

GLOBAL SOLVABILITY FOR THE HEAT EQUATION WITH BOUNDARY FLUX GOVERNED BY NONLINEAR MEMORY

BY

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Abstract. We introduce the study of global existence and blowup in finite time for the heat equation with flux at the boundary governed by a nonlinear memory term. Via a simple transformation, the model may be written in a form which has been introduced in previous studies of tumor-induced angiogenesis. The present study is also in the spirit of extending work on models of the heat equation with local, nonlocal, and delay nonlinearities present in the boundary flux. Additionally, we provide a brief summary of related studies regarding heat equation models where memory terms are incorporated within reaction or diffusion.

1. Introduction. We investigate the global solvability and blowup in finite time for the heat equation with flux at the boundary governed by a nonlinear memory law

$$\begin{aligned} u_t &= \Delta u && \text{on } \Omega_T, \\ \nabla u \cdot \mathbf{n} &= u^q \int_0^t u^p(\cdot, s) ds && \text{on } (\partial\Omega)_T, \\ u &= u_0 && \text{on } \overline{\Omega} \times \{0\}. \end{aligned} \tag{1}$$

Here, $p \geq 0$, $q \geq 0$, and $\Omega_T = \Omega \times (0, T)$, where Ω is a bounded domain in \mathbf{R}^N having piecewise smooth boundary $\partial\Omega$ with outward pointing unit normal \mathbf{n} . The initial condition u_0 is a nonnegative, continuous function on $\overline{\Omega}$. Our primary result is that if

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$0 \leq p + q \leq 1$, then every solution of (1) is global. On the other hand, if $p + q > 1$, then all nontrivial, nonnegative solutions blow up in finite time.

There are two main sources of motivation for the present study of (1). First, we conduct our investigation in parallel to previous studies of the semilinear heat equation model with memory,

$$\begin{aligned} u_t &= \Delta u + u^q \int_0^t u^p(\cdot, s) ds && \text{on } \Omega_T, \\ u &= 0 && \text{on } (\partial\Omega)_T, \\ u &= u_0 && \text{on } \overline{\Omega} \times \{0\}, \end{aligned} \tag{2}$$

and the localized boundary flux version of (1),

$$\begin{aligned} u_t &= \Delta u && \text{on } \Omega_T, \\ \nabla u \cdot \mathbf{n} &= u^{q+p} && \text{on } (\partial\Omega)_T, \\ u &= u_0 && \text{on } \overline{\Omega} \times \{0\}. \end{aligned} \tag{3}$$

It is known that all solutions of (2) are global in the case $0 \leq p + q \leq 1$ [20, 29, 35, 37], and the same result is also true for solutions of (3); e.g., see [26]. On the other hand, while all nontrivial solutions of (3) blow up in finite time if $p + q > 1$ [28, 34], the situation is slightly different for the model with memory terms in the reaction. In the case $p + q > 1$ with $q < 1$, all nontrivial solutions of (2) blow up in finite time, while in the case $p + q > 1$, $q \geq 1$, solutions of (2) which blow up in finite time and which are global both exist depending upon the size of u_0 [20, 29].

Our second source of motivation for the study of (1) originates with a model of capillary growth in solid tumors as initiated by angiogenic growth factors [27]. In this reference a model for the transmission of growth factors across a capillary wall is introduced in the following form:

$$\begin{aligned} u_t &= \nabla \cdot (\nabla \phi(x, t, u) + \mathbf{f}(x, t, u)) + h(x, t, u) && \text{on } \Omega_T, \\ (\nabla \phi(x, t, u) + \mathbf{f}(x, t, u)) \cdot \mathbf{n} &= g(x, t, u, v) && \text{on } (\partial\Omega)_T, \\ u &= u_0 && \text{on } \overline{\Omega} \times \{0\} \end{aligned} \tag{4}$$

with $g_v \geq 0$ on $(\partial\Omega)_T$; $g_v \equiv 0$ on $(\partial\Omega \setminus \Sigma)_T$, and

$$\begin{aligned} v_t &= F(x, t, u, v) + G(u)_t && \text{on } \Sigma_T, \\ v &= v_0 && \text{on } \Sigma \times \{0\}. \end{aligned} \tag{5}$$

Here, Σ is a relatively open subset of $\partial\Omega$, which represents the capillary wall.

In (1), if we let

$$v(x, t) \equiv \int_0^t u^p(x, s) ds$$

and $\Sigma = \partial\Omega$, then the model may be seen to be of the type (4)-(5), with $g = u^q v$, $F = u^p$, $G = 0$, and $v_0 = 0$. In previous works, we have studied existence, uniqueness, and blowup in finite time for various forms of (4)-(5) [2, 3]. However, blowup results in these references are only those which may be achieved through comparison with a related localized model. With the present study, we intend to provide a more complete analysis,

albeit for a specialized version of (4)-(5), via identification with a heat equation problem involving memory flux at the boundary (1).

Considering the nonlinear time integral condition governing flux through the boundary, the model (1) involves a continuous time delay, often referred to as an integrodifferential system or a memory condition in the literature. Models involving memory terms in reaction have arisen in studies of nuclear reactor dynamics [24, 32] and population dynamics [25], specifically in the case of logistic growth models involving both nondelayed and hereditary effects [7, 40]. Local/global existence, stabilization, and blowup in finite time of solutions for a variety of generalizations of such models have subsequently been investigated within a number of previous works, e.g., [6, 12, 36, 38, 39].

Memory terms in diffusion have been studied as well, arising in models of viscoelastic forces in non-Newtonian fluids [18, 31] and resulting from a modified Fourier law applied to an anisotropic, nonhomogeneous media [41]. Numerical studies of closely related models and also those involving a form of the Fisher equation with memory may be found in [4, 5, 8]. It is useful to note here that fractional order time derivatives, as memory operators, have been studied in a memory formalism of D'Arcy's law and molecule transport across a biological membrane. See [9, 10] and the references therein. Memory terms in diffusion and reaction have additionally been introduced in the context of climate models [15, 19].

Despite the volume of work done on models incorporating memory in reaction, diffusion, or both, there appear to be very few appearances in the literature of diffusion models in which such terms are present in the boundary flux. In [11], a linear memory term of the type in (1), with $q = 0$, $p = 1$, is introduced as arising in hereditary boundary conditions of Newtonian radiation and calorimetry. Similar hereditary boundary conditions have been employed in models of time-dependent electromagnetic fields at dissipative boundaries [16]. Boundary conditions incorporating convolution in time have been studied more extensively in the case of wave equation models arising in shock problems involving a linear viscoelastic bar; see, e.g., [30, 33] and the references therein. We have been able to locate only one general study of existence of weak solutions for a quasilinear parabolic system in which memory terms are incorporated in diffusion, reaction, and boundary flux conditions [1]. Owing to our need for comparison results for classical solutions, we have developed these fundamental results for (1) herein.

The plan of the paper is as follows. In Section 2 we present the local existence and the comparison principle. In Section 3 we prove that for $p + q \leq 1$, every solution of (1) is global, while for $p + q > 1$, all solutions blow up in finite time. In Section 4 we show that for $p > 1$ or $q > 1$, blowup can occur only on the boundary.

2. Local existence and comparison results. In this section, we establish the existence of classical solutions, along with sub- and super-solution comparisons for (1). While these arguments are fairly standard, we provide them here as a way to introduce the Green's function formulation, to be utilized later, and to present comparison results as applicable even in the non-Lipschitz cases $p, q < 1$. For such purposes, we begin by defining $G_N(x, y, t, \tau)$ as the Green's function for the heat equation with homogeneous Neumann boundary condition. $u(x, t) \in C^{2,1}(\Omega_T) \cap C(\overline{\Omega_T})$ is a classical solution of (1);

i.e., u is C^1 in t , C^2 in x on $\overline{\Omega} \times (0, T)$ and continuous on $\overline{\Omega} \times [0, T]$ if and only if

$$\begin{aligned} u(x, t) = & \int_{\Omega} G_N(x, y, t, 0)u_0(y)dy \\ & + \int_0^t \int_{\partial\Omega} G_N(x, y, t, \tau)u^q(y, \tau) \int_0^\tau u^p(y, s)dsdS_yd\tau \end{aligned} \quad (6)$$

for all $x \in \overline{\Omega}$, $0 \leq t \leq T$.

We prove local existence of solutions of (1) via a fixed point argument. Despite the presence of a memory term, the necessary steps are nearly identical to well-known results for corresponding localized problems [17]. Thus, we do not repeat all of the details here. Although presented in the context of (1), the work below may be applied to more general uniformly parabolic models and nonlinear memory terms. We restrict this discussion to nonnegative solutions; however, such a condition may also be removed.

Given $u \in C(\overline{\Omega} \times [0, T])$, define

$$\begin{aligned} \mathcal{T}[u](x, t) \equiv & \int_{\Omega} G_N(x, y, t, 0)u_0(y)dy \\ & + \int_0^t \int_{\partial\Omega} G_N(x, y, t, \tau)u^q(y, \tau) \int_0^\tau u^p(y, s)dsdS_yd\tau \end{aligned}$$

and

$$M_0 \equiv \sup_{\overline{\Omega}} |u_0|.$$

Continuing in a manner also similar to [13], let

$$\kappa(t) \equiv \sup_{x \in \overline{\Omega}, 0 \leq \tau \leq t} \int_0^\tau \int_{\partial\Omega} G_N(x, y, \tau, \eta)dS_yd\eta.$$

We note that $\kappa(t) \leq 2C_0\sqrt{t}$ for constants $\varepsilon_0, C_0 > 0$ with $t < \varepsilon_0$ [22]. Fix $\hat{T} < \varepsilon_0$ and $M > M_0$ so that

$$M^{q+p}\hat{T}\kappa(\hat{T}) \leq M - M_0.$$

If $|u| \leq M$ and $0 \leq t \leq \hat{T}$, then

$$\begin{aligned} |\mathcal{T}[u]| & \leq M_0 \int_{\Omega} G_N(x, y, t, 0)dy + M^{q+p}\hat{T}\kappa(\hat{T}) \\ & \leq M_0 + (M - M_0). \end{aligned}$$

So, $\mathcal{T} : K \rightarrow K$, where

$$K \equiv \left\{ u \in C(\overline{\Omega} \times [0, \hat{T}]) : \|u\|_\infty \leq M \right\}.$$

The continuity and compactness of \mathcal{T} follow from the same arguments as in [17]. Therefore, \mathcal{T} has a fixed point, which is a solution of (1), and, by standard results, is a classical solution and may be continued unless $\|u\|_\infty$ becomes unbounded.

Uniqueness and comparison results for (1) may be developed using the same techniques as used in [3] for the transmission model (4) – (5). Nevertheless, we herein provide an adaptation of this more general work to the case of classical sub- and super-solutions, towards a more complete comparison theory than is currently known for the case of nonlinear diffusion models addressed in this reference. In such respect, our treatment

here actually bears a great deal more similarity to that of [29] in the use of classical maximum principles and development of the comparison results.

Consider $w(x, t) \in C^{2,1}(\Omega_T) \cap C(\overline{\Omega_T})$ such that

$$\begin{aligned} w_t &\geq \Delta w && \text{on } \Omega_T, \\ \nabla w \cdot \mathbf{n} &\geq aw + b \int_0^t cw(\cdot, s) ds && \text{on } (\partial\Omega)_T, \\ w &\geq 0 && \text{on } \overline{\Omega} \times \{0\} \end{aligned} \quad (7)$$

with bounded continuous functions a and $b, c \geq 0$. For a positive smooth function ξ , which satisfies $\nabla \xi \cdot \mathbf{n} \geq \alpha \xi$ on $\partial\Omega$, let

$$w(x, t) = e^{\lambda t} \xi(x) W(x, t).$$

The positive constants α, λ are chosen to satisfy

$$\alpha > 1 + T \|b\|_\infty \|c\|_\infty$$

and

$$\lambda > \frac{\Delta \xi}{\xi}$$

on $\overline{\Omega}$.

Standard calculations show that

$$W_t \geq \Delta W + \frac{2\nabla \xi}{\xi} \cdot \nabla W + \left(\frac{\Delta \xi}{\xi} - \lambda \right) W$$

in Ω_T . Therefore, W cannot assume a negative minimum anywhere within $\Omega \times (0, T]$. Similarly, if $W(x_0, t_0) = \min W < 0$ for some $x_0 \in \partial\Omega$, at the point (x_0, t_0) it follows that

$$\nabla W \cdot \mathbf{n} \geq -W(x_0, t_0) \left[\nabla \xi \cdot \mathbf{n} - \xi - e^{-\lambda t} \xi b \int_0^t c e^{\lambda s} ds \right] > 0,$$

which is a contradiction. Therefore, $W \geq 0$ on $\overline{\Omega_T}$, and hence $w \geq 0$ also.

Now, for a (classical) supersolution \bar{u} and a subsolution \underline{u} of (1), we have that $w \equiv \bar{u} - \underline{u}$ satisfies (7), where

$$\begin{aligned} [\bar{u}^q(x, t) - \underline{u}^q(x, t)] \int_0^t \bar{u}^p(x, s) ds &\equiv a(x, t) [\bar{u}(x, t) - \underline{u}(x, t)], \\ b(x, t) &\equiv \underline{u}^q(x, t), \end{aligned}$$

and

$$\bar{u}^p(x, t) - \underline{u}^p(x, t) \equiv c(x, t) [\bar{u}(x, t) - \underline{u}(x, t)].$$

Provided $\bar{u} \geq \delta > 0$ and $\underline{u} \geq 0$, a, b, c will be bounded for all cases $p \geq 0, q \geq 0$. It follows that if $\bar{u}(\cdot, 0) \geq \underline{u}(\cdot, 0)$, then $\bar{u} \geq \underline{u}$.

Considering a solution u_δ of (1) with initial condition $u_0 + \delta$, we see that the limit

$$U(x, t) \equiv \lim_{\delta \rightarrow 0^+} u_\delta(x, t)$$

exists and yields a maximal solution of (1). In the case $p, q \geq 1$, this is, in fact, the unique solution of the problem. Furthermore, if \bar{u} is a supersolution and \underline{u} is a subsolution of (1) such that $\bar{u} \geq \delta$ for some positive constant δ and $\underline{u} \geq 0$, then $\bar{u}(\cdot, 0) \geq u_0 \geq \underline{u}(\cdot, 0)$ implies $\bar{u} \geq U \geq \underline{u}$ on $\overline{\Omega_T}$.

3. Global existence and finite time blowup. In this section, we begin with the global existence for solutions of (1), that is, for any $T > 0$, $u(x, t) \leq C(T) < \infty$ in $\overline{\Omega_T}$ with some positive constant $C = C(T)$.

THEOREM 3.1. If $p + q \leq 1$, then every solution of (1) is global.

Proof. We seek a global supersolution \bar{u} of (1). From [14], there exists a function $\varphi(x) \in C^2(\overline{\Omega})$ satisfying

$$0 < \varphi(x) \leq 1 \quad \text{in } \Omega \quad \text{and} \quad \nabla \varphi \cdot \mathbf{n} \geq 1 \quad \text{on } \partial\Omega.$$

Let $m_1 = \max_{\overline{\Omega}} |\nabla \varphi|$ and $m_2 = \max_{\overline{\Omega}} |\Delta \varphi|$. We define

$$\bar{u} = M e^{\lambda t + \varphi},$$

where

$$\begin{aligned} M &= \max \{ \|u_0\|_{L^\infty(\Omega)}, 1 \}, \\ \lambda &= \max \{ m_1^2 + m_2, 1/p \}. \end{aligned}$$

Then \bar{u} satisfies

$$\begin{aligned} \bar{u}_t &\geq \Delta \bar{u} && \text{on } \Omega_T, \\ \nabla \bar{u} \cdot \mathbf{n} &\geq \bar{u}^q \int_0^t \bar{u}^p(x, \tau) d\tau && \text{on } (\partial\Omega)_T, \\ \bar{u}(x, 0) &\geq u_0(x) && \text{on } \overline{\Omega} \times \{0\}. \end{aligned} \tag{8}$$

Hence, \bar{u} is a desired supersolution. \square

We then establish the blowup result. To this end, we first consider the following problem:

$$\begin{aligned} u_t &= \Delta u && \text{on } \Omega_T, \\ \nabla u \cdot \mathbf{n} &= c \int_0^t u^\mu(\cdot, s) ds && \text{on } (\partial\Omega)_T, \\ u &= u_0 && \text{on } \overline{\Omega} \times \{0\}, \end{aligned} \tag{9}$$

where c, μ are positive constants. From now on, without causing any confusion, we may use c_i, C_i or \tilde{c}_i ($i = 0, 1, 2, \dots$) to denote various positive constants.

LEMMA 3.2. For any $c > 0$, if $\mu > 1$, all nonnegative solutions of (9) blow up in finite time.

Proof. The solution u of (9) satisfies

$$\begin{aligned} u(x, t) &= \int_{\Omega} G_N(x, y, t, 0) u_0(y) dy \\ &\quad + c \int_0^t \int_{\partial\Omega} G_N(x, y, t, \tau) \int_0^\tau u^\mu(y, s) ds dS_y d\tau \end{aligned} \tag{10}$$

for all $x \in \overline{\Omega}$, $0 \leq t \leq T$. As in [22], one can show that

$$\int_{\partial\Omega} G_N(x, y, t, \tau) dS_x \geq c_0 > 0 \quad \text{for } y \in \overline{\Omega}, T > t > \tau \geq 0. \tag{11}$$

By (10), (11) and Jensen's inequality, we have

$$\begin{aligned} \int_{\partial\Omega} u(x, t) dS_x &\geq c \int_{\partial\Omega} \left(\int_0^t \int_{\partial\Omega} G_N(x, y, t, \tau) \int_0^\tau u^\mu(y, s) ds dS_y d\tau \right) dS_x \\ &\geq cc_0 \int_0^t \int_0^\tau \int_{\partial\Omega} u^\mu(y, s) dS_y ds d\tau \\ &\geq c_1 \int_0^t \tau^{1-\mu} \left(\int_0^\tau \int_{\partial\Omega} u(y, s) dS_y ds \right)^\mu d\tau. \end{aligned} \quad (12)$$

On the other hand, by (10) and (11),

$$\begin{aligned} \int_{\partial\Omega} u(x, t) dS_x &\geq \int_{\partial\Omega} \int_{\Omega} G_N(x, y, t, 0) u_0(y) dy dS_x \\ &= \int_{\Omega} u_0(y) \int_{\partial\Omega} G_N(x, y, t, 0) dS_x dy \\ &\geq c_0 \int_{\Omega} u_0(y) dy \\ &\geq c_2 > 0. \end{aligned} \quad (13)$$

Set

$$F(t) = \int_0^t \int_{\partial\Omega} u(x, \tau) dS_x d\tau \quad \text{for } t > 0.$$

From (12) and (13), it follows that $F(t)$ satisfies

$$F'(t) \geq c_4 + c_3 \int_0^t \tau^{1-\mu} F^\mu(\tau) d\tau \quad \text{for } t > 0. \quad (14)$$

Integration of the above inequality over $(0, t)$ then yields

$$\begin{aligned} F(t) &\geq c_4 t + c_3 \int_0^t \int_0^\tau \zeta^{1-\mu} F^\mu(\zeta) d\zeta d\tau \\ &\geq c_4 t + c_3 \int_0^t (t - \zeta) \zeta^{1-\mu} F^\mu(\zeta) d\zeta \\ &\geq c_4 t + c_3 t^{1-\mu} \int_0^t (t - \zeta) F^\mu(\zeta) d\zeta \end{aligned} \quad (15)$$

for $t > 0$. Assume to the contrary that (9) has a global solution u . Then for any positive number T , we have

$$F(t) \geq c_4 T + c_5 T^{1-\mu} \int_T^t (t - \zeta) F^\mu(\zeta) d\zeta \quad \text{for } T \leq t \leq 2T.$$

Thus, by comparison, $F(t) \geq H(t)$ on $[T, 2T]$, where

$$H(t) = c_4 T + c_5 T^{1-\mu} \int_T^t (t - \zeta) H^\mu(\zeta) d\zeta \quad \text{for } T \leq t \leq 2T.$$

Clearly, $H(t)$ satisfies

$$\begin{aligned} H''(t) &= c_5 T^{1-\mu} H^\mu(t), \quad T < t < 2T, \\ H(T) &= c_4 T, \\ H'(T) &= 0. \end{aligned} \quad (16)$$

Multiplying the equation in (16) by $H'(t)$ and integrating from T to t , we obtain

$$H'(t) = c_6 T^{(1-\mu)/2} (H^{\mu+1}(t) - H^{\mu+1}(T))^{1/2}.$$

Integration of this relation over $(T, 2T)$ then leads to

$$\begin{aligned} c_6 T^{(3-\mu)/2} &= \int_{H(T)}^{H(2T)} (z^{\mu+1} - H^{\mu+1}(T))^{-1/2} dz \\ &\leq (\mu + 1)^{-1/2} H^{-\mu/2}(T) \int_{H(T)}^{2H(T)} (z - H(T))^{-1/2} dz \\ &\quad + 2^{(\mu+1)/2} \int_{2H(T)}^{\infty} z^{-(\mu+1)/2} dz \\ &= 2[(\mu + 1)^{-1/2} + 2(\mu - 1)^{-1}] c_4^{(1-\mu)/2} T^{(1-\mu)/2}, \end{aligned}$$

which is equivalent to

$$T \leq 2[(\mu + 1)^{-1/2} + 2(\mu - 1)^{-1}] c_4^{(1-\mu)/2} / c_6. \tag{17}$$

For sufficiently large T , inequality (17) yields a contradiction, which completes the proof. \square

We now consider the general case that $p + q > 1$.

THEOREM 3.3. If $p + q > 1$, then all nonnegative solutions of (1) blow up in finite time.

Proof. We consider two cases.

Case 1 ($q < 1$). By the maximum principle, $u(x, t) > 0$ for $x \in \bar{\Omega}$ and $t > 0$. Let

$$z = Mu^{1-q},$$

where $M = \|u_0\|_{L^\infty(\Omega)}^q$. Then z satisfies

$$\begin{aligned} z_t &\geq \Delta z && \text{on } \Omega_T, \\ \nabla z \cdot \mathbf{n} &= M^{\frac{1-p-q}{1-q}} (1-q) \int_0^t z^{\frac{p}{1-q}}(x, \tau) d\tau && \text{on } (\partial\Omega)_T, \\ z(x, 0) &\geq u_0(x) && \text{on } \bar{\Omega} \times \{0\}. \end{aligned} \tag{18}$$

z is a supersolution of problem (9). Since $p/(1-q) > 1$, by Lemma 3.2, z blows up in finite time, and so does u .

Case 2 ($q \geq 1$). Because the solution of the heat equation with homogenous Neumann boundary condition is a subsolution of (1), for a small $\sigma > 0$, there exists a constant $c_\sigma > 0$ such that the solution u of (1) satisfies

$$u(x, t) \geq c_\sigma \quad \text{for } x \in \partial\Omega, t \geq \sigma > 0. \tag{19}$$

Let $q = \alpha + \gamma$, where $\gamma < 1$ and $\gamma + p > 1$. Then we have

$$u^q = u^{\alpha+\gamma} \geq c_\sigma^\alpha u^\gamma \quad \text{for } x \in \partial\Omega, t \geq \sigma > 0.$$

Consider the following problem:

$$\begin{aligned} \underline{u}_t &= \Delta \underline{u} && \text{on } \Omega \times (\sigma, T), \\ \nabla \underline{u} \cdot \mathbf{n} &= c_\sigma^\alpha \underline{u}^\gamma \int_\sigma^t \underline{u}^p(x, s) ds && \text{on } \partial\Omega \times (\sigma, T), \\ \underline{u}(x, \sigma) &\leq u(x, \sigma) && \text{on } \bar{\Omega}. \end{aligned} \tag{20}$$

Proceeding analogously as in the proofs of Lemma 3.2 and case 1 of Theorem 3.3, we can show that \underline{u} blows up in finite time. Since u is a supersolution of (20), u cannot exist globally. \square

4. Blowup on the boundary. In this section, we show that for problem (1) in the case $p > 1$ or $q > 1$, blowup cannot occur at the interior domain. For definiteness, we may assume that T is the blowup time.

THEOREM 4.1. If $p > 1, q = 0$ or $p \geq 0, q > 1$, then blowup can occur only on the boundary.

Proof. We consider two cases.

Case 1 ($p > 1, q = 0$). Following the general idea of [21], we set

$$J(t) = \int_0^t \int_0^\tau \int_{\partial\Omega} u^p(y, s) dS_y ds d\tau.$$

By (6), (11) and Jensen’s inequality, we find

$$J''(t) \geq c_0^p |\partial\Omega|^{1-p} J^p(t).$$

Multiplying both sides of the above inequality by $J'(t)$ and integrating over $(0, t)$, we obtain

$$J'(t) \geq c_7 J^{\frac{p+1}{2}}(t),$$

which, upon integration over (t, T) , yields

$$\int_{J(t)}^\infty s^{-\frac{p+1}{2}} ds \geq c_8(T - t),$$

or equivalently,

$$J(t) \leq c_9(T - t)^{-\frac{2}{p-1}} \quad \text{for } t \in [0, T]. \tag{21}$$

We now take an arbitrary $\Omega' \subset\subset \Omega$ with $\text{dist}(\partial\Omega, \Omega') = \varepsilon > 0$. For this Ω' , we then take $\Omega'' \subset\subset \Omega$ such that $\Omega' \subset\subset \Omega''$, $\text{dist}(\partial\Omega'', \Omega') \geq \varepsilon/3$ and $\text{dist}(\partial\Omega, \Omega'') \geq \varepsilon/3$. It is well known that for any $\varepsilon > 0$,

$$0 \leq G_N(x, y, t, \tau) \leq C_\varepsilon \quad \text{for } |x - y| \geq \frac{\varepsilon}{3}, x, y \in \bar{\Omega}', 0 < \tau < t < T, \tag{22}$$

where C_ε is a positive constant depending on ε . Then by (6), (21) and (22),

$$\max_{\bar{\Omega}'} u(x, t) \leq C_0 + C_\varepsilon J(t) \leq C_1(T - t)^{-\frac{2}{p-1}}.$$

Proceeding similarly as in the proof of Theorem 4.1 of [23], we have

$$u(x, t) \leq \frac{C_3}{[\psi(x) + (C_2 + 1)(T - t)]^{\frac{2}{p-1}}} \quad \text{in } \bar{\Omega}' \times [0, T], \tag{23}$$

where $\psi(x) \in C^2(\overline{\Omega}')$ satisfies

$$\begin{aligned} \psi(x) &> 0 \quad \text{in } \Omega', \quad \psi(x) = 0 \quad \text{on } \partial\Omega', \\ \Delta\psi - \frac{p+1}{p-1} \frac{|\nabla\psi|^2}{\psi} &\geq -C_2 \quad \text{in } \Omega' \end{aligned}$$

for $C_2 > 0$. Inequality (23) shows that $u(x, t)$ cannot blow up in $\Omega' \times (0, T)$. Case 2 ($p \geq 0, q > 1$). We now set

$$\tilde{J}(t) = \int_{\sigma}^t \int_{\partial\Omega} u^q(y, \tau) \int_{\sigma}^{\tau} u^p(y, s) ds dS_y d\tau,$$

where $0 < \sigma < T/4$. In view of Case 1, it suffices to prove that $\tilde{J}(t)$ satisfies a similar estimate as (21). More precisely, if we can show that

$$J(t) \leq \tilde{c}(T-t)^{-\beta} \quad \text{for } T/2 \leq t < T, \tag{24}$$

where $\tilde{c}, \beta > 0$, then $u(x, t)$ must blow up on the boundary. By (19), we have that for $y \in \partial\Omega$ and $t \in [T/2, T)$,

$$\int_{\sigma}^t u^p(y, s) ds \geq c_{\sigma}^p(t - \sigma) \geq c_{\sigma}^p T/4.$$

Let $\tilde{c}_0 = c_{\sigma}^p T/4$. Then

$$\int_{\sigma}^t u^p(y, s) ds \geq \tilde{c}_0. \tag{25}$$

On the other hand, by (11) we find

$$\begin{aligned} \int_{\partial\Omega} u(x, t) dS_x &\geq \int_{\partial\Omega} \left(\int_0^t \int_{\partial\Omega} G_N(x, y, t, \tau) u^q(y, \tau) \int_0^{\tau} u^p(y, s) ds dS_y d\tau \right) dS_x \\ &\geq c_0 \int_0^t \int_{\partial\Omega} u^q(y, \tau) \int_0^{\tau} u^p(y, s) ds dS_y d\tau \\ &> c_0 \int_{\sigma}^t \int_{\partial\Omega} u^q(y, \tau) \int_{\sigma}^{\tau} u^p(y, s) ds dS_y d\tau \\ &= c_0 \tilde{J}(t). \end{aligned} \tag{26}$$

Combining (25), (26) and applying Jensen's inequality, we further find that

$$\begin{aligned} \tilde{J}'(t) &= \int_{\partial\Omega} u^q(y, t) \int_{\sigma}^t u^p(y, s) ds dS_y \\ &\geq \tilde{c}_0 \int_{\partial\Omega} u^q(y, t) dS_y \\ &\geq \tilde{c}_1 \left(\int_{\partial\Omega} u(y, t) dS_y \right)^q \\ &\geq \tilde{c}_2 \tilde{J}^q(t). \end{aligned}$$

Integration of the above inequality over (t, T) yields

$$\int_{\tilde{J}(t)}^{\infty} s^{-q} ds > \tilde{c}_2(T-t),$$

which gives

$$\tilde{J}(t) \leq \tilde{c}_3(T-t)^{-\frac{1}{q-1}} \quad \text{for } t \in [T/2, T).$$

Hence, the proof is completed. \square

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