respect to the \( L^2 \) inner product of the Riemannian manifold. Indeed, for any two smooth functions \( f \) and \( h \), we have

\[ \langle L(f), h \rangle = \int_M L(f) h dv \]

\[ = \int_M \left[ \frac{B_{ij} f_{ij} h}{24(n-4)} + \frac{1}{8} T_{2ij} f_{ij} h \right] dv \]

\[ = \int_M \left[ - \frac{B_{ij} f_{ij} h_j}{24(n-4)} - \frac{1}{8} T_{2ij} f_{ij} h_j - \frac{B_{ij, j} f_{ij}}{24(n-4)} - \frac{1}{8} T_{2ij, j} f_{ij} h - \frac{1}{12} P_{ij} C_{ijk} f_k h \right] dv \]

\[ = \int_M \left[ - \frac{B_{ij} f_{ij} h_j}{24(n-4)} - \frac{T_{2ij, j} f_{ij}}{8} \right] dv \]

where we have used (2) and (3) in Lemma 2.1 and integration by parts. Denote \( \phi - V^{-1} \int \phi \) by \( \bar{\phi} \). From (4.20), we see that

\[ \frac{d^2}{dt^2} \bigg|_{t=0} F_3(g_t) \]

\[ = (n-6) V^{-\frac{n-6}{n}} \left\{ \int \left[ -6v^{(6)}(g) \bar{\phi}^2 + L(\bar{\phi}) \phi \right] dv \right\} \]

\[ = (n-6) V^{-\frac{n-6}{n}} \left\{ \int \left[ -6v^{(6)}(g) \bar{\phi}^2 + L(\bar{\phi}) \phi \right] dv \right\}. \]

Thus we complete the proof of Theorem 1.2. \( \square \)

To prove Theorem 1.3, we need the following famous theorem.

**Theorem 4.1** (Lichnerowicz and Obata; see e.g. [5]). Let \( M \) be an \( n \)-dimensional compact manifold. Suppose the Ricci curvature of \( M \) is bounded from below by

\[ \text{Ric} \geq (n-1)K \]

for some positive constant \( K \). Then the first nonzero eigenvalue of the Laplacian on \( M \) must satisfy

\[ \lambda_1 \geq nK. \]
Moreover, equality holds if and only if $M$ is isometric to a standard sphere of radius $\frac{1}{\sqrt{K}}$.

By the min-max principle, for the first nonzero eigenvalue $\lambda_1$ of the Laplacian, it follows that
\begin{equation}
\lambda_1 \int_M f^2 \, dv \leq \int_M |\nabla f|^2 \, dv,
\end{equation}
for any $f \in C^\infty(M)$ satisfying $\int_M f \, dv = 0$.

**Proof of Theorem 1.3.** Note that an Einstein manifold $(M^n, g)$ is a critical metric in $[g]$; i.e., it satisfies (1.2). Now let $(M^n, g)$ be an Einstein manifold with positive scalar curvature. Then it follows from Theorem 4.1 and (4.22) that
\begin{equation}
R^n - \frac{1}{n-1} \int_M \phi^2 \, dv \leq \int_M |\nabla \phi|^2 \, dv.
\end{equation}

Note that for an Einstein manifold, $v(6)(g) = -\frac{(n-2)R^n}{386n^2(n-1)}$, $\mathcal{L}(\phi) = \frac{(n-2)R^2}{64n^2(n-1)} \Delta \phi$. Hence, we see that
\begin{equation}
\frac{d^2}{dt^2} \bigg|_{t=0} \mathcal{F}_3[g_t] = (n-6)V^{- \frac{n-a}{n-1}} \int_M \left[ \frac{(n-2)R^3}{64n^2(n-1)^2} \phi^2 + \frac{(n-2)R^2}{64n^2(n-1)} \Delta \phi \right] \, dv
\end{equation}
\begin{align*}
&= \frac{(n-2)(n-6)R^2}{64n^2(n-1)} \int_M \left[ R\phi^2 - \frac{1}{n-1} - |\nabla \phi|^2 \right] \, dv \\
&\leq (n-2)(n-6)R^2 \frac{n-a}{n-1} \left( \frac{R}{n-1} - \lambda_1 \right) \int_M \phi^2 \, dv \\
&\leq 0,
\end{align*}
with equality if and only if $\lambda_1 = \frac{R}{n-1}$. Hence, by Theorem 4.1 in this case, $(M, g)$ is isometric to the standard sphere $S^n$.

When $(M^n, g)$ is an $n$-dimensional Einstein manifold with nonpositive scalar curvature, then we have from (4.24) that
\begin{equation}
\frac{d^2}{dt^2} \bigg|_{t=0} \mathcal{F}_3[g_t] < 0.
\end{equation}
Hence we complete the proof of Theorem 1.3. \hfill $\square$

**Remark 4.2.** When $M^n$ is an Einstein manifold with dimension $n = 5$, we see from (4.24) that
\begin{equation}
\frac{d^2}{dt^2} \bigg|_{t=0} \mathcal{F}_3[g_t] \geq 0,
\end{equation}
with equality if and only if $\lambda_1 = \frac{R}{4}$. Theorem 4.1 shows that in this case $(M^n, g)$ is isometric to the sphere $S^5$ with the standard metric up to a multiple of a constant. We see that this Einstein metric is a strict local minimum of the functional $\mathcal{F}_3$ within its conformal class if the equality does not hold in (4.25).

**Remark 4.3.** Letting $T_{ij}(g) = T_{2ij}(g) + \frac{1}{n-4}(B_g)_{ij}$, we have
\begin{equation}
\sum_j \nabla^j T_{ij} = 0;
\end{equation}
that is, $T_{ij}$ is a divergence-free tensor. We note that $v(6)(g) = -\frac{1}{27} \sum_{ij} T_{ij}(g)(P_g)_{ij}$. 

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THE SECOND VARIATIONAL FORMULA FOR $\int v^{(6)}(g) dV_g$

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