# FISHER-KOLMOGOROV TYPE PERTURBATIONS OF THE RELATIVISTIC OPERATOR: DIFFERENTIAL VS. DIFFERENCE 

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(Communicated by Joachim Krieger)
Dedicated to Jean Mawhin for his 75th anniversary
Abstract. We are concerned with the existence of multiple periodic solutions for differential equations involving Fisher-Kolmogorov perturbations of the relativistic operator of the form

$$
-\left[\phi\left(u^{\prime}\right)\right]^{\prime}=\lambda u\left(1-|u|^{q}\right)
$$

as well as for difference equations, of type

$$
-\Delta[\phi(\Delta u(n-1))]=\lambda u(n)\left(1-|u(n)|^{q}\right) ;
$$

here $q>0$ is fixed, $\Delta$ is the forward difference operator, $\lambda>0$ is a real parameter and

$$
\phi(y)=\frac{y}{\sqrt{1-y^{2}}} \quad(y \in(-1,1)) .
$$

The approach is variational and relies on critical point theory for convex, lower semicontinuous perturbations of $C^{1}$-functionals.

## 1. Introduction

This paper is concerned with problems involving Fisher-Kolmogorov nonlinearities of the type:

$$
\begin{equation*}
-\left[\phi\left(u^{\prime}\right)\right]^{\prime}=\lambda u\left(1-|u|^{q}\right), \quad u(0)-u(T)=0=u^{\prime}(0)-u^{\prime}(T), \tag{1.1}
\end{equation*}
$$

respectively,

$$
\begin{equation*}
-\Delta[\phi(\Delta u(n-1))]=\lambda u(n)\left(1-|u(n)|^{q}\right), \quad u(n)=u(n+T) \quad(n \in \mathbb{Z}) \tag{1.2}
\end{equation*}
$$

where $q>0$ is fixed, $\Delta u(n)=u(n+1)-u(n)$ is the usual forward difference operator, $\lambda>0$ is a real parameter and

$$
\phi(y)=\frac{y}{\sqrt{1-y^{2}}} \quad(y \in(-1,1)) .
$$

Notice that besides the trivial solution, problems (1.1) and (1.2) always have the pair of constant solutions $u \equiv \pm 1$ and these are the only constant nontrivial solutions of (1.1) and (1.2). Here we are interested in the multiplicity of pairs of nonconstant solutions of (1.1) and (1.2).

[^0]The typical example which involves the above type of nonlinearities was originally motivated by models in biological population dynamics and led to the scalar reaction-diffusion equation

$$
\frac{\partial u}{\partial t}-\frac{\partial^{2} u}{\partial x^{2}}=u\left(1-u^{2}\right)
$$

referred to as the classical Fisher-Kolmogorov (FK) equation (14, [15], [18]). In the last years interest has turned to higher-order equations of type

$$
u^{i v}-p u^{\prime \prime}=u\left(1-u^{2}\right),
$$

which corresponds, if $p>0$, to the extended Fisher-Kolmogorov (EFK) equations; these are models for phase transitions and other bistable phenomena. In this direction we refer the reader to [11], [12], [21] - [24], [28] where existence of solutions was studied by a variety of methods such as topological shooting methods, phase-plane analysis and variational methods. Also, a difference equation related to the FK equation was considered in [1], 10].

A multiplicity result as the one in this paper (Theorem 2.1) was obtained in [11], [12], [28] for EFK equations. We notice also the paper [8], where a multiplicity result is given for periodic problems involving the discrete $p$-Laplacian operator.

On the other hand, in recent years special attention was paid to various qualitative aspects for boundary value problems involving the so-called relativistic operator: $u \mapsto\left[\phi\left(u^{\prime}\right)\right]^{\prime}$. Among others and far from being exhausted, related to existence and multiplicity of periodic solutions for such problems, we refer the reader to 3] - 5], [7] [9, [16], respectively to [2] [17], 19, [20] for discrete versions.

It is the aim of this paper to obtain multiplicity of nonconstant solutions for problems (1.1) and (1.2). First, let us note that both of them can be seen as eigenvalue problems. In this view, we prove in Theorem 2.1 (resp. Theorem 3.1) that (1.1) (resp. (1.2)) has a prescribed number of distinct pairs of nonconstant solutions for large enough values of the parameter $\lambda$. On the other hand, for any $\lambda>0$ fixed, we obtain that a prescribed number of distinct pairs of nonconstant solutions can be obtained for (1.1), provided that the period $T$ is sufficiently large (Theorem 2.1). Our approach is a variational one and relies on a generalization of a result for smooth functionals due to Clark [13] to convex, lower semicontinuous perturbations of $C^{1}$-functionals.

Before concluding this introductory part, we briefly recall some topics in the frame of Szulkin's critical point theory [27, which will be needed in the sequel.

Let $(Y,\|\cdot\|)$ be a real Banach space and let $\mathcal{I}: Y \rightarrow(-\infty,+\infty]$ be a functional of the type

$$
\begin{equation*}
\mathcal{I}=\mathcal{F}+\psi \tag{1.3}
\end{equation*}
$$

where $\mathcal{F} \in C^{1}(Y, \mathbb{R})$ and $\psi: Y \rightarrow(-\infty,+\infty]$ is convex, lower semicontinuous and proper (i.e., $D(\psi):=\{u \in Y: \psi(u)<+\infty\} \neq \emptyset$ ). A point $u \in Y$ is said to be $a$ critical point of $\mathcal{I}$ if $u \in D(\psi)$ and satisfies the inequality

$$
\left\langle\mathcal{F}^{\prime}(u), v-u\right\rangle+\psi(v)-\psi(u) \geq 0 \quad \forall v \in D(\psi) .
$$

A sequence $\left\{u_{n}\right\} \subset D(\psi)$ is called a (PS)-sequence if $\mathcal{I}\left(u_{n}\right) \rightarrow c \in \mathbb{R}$ and

$$
\left\langle\mathcal{F}^{\prime}\left(u_{n}\right), v-u_{n}\right\rangle+\psi(v)-\psi\left(u_{n}\right) \geq-\varepsilon_{n}\left\|v-u_{n}\right\| \quad \forall v \in D(\psi),
$$

where $\varepsilon_{n} \rightarrow 0$. The functional $\mathcal{I}$ is said to satisfy the (PS) condition if any (PS)sequence has a convergent subsequence in $Y$.

Let $\Sigma$ be the collection of all symmetric subsets of $Y \backslash\{0\}$ which are closed in $Y$. A nonempty set $A \in \Sigma$ is said to have genus $k$ (denoted $\gamma(A)=k$ ) if $k$ is the smallest integer with the property that there exists an odd continuous mapping $h: A \rightarrow \mathbb{R}^{k} \backslash\{0\}$. If such an integer does not exist, $\gamma(A)=+\infty$. For properties and more details of the notion of genus we refer the reader to [25, 26]. Let $\Gamma \subset 2^{Y}$ be the collection of all nonempty compact symmetric subsets of $Y$, considered with the Hausdorff-Pompeiu distance and

$$
\Gamma_{j}:=\operatorname{cl}\{A \in \Gamma: 0 \notin A, \gamma(A) \geq j\}
$$

( $c l$ is the closure in $\Gamma$ ). The following is a generalization of the result for smooth functions in [25, Theorem 5.19] to functionals of type (1.3) and it is proved in [27, Theorem 4.3].
Theorem 1.1. Let $\mathcal{I}$ be of type (1.3) with $\mathcal{F}$ and $\psi$ even. Also, suppose that $\mathcal{I}$ is bounded from below, satisfies the (PS) condition and $\mathcal{I}(0)=0$. If

$$
c_{m}:=\inf _{A \in \Gamma_{m}} \sup _{v \in A} \mathcal{I}(v)<0,
$$

then the functional $\mathcal{I}$ has at least $m$ distinct pairs of nontrivial critical points.

## 2. The differential problem (1.1)

Using the ideas from [4, we introduce a variational formulation for problem (1.1). With this aim let $C:=C[0, T]$ be endowed with the usual supremum norm $\|\cdot\|_{\infty}$ and $W^{1, \infty}:=W^{1, \infty}(0, T)$. For each $v \in C$ we set $\bar{v}:=\frac{1}{T} \int_{0}^{T} v(t) d t$ and we write $v(t)=\bar{v}+\tilde{v}(t)$, where $\overline{\tilde{v}}=0$. If $v \in W^{1, \infty}$, then $\tilde{v}$ vanishes at some $t_{0} \in(0, T)$ and so

$$
|\tilde{v}(t)|=\left|\tilde{v}(t)-\tilde{v}\left(t_{0}\right)\right| \leq \int_{0}^{T}\left|v^{\prime}(s)\right| d s \leq T\left\|v^{\prime}\right\|_{\infty}
$$

Next, denoting

$$
K:=\left\{v \in W^{1, \infty}:\left\|v^{\prime}\right\|_{\infty} \leq 1, v(0)=v(T)\right\}
$$

it is clear that

$$
\begin{equation*}
\|\tilde{v}\|_{\infty} \leq T \quad \forall v \in K \tag{2.1}
\end{equation*}
$$

Also, it is not difficult to show that (see [4, Lemma 4])

$$
\begin{equation*}
|v(t)|^{p} \geq|\bar{v}|^{p}-p T|\bar{v}|^{p-1} \quad \forall v \in K, \forall t \in[0, T] \text { and } p \geq 1 . \tag{2.2}
\end{equation*}
$$

From [6] we know that the even functional $\Psi: C \rightarrow(-\infty,+\infty$ ],

$$
\Psi(v)= \begin{cases}\int_{0}^{T}\left[1-\sqrt{1-v^{\prime 2}}\right] \quad \text { if } v \in K, \\ +\infty \text { otherwise }\end{cases}
$$

is proper, convex and lower semicontinuous on $C$ and it is easy to see that

$$
\begin{equation*}
\Psi(v) \leq \int_{0}^{T}\left|v^{\prime}\right|^{2} \quad \forall v \in K \tag{2.3}
\end{equation*}
$$

Next, let $\mathcal{G}_{\lambda}: C \rightarrow \mathbb{R}$ be defined by

$$
\mathcal{G}_{\lambda}(u)=\lambda \int_{0}^{T}\left[\frac{|u|^{q+2}}{q+2}-\frac{u^{2}}{2}\right] .
$$

Notice that $\mathcal{G}_{\lambda}$ is even, of class $C^{1}$ on $C$ and its derivative is given by

$$
\left\langle\mathcal{G}_{\lambda}^{\prime}(u), v\right\rangle=\lambda \int_{0}^{T}\left(|u|^{q}-1\right) u v, \quad u, v \in C .
$$

Then the energy functional $I_{\lambda}: C \rightarrow(-\infty,+\infty]$ associated to problem (1.1) is given by

$$
I_{\lambda}=\Psi+\mathcal{G}_{\lambda}
$$

and it has the structure required by Szulkin's critical point theory.
Recall, by a solution of (1.1) we mean a function $u \in C^{1}[0, T]$, with $\left\|u^{\prime}\right\|_{\infty}<1$ and $\phi\left(u^{\prime}\right)$ differentiable, which satisfies (1.1).

From Proposition 2 in [4, one has the following:
Proposition 2.1. If $u \in K$ is a critical point of $I_{\lambda}$, then $u$ is a solution of problem (1.1).

Lemma 2.1. $I_{\lambda}$ is bounded from below and satisfies the (PS) condition.
Proof. Let $u \in K=D(\Psi)$. From (2.1) we have

$$
\int_{0}^{T} \frac{u^{2}}{2}=\int_{0}^{T} \frac{(\bar{u}+\tilde{u})^{2}}{2} \leq \frac{T^{3}}{2}+\frac{T}{2}|\bar{u}|^{2} .
$$

Also, on account of (2.2) we obtain

$$
\mathcal{G}_{\lambda}(u) \geq \frac{\lambda T}{q+2}|\bar{u}|^{q+2}-\lambda T^{2}|\bar{u}|^{q+1}-\lambda \int_{0}^{T} \frac{u^{2}}{2} .
$$

It follows

$$
\begin{equation*}
I_{\lambda}(u) \geq \mathcal{G}_{\lambda}(u) \geq \frac{\lambda T}{q+2}|\bar{u}|^{q+2}-\lambda T^{2}|\bar{u}|^{q+1}-\frac{\lambda T}{2}|\bar{u}|^{2}-\frac{\lambda T^{3}}{2} \tag{2.4}
\end{equation*}
$$

which clearly shows that $I_{\lambda}$ is bounded from below. To see that $I_{\lambda}$ satisfies the (PS) condition, let $\left\{u_{n}\right\} \subset K=D(\Psi)$ be a (PS)-sequence. We write (2.4) with $u_{n}$ instead of $u$ and, from the fact that $\left\{I_{\lambda}\left(u_{n}\right)\right\}$ is bounded, we get that $\left\{\bar{u}_{n}\right\}$ is bounded. Then, Lemma 3 ii) in [4] ensures that $\left\{u_{n}\right\}$ has a convergent subsequence in $C$.

Theorem 2.1. If $\lambda>4 \pi^{2} m^{3} / T^{2}$ for some $m \in \mathbb{N}, m \geq 2$, then problem (1.1) has at least $m-1$ distinct pairs of nonconstant solutions.

Proof. Using that $u= \pm 1$ is the only pair of nontrivial constant solutions for (1.1), it suffices to prove that (1.1) has at least $m$ distinct pairs of nontrivial solutions. From Theorem 1.1. Proposition 2.1] and Lemma 2.1) this can be reduced to showing that there is some $A_{m} \in \Gamma_{m} \subset 2^{C}$ such that

$$
\begin{equation*}
\sup _{v \in A_{m}} I_{\lambda}(v)<0 . \tag{2.5}
\end{equation*}
$$

With this aim, we consider the finite dimensional space

$$
X_{m}:=\operatorname{span}\left\{\sin \frac{\pi x}{T}, \sin \frac{2 \pi x}{T}, \ldots, \sin \frac{m \pi x}{T}\right\}
$$

equipped with the norm

$$
\left\|\alpha_{1} \sin \frac{\pi x}{T}+\cdots+\alpha_{m} \sin \frac{m \pi x}{T}\right\|_{X_{m}}^{2}=\alpha_{1}^{2}+\cdots+\alpha_{m}^{2} .
$$

Since the norms $\|\cdot\|_{X_{m}}$ and $\|\cdot\|_{L^{q+2}}$ are equivalent on $X_{m}$, there exists a positive constant $c(m)$ such that

$$
\begin{equation*}
\|v\|_{L^{q+2}} \leq c(m)\|v\|_{X_{m}} \tag{2.6}
\end{equation*}
$$

Next, as in e.g. [24], [28], we introduce the subset $A_{m}$ of $C$ by

$$
A_{m}=\left\{\sum_{k=1}^{m} \alpha_{k} \sin \frac{k \pi x}{T}: \alpha_{1}^{2}+\cdots+\alpha_{m}^{2}=\rho^{2}\right\}\left(\subset X_{m}\right)
$$

where, since $\lambda>4 \pi^{2} m^{3} / T^{2}$, the positive number $\rho$ can be chosen $\leq 2 / \sqrt{\lambda}$ and such that

$$
\frac{m^{3} \pi^{2}}{T}-\frac{\lambda T}{4}+\frac{\lambda(c(m))^{q+2}}{q+2} \rho^{q}<0 .
$$

It is easy to see that the odd mapping $H: A_{m} \rightarrow S^{m-1}(m-1$ dimension unit sphere in the Euclidean space $\mathbb{R}^{m}$ ) defined by

$$
H\left(\sum_{k=1}^{m} \alpha_{k} \sin \frac{k \pi x}{T}\right)=\left(\frac{\alpha_{1}}{\rho}, \ldots, \frac{\alpha_{m}}{\rho}\right)
$$

is a homeomorphism between $A_{m}$ and $S^{m-1}$. According to [26, Corrolary 5.5], $\gamma\left(A_{m}\right)=m$ and so, $A_{m} \in \Gamma_{m}$.

Let $v \in A_{m}$. Clearly, $v(0)=v(T)$ and we have

$$
\begin{align*}
\left|v^{\prime}\right| & \leq \sum_{k=1}^{m}\left|\alpha_{k} \frac{k \pi}{T} \cos \frac{k \pi x}{T}\right| \leq \frac{m \pi}{T} \sum_{k=1}^{m}\left|\alpha_{k}\right| \\
& \leq \frac{m^{3 / 2} \pi}{T}\left(\sum_{k=1}^{m} \alpha_{k}^{2}\right)^{1 / 2}=\frac{m^{3 / 2} \pi}{T} \rho . \tag{2.7}
\end{align*}
$$

Therefore, as $T>2 \pi m \sqrt{m / \lambda} \geq \pi m^{3 / 2} \rho$, one has $\left\|v^{\prime}\right\|_{\infty}<1$, meaning that $v \in K$.
On the other hand, we compute

$$
\begin{align*}
\int_{0}^{T} v^{2} & =\int_{0}^{T}\left(\sum_{k=1}^{m} \alpha_{k} \sin \frac{k \pi x}{T}\right)^{2}=\sum_{k=1}^{m} \alpha_{k}^{2} \int_{0}^{T} \sin ^{2} \frac{k \pi x}{T} \\
& =\frac{1}{2} \sum_{k=1}^{m} \alpha_{k}^{2} \int_{0}^{T}\left(1-\cos \frac{2 k \pi x}{T}\right)=\frac{T}{2} \rho^{2} . \tag{2.8}
\end{align*}
$$

Then, using (2.3), (2.6) - (2.8), we estimate $I_{\lambda}$ as follows:

$$
\begin{aligned}
I_{\lambda}(v) & =\Psi(v)+\lambda \int_{0}^{T} \frac{|v|^{q+2}}{q+2}-\lambda \int_{0}^{T} \frac{v^{2}}{2} \leq \int_{0}^{T}\left|v^{\prime}\right|^{2}+\frac{\lambda(c(m) \rho)^{q+2}}{q+2}-\frac{\lambda T}{4} \rho^{2} \\
& \leq \rho^{2}\left[\frac{m^{3} \pi^{2}}{T}-\frac{\lambda T}{4}+\frac{\lambda(c(m))^{q+2} \rho^{q}}{q+2}\right](<0-\text { from the choice of } \rho),
\end{aligned}
$$

which shows that (2.5) holds true.

## 3. The difference problem (1.2)

Analogously to the previous section, we first give the variational formulation for problem (1.2). Let $H_{T}$ be the space of all $T$-periodic $\mathbb{Z}$-sequences in $\mathbb{R}$, i.e., of mappings $u: \mathbb{Z} \rightarrow \mathbb{R}$ such that $u(n)=u(n+T)$ for all $n \in \mathbb{Z}$. On $H_{T}$ we shall refer to the following two (equivalent) norms:

$$
\|u\|=\left(\sum_{j=1}^{T}|u(j)|^{2}\right)^{1 / 2} \quad \text { and } \quad\|u\|_{q+2}=\left(\sum_{j=1}^{T}|u(j)|^{q+2}\right)^{\frac{1}{q+2}}
$$

Also, for each $u \in H_{T}$ we set

$$
\bar{u}:=\frac{1}{T} \sum_{j=1}^{T} u(j), \quad \tilde{u}:=u-\bar{u} .
$$

Let the closed convex subset $\mathbf{K}$ of $H_{T}$ be defined by

$$
\mathbf{K}:=\left\{u \in H_{T}:|\Delta u|_{\infty} \leq 1\right\},
$$

where $|\Delta u|_{\infty}:=\max _{j=1, \ldots, T}|\Delta u(j)|$. We introduce the (even) functions

$$
\boldsymbol{\Psi}(u)=\left\{\begin{array}{l}
\sum_{j=1}^{T} \Phi[\Delta u(j)] \quad \text { if } u \in \mathbf{K} \\
+\infty \quad \text { otherwise }
\end{array}\right.
$$

where $\Phi(y)=1-\sqrt{1-y^{2}}(y \in[-1,1])$, respectively

$$
\mathbf{G}_{\lambda}(u)=\lambda \sum_{j=1}^{T}\left[\frac{|u(j)|^{q+2}}{q+2}-\frac{u(j)^{2}}{2}\right]
$$

Then the functional $\mathbf{I}_{\lambda}: H_{T} \rightarrow(-\infty,+\infty]$ associated to problem (1.2) will be given by

$$
\mathbf{I}_{\lambda}=\boldsymbol{\Psi}+\mathbf{G}_{\lambda}
$$

and it is not difficult to see that it has the structure required by Szulkin's critical point theory, the derivative of $\mathbf{G}_{\lambda}$ being given by

$$
\left\langle\mathbf{G}_{\lambda}^{\prime}(u), v\right\rangle=\lambda \sum_{j=1}^{T}\left(|u(j)|^{q}-1\right) u(j) v(j), \quad\left(u, v \in H_{T}\right)
$$

A solution of problem (1.2) is an element $u \in H_{T}$ such that $|\Delta u(n)|<1$, for all $n \in \mathbb{Z}$, which satisfies the equation in (1.2).

Proposition 3.1. Any critical point of $\mathbf{I}_{\lambda}$ is a solution of problem (1.2).
Proof. For any $e \in H_{T}$, on account of Lemmas 5 and 6 in [19], problem

$$
\begin{equation*}
\Delta[\phi(\Delta u(n-1))]=\bar{u}+e(n), \quad u(n)=u(n+T) \quad(n \in \mathbb{Z}) \tag{3.1}
\end{equation*}
$$

has a unique solution $u_{e}$, which is also a solution of the variational inequality

$$
\begin{equation*}
\sum_{j=1}^{T}\{\Phi[\Delta v(j)]-\Phi[\Delta u(j)]+\bar{u}(\bar{v}-\bar{u})+e(j)(v(j)-u(j))\} \geq 0 \forall v \in \mathbf{K} \tag{3.2}
\end{equation*}
$$

We show that $u_{e}$ is actually the unique solution of (3.2). With this aim, let $J$ : $\mathbf{K} \rightarrow \mathbb{R}$ be defined by

$$
J(u)=\sum_{j=1}^{T}\left\{\Phi[\Delta u(j)]+\frac{\bar{u}^{2}}{2}+e(j) u(j)\right\} .
$$

If $u$ is a solution of (3.2), then the inequality $\frac{\bar{v}^{2}}{2}-\frac{\bar{u}^{2}}{2} \geq \bar{u}(\bar{v}-\bar{u})$ plugged in (3.2) implies that

$$
\sum_{j=1}^{T}\left\{\Phi[\Delta v(j)]-\Phi[\Delta u(j)]+\frac{\bar{v}^{2}}{2}+e(j) v(j)-\frac{\bar{u}^{2}}{2}-e(j) u(j)\right\} \geq 0 \forall v \in \mathbf{K}
$$

which shows that $J$ has a minimum at $u$. Then the uniqueness of $u_{e}$ as a solution of (3.2) follows by the strict convexity of $J$.

Next, let $w$ be a critical point of $\mathbf{I}_{\lambda}$. Then, for all $v \in \mathbf{K}$, one has

$$
\sum_{j=1}^{T}\left\{\Phi[\Delta v(j)]-\Phi[\Delta w(j)]+\lambda\left(|w(j)|^{q}-1\right) w(j)(v(j)-w(j))\right\} \geq 0
$$

which can be written

$$
\begin{aligned}
& \sum_{j=1}^{T}\{\Phi[\Delta v(j)]-\Phi[\Delta w(j)]+\bar{w}(v(j)-w(j))\} \\
+ & \sum_{j=1}^{T}\left[\lambda\left(|w(j)|^{q}-1\right) w(j)-\bar{w}\right](v(j)-w(j)) \geq 0
\end{aligned}
$$

for all $v \in \mathbf{K}$. Hence, $w$ is a solution of the variational inequality

$$
\begin{equation*}
\sum_{j=1}^{T}\left\{\Phi[\Delta v(j)]-\Phi[\Delta w(j)]+\bar{w}(\bar{v}-\bar{w})+e_{w}(j)(v(j)-w(j))\right\} \geq 0 \quad \forall v \in \mathbf{K} \tag{3.3}
\end{equation*}
$$

with $e_{w} \in H_{T}$ given by $e_{w}(n)=\lambda\left(|w(n)|^{q}-1\right) w(n)-\bar{w}(n \in \mathbb{Z})$.
Therefore, by (3.3) and the uniqueness of the solution of (3.2), one can see that actually $w$ solves problem (1.2).

Lemma 3.1. $\mathbf{I}_{\lambda}$ is bounded from below and satisfies the (PS) condition.
Proof. Let $u \in \mathbf{K}=D(\boldsymbol{\Psi})$. By the equivalence of the norms in $H_{T}$, there exists a positive constant $C_{1}$ such that

$$
\|u\|_{q+2}^{q+2} \geq C_{1}\|u\|^{q+2}
$$

Then, we have

$$
\begin{equation*}
\mathbf{I}_{\lambda}(u) \geq \mathbf{G}_{\lambda}(u) \geq \frac{\lambda C_{1}}{q+2}\|u\|^{q+2}-\frac{\lambda}{2}\|u\|^{2} \tag{3.4}
\end{equation*}
$$

which clearly shows that $\mathbf{I}_{\lambda}$ is bounded from below. Now, if $\left\{u_{k}\right\}$ is a sequence in $\mathbf{K}$ such that $\left\{\mathbf{I}_{\lambda}\left(u_{k}\right)\right\}$ is bounded, then one has from (3.4) that $\left\{u_{k}\right\}$ is bounded in $H_{T}$ and hence it contains a convergent subsequence.

Theorem 3.1. If $\lambda>8 m T$ for some $m \in \mathbb{N}$ with $2 \leq m \leq T$, then problem (1.2) has at least $m-1$ distinct pairs of nonconstant solutions.

Proof. Similar to the proof of Theorem 2.1] since $u= \pm 1$ is the only pair of nontrivial constant solutions for (1.2), it is sufficient to show that (1.2) has at least $m$ distinct pairs of nontrivial solutions. By virtue of Theorem 1.1 Proposition 3.1 and Lemma 3.1 we have to prove that there is some $A_{m} \in \Gamma_{m} \subset 2^{H_{T}}$ such that

$$
\begin{equation*}
\sup _{v \in A_{m}} \mathbf{I}_{\lambda}(v)<0 \tag{3.5}
\end{equation*}
$$

For this, let $w^{1}, w^{2}, \ldots, w^{T}$ be an orthonormal basis in the space $H_{T}$ endowed with the Euclidean norm $\|\cdot\|$. We consider the set

$$
A_{m}:=\left\{\sum_{k=1}^{m} \alpha_{k} w^{k}: \alpha_{1}^{2}+\cdots+\alpha_{m}^{2}=\rho^{2}\right\}
$$

where, since $\lambda>8 m T$, the positive number $\rho$ can be chosen $\leq 1 /(2 \sqrt{m})$ and such that

$$
4 m T-\frac{\lambda}{2}+\frac{\lambda T \sqrt{m^{q+2}}}{q+2} \rho^{q}<0
$$

Since the mapping $\mathbf{H}: A_{m} \rightarrow S^{m-1}$ defined by

$$
\mathbf{H}\left(\sum_{k=1}^{m} \alpha_{k} w^{k}\right)=\left(\frac{\alpha_{1}}{\rho}, \ldots, \frac{\alpha_{m}}{\rho}\right)
$$

is an odd homeomorphism between $A_{m}$ and $S^{m-1}$, then we have $\gamma\left(A_{m}\right)=m$. Hence, $A_{m} \in \Gamma_{m}$.

Now, let $v=\sum_{k=1}^{m} \alpha_{k} w^{k} \in A_{m}$. Then, for each $j=1, \ldots, T$, we obtain

$$
\begin{align*}
|\Delta v(j)| & \leq \sum_{k=1}^{m}\left|\alpha_{k} w^{k}(j+1)\right|+\sum_{k=1}^{m}\left|\alpha_{k} w^{k}(j)\right| \leq 2 \sum_{k=1}^{m}\left|\alpha_{k}\right| \\
& \leq 2 \sqrt{m}\left(\sum_{k=1}^{m} \alpha_{k}^{2}\right)^{1 / 2}=2 \rho \sqrt{m} \tag{3.6}
\end{align*}
$$

and since $\rho \leq 1 /(2 \sqrt{m})$, one has $|\Delta v|_{\infty} \leq 1$, which shows that $v \in \mathbf{K}$. On the other hand, we have

$$
\begin{equation*}
\sum_{j=1}^{T} v(j)^{2}=\|v\|^{2}=\sum_{k=1}^{m} \alpha_{k}^{2}=\rho^{2} \tag{3.7}
\end{equation*}
$$

and

$$
\begin{align*}
\sum_{j=1}^{T}|v(j)|^{q+2} & \leq \sum_{j=1}^{T}\left(\sum_{k=1}^{m}\left|\alpha_{k} \| w^{k}(j)\right|\right)^{q+2} \leq T\left(\sum_{k=1}^{m}\left|\alpha_{k}\right|\left\|w^{k}\right\|\right)^{q+2} \\
& =T\left(\sum_{k=1}^{m}\left|\alpha_{k}\right|\right)^{q+2} \leq T(\rho \sqrt{m})^{q+2} \tag{3.8}
\end{align*}
$$

Then, using (3.6) - (3.8), it follows

$$
\begin{aligned}
\mathbf{I}_{\lambda}(v) & =\boldsymbol{\Psi}(v)+\frac{\lambda}{q+2} \sum_{j=1}^{T}|v(j)|^{q+2}-\frac{\lambda}{2} \sum_{j=1}^{T} v(j)^{2} \\
& \leq \sum_{j=1}^{T}|\Delta v(j)|^{2}+\frac{\lambda T(\rho \sqrt{m})^{q+2}}{q+2}-\frac{\lambda \rho^{2}}{2} \\
& \leq \rho^{2}\left[4 m T-\frac{\lambda}{2}+\frac{\lambda T \sqrt{m^{q+2}}}{q+2} \rho^{q}\right](<0-\text { from the choice of } \rho) .
\end{aligned}
$$

Therefore, (3.5) holds true and the proof is complete.

## References

[1] Meline Aprahamian, Diko Souroujon, and Stepan Tersian, Decreasing and fast solutions for a second-order difference equation related to Fisher-Kolmogorov's equation, J. Math. Anal. Appl. 363 (2010), no. 1, 97-110, DOI 10.1016/j.jmaa.2009.08.009. MR2559044
[2] Cristian Bereanu and Dana Gheorghe, Topological methods for boundary value problems involving discrete vector $\phi$-Laplacians, Topol. Methods Nonlinear Anal. 38 (2011), no. 2, 265276. MR2932036
[3] Cristian Bereanu, Dana Gheorghe, and Manuel Zamora, Periodic solutions for singular perturbations of the singular $\phi$-Laplacian operator, Commun. Contemp. Math. 15 (2013), no. 4, 1250063, 22, DOI 10.1142/S0219199712500630. MR3073446
[4] Cristian Bereanu, Petru Jebelean, and Jean Mawhin, Variational methods for nonlinear perturbations of singular $\phi$-Laplacians, Atti Accad. Naz. Lincei Rend. Lincei Mat. Appl. 22 (2011), no. 1, 89-111, DOI 10.4171/RLM/589. MR2799910
[5] Cristian Bereanu, Petru Jebelean, and Jean Mawhin, Multiple solutions for Neumann and periodic problems with singular $\phi$-Laplacian, J. Funct. Anal. 261 (2011), no. 11, 3226-3246, DOI 10.1016/j.jfa.2011.07.027. MR 2835997
[6] Cristian Bereanu, Petru Jebelean, and Jean Mawhin, Radial solutions of Neumann problems involving mean extrinsic curvature and periodic nonlinearities, Calc. Var. Partial Differential Equations 46 (2013), no. 1-2, 113-122, DOI 10.1007/s00526-011-0476-x. MR3016504
[7] C. Bereanu and J. Mawhin, Existence and multiplicity results for some nonlinear problems with singular $\phi$-Laplacian, J. Differential Equations 243 (2007), no. 2, 536-557, DOI 10.1016/j.jde.2007.05.014. MR2371799
[8] Li-Hua Bian, Hong-Rui Sun, and Quan-Guo Zhang, Solutions for discrete p-Laplacian periodic boundary value problems via critical point theory, J. Difference Equ. Appl. 18 (2012), no. 3, 345-355, DOI 10.1080/10236198.2010.491825. MR 2901826
[9] Haïm Brezis and Jean Mawhin, Periodic solutions of the forced relativistic pendulum, Differential Integral Equations 23 (2010), no. 9-10, 801-810. MR2675583
[10] Alberto Cabada, Diko Souroujon, and Stepan Tersian, Heteroclinic solutions of a secondorder difference equation related to the Fisher-Kolmogorov's equation, Appl. Math. Comput. 218 (2012), no. 18, 9442-9450, DOI 10.1016/j.amc.2012.03.032. MR 2923041
[11] Julia Chaparova, Existence and numerical approximations of periodic solutions of semilinear fourth-order differential equations, J. Math. Anal. Appl. 273 (2002), no. 1, 121-136, DOI 10.1016/S0022-247X(02)00216-0. MR 1933020
[12] J. V. Chaparova, L. A. Peletier, and S. A. Tersian, Existence and nonexistence of nontrivial solutions of semilinear sixth-order ordinary differential equations, Appl. Math. Lett. 17 (2004), no. 10, 1207-1212, DOI 10.1016/j.aml.2003.05.014. MR2091859
[13] David C. Clark, A variant of the Lusternik-Schnirelman theory, Indiana Univ. Math. J. 22 (1972/73), 65-74, DOI 10.1512/iumj.1972.22.22008. MR0296777
[14] Paul C. Fife and J. B. McLeod, The approach of solutions of nonlinear diffusion equations to travelling front solutions, Arch. Ration. Mech. Anal. 65 (1977), no. 4, 335-361, DOI 10.1007/BF00250432. MR0442480
[15] R.A. Fisher, The advance of advantageous genes, Ann. Eugen. 7 (1937), 335-369.
[16] Petru Jebelean, Jean Mawhin, and Călin Şerban, Multiple periodic solutions for perturbed relativistic pendulum systems, Proc. Amer. Math. Soc. 143 (2015), no. 7, 3029-3039, DOI 10.1090/S0002-9939-2015-12542-7. MR 3336627
[17] Petru Jebelean, Jean Mawhin, and Călin Şerban, Morse theory and multiple periodic solutions of some quasilinear difference systems with periodic nonlinearities, Georgian Math. J. 24 (2017), no. 1, 103-112, DOI 10.1515/gmj-2016-0075. MR3607244
[18] A. Kolmogorov, I. Petrovski and N. Piscounov, Étude de l'équation de la diffusion avec croissance de la quantité de matière at son application á un probléme biologique, Bull. Univ. d'État á Moscou, Sér. Int., Sec.A, 1 (1937), 1-25.
[19] Jean Mawhin, Periodic solutions of second order nonlinear difference systems with $\phi$ Laplacian: a variational approach, Nonlinear Anal. 75 (2012), no. 12, 4672-4687, DOI 10.1016/j.na.2011.11.018. MR2927127
[20] J. Mawhin, A simple proof of multiplicity for periodic solutions of Lagrangian difference systems with relativistic operator and periodic potential, J. Difference Equ. Appl. 22 (2016), no. 2, 306-315, DOI 10.1080/10236198.2015.1089867. MR3474984
[21] L. A. Peletier and W. C. Troy, A topological shooting method and the existence of kinks of the extended Fisher-Kolmogorov equation, Topol. Methods Nonlinear Anal. 6 (1995), no. 2, 331-355, DOI 10.12775/TMNA.1995.049. MR1399544
[22] L. A. Peletier and W. C. Troy, Spatial patterns described by the extended Fisher-Kolmogorov equation: periodic solutions, SIAM J. Math. Anal. 28 (1997), no. 6, 1317-1353, DOI 10.1137/S0036141095280955. MR1474217
[23] L. A. Peletier and W. C. Troy, Spatial patterns, Progress in Nonlinear Differential Equations and their Applications, vol. 45, Birkhäuser Boston, Inc., Boston, MA, 2001. Higher order models in physics and mechanics. MR 1839555
[24] L. A. Pelet'e, R. K. A. M. Van der Vorst, and V. K. Troŭ, Stationary solutions of a fourth-order nonlinear diffusion equation (Russian, with Russian summary), Differentsial'nye Uravneniya 31 (1995), no. 2, 327-337, 367; English transl., Differential Equations 31 (1995), no. 2, 301314. MR1373793
[25] P. H. Rabinowitz, Variational methods for nonlinear eigenvalue problems, Eigenvalues of non-linear problems (Centro Internaz. Mat. Estivo (C.I.M.E.), III Ciclo, Varenna, 1974), Edizioni Cremonese, Rome, 1974, pp. 139-195. MR0464299
[26] P.H. Rabinowitz, Some aspects of critical point theory, MRC Tech. Rep. \#2465, Madison, Wisconsin, 1983.
[27] Andrzej Szulkin, Minimax principles for lower semicontinuous functions and applications to nonlinear boundary value problems (English, with French summary), Ann. Inst. H. Poincaré Anal. Non Linéaire 3 (1986), no. 2, 77-109. MR837231
[28] Stepan Tersian and Julia Chaparova, Periodic and homoclinic solutions of extended FisherKolmogorov equations (English, with English and French summaries), C. R. Acad. Sci. Paris Sér. I Math. 331 (2000), no. 4, 287-292, DOI 10.1016/S0764-4442(00)01629-3. MR1787196

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[^0]:    Received by the editors June 24, 2017.
    2010 Mathematics Subject Classification. Primary 34B15, 34C25, 39A10, 39A23.
    Key words and phrases. Relativistic operator, Fisher-Kolmogorov nonlinearities, difference equations, periodic solution, critical point, Palais-Smale condition, Krasnoselskii genus.

