# GLOBAL-IN-TIME SMOOTHING EFFECTS FOR SCHRÖDINGER EQUATIONS WITH INVERSE-SQUARE POTENTIALS 

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#### Abstract

The purpose of this note is to prove global-in-time smoothing effects for the Schrödinger equation with potentials exhibiting critical singularity. A typical example of admissible potentials is the inverse-square potential $a|x|^{-2}$ with $a>-(n-2)^{2} / 4$. This particularly gives an affirmative answer to a question raised by T. A. Bui et al. (J. Differential Equations 262 (2017), 2771-2807). The proof employs a uniform resolvent estimate proved by Barceló, Vega, and Zubeldia (Adv. Math. 240 (2013), 636-671) an abstract perturbation method by Bouclet and Mizutani (preprint).


## 1. Introduction

This note is concerned with smoothing properties of the time-dependent Schrödinger equation

$$
\begin{equation*}
i \partial_{t} u(t, x)=H u(t, x)+F(t, x) ; \quad u(0, x)=\psi(x), \tag{1.1}
\end{equation*}
$$

with given data $\psi \in L^{2}\left(\mathbb{R}^{n}\right)$ and $F \in L_{\text {loc }}^{1}\left(\mathbb{R} ; L^{2}\left(\mathbb{R}^{n}\right)\right.$ ), where $H=-\Delta+V(x)$ is a Schrödinger operator on $\mathbb{R}^{n}, n \geq 3$, with a real-valued function $V$ which decays at spatial infinity in a suitable sense and has a critical singularity at the origin. A typical example of potentials we have in mind is the inverse-square potential $V(x)=a|x|^{-2}$ satisfying $a>-(n-2)^{2} / 4$.

Let us first recall several known results for the free case, describing the motivation of this paper. It is well-known that the solution $u=e^{i t \Delta} \psi$ to the free Schrödinger equation

$$
i \partial_{t} u(t, x)=-\Delta u(t, x) ;\left.\quad u\right|_{t=0}=\psi \in L^{2}\left(\mathbb{R}^{n}\right),
$$

satisfies the global-in-time smoothing effect

$$
\begin{equation*}
\left\|\langle x\rangle^{-\rho}|D|^{1 / 2} e^{i t \Delta} \psi\right\|_{L^{2}\left(\mathbb{R}^{1+n}\right)} \leq C\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)} \tag{1.2}
\end{equation*}
$$

where $\langle x\rangle=\left(1+|x|^{2}\right)^{1 / 2}, \rho>1 / 2$ and $|D|=(-\Delta)^{1 / 2}$ (see Ben-Artzi and Klainerman [2] for $n \geq 3$ and Chihara [8] for $n=2$ ). When $n \geq 3$, the estimate of the form

$$
\begin{equation*}
\left\|w(x) e^{i t \Delta} \psi\right\|_{L^{2}\left(\mathbb{R}^{1+n}\right)} \leq C\|w\|_{L^{n}\left(\mathbb{R}^{n}\right)}\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)} \tag{1.3}
\end{equation*}
$$

[^0]was proved by Kato and Yajima [18]. The estimate (1.3) also follows from Hölder's inequality and the endpoint Strichartz estimate proved by Keel and Tao [19]:
\[

$$
\begin{equation*}
\left\|e^{i t \Delta} \psi\right\|_{L^{2}\left(\mathbb{R} ; L^{\frac{2 n}{n-2}}\left(\mathbb{R}^{n}\right)\right)} \leq C\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)} \tag{1.4}
\end{equation*}
$$

\]

which can also be regarded as a smoothing property in $L^{p}$-spaces. All of these three estimates are fundamental tools in the study of the Cauchy problem and scattering theory for both linear and nonlinear Schrödinger equations (see [7, 17, 18, 27, 28] and references therein). It is also worth noting that (1.2) and (1.3) are closely connected with uniform estimates for the resolvent $(-\Delta-z)^{-1}$ with respect to $z \in \mathbb{C} \backslash[0, \infty)$ (see the next section for more details).

There is a vast literature on extending estimates (1.2)-(1.4) to the Schrödinger operator $H=-\Delta+V$ with potential $V(x)$. For the case when $V$ has enough regularity and decays sufficiently fast at spatial infinity, we refer to [12, 13, 22, 24, 26] and references therein. There are also several results in the case when $V$ has critical singularity. In particular, the Schrödinger operator with the inverse-square potential of the form

$$
H_{a}=-\Delta+a|x|^{-2}, \quad a>-\frac{(n-2)^{2}}{4}
$$

has attracted increasing attention since it represents a borderline case for the validity of (1.2)-(1.4) (11,14), where we note that $(n-2)^{2} / 4$ is the best constant in Hardy's inequality

$$
\begin{equation*}
\frac{(n-2)^{2}}{4} \int|x|^{-2}|u|^{2} d x \leq \int|\nabla u|^{2} d x, \quad u \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right) \tag{1.5}
\end{equation*}
$$

We refer to [1, 3, 5, 6, for Kato-Yajima type estimates (1.3) and to [3, 5, 6, 23] for Strichartz estimates (1.4). Concerning the estimate (1.2), in a recent paper [4], the authors showed, among others, the following theorem.

Theorem 1.1 ([4, Theorem 1.2]). Let $n \geq 3, a \geq-(n-2)^{2} / 4+1 / 4$ and $\varepsilon>0$. Then there exists $C_{\varepsilon}>0$ such that for all $\psi \in L^{2}\left(\mathbb{R}^{n}\right)$, $e^{-i t H_{a}} \psi$ satisfies

$$
\begin{equation*}
\left\|w(|x|)|D|^{1 / 2} e^{-i t H_{a}} \psi\right\|_{L^{2}\left(\mathbb{R}^{1+n}\right)} \leq C_{\varepsilon}\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)} \tag{1.6}
\end{equation*}
$$

where $w(r)=r^{(\varepsilon-1) / 2}\left(1+r^{\varepsilon}\right)^{-1}$.
The condition $a \geq-(n-2)^{2} / 4+1 / 4$ was used to ensure that $w$ satisfies some conditions for two-sided weighted norm estimates

$$
C_{1}\left\|w(x)(-\Delta)^{s / 2} f\right\|_{L^{2}} \leq\left\|w(x) H_{a}^{s / 2} f\right\|_{L^{2}} \leq C_{2}\left\|w(x)(-\Delta)^{s / 2} f\right\|_{L^{2}}
$$

established by the same paper (see [4, Theorem 1.1]). Then the authors raised a question whether (1.6) holds under the condition $a>-(n-2)^{2} / 4$. The main purpose of the present note is to give an affirmative answer to this question. More precisely, we prove global-in-time smoothing effects of the form (1.6) for Schrödinger operators $H=-\Delta+V(x)$ with a large class of real-valued potentials which particularly includes the inverse-square potential with $a>-(n-2)^{2} / 4$. Furthermore, global-in-time smoothing effects for the solution to (1.1) with the inhomogeneous term $F$ are also studied for the same class of potentials. The proofs are based on an abstract perturbation method from our previous work [3] and a uniform estimate proved by [1] for the weighted resolvent $|x|^{-1}(H-z)^{-1}|x|^{-1}$ with respect to $z \in \mathbb{C} \backslash[0, \infty)$.

In order to state the main results, we introduce some notation. From now on we let $n \geq 3$ and impose the following condition:

Assumption A. Let $V(x)$ be a real-valued function on $\mathbb{R}^{n}$ such that $|x| V \in$ $L^{n, \infty}\left(\mathbb{R}^{n}\right)$ and $x \cdot \nabla V \in L^{n / 2, \infty}\left(\mathbb{R}^{n}\right)$. Moreover, there exists $\delta>0$ such that $-\Delta+V \geq-\delta \Delta$ and $-\Delta-V-x \cdot \nabla V \geq-\delta \Delta$ in the sense of forms; that is, for all $u \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\langle(-\Delta+V) u, u\rangle \geq \delta\|\nabla u\|_{L^{2}}, \quad\langle(-\Delta-V-x \cdot(\nabla V)) u, u\rangle \geq \delta\|\nabla u\|_{L^{2}} \tag{1.7}
\end{equation*}
$$

Here $\langle f, g\rangle=\int f \bar{g} d x$ is the inner product in $L^{2}\left(\mathbb{R}^{n}\right)$ and $L^{p, q}$ is the Lorentz space (see the end of this section). A typical example satisfying Assumption A is the inverse-square potential $V(x)=a|x|^{-2}$ with $a>-(n-2)^{2} / 4$. In this case, it follows from Hardy's inequality (1.5) that (1.7) is satisfied with $\delta=1-4|a| /(n-2)^{2}>0$ if $a<0$ or $\delta=1$ if $a \geq 0$. Moreover, Assumption A is general enough to include some potentials such that $|x|^{2} V \notin L^{\infty}$. For instance, we let $c_{1}, c_{2}>0, \alpha \in \mathbb{R}^{n}$ and $\chi \in C^{1}(\mathbb{R})$ such that $0 \leq \chi \leq 1$ and $\left|\chi^{(k)}(t)\right| \leq|t|^{-k-1}$ for $|t| \geq 1$. Then

$$
V(x)=\frac{-(n-2)^{2} / 4+c_{1}}{|x|^{2}}-\frac{c_{2} \chi(|x-\alpha|)}{|x-\alpha|}
$$

satisfies Assumption A with $\delta=c_{1}-c_{2}\left(2+\sup \left|\chi^{\prime}\right|\right)(|\alpha|+1)$ if

$$
0<c_{2}<\frac{c_{1}}{\left(2+\sup \left|\chi^{\prime}\right|\right)(|\alpha|+1)}
$$

Under Assumption A. Hardy's inequality (1.5) implies that the sesquilinear form

$$
Q_{H}(u, v)=\langle(-\Delta+V) u, v\rangle, \quad u, v \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)
$$

is symmetric, non-negative and closable such that the domain of its closure $\bar{Q}_{H}$ satisfies $D\left(\bar{Q}_{H}\right)=\mathcal{H}^{1}\left(\mathbb{R}^{n}\right)$. Let $H$ be the Friedrichs extension of $\bar{Q}_{H}$, $e^{-i t H}$ the unitary group on $L^{2}\left(\mathbb{R}^{n}\right)$ generated by $H$ via Stone's theorem and $\Gamma_{H}$ the inhomogeneous propagator defined by

$$
\begin{equation*}
\Gamma_{H} F(t)=\int_{0}^{t} e^{-i(t-s) H} F(s) d s, \quad F \in L_{\mathrm{loc}}^{1}\left(\mathbb{R} ; L^{2}\left(\mathbb{R}^{n}\right)\right) \tag{1.8}
\end{equation*}
$$

Then a unique (mild) solution to the Schrödinger equation (1.1) is given by

$$
\begin{equation*}
u(t)=e^{-i t H} \psi-i \Gamma_{H} F(t) \tag{1.9}
\end{equation*}
$$

We say that $v(x)$ belongs to the Muckenhoupt $A_{2}$ class if $v, 1 / v \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{n}\right), v \geq 0$ and

$$
\left(\frac{1}{|B|} \int_{B} v(x) d x\right)\left(\frac{1}{|B|} \int_{B} \frac{1}{v(x)} d x\right) \leq C
$$

for all balls $B \subset \mathbb{R}^{n}$ with some constant $C>0$ independent of $B$.
The main result in this paper then is as follows.
Theorem 1.2. Let $n \geq 3, V$ satisfy Assumption A and $w \in L^{2}(\mathbb{R})$. Suppose $w(|x|)^{2} \in A_{2}$ and, for any $j=1,2, \ldots, n$, there exists $C_{j}>0$ such that

$$
w(|x|) \leq C_{j} w\left(x_{j}\right), \quad x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}
$$

Then there exists $C>0$, independent of $w$, such that $e^{-i t H}$ satisfies

$$
\left\|w(|x|)|D|^{1 / 2} e^{-i t H} \psi\right\|_{L^{2}\left(\mathbb{R}^{1+n}\right)} \leq C\|w\|_{L^{2}(\mathbb{R})}\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)}, \quad \psi \in L^{2}\left(\mathbb{R}^{n}\right)
$$

Assuming $0<\varepsilon<1$ without loss of generality, it is easy to see that $w(r)=$ $r^{(\varepsilon-1) / 2}\left(1+r^{\varepsilon}\right)^{-1}$ fulfills the above conditions. Another typical example of $w$ is $\langle x\rangle^{-\rho}$, in which case we have
Theorem 1.3. Let $n \geq 3, \rho>1 / 2$ and $\mathcal{A}, \mathcal{B} \in\left\{\dot{\mathcal{H}}^{-1 / 2, \rho}\left(\mathbb{R}^{n}\right)\right.$, $\left.L^{\frac{2 n}{n+2}, 2}\left(\mathbb{R}^{n}\right)\right\}$. Suppose $V$ satisfies Assumption A. Then the solution $u$ to (1.1) given by (1.9) satisfies

$$
\begin{equation*}
\|u\|_{L^{2}\left(\mathbb{R} ; \mathcal{B}^{*}\right)} \leq C\|\psi\|_{L^{2}\left(\mathbb{R}^{n}\right)}+C\|F\|_{L^{2}(\mathbb{R} ; \mathcal{A})} \tag{1.10}
\end{equation*}
$$

for all $\psi \in L^{2}\left(\mathbb{R}^{n}\right)$ and $F \in L_{\mathrm{loc}}^{1}\left(\mathbb{R} ; L^{2}\left(\mathbb{R}^{n}\right)\right) \cap L^{2}(\mathbb{R} ; \mathcal{A})$.
Here $\dot{\mathcal{H}}^{s, \mu}\left(\mathbb{R}^{n}\right)$ denotes the weighted homogeneous Sobolev space equipped with the norm $\|f\|_{\dot{\mathcal{H}}^{s}, \mu}=\left\|\langle x\rangle^{\mu}|D|^{s} f\right\|_{L^{2}}$. Note that $\left(L^{\frac{2 n}{n+2}, 2}\left(\mathbb{R}^{n}\right)\right)^{*}=L^{\frac{2 n}{n-2}, 2}\left(\mathbb{R}^{n}\right)$ and $\left(\dot{\mathcal{G}}^{-1 / 2, \rho}\left(\mathbb{R}^{n}\right)\right)^{*}=\dot{\mathcal{G}}^{1 / 2,-\rho}\left(\mathbb{R}^{n}\right)$. If $\mathcal{A}=\mathcal{B}=L^{\frac{2 n}{n+2}, 2}\left(\mathbb{R}^{n}\right)$, (1.10) becomes the endpoint Strichartz estimate and was proved by our previous work [3]. If $\mathcal{A}=$ $\mathcal{B}=\dot{\mathscr{H}}^{-1 / 2, \rho}\left(\mathbb{R}^{n}\right)$, (1.10) is a generalization of (1.2) and seems to be new under Assumption ( Here we stress that $\mathcal{A}$ and $\mathcal{B}$ do not have to coincide.

Notation. Throughout the paper we use the following notation. For $T>0$ and a Banach space $X$, we denote $\|F\|_{L_{T}^{p} x}=\|F\|_{L^{p}([-T, T] ; x)} . L^{p, q}\left(\mathbb{R}^{n}\right)$ denotes the Lorentz space equipped with the norm $\|\cdot\|_{L^{p, q}\left(\mathbb{R}^{n}\right)}$ satisfying

$$
\|f\|_{L^{p, q}\left(\mathbb{R}^{n}\right)} \sim\left\|t d_{f}(t)^{1 / p}\right\|_{L^{q}\left(\mathbb{R}_{+}, t^{-1} d t\right)},
$$

where $d_{f}(t):=\mu\left(\left\{x \in \mathbb{R}^{n}| | f(x) \mid>t\right\}\right)$ is the distribution function of $f$. We use the convention $L^{\infty, \infty}=L^{\infty}$. For $1 \leq p, p_{1}, p_{2}<\infty$ and $1 \leq q, q_{1}, q_{2} \leq \infty$ satisfying $1 / p=1 / p_{1}+1 / p_{2}, 1 / q=1 / q_{1}+1 / q_{2}$, we have Hölder's inequality for Lorentz spaces:

$$
\begin{equation*}
\|f g\|_{L^{p, q}} \leq C\|f\|_{L^{p_{1}, q_{1}}}\|g\|_{L^{p_{2}, q_{2}}}, \quad\|f g\|_{L^{p, q}} \leq C\|f\|_{L^{\infty}}\|g\|_{L^{p, q}} . \tag{1.11}
\end{equation*}
$$

When $n \geq 3$, we also have Sobolev's inequality in Lorentz spaces:

$$
\begin{equation*}
\|f\|_{L^{\frac{2 n}{n-2}, 2}} \leq C\|\nabla f\|_{L^{2}}, \quad f \in \mathcal{H}^{1} \tag{1.12}
\end{equation*}
$$

We refer to [15] for more details on Lorentz spaces. In what follows we often omit $\mathbb{R}^{n}$ from $L^{p}\left(\mathbb{R}^{n}\right)$ and so on, if there is no confusion.

The rest of the paper is organized as follows. We first recall in the next section the abstract perturbation method developed in [3], which plays an important role in the proof of the main theorems. The proof of Theorems 1.2 and 1.3 is given in section 3.

## 2. An abstract perturbation method

Here we recall the abstract method developed in [3. We begin by recalling the notion of the (super)smoothness in the sense of Kato [17] and Kato-Yajima [18]. Let $\mathcal{H}$ be a Hilbert space with inner product $\langle\cdot, \cdot\rangle$ and norm $\|\cdot\|, H$ a self-adjoint operator on $\mathcal{H}$ and $A$ a densely defined closed operator on $\mathcal{H}$. Note that $A^{*}$ is also a densely defined closed operator (see [25, Theorem VIII.1]). Let $R_{H}(z):=(H-z)^{-1}$, $z \notin \sigma(H)$. Then we say that $A$ is $H$-smooth with bound $a$ if

$$
\sup _{z \in \mathbb{C} \backslash \mathbb{R}}\left|\left\langle\left(R_{H}(z)-R_{H}(\bar{z})\right) A^{*} \psi, A^{*} \psi\right\rangle\right| \leq \frac{a^{2}}{2}\|\psi\|^{2}, \quad \psi \in D\left(A^{*}\right) .
$$

We say that $A$ is $H$-supersmooth with bound $a$ if

$$
\sup _{z \in \mathbb{C} \backslash \mathbb{R}}\left|\left\langle R_{H}(z) A^{*} \psi, A^{*} \psi\right\rangle\right| \leq \frac{a}{2}\|\psi\|^{2}, \quad \psi \in D\left(A^{*}\right) .
$$

Note that if $A$ is $H$-supersmooth with bound $a$, then $A$ is $H$-smooth with bound $(2 a)^{1 / 2}$. The $H$-(super)smoothness is closely connected with smoothing effects.

Proposition 2.1. (1) $A$ is $H$-smooth with bound a if and only if, for any $\psi \in \mathcal{H}$, $e^{-i t H} \psi$ belongs to $D(A)$ for a.e. $t \in \mathbb{R}$ and

$$
\begin{equation*}
\left\|A e^{-i t H} \psi\right\|_{L^{2}(\mathbb{R} ; \mathcal{H})} \leq a\|\psi\| \tag{2.1}
\end{equation*}
$$

(2) Suppose $A$ is $H$-supersmooth with bound $a$. Then, for any simple function $F: \mathbb{R} \rightarrow D\left(A^{*}\right)$ and $t \in \mathbb{R}, A e^{-i(t-s) H} A^{*} F(s)$ is Bochner integrable in $s$ over $[0, t]$ (or $[t, 0]$ ) and satisfies

$$
\begin{equation*}
\left\|e^{-|\varepsilon t|} \int_{0}^{t} A e^{-i(t-s) H} A^{*} F(s) d s\right\|_{L^{2}(\mathbb{R} ; \mathcal{H})} \leq a\left\|e^{-|\varepsilon t|} F\right\|_{L^{2}(\mathbb{R} ; \mathcal{H})} \tag{2.2}
\end{equation*}
$$

for all $\varepsilon \in \mathbb{R}$. Conversely, if the estimate (2.2) holds for all simple functions $F: \mathbb{R} \rightarrow D\left(A^{*}\right)$ and $|\varepsilon|<\varepsilon_{0}$ with some $\varepsilon_{0}>0$, then $A$ is $H$-supersmooth with bound $a$.

Proof. The first statement is due to [17, Lemma 3.6 and Theorem 5.1] (see also [25, Theorem XIII.25]). The second assertion was proved by [10, Theorem 2.3].

Note that if $A$ is $H$-smooth, then $A$ is infinitesimally $H$-bounded ([25), Theorem XIII.22]). Also note that the estimate (2.2) can be replaced by

$$
\begin{equation*}
\left\|e^{-|\varepsilon t|} \int_{0}^{t} A e^{-i(t-s) H} A^{*} F(s) d s\right\|_{L_{T}^{2} \mathcal{H}} \leq a\left\|e^{-|\varepsilon t|} F\right\|_{L_{T}^{2} \mathcal{H}} . \tag{2.3}
\end{equation*}
$$

Indeed, (2.2) implies (2.3) since $s \in[-T, T]$ if $t \in[-T, T]$ and $s \in[0, t]$ (or $s \in$ $[-t, 0]$ ). Conversely, since (2.3) implies that

$$
\left\|e^{-|\varepsilon t|} \int_{0}^{t} A e^{-i(t-s) H} A^{*} F(s) d s\right\|_{L_{T}^{2} \mathcal{H}} \leq a\left\|e^{-|\varepsilon t|} F\right\|_{L^{2}(\mathbb{R} ; \mathcal{H})}
$$

and $a$ is independent of $T$, one has (2.2) by letting $T \rightarrow \infty$. Let $\Gamma_{H}$ be the inhomogeneous propagator defined by the formula (1.8) and set

$$
\begin{equation*}
\Gamma_{H}^{*} F(t)=\mathbb{1}_{[0, \infty)}(t) \int_{t}^{T} e^{-i(t-s) H} F(s) d s-\mathbb{1}_{(-\infty, 0]}(t) \int_{-T}^{t} e^{-i(t-s) H} F(s) d s \tag{2.4}
\end{equation*}
$$

By a direct calculation, $\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}=\left\langle\left\langle F, \Gamma_{H}^{*} G\right\rangle\right\rangle_{T}$ for $F, G \in L_{\text {loc }}^{1} \mathcal{H}$, where

$$
\langle\langle F, G\rangle\rangle_{T}:=\int_{-T}^{T}\langle F(t), G(t)\rangle d t
$$

Hence $\Gamma_{H}^{*}$ is the adjoint of $\Gamma_{H}$ in $L_{T}^{2} \mathcal{H}$. In the abstract theorem below, the operators $A \Gamma_{H}$ and $A \Gamma_{H}^{*}$ for some $H$-smooth operator $A$ play important roles. These operators are a priori well-defined on $L_{\mathrm{loc}}^{1}(\mathbb{R} ; D(H))$ since, for some $z \notin \sigma(H)$ and each $T>0$,

$$
\left\|A \Gamma_{H} F\right\|_{L_{T}^{2} \mathcal{H}}+\left\|A \Gamma_{H}^{*} F\right\|_{L_{T}^{2} \mathcal{H}} \leq C_{T}\left\|A R_{0}(z)\right\|_{\mathbb{B}(\mathcal{H})}\left\|R_{0}(z) F\right\|_{L_{T}^{1} \mathcal{H}}<\infty
$$

The next lemma provides a rigorous definition of $A \Gamma_{H} F(t)$ for $F \in L_{\mathrm{loc}}^{1}(\mathbb{R} ; \mathcal{H})$.

Lemma 2.2 (3, Lemma 4.3 and Lemma 4.5]). Suppose that $A$ is $H$-smooth with bound a. Then $A \Gamma_{H}$ and $A \Gamma_{H}^{*}$ extend to bounded operators from $L_{T}^{1} \mathcal{H}$ to $L_{T}^{2} \mathcal{H}$ such that

$$
\begin{equation*}
\left\|A \Gamma_{H} F\right\|_{L_{T}^{2} \mathcal{H}} \leq C a\|F\|_{L_{T}^{1} \mathcal{H}}, \quad\left\|A \Gamma_{H}^{*} F\right\|_{L_{T}^{2} \mathcal{H}} \leq C a\|F\|_{L_{T}^{1} \mathcal{H}} \tag{2.5}
\end{equation*}
$$

for any $T>0$ and $F \in L_{T}^{1} \mathcal{H}$ with some $C>0$ independent of $a, T$ and $F$. Moreover, we have

$$
\begin{equation*}
A \Gamma_{H} F(t)=\int_{0}^{t} A e^{-i(t-s) H} F(s) d s \tag{2.6}
\end{equation*}
$$

for all simple functions $F:[-T, T] \rightarrow \mathcal{H}$ and a.e. $t \in[-T, T]$. In particular $A \Gamma_{H} F(t)$ and $\int_{0}^{t} A e^{-i(t-s) H} F(s) d s$ coincide in $L_{T}^{2} \mathcal{H}$.
Proof. For the sake of self-containedness, we give the proof for $\Gamma_{H}$ in detail. The proof for $\Gamma_{H}^{*}$ is analogous. Let $\chi \in C_{0}^{\infty}(\mathbb{R})$ be such that $\chi \equiv 1$ near 0 and $0 \leq \chi \leq 1$ and set $\chi_{\varepsilon}(t)=\chi(\varepsilon t)$. Note that $A \Gamma_{H} \chi_{\varepsilon}(H)=A \chi_{\varepsilon}(H) \Gamma_{H}$ is well-defined on $L_{T}^{1} \mathcal{H}$ since $A$ is $H$-bounded. Moreover, the $H$-smoothness of $A$ shows that

$$
\left\|A e^{-i t H} \int_{[0, T]} e^{i s H} \chi_{\varepsilon}(H) F(s) d s\right\|_{L_{T}^{2} \mathcal{H}} \leq a| | F \|_{L_{T}^{1} \mathcal{H}}
$$

Using the Christ-Kiselev lemma [9, one can replace $[0, T]$ by $[0, t]$ in the left hand side to obtain

$$
\left\|A \chi_{\varepsilon}(H) \Gamma_{H} F\right\|_{L_{T}^{2} \mathcal{H}} \leq C a\|F\|_{L_{T}^{1} \mathcal{H}}
$$

with some universal constant $C>0$ independent of $F$. This estimate implies that, for all $F \in L_{T}^{1} \mathcal{H}, A \chi_{\varepsilon}(H) \Gamma_{H} F$ converges in $L_{T}^{2} \mathcal{H}$ as $\varepsilon \rightarrow 0$ and the limit denoted by the same symbol $A \Gamma_{H} F$ satisfies the first estimate in (2.5). This proves the first half of this lemma.

Next, let $F:[-T, T] \rightarrow \mathcal{H}$ be a simple function. As we have shown, $A \chi_{\varepsilon_{n}}(H) \Gamma_{H}$ converges to $A \Gamma_{H} F$ in $L_{T}^{2} \mathcal{H}$ as $n \rightarrow \infty$ for any sequence $\varepsilon_{n}>0$ with $\varepsilon_{n} \rightarrow$ 0 . Then one can find a subsequence $\varepsilon_{n_{k}}$ and a null set $\mathcal{N} \subset[-T, T]$ such that $A \chi_{\varepsilon_{n_{k}}}(H) \Gamma_{H} F \rightarrow A \Gamma_{H} F(t)$ in $\mathcal{H}$ as $k \rightarrow \infty$ for all $t \in[-T, T] \backslash \mathcal{N}$. Hence, in order to show (2.6), it suffices to check that

$$
\left\|\int_{0}^{t} A e^{-i(t-s) H} F(s) d s-A \chi_{\varepsilon_{n_{k}}}(H) \Gamma_{H} F(t)\right\|_{\mathcal{H}} \rightarrow 0
$$

as $k \rightarrow \infty$ for all $t \in[-T, T] \backslash \mathcal{N}$. Let us write $F=\sum_{j=1}^{N} \mathbb{1}_{E_{j}}(t) f_{j}$ with $f_{j} \in \mathcal{H}$ and measurable sets $E_{j} \subset[-T, T]$. Then we have for all $t \in[-T, T] \backslash \mathcal{N}$,

$$
\begin{aligned}
& \left\|\int_{0}^{t} A e^{-i(t-s) H} F(s) d s-A \chi_{\varepsilon_{n_{k}}}(H) \Gamma_{H} F(t)\right\|_{\mathcal{H}} \\
& \leq \sum_{j=1}^{N}\left\|\int_{[0, t] \cap M_{j}} A e^{i s H}\left(1-\chi_{\varepsilon_{n_{k}}}(H)\right) e^{-i t H} f_{j} d s\right\|_{\mathcal{H}} \\
& \leq \sum_{j=1}^{N} T^{1 / 2}\left(\int_{\mathbb{R}}\left\|A e^{i s H} e^{-i t H}\left(1-\chi_{\varepsilon_{n_{k}}}(H)\right) f_{j}\right\|_{\mathcal{H}}^{2} d s\right)^{1 / 2} \\
& \leq C T^{1 / 2} \sum_{j=1}^{N}\left\|\left(1-\chi_{\varepsilon_{n_{k}}}(H)\right) f_{j}\right\|_{\mathcal{H}} \rightarrow 0
\end{aligned}
$$

as $k \rightarrow \infty$, where we have used Hölder's inequality in the third line and the $H$ smoothness of $A$ and the unitarity of $e^{-i t H}$ in the last line. This completes the proof.

In what follows, $A \Gamma_{H}$ and $A \Gamma_{H}^{*}$ denote such extensions. By Proposition [2.1(1), (2.3) and this lemma, if $A$ is $H$-supersmooth with bound $a$, then $A \Gamma_{H} A^{*}$ is welldefined on $L_{\text {loc }}^{1}\left(\mathbb{R} ;\left(D\left(A^{*}\right)\right)\right)$ and satisfies $\left\|A \Gamma_{H} A^{*} F\right\|_{L_{T}^{2} \mathcal{H}} \leq a\|F\|_{L_{T}^{2} \mathcal{H}}$ for all simple functions $F:[-T, T] \rightarrow D\left(A^{*}\right)$.

Now we are ready to recall an abstract theorem from 3]. Let $\left(H_{0}, H\right)$ be a pair of self-adjoint operators on $\mathcal{H}$ such that $H=H_{0}+V_{1}^{*} V_{2}$ in the following sense:

- $V_{1}, V_{2}$ are densely defined closed operators on $\mathcal{H}$ such that we have continuous embeddings $D\left(H_{0}\right) \subset D\left(V_{1}\right), D(H) \subset D\left(V_{1}\right)$ and $D(H) \subset D\left(V_{2}\right)$.
- $\langle H f, g\rangle=\left\langle f, H_{0} g\right\rangle+\left\langle V_{2} f, V_{1} g\right\rangle$ for $f \in D(H)$ and $g \in D\left(H_{0}\right)$.

Note that, under the above conditions, $V_{1}^{*}, V_{2}^{*}$ are also densely defined. Recall that two Banach spaces $(\mathcal{A}, \mathcal{B})$ are said to be a Banach couple if both $\mathcal{A}$ and $\mathcal{B}$ are algebraically and topologically embedded in a Hausdorff topological vector space e.

Proposition 2.3 ( 3 , Theorem 4.7]). Let $\mathcal{A}, \mathcal{B}$ be two Banach spaces such that $(\mathcal{A}, \mathcal{H})$ and $(\mathcal{B}, \mathcal{H})$ are Banach couples. Suppose that $V_{1}$ is $H_{0}$-smooth and $V_{2}$ is both $H_{0}$-smooth and $H$-smooth. Consider the following series of estimates:

$$
\begin{align*}
\left|\left\langle\left\langle e^{-i t H_{0}} \psi, G\right\rangle\right\rangle_{T}\right| & \leq s_{1}\|\psi\|_{\mathcal{H}}\|G\|_{L_{T}^{2} \mathcal{B}},  \tag{2.7}\\
\left|\left\langle\left\langle\Gamma_{H_{0}} F, G\right\rangle\right\rangle_{T}\right| & \leq s_{2}\|F\|_{L_{T}^{2} \mathcal{A}}\|G\|_{L_{T}^{2} \mathcal{B}},  \tag{2.8}\\
\left\|V_{1} \Gamma_{H_{0}}^{*} G\right\|_{L_{T}^{2} \mathcal{H}} & \leq s_{3}\|G\|_{L_{T}^{2} \mathcal{B}},  \tag{2.9}\\
\left\|V_{1} \Gamma_{H_{0}} F\right\|_{L_{T}^{2} \mathcal{H}} & \leq s_{4}\|F\|_{L_{T}^{2} \mathcal{A}}  \tag{2.10}\\
\left\|V_{2} \Gamma_{H_{0}} F\right\|_{L_{T}^{2} \mathcal{H}} & \leq s_{5}\|F\|_{L_{T}^{2} \cdot \mathcal{A}}  \tag{2.11}\\
\left\|V_{2} e^{-i t H} \psi\right\|_{L_{T}^{2} \mathcal{H}} & \leq s_{6}\|\psi\|_{\mathcal{H}},  \tag{2.12}\\
\left\|V_{2} \Gamma_{H} V_{2}^{*} \widetilde{G}\right\|_{L_{T}^{2} \mathcal{H}} & \leq s_{7}\|\widetilde{G}\|_{L_{T}^{2} \mathcal{H}} . \tag{2.13}
\end{align*}
$$

(1) Suppose there exist constants $s_{1}, s_{3}, s_{6}>0$ such that (2.7), (2.9) and (2.12) are satisfied for all $\psi \in \mathscr{H}$ and simple function $G:[-T, T] \rightarrow \mathcal{H} \cap \mathcal{B}$. Then one has

$$
\left|\left\langle\left\langle e^{-i t H} \psi, G\right\rangle\right\rangle_{T}\right| \leq\left(s_{1}+s_{3} s_{6}\right)\|\psi\|_{\mathcal{H}}\|G\|_{L_{T}^{2} \mathcal{B}}
$$

for all $\psi \in \mathcal{H}$ and simple function $G:[-T, T] \rightarrow \mathcal{H} \cap \mathcal{B}$.
(2) Suppose that there exist constants $s_{2}, s_{3}, s_{4}, s_{5}, s_{7}>0$ such that (2.8), (2.9), (2.10), (2.11) and (2.13) hold for all simple functions $F:[-T, T] \rightarrow \mathcal{H} \cap \mathcal{A}, G:$ $[-T, T] \rightarrow \mathcal{H} \cap \mathcal{B}$ and $\widetilde{G}:[-T, T] \rightarrow D(H)$. Then one has

$$
\left|\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}\right| \leq\left(s_{2}+s_{3} s_{5}+s_{3} s_{4} s_{7}\right)\|F\|_{L_{T}^{2} \mathcal{A}}\|G\|_{L_{T}^{2} \mathcal{B}}
$$

for all simple functions $F:[-T, T] \rightarrow \mathcal{H} \cap \mathcal{A}$ and $G:[-T, T] \rightarrow \mathcal{H} \cap \mathcal{B}$.
Proof. We give the proof in detail for the sake of self-containedness. Let us first show the first statement. The Duhamel formula implies that

$$
\begin{equation*}
\left\langle e^{-i t H} \psi, \varphi\right\rangle=\left\langle e^{-i t H_{0}} \psi, \varphi\right\rangle-i \int_{0}^{t}\left\langle V_{2} e^{-i r H} \psi, V_{1} e^{-i(r-t) H_{0}} \varphi\right\rangle d r, \quad \psi, \varphi \in L^{2} \tag{2.14}
\end{equation*}
$$

Plugging in $\varphi=G(t)$, integrating over $t \in[-T, T]$ and using Fubini's theorem, we learn from the formula (2.4) of the adjoint $\Gamma_{H_{0}}^{*}$ that

$$
\left\langle\left\langle e^{-i t H} \psi, G\right\rangle\right\rangle_{T}=\left\langle\left\langle e^{-i t H_{0}} \psi, G\right\rangle\right\rangle_{T}-i\left\langle\left\langle V_{2} e^{-i t H} \psi, V_{1} \Gamma_{H_{0}}^{*} G\right\rangle\right\rangle_{T}
$$

Applying (2.7), (2.9) and (2.12), we then obtain the first assertion (1).
In order to prove the second assertion (2), we replace $t$ by $t-s$ and plug in $\psi=F(s), \varphi=G(t)$ and integrate over $s \in[0, t]$ in (2.14) to obtain

$$
\begin{aligned}
&\left\langle\Gamma_{H}\right.F(t), G(t)\rangle \\
& \quad=\left\langle\Gamma_{H_{0}} F(t), G(t)\right\rangle-i \int_{0}^{t} \int_{s}^{t}\left\langle V_{2} e^{-i(\tau-s) H} F(s), V_{1} e^{-i(\tau-t) H_{0}} G(t)\right\rangle d \tau d s \\
& \quad=\left\langle\Gamma_{H_{0}} F(t), G(t)\right\rangle-i \int_{0}^{t}\left\langle V_{2} \Gamma_{H} F(\tau), V_{1} e^{-i(\tau-t) H_{0}} G(t)\right\rangle d \tau
\end{aligned}
$$

As above, integrating in $t \in[-T, T]$ and using (2.4) implies that

$$
\begin{equation*}
\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}=\left\langle\left\langle\Gamma_{H_{0}} F, G\right\rangle\right\rangle_{T}-i\left\langle\left\langle V_{2} \Gamma_{H} F, V_{1} \Gamma_{H_{0}}^{*} G\right\rangle\right\rangle_{T} \tag{2.15}
\end{equation*}
$$

Exchanging the roles of $H$ and $H_{0}$, we also obtain

$$
\begin{equation*}
\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}=\left\langle\left\langle\Gamma_{H_{0}} F, G\right\rangle\right\rangle_{T}-i\left\langle\left\langle V_{1} \Gamma_{H_{0}} F, V_{2} \Gamma_{H}^{*} G\right\rangle\right\rangle_{T} . \tag{2.16}
\end{equation*}
$$

Now applying (2.8), (2.9) to (2.15) implies that

$$
\begin{equation*}
\left|\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}\right| \leq s_{2}\|F\|_{L_{T}^{2} \mathcal{A}}\|G\|_{L_{T}^{2} \mathcal{B}}+s_{3}\left\|V_{2} \Gamma_{H} F\right\|_{L_{T}^{2} \mathcal{H}}\|G\|_{L_{T}^{2} \mathcal{B}} \tag{2.17}
\end{equation*}
$$

It remains to deal with $\left\|V_{2} \Gamma_{H} F\right\|_{L_{T}^{2} \mathcal{H}}=\sup _{\|\widetilde{G}\|_{L_{T}^{2} \mathscr{H}}=1}\left|\left\langle\left\langle V_{2} \Gamma_{H} F, \widetilde{G}\right\rangle\right\rangle_{T}\right|$. Since $D(H)$ is dense in $\mathcal{H}$, we may assume $\widetilde{G}(t) \in D(H)$. Then, taking $D(H) \subset D\left(V_{2}^{*}\right)$ into account, we use (2.16) with $G=V_{2}^{*} \widetilde{G}$, (2.10), (2.11) and (2.13) to obtain $\left|\left\langle\left\langle V_{2} \Gamma_{H} F, \widetilde{G}\right\rangle\right\rangle_{T}\right| \leq\left.\left(s_{5}+s_{4} s_{7}\right)| | F\right|_{L_{T}^{2} \mathcal{A}}$, which, together with (2.16), gives us the second assertion. This completes the proof.

## 3. Proof of Theorems 1.2 and 1.3

Let $H=-\Delta+V(x)$ be as in Theorem 1.2. This section is devoted to the proof of the main theorems. In what follows we use a standard notation $2^{*}=2 n /(n-2)$, $2_{*}=2 n /(n+2)$. We write $\Gamma_{0}=\Gamma_{-\Delta}$. Let us first recall various estimates for the free Schrödinger equation.

Lemma 3.1. There exists $C>0$ such that, for any $v \in L^{n, \infty}\left(\mathbb{R}^{n}\right)$ and $T>0$,

$$
\begin{align*}
\left\|e^{i t \Delta} \psi\right\|_{L_{T}^{2} L^{2^{*}, 2}} & \leq C\|\psi\|_{L^{2}},  \tag{3.1}\\
\left\|\Gamma_{0} F\right\|_{L_{T}^{2} L^{2^{*}, 2}} & \leq C\|F\|_{L_{T}^{2} L^{2 *}, 2}  \tag{3.2}\\
\left\|v(x) \Gamma_{0} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|v\|_{L^{n, \infty}}\|F\|_{L_{T}^{2} L^{2 *, 2}}  \tag{3.3}\\
\left\|v(x) \Gamma_{0}^{*} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|v\|_{L^{n, \infty}}\|F\|_{L_{T}^{2} L^{2 *, 2}} \tag{3.4}
\end{align*}
$$

Proof. (3.1) and (3.2) are endpoint Strichartz estimates proved by (19, Theorem 10.1]. The latter two estimates follow from (3.2), Hölder's inequality (1.11) and the duality.

Lemma 3.2. Let $w \in L^{2}(\mathbb{R})$ be as in Theorem 1.2, $\rho>1 / 2$ and $v \in L^{n, \infty}\left(\mathbb{R}^{n}\right)$. Then there exists $C>0$, independent of $w, v$, and $T>0$, such that, for all $\psi \in L^{2}$ and simple function $F: \mathbb{R} \rightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)$, one has

$$
\begin{align*}
\left\|w(|x|)|D|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{T}^{2} L^{2}} & \leq C\|w\|_{L^{2}(\mathbb{R})}\|\psi\|_{L^{2}},  \tag{3.5}\\
\left\|w(|x|)|D|^{1 / 2} \Gamma_{0} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|w\|_{L^{2}(\mathbb{R})}\|F\|_{L_{T}^{2} L^{2 *}, 2}  \tag{3.6}\\
\left\|v(x) \Gamma_{0}^{*} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|v\|_{L^{n, \infty}}\|w\|_{L^{2}(\mathbb{R})}\left\|w(|x|)^{-1}|D|^{-1 / 2} F\right\|_{L_{T}^{2} L^{2}},  \tag{3.7}\\
\left\|\langle x\rangle^{-\rho}|D|^{1 / 2} \Gamma_{0} F\right\|_{L_{T}^{2} L^{2}} & \leq C\left\|\langle x\rangle^{\rho}|D|^{-1 / 2} F\right\|_{L_{T}^{2} L^{2}},  \tag{3.8}\\
\left\|\langle x\rangle^{-\rho}|D|^{1 / 2} \Gamma_{0} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|F\|_{L_{T}^{2} L^{2 *}, 2}  \tag{3.9}\\
\left\|v(x) \Gamma_{0}^{*} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|v\|_{L^{n, \infty}}\left\|\langle x\rangle^{\rho}|D|^{-1 / 2} F\right\|_{L_{T}^{2} L^{2}},  \tag{3.10}\\
\left\|v(x) \Gamma_{0} F\right\|_{L_{T}^{2} L^{2}} & \leq C\|v\|_{L^{n, \infty}}\left\|\langle x\rangle^{\rho}|D|^{-1 / 2} F\right\|_{L_{T}^{2} L^{2}} . \tag{3.11}
\end{align*}
$$

Proof. Let us first consider (3.5). When $n=1$, it was proved by [20] that

$$
\sup _{x \in \mathbb{R}}\left|\left\|\left.D_{x}\right|^{1 / 2} e^{i t \partial_{x}^{2}} f\right\|_{L^{2}\left(\mathbb{R}_{t}\right)} \leq C\|f\|_{L^{2}\left(\mathbb{R}_{x}\right)}\right.
$$

which, together with the unitarity of $e^{i t \Delta_{\widehat{x}_{j}}}$ in $L^{2}\left(\mathbb{R}^{n-1}\right)$, implies that

$$
\left\|\left|D_{j}\right|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{x_{j}}^{\infty} L_{T}^{2} L_{\hat{x}_{j}}^{2}} \leq C\|\psi\|_{L^{2}}, \quad j=1,2, \ldots, n
$$

uniformly in $T>0$, where $\widehat{x}_{j}=\left(x_{1}, \ldots, x_{j-1}, x_{j+1}, \ldots, x_{n}\right) \in \mathbb{R}^{n-1}$ and $D_{j}=-i \partial_{x_{j}}$. (3.5) then is derived from this estimate as follows. Let $\left\{C_{j}(\xi)\right\}$ be a conical partition of unity on $\mathbb{R}^{n}$ so that $I=\sum_{j=1}^{n} C_{j}(\xi)$, where $C_{j} \in$ $C^{\infty}\left(\mathbb{R}^{n} \backslash\{0\}\right)$ such that $\operatorname{supp} C_{j} \subset\left\{2\left|\xi_{j}\right|>|\xi|\right\}$ and $\partial_{\xi}^{\alpha} C_{j}(\xi)=O\left(|\xi|^{-|\alpha|}\right)$. If we set $\widetilde{C}_{j}(\xi)=C_{j}(\xi)|\xi|^{1 / 2}\left|\xi_{j}\right|^{-1 / 2}$, then $\widetilde{C}_{j}$ also satisfies $\partial_{\xi}^{\alpha} C_{j}(\xi)=O\left(|\xi|^{-|\alpha|}\right)$ and $|\xi|^{1 / 2}=\sum_{j=1}^{n} \widetilde{C}_{j}(\xi)\left|\xi_{j}\right|^{1 / 2}$. Since $w(|x|)^{2}$ belongs to the Muckenhoupt $A_{2}$-class, $\widetilde{C}_{j}(D)$ is bounded on a weighted space $L^{2}\left(\mathbb{R}^{n}, w(|x|)^{2} d x\right)$ by weighted Mikhlin's multiplier theorem (see 21). Thus we conclude that

$$
\begin{aligned}
\left\|w(|x|)|D|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{T}^{2} L^{2}}^{2} & \leq \sum_{j=1}^{n}\left\|w(|x|) \widetilde{C}_{j}(D)\left|D_{j}\right|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{T}^{2} L^{2}}^{2} \\
& \leq C \sum_{j=1}^{n}\left\|w(|x|)\left|D_{j}\right|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{T}^{2} L^{2}}^{2} \\
& \leq C \sum_{j=1}^{n}\|w\|_{L^{2}(\mathbb{R})}^{2}\left|\left\|\left.D_{j}\right|^{1 / 2} e^{i t \Delta} \psi\right\|_{L_{x_{j}}^{\infty} L_{T}^{2} L_{\widehat{x}_{j}}^{2}}^{2}\right. \\
& \leq C\|w\|_{L^{2}(\mathbb{R})}^{2}\|\psi\|_{L^{2}}^{2}
\end{aligned}
$$

uniformly in $T>0$, where we used the properties $w(|x|) \leq C_{j} w\left(x_{j}\right)$ and $w \in L^{2}(\mathbb{R})$ in the third line. Next, by the same argument as above, (3.6) follows from the estimate

$$
\left|\left\|\left.D_{j}\right|^{1 / 2} \Gamma_{0} F\right\|_{L_{x_{j}}^{\infty} L_{T}^{2} L_{\bar{x}_{j}}^{2}} \leq C\|F\|_{L_{T}^{2} L^{2} *, 2},\right.
$$

which is a slight generalization of [16, Lemma 4], in which the same estimate with $L^{2_{*}, 2}$ replaced by $L^{2_{*}}$ was proved. Although the proof is essentially the same as
that of [16, Lemma 4], we briefly recall its strategy for the reader's convenience. Without loss of generality, we may assume $j=1$. Then it suffices to show that

$$
\begin{equation*}
\sup _{x_{1}}| |\left|D_{j}\right|^{1 / 2} \widetilde{\Gamma}_{0} F\left\|_{L_{T}^{2} L_{\widehat{x}_{1}}^{2}} \leq C\right\| F \|_{L_{T}^{2} L^{2 *, 2}} \tag{3.12}
\end{equation*}
$$

where $\widetilde{\Gamma}_{0}$ is defined by

$$
\widetilde{\Gamma}_{0} F(t):=\int_{-\infty}^{t} e^{i(t-s) \Delta} F(s) d s
$$

Indeed, the corresponding estimate for $\Gamma_{0}-\widetilde{\Gamma}_{0}$ follows from (3.5) and the dual estimate of (3.1). The only difference from the proof of [16, Lemma 4] is an interpolation step. While they used the complex interpolation, we will use a real interpolation technique as in [19, section 6]. Let $I_{ \pm}=1_{ \pm}\left(D_{1}\right)\left|D_{1}\right|^{1 / 2} \widetilde{\Gamma}_{0}$, where $\mathbf{1}_{ \pm}(t)=1$ for $\pm t \geq 0$ and $\mathbf{1}_{ \pm}(t)=0$ for $\mp t \geq 0$. It suffices to show that $I_{+}$is bounded from $L_{T}^{2} L^{2 *, 2}$ to $L_{T}^{2} L_{\widehat{x}_{1}}^{2}$ uniformly in $x_{1}$ since the proof for $I_{-}$is analogous. By the $T T^{*}$ argument, $I_{+} \in \mathbb{B}\left(L_{T}^{2} L^{2_{*}, 2}, L_{T}^{2} L_{\widehat{x}_{1}}^{2}\right)$ if $I_{+}^{*} I_{+}$is bounded from $L_{T}^{2} L^{2_{*}, 2}$ to $L_{T}^{2} L^{2^{*}, 2}$. Hence, if we define a bilinear form $I$ by

$$
I(F, G):=\iint\left\langle I_{+} F(s, \cdot), I_{+} G(t, \cdot)\right\rangle d s d t
$$

then it suffices to show that

$$
\begin{equation*}
|I(F, G)| \leq C\|F\|_{L_{T}^{2} L^{2 *}, 2}\|G\|_{L_{T}^{2} L^{2 *, 2}} \tag{3.13}
\end{equation*}
$$

uniformly in $x_{1}$ and $T>0$. It was shown by 16 that $I_{+}^{*} I_{+}$is bounded on $L^{2}\left(\mathbb{R}^{n}\right)$ and the kernel of $I_{+}^{*} I_{+}$, denoted by $K_{+}(t, s, x, y)$, satisfies the dispersive estimate:

$$
\left|K_{+}(t, s, x, y)\right| \leq C|t-s|^{-n / 2}, \quad t \neq s
$$

We then decompose $I(F, G)$ as

$$
I(F, G)=\sum_{k \in \mathbb{Z}} I_{k}(F, G), I_{k}(F, G):=\iint_{t-2^{k+1}}^{t-2^{k}}\left\langle I_{+} F(s, \cdot), I_{+} G(t, \cdot)\right\rangle d s d t
$$

By using the same argument as in [19, Lemma 4.1], we see that

$$
\left|I_{k}(F, G)\right| \leq C 2^{-k \beta(a, b)}| | F\left\|_{L_{T}^{2} L^{a^{\prime}}}\right\| G \|_{L_{T}^{2} L^{b^{\prime}}}, \quad \beta(a, b)=\frac{n}{2}-1-\frac{n}{2}\left(\frac{1}{a}-\frac{1}{b}\right)
$$

uniformly in $k \in \mathbb{Z}$, where $(a, b)$ satisfies one of the following conditions:

$$
\begin{gathered}
\text { (i) } \frac{1}{a}=\frac{1}{b}=0 ; \quad \text { (ii) } \frac{n-1}{2 n} \leq \frac{1}{a} \leq \frac{1}{2} \text { and } \frac{1}{b}=\frac{1}{2} \text {; } \\
\text { (iii) } \frac{n-1}{2 n} \leq \frac{1}{b} \leq \frac{1}{2} \text { and } \frac{1}{a}=\frac{1}{2} .
\end{gathered}
$$

In other words, a vector valued sequence $\left(I_{k}\right)_{k \in \mathbb{Z}}$ is bounded from $L_{T}^{2} L_{x}^{a^{\prime}} \times L_{T}^{2} L_{x}^{b^{\prime}}$ to $\ell_{\beta(a, b)}^{\infty}$, where $\ell_{s}^{p}=L^{p}\left(\mathbb{Z}, 2^{j s} d j\right)$ is a weighted $\ell^{p}$ space with the counting measure $d j$. Then (3.13) follows from the technique by [19, Section 6] based on a bilinear real interpolation.

The estimate (3.7) follows from the dual estimate of (3.6) and Hölder's inequality (1.11). For (3.8), we refer to [8]. (3.9) and (3.10) follow from (3.6) and (3.7) since
$\langle x\rangle^{-\rho}$ satisfies the condition on $w$ in Theorem 1.2, In order to derive (3.11), we observe from the formula (2.4) that $\Gamma_{0}$ can be brought to the form

$$
\Gamma_{0} F(t)=-\Gamma_{0}^{*} F(t) \pm \Gamma_{0}^{0} F(t)+\Gamma_{0}^{\mp} F(t)
$$

for $\pm t \geq 0$, where

$$
\Gamma_{0}^{0} F(t)=\int_{-T}^{T} e^{-i(t-s) H} F(s) d s, \quad \Gamma_{0}^{ \pm} F(t)=\int_{0}^{ \pm T} e^{-i(t-s) H} F(s) d s
$$

Then the desired estimate for $\Gamma_{0}^{*}$ is nothing but (3.9); the desired estimates for $\Gamma_{0}^{0}$ and $\Gamma_{0}^{\mp}$ follow from (3.1), the dual estimate of (3.5) with $w=\langle x\rangle^{-\rho}$ and Hölder's inequality (1.11).

The following fact, proved by [1, Theorem 1.6 and (1.23)] (see also [3, Theorem 6.1 and Appendix B]), also plays an important role.

Proposition 3.3. Let $n \geq 3$ and $V$ satisfy Assumption A. Then $|x|^{-1}$ is H supersmooth.

We are in a position to show the main theorems.
Proof of Theorem 1.2. Let us set $V_{1}=|x| V$ and $V_{2}=|x|^{-1}$. By Sobolev's inequality (1.12),

$$
\left\|V_{j} f\right\|_{L^{2}} \leq C\left\|\mid V_{j}\right\|_{L^{n, \infty}}\|f\|_{L^{2^{*}, 2}} \leq C\left\|V_{j}\right\|_{L^{n, \infty}}\|\nabla f\|_{L^{2}}
$$

and hence $D\left(V_{j}\right) \supset \mathcal{H}^{1} \supset D(\Delta) \cup D(H)$. Moreover, (3.1) and Proposition 2.1(1) show that both $V_{1}$ and $V_{2}$ are $\Delta$-smooth. On the other hand, Propositions 2.1 and 3.3 show that

$$
\begin{equation*}
\left\|V_{2} e^{-i t H} \psi\right\|_{L_{T}^{2} L^{2}} \leq C\|\psi\|_{L^{2}} \tag{3.14}
\end{equation*}
$$

uniformly in $T>0$. Let $\mathcal{B}$ be the completion of $C_{0}^{\infty}$ with respect to the norm $\|\left. w(|x|)^{-1}|D|^{-1 / 2} f\right|_{L^{2}}$. By virtue of (3.5), (3.7) and (3.14), one can use Proposition 2.3 with $H_{0}=-\Delta, H=-\Delta+V$ and this $\mathcal{B}$ to obtain

$$
\left|\left\langle\left\langle e^{-i t H} \psi, G\right\rangle\right\rangle_{T}\right| \leq C\|w\|_{L^{2}(\mathbb{R})}\|\psi\|_{L^{2}}\|G\|_{\mathcal{B}}
$$

for all $\psi \in L^{2}$ and simple function $G:[-T, T] \rightarrow \mathcal{S}$ uniformly in $T>0$. Then the desired estimate follows from density and duality arguments.

Proof of Theorem 1.3. We use the same decomposition $V=V_{1} V_{2}$ as above. Since $V_{2}$ is $H$-supersmooth, we learn by Proposition 2.1 and a remark after Lemma 2.2 that

$$
\begin{equation*}
\left\|V_{2} \Gamma_{H} V_{2} \widetilde{G}\right\|_{L_{T}^{2} L^{2}} \leq C\|\widetilde{G}\|_{L_{T}^{2} L^{2}} \tag{3.15}
\end{equation*}
$$

for all simple functions $\widetilde{G}: \mathbb{R} \rightarrow D\left(V_{2}\right)$ with the constant $C$ independent of $T$ and $\widetilde{G}$. By virtue of (3.1)-(3.4), (3.5) with $w=\langle x\rangle^{-\rho}$, (3.8)-(3.11) with $v \in\left\{V_{1}, V_{2}\right\}$, (3.14) and (3.15), we can use Proposition 2.3 with $\left.\mathcal{A}, \mathcal{B} \in\left\{\dot{\mathcal{H}}^{-1 / 2, \rho}, L^{2 *}\right)^{2}\right\}$ to obtain

$$
\left|\left\langle\left\langle e^{-i t H} \psi, G\right\rangle\right\rangle_{T}\right| \leq C\|\psi\|_{L^{2}}\|G\|_{L_{T}^{2} \mathcal{B}}, \quad\left|\left\langle\left\langle\Gamma_{H} F, G\right\rangle\right\rangle_{T}\right| \leq C\|F\|_{L_{T}^{2} \mathcal{A}}\|G\|_{L_{T}^{2} \mathcal{B}}
$$

uniformly in $T>0, \psi \in L^{2}$ and simple functions $F, G: \mathbb{R} \rightarrow \mathcal{S}$. Then the assertion follows from density of simple functions $F: \mathbb{R} \rightarrow \mathcal{S}$ in $L_{T}^{2} \mathcal{A}$ and $L_{T}^{2} \mathcal{B}$ and the formula (1.9).

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