THE LITTLEWOOD-PALEY THEORY FOR
JACOBI EXPANSIONS

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Abstract. The machinery for harmonic analysis utilizing Jacobi polynomial
expansions is developed using the explicit form of the convolution kernel
discovered by Gasper. Various maximal functions, and the standard Lit-
tlewood-Paley functionals are studied and an application is given to
multiplier theorems.

1. Introduction. The Littlewood-Paley theory involves certain statements
about the $L^p$ norm of various functionals constructed out of derivatives of
Abel means and differences of Cesàro sums. The theory was originally used
to study some of the subtler properties of trigonometric series, but later it was
extended to the Fourier transform in $\mathbb{R}^n$ and then to expansions related to
various Lie groups, symmetric spaces, and some weighted measure spaces (see
for example [19], [18], and [8]).

In this paper the theory is extended to Jacobi series. The major tool for this
extension is the convolution kernel studied by Gasper in [9] and [10]. The
major part of what follows is the development of the interplay between the
convolution kernel and the Poisson kernel or the Cesàro kernel, or the role of
the convolution in various maximal theorems. The paper is divided as
follows: In §2 the convolution structure is introduced and the various
operators $g_\alpha$, $g^*_\alpha$, $S$, $\gamma_\alpha$, and $\gamma^*_\alpha$ are defined. Theorem 1 describes the range of
validity of the standard inequalities. The proof of this theorem uses a number
of other results which are established in §§2 and 3. In particular Theorems 2
and 3 are results about Cesàro summability. Theorems 4, 5, and 6 study the
interplay of the convolution structure and the various maximal functions.

In §4 the $g_\alpha$, $g^*_\alpha$, and $S$ functionals are used to prove a multiplier theorem
for the Banach space $S(2, \alpha + 1)$ acting on Jacobi series (see §4 for defini-
tions). This theorem of Hörmander type extends and improves the earlier
results in [7] and [8] and is then used as an end point result along with the
results of Askey [1] and Gasper and Trebels [12] to obtain in §4 a general
multiplier theorem which describes a set of parameters $(1/p, 1/q, \alpha)$ for
which $S(q, \alpha)$ is an $L^p$ multiplier for Jacobi expansions.

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2. The Littlewood-Paley theory for Jacobi expansions. In all that follows, assume \( \alpha > \beta > -1 \). Define \( L^p \) as the collection of all \( f \) for which

\[
\|f\|_p = \left[ \int_0^{\pi/2} |f(\theta)|^p \, d\mu \right]^{1/p} < \infty
\]

where \( d\mu(\theta) = 2^{\alpha+\beta+2} \sin^{\alpha+1} \theta \cos^{\beta+1} \theta \, d\theta \). The Jacobi polynomials are the polynomials with argument \( \cos 2\theta \), \( P_n^{(\alpha,\beta)}(\cos 2\theta) \) orthogonal with respect to \( d\mu \), normalized so that \( P_n^{(\alpha,\beta)}(1) = (\alpha+\beta)_n \). Define

\[
R_n(\theta) = \frac{P_n^{(\alpha,\beta)}(\cos 2\theta)}{P_n^{(\alpha,\beta)}(1)}
\]

so that (cf. Szegö [20, 7.32.2])

\[-1 < R_n(\theta) < R_n(0) = 1 \quad (\alpha > -\frac{1}{2}).\]

The Jacobi series of a function \( f \) in \( L^1 \) is given by

\[
f(\theta) \sim \sum_{n=0}^{\infty} \hat{f}(n) h_n R_n(\theta)
\]

where

\[
\hat{f}(n) = \int_0^{\pi/2} f(\theta) R_n(\theta) \, d\mu(\theta)
\]

and

\[
h_n = \left[ \int_0^{\pi/2} R_n^2(\theta) \, d\mu(\theta) \right]^{-1} \approx C n^{2\alpha+1}.
\]

The convolution structure is based on the work of Gasper [10]. A necessary and sufficient condition that there exist a unique measure \( m_{\theta,\phi}(\psi) \) such that

(i) \( R_n(\theta) R_n(\phi) = \int_0^{\pi/2} R_n(\psi) \, dm_{\theta,\phi}(\psi) \),

(ii) \( \|dm_{\theta,\phi}(\psi)\| < K \) (uniformly in \( \theta, \phi \))

is that \( \alpha + \beta > -1 \). The measure is nonnegative if and only if \( \beta > -\frac{1}{2} \) or \( \alpha + \beta > 0 \). An important property of this measure is

(iii) \( \text{supp } m_{\theta,\phi}(\psi) \subset [\theta - \phi, \theta + \phi] \).

Using this measure it is possible to define a convolution

\[
f \ast g(\theta) = \int_0^{\pi/2} f(\phi) g(\psi) \, dm_{\theta,\phi}(\psi) \, d\mu(\phi)
\]

and it follows that, if \( f, g \in L^1 \),

\[
(f \ast g)(n) = \hat{f}(n) \hat{g}(n)
\]

and if \( 1 < r, p, q < \infty \) with \( 1/r = 1/p + 1/q - 1 \). Then

\[
\|f \ast g\|_r \leq K \|f\|_p \|g\|_q.
\]

(2.1)

If \( \delta_\phi \) is the unit mass concentrated at \( \phi \), then a generalized translation is
defined by
\[ f \ast \delta_\phi(\theta) = \int f(\psi) \, dm_{\psi, \phi}(\psi). \]

The Littlewood-Paley theory depends on manipulating certain derivatives or differences of “mollified” functions. The \( g \)-functions are defined in terms of Abel means and their derivatives. The function
\[ P_r(\theta) = \sum r^n h_n R_n(\theta) \]
is called the untranslated Poisson kernel, the function
\[ P_r(\theta, \phi) = \sum r^n h_n R_n(\theta) R_n(\phi) \]
is called the translated Poisson kernel. It is clear that \( P_r(\theta, 0) = P_r(\theta) \), and, if \( \alpha + \beta > -1 \), \( P_r(\theta, \phi) = P_r(\theta) \ast \delta_\phi(\theta) \). The Abel mean of \( f \) is given by
\[ f(r, \theta) = \frac{1}{\Gamma(\beta)} \int_0^r f(\phi) P_r(\theta, \phi) \, d\mu(\phi) = \sum r^n \hat{f}(n) h_n R_n(\theta). \]
Again if \( \alpha + \beta > -1 \), \( f(r, \theta) = f \ast P_r(\theta) \). The radial derivative is defined by
\[ d_\delta f(r, \theta) = \left( r \frac{\partial}{\partial r} \right)^\delta f(r, \theta) \]
if \( \delta \) is an integer, and if \( \delta = k - \beta, 0 < \beta < 1 \),
\[ d_\delta f(r, \theta) = \frac{1}{\Gamma(\beta)} \int_0^r d_\delta f(rt, \theta) |\log r|^{\beta - 1} \frac{dt}{t}. \]
A direct calculation then shows that
\[ d_\delta f(r, \theta) = \sum n^\delta r^n \hat{f}(n) h_n R_n(\theta) \]
and \( d_\delta d_n f = d_{\delta + n} f \).
For all \( \alpha, \beta > -1 \), define
\[ g_\delta(f, \theta) = \left\{ \int_0^1 \left| d_\delta f(r, \theta) \right|^2 |\log r|^{2\alpha - 1} \frac{dr}{r} \right\}^{1/2} \]
and for \( \alpha + \beta > -1 \), define
\[ S(f, \theta) = \left\{ \int_0^1 \left| d_1 f(r, \cdot) \right|^2 \chi_{1-r}(\theta)(1 - r)^{-2\alpha - 1} \, dr \right\}^{1/2}, \]
\[ g_\eta^*(f, \theta) = \left\{ \int_0^1 \left| d_1 f(r, \cdot) \right|^2 \kappa_{\eta, r}(\theta)(1 - r)^{-2\alpha - 1} \, dr \right\}^{1/2}, \]
where \( \chi_{1-r} \) is the characteristic function of the interval \([0, 1 - r]\) and
\[ k_{\eta, r}(\theta) = \left( \frac{1 - r}{1 - r + \theta} \right)^{\eta(2\alpha + 2)}. \]
The \( \gamma \)-functions of Bonami and Clerc [4] are defined in terms of the Cesàro sums:

\[
S^\delta f = \frac{1}{A_L^\delta} \sum_{l=0}^{L} A_{L-l}^\delta h_l f(l) R_l,
\]

\[
A_L^\delta = \frac{\Gamma(L + \delta + 1)}{\Gamma(L + 1)\Gamma(\delta + 1)}.
\]

These Cesàro sums can be represented by integration against a kernel. Define

\[
K_L^\delta (\theta, \phi) = \frac{1}{A_L^\delta} \sum_{l=0}^{L} A_{L-l}^\delta h_l R_l(\theta) R_l(\phi),
\]

and it follows that

\[
S^\delta f(\theta) = \int f(\phi) K^\delta_L (\theta, \phi) \, d\mu(\phi).
\]

If \( \alpha + \beta > -1 \), define

\[
K^\delta_L (\theta, \phi) = K^\delta_L (\theta, 0); \quad S^\delta_L f(\theta) = f \star K^\delta_L (0).
\]

The \( \gamma \)-functions are defined in terms of certain differences of these sums:

\[
\gamma_\delta(f, \theta) = \left\{ \sum_{L=0}^{\infty} \left| S_L^{\delta+1} f(\theta) - S_L^\delta f(\theta) \right|^2 V_L^{-1} \right\}^{1/2},
\]

\[
\gamma^\delta(f, \theta) = \left\{ \sum_{L=0}^{\infty} \left| S_L^{\delta+1} f(\theta) - S_L^\delta f(\theta) \right|^2 V_L^{-1} \right\}^{1/2}
\]

where it is assumed that the sequence \( V_L \) used in defining \( \gamma^\delta \) always satisfies:

(a) \( V_L > 0 \), (b) \( \sup_{L>1} V_L^{-1} \sum_{k=1}^{L-1} V_K = M < \infty \). (N.B. Bonami and Clerc call these functions \( g_\delta \) and \( g^\delta \), respectively. The older notation using \( \gamma \)'s is from Zygmund [22, II, p. 222].)

The Littlewood-Paley theory for the Jacobi expansions then consists of the following statements:

**Theorem 1.** If \( f \in L^2 \cap L^p \) and \( \int f \, d\mu = 0 \), then

(a) \( \| g_\delta(f, \cdot) \|_p < A_p \| f \|_p \), \( 1 < p < \infty \), \( \alpha, \beta > -1 \);

(b) \( \| f \|_p < A_p \| g_\delta(f, \cdot) \|_p \), \( 1 < p < \infty \), \( \alpha, \beta > -1 \);

(c) \( \| g^\delta(f, \cdot) \|_p < A_p \| g_\delta(f, \cdot) \|_p \), \( 2 < p < \infty \), \( \eta > 1 \), \( \alpha + \beta > 1 \);

(d) \( g_\delta(f, \theta) < A_{\delta, \eta} g_\delta(f, \theta) \), \( \eta > \delta \), \( \alpha + \beta > -1 \);

(e) \( S(f, \theta) < A_{\delta} g^\delta_\theta(f, \theta) \), \( \eta > 0 \), \( \alpha + \beta > -1 \);

(f) \( g_1(f, \theta) < C g_\delta(f, \theta) \), \( \delta > 0 \), \( \alpha + \beta > 0 \) or \( \beta > -\frac{1}{2} \).
(g) if $\alpha + \beta > 0$ or $\beta > -\frac{1}{2}$, and if $\delta > 0$ and $p$ satisfy $|1/p - \frac{1}{2}| < (2\delta + 1)/(4\alpha + 1)$, then $\|\gamma_{\delta}(f, \cdot)\|_p < C\|f\|_p$;

(h) under the assumptions as in (g), $\|\gamma_{\delta}^*(f, \cdot)\|_p < \|g_1(f, \cdot)\|_p$.

The relations (a) and (b) are demonstrated by defining the Poisson semigroup

$$T'f(\theta) = \int_0^{\pi/2} f(\phi)P_{e^{-r}}(\theta, \phi)\,d\mu(\phi)$$

and verifying that $T'$ satisfies the four axioms for a symmetric diffusion semigroup given by Stein [18] as:

I. Contraction: $\|T'f\|_p < \|f\|_p$, $1 < p < \infty$;

II. Symmetry: each $T'$ is self-adjoint on $L^2$;

III. Positivity: $T'f > 0$ if $f > 0$;

IV. Conservation: $T'1 = 1$.

Then (a) and (b) follow from [18, pp. 111, 120].

Property III follows from the fact that $P_r(\theta, \phi)$ is always positive, $\alpha + \beta > -2$, even if the convolution measure is not positive. This can be seen from the formula

$$P_r(\theta, \phi) = \frac{\Gamma(\alpha + \beta + 2)(1 - r)}{2^{\alpha + \beta + 2}\Gamma(\alpha + 1)\Gamma(\beta + 1)(1 + r)^{\alpha + \beta + 2}}$$

$$\cdot F_4\left(\frac{1}{2}(\alpha + \beta + 2), \frac{1}{2}(\alpha + \beta + 3); \alpha + 1, \beta + 1; \frac{a^2}{k^2}, \frac{b^2}{k^2}\right),$$

(2.2)

where $a = \sin \theta \sin \phi$, $b = \cos \theta \cos \phi$, and $k = \frac{1}{2}(r^{-1/2} + r^{1/2})$ (Bailey [2, p. 102]). The definition of the Poisson kernel implies $\int P_r(\theta, \phi)\,d\mu(\theta) = r^0 R_0^2 R_0(\phi) = 1$ which is IV.

Property I follows from the inequality

$$\|P_r f(\cdot)\|_p < \|f\|_p\|P_r(\cdot, \cdot)\|_1 < \|f\|_p,$$

by an adaption of the argument for Young’s inequality [16, p. 7], and II follows by a simple change in the order of integration.

The relation (c) is established in §3 and it holds in the region where the convolution is bounded on $L^1$, in particular for $\alpha + \beta > -1$.

Relations (d) and (e) were proved in [8, p. 66] for the ultraspherical setting and that proof carries over directly to all expansions $\alpha + \beta > -1$.

Relations (f), (g), and (h) are given in [4] for $\alpha = \beta$ and for a few geometric cases. Gasper’s work on the Jacobi convolution can be used to extend these inequalities to the $(\alpha, \beta)$ region for which the convolution kernel is positive.

Bonami and Clerc [4] have shown that (g) and (h) follow from Zygmund’s
inequality (see Theorem 3 infra), thus it suffices to establish that inequality.

Szegö showed [20, p. 258] that for all $\alpha, \beta > -1, \delta > \alpha + \frac{1}{2}$,
\[
\|K_L^\delta\|_1 < B \quad \text{(uniformly in } L). \tag{2.3}
\]
Therefore, for $\alpha + \beta > -1$, relation (2.1) yields
\[
\|S_L^\delta f\|_p = \|f \ast K_L^\delta\|_p < K\|K_L^\delta\|_1\|f\|_p < KB\|f\|_p.
\]
The constant $KB$ is independent of $L$, and when such an inequality holds $L^p$ is said to be $(C, \delta)$ summable. This is the full range result. When $\delta = 0$, Pollard [17] showed that $L^p$ is $(C, 0)$ summable if
\[
\left| \frac{1}{p} - \frac{1}{2} \right| < \frac{1}{4(\alpha + 1)} < \frac{1}{4(\beta + 1)}.
\]
Therefore by the interpolation of analytic families of operators [19, p. 239], this proves

**Theorem 2.** If $\alpha + \beta > -1$, then $L^p$ is $(C, \delta)$ summable for all $p$ such that
\[
\left| \frac{1}{p} - \frac{1}{2} \right| < \frac{2\delta + 1}{4(\alpha + 1)} \quad (\delta > 0).
\]

That is,
\[
\|S_L^\delta f\|_p < C\|f\|_p
\]
where the constant depends on $\alpha, p, \delta$ but not $L$.

The proof of the above theorem merely depends on the existence of a bounded convolution. In order to prove Zygmund's inequality, it is necessary to restrict $\alpha$ and $\beta$ so that the convolution measure becomes positive. In particular,

**Theorem 3.** For all sequences of functions $\{f_k\}$ and all sequences of integers $\{n_k\}$,
\[
\left\| \left( \sum_0^\infty |S_L^\delta f_k|^2 \right)^{1/2} \right\|_p < A_p \left( \left( \sum_0^\infty |f_k|^2 \right)^{1/2} \right)_p
\]
if $\alpha + \beta > 0 \text{ or } \beta > -\frac{1}{2}$, and $\delta$ and $p$ satisfy
\[
\left| \frac{1}{p} - \frac{1}{2} \right| < \frac{2\delta + 1}{4(\alpha + 1)} \quad (\delta > 0).
\]
The proof is adapted from Bonami and Clerc [4]. A full range result is established for $\delta > \alpha + \frac{1}{2}$ by showing that the maximal operator
\[
\sup_L S_L^\delta f(x) = S_L^\delta f(x)
\]
is a bounded $L^p$ operator $1 < p < \infty$ (this is proved in §3, Theorem 7). The rest of the proof is the same as in [4].
If $0 > \delta > \alpha + \frac{1}{2}$, then $L^2$ is not $(C, \delta)$ summable even though (2.3) still holds. To see this observe that $S^2 = \frac{A^2}{A}$ is given by a sequence of multipliers $m_L^* = \frac{A_{L-1}}{A_L}$, and $m_L^* \sim L^{-\delta}$ which is unbounded.

If $-\frac{1}{2} < \alpha, \alpha + \beta < -1$, then this objection does not apply, and it is conceivable that, as in the Poisson case, $\int |K_\phi^L(\theta, \phi)| \, d\mu(\theta) \, d\mu(\phi) < K$ uniformly in $L$, even though the convolution measure is unbounded (for this see Gasper [10, pp. 263, 278 (6.13)]). In any event, the range of validity for (g) and (h) is not as extensive as the range for (a) and (b).

3. Maximal functions and the harmonic structure of Jacobi expansions. If $f(\theta) > 0$ is periodic on $[-\pi, \pi]$, then the maximal function of Hardy-Littlewood is defined by

$$f^*(\phi) = \sup_{\varepsilon > 0} \frac{1}{2\varepsilon} \int_{\phi - \varepsilon}^{\phi + \varepsilon} f(\theta) \, d\theta.$$  \hspace{1cm} (3.1)

Define the function $h_\varepsilon(\theta) = (2\varepsilon)^{-1} \chi_{[-\varepsilon, \varepsilon]}(\theta)$, where $\chi_{[-\varepsilon, \varepsilon]}$ is the characteristic function of $[-\varepsilon, \varepsilon]$, and extend it periodically, then (3.1) can be rewritten as a convolution:

$$f^*(\phi) = \sup_{\varepsilon > 0} h_\varepsilon \ast f(\phi) = \sup_{\varepsilon > 0} \int_{\phi - \varepsilon}^{\phi + \varepsilon} h_\varepsilon(\phi - \theta) f(\theta) \, d\theta.$$  \hspace{1cm} (3.2)

If $f(r, \theta)$ is the Abel mean of $f(\theta)$, then a second maximal function can be defined by

$$P_\phi f(\phi) = \sup_{0 < r < 1} f(r, \phi).$$ \hspace{1cm} (3.3)

It is possible to show that there exist positive constants $A$ and $B$ such that

$$Af^*(\phi) < P_\phi f(\phi) < Bf^*(\phi)$$

(see for example [19, p. 59]).

The generalization of these ideas to the Jacobi setting leads to three different maximal functions, two of which are equivalent when the convolution is nonnegative. The first maximal function is

$$Mf(\phi) = \sup_{\varepsilon > 0} \frac{1}{\mu[\phi - \varepsilon, \phi + \varepsilon]} \int_{\phi - \varepsilon}^{\phi + \varepsilon} |f(\theta)| \, d\mu(\theta);$$ \hspace{1cm} (3.4)

the second is defined where there is a bounded convolution, i.e. when $\alpha + \beta > -1$:

$$\mathfrak{M}f(\phi) = \sup_{\varepsilon > 0} \int h_\varepsilon(\psi) |f(\theta)| \, dm_{\phi, \theta}(\psi) \, d\mu(\theta)$$ \hspace{1cm} (3.5)

where now $h_\varepsilon(\theta) = \chi_{[0, \varepsilon]}(\theta)/\mu[0, \varepsilon]$. If $\alpha + \beta > 0$ or $\beta > -\frac{1}{2}$, $dm_{\phi, \theta}(\psi)$ is positive and this integral can be written $h_\varepsilon \ast |f(\phi)|$, and in any case $h_\varepsilon \ast |f(\phi)|$ is bounded by the double integral. The third maximal function is analogous to (3.3):
\[ P_{\phi_0}(\phi) = \sup_{0 < r < 1} |f(r, \phi)| = \sup_{0 < r < 1} \left| \int_0^{\pi/2} f(\theta) \, P_r(\theta, \phi) \, d\mu(\theta) \right|. \quad (3.6) \]

\( \mathfrak{M}_f \) and \( M_f \) are not obviously related since \( \mu[\phi - \epsilon, \phi + \epsilon]/\mu[0, \epsilon] \) is not a bounded ratio; however, the \( L^p \) inequalities for \( M_f \) will follow from

**Theorem 4.** If \( \alpha + \beta > -1 \), then there exists \( K > 0 \) such that

\[ \mathfrak{M}_f(\phi) < KM_f(\phi). \]

**Proof.** The special case \( \alpha = \beta \) is equivalent to the lemma in [6, p. 63], thus it will be assumed below that \( \alpha > \beta \). The proof in the Jacobi setting requires even more knowledge of special functions to carry through the details, in particular a formula due to Macdonald [15] for the representation of a triple integral of Bessel functions as an associated Legendre function (this is also in Watson [21, p. 412]; see Gasper [10, p. 268] and Koornwinder [14, p. 135] for more modern presentations of this formula and its application to Jacobi expansions).

Let \( f > 0 \); define

\[ \mathfrak{M}_\epsilon f(\phi) = \frac{1}{\mu[0, \epsilon]} \int_{\phi - \epsilon}^{\phi + \epsilon} |f(\theta)| \, d\mu(\phi) \]

and

\[ M_\epsilon f(\phi) = \frac{1}{\mu[\phi - \epsilon, \phi + \epsilon]} \int_{\phi - \epsilon}^{\phi + \epsilon} f(\theta) \, d\mu(\theta). \]

Thus it suffices to show that \( \mathfrak{M}_\epsilon f(\phi) < KM_\epsilon f(\phi) \), with the constant independent of \( \epsilon \). This will follow immediately from

\[ \mathcal{J} = \mu(\phi - \epsilon, \phi + \epsilon) \int_{\phi - \epsilon}^{\phi + \epsilon} |d\mu_{\phi, \epsilon}(\psi)| < K\mu[0, \epsilon] \simeq Ce^{2\alpha + 2}. \quad (3.7) \]

There are a number of cases to consider:

(i) \( \alpha + \beta > -1, \quad \alpha < -\frac{1}{2} \),

(ii) \( \alpha + \beta > -1, \quad \alpha = -\frac{1}{2} \),

(iii) \( \alpha + \beta > -1, \quad \alpha > -\frac{1}{2} \),

(iv) \( \alpha + \beta = -1 \) (the singular case).

In each case, the formula for \( d\mu_{\phi, \epsilon}(\psi) \) is given explicitly by Gasper [10, formulas (1.12), (3.3), (3.4), (6.4)] and is different. In each case the evaluation of the integral in (3.7) involves considering typical values of the variable: \( \phi < 3\epsilon, \quad 3\epsilon < \phi < \frac{\pi}{2} - 3\epsilon, \quad \phi > \frac{\pi}{2} - 3\epsilon \), and for each of these ranges of the variables, the argument is different. Since there are many cases, and so many of the computations are similar, only the computations needed to verify case (ii) will be given in detail here—it was chosen because of a logarithmic singularity which is not present in the other cases.
Let \( \alpha = \frac{1}{2}, \ a = \sin \phi \sin \theta, \ b = \cos \phi \cos \theta, \ c = \cos \psi, \ B = (b^2 + c^2 - a^2)/2bc, \) then Gasper shows that \( d\mu_{\theta, \phi}(\psi) = G(\theta, \phi, \psi) \, d\mu(\psi) \) with

\[
G(\theta, \phi, \psi) = \frac{\Gamma\left(\frac{1}{2}\right) a^{-1}(1 - c^2)^{-1/2}(bc)^{-\beta - 1/2}}{\Gamma\left(\frac{1}{2}\right)2^{\beta + 5/2}} \cdot F\left[\frac{1}{2} + \beta, \frac{1}{2} - \beta; 1; \frac{1}{2}(1 - B)\right]
\]

(3.8)

if \( |a - b| < c < a + b, \) and

\[
G(\theta, \phi, \psi) = \frac{\Gamma\left(\frac{1}{2}\right) a^{-1}(1 - c^2)^{-1/2}(a^2 - b^2 - c^2)^{-\beta - 1/2}}{\Gamma\left(\frac{1}{2} - \beta\right)\Gamma(\beta + 1)2^{\beta + 3/2}} \cdot F\left[\frac{1}{2} + \beta + \frac{1}{2}, \frac{1}{2}(\beta + \frac{3}{2}); \beta + 1; B^{-2}\right]
\]

(3.9)

if \( c < a - b, \) and \( G(\theta, \phi, \psi) = 0 \) if \( c < b - a \) or \( c > a + b. \)

In the other three cases a hypergeometric function \( F[\lambda, \gamma; \nu; s] \) appears, but because of a theorem of Gauss, \( F \) is uniformly bounded for \( |s| < 1 \) provided \( \nu - \lambda - \gamma > 0. \) In the case being considered \( \nu - \lambda - \gamma = 0, \) and by a theorem of Appell (see, for example, Bartle [3, p. 419]),

\[
\lim_{s \to 1^-}(F[\lambda, \gamma; \nu; s]/\log(1/(1 - s))) \text{ exists and is finite.}
\]

The support of \( G \) requires that \( b - a < c < a + b \) which by a change of variables becomes \( |\theta - \phi| < \psi < \theta + \phi. \) The proof of case (ii) proceeds by considering four cases.

Case (1), \( \phi < 3\epsilon. \) This is immediate since

\[
J < \mu[0, 4\epsilon] \int_0^{\pi/2} |d\mu_{\theta, \phi}(\psi)| < Ce^{2\epsilon + 2}.
\]

In all of the cases that follow, \( C \) will indicate a constant that does not depend on \( \epsilon \) or \( \phi, \) but may change from place to place.

The change of variables \( c = \cos \psi \) yields

\[
J = \mu[\phi - \epsilon, \phi + \epsilon] \int_{\cos \epsilon}^{\cos(\theta - \phi)} G(\theta, \phi, \psi)(1 - c^2)^{1/2}c^{2\beta + 1} dc.
\]

(3.10)

In all of the remaining cases, \( 3\epsilon < \phi < \pi/2, \) so the following inequalities hold:

\[
a + b - \cos \epsilon = \cos(\theta - \phi) - \cos \epsilon < \epsilon^2,
\]

\[
a + b > c > \cos \epsilon > \sqrt{3}/2,
\]

\[
a^2 - b^2 - c^2 < (a + b + c)(a + b - c) < 2(a + b - c)
\]

\[
> (a - b - c)(a - b + c) > \sqrt{3}/2(a - b - c) \quad \text{(if} \ a > b)\]

Case (2), $3\pi < \phi < \pi/2$ and $a < b$. For this case the following inequalities hold:

$$b > \sqrt{3}/4, \quad \mu[\phi - e, \phi + e] \sim ea, \quad 1 + B \sim 1$$

(the symbol $x \sim y$ means $x/y$ and $y/x$ are bounded). Now substituting (3.8) in (3.10), assuming $b - a < \cos e$ and using Appell's theorem, obtain

$$J \leq Ce(1 - c^2)^{1/2}c^{2e+1} \int_{\cos e}^{a+b} \left( a^{-1}(1 - c^2)^{-1/2}(bc)^{-\beta-1/2} \cdot \log \left( \frac{1}{1 - \frac{1}{2}(1 - B)} \right) \right) (1 - c^2)^{1/2}c^{2e+1} \, dc$$

$$\leq Ce\int_{\cos e}^{a+b} \log(2(1 + B)^{-1}) \, dc$$

$$\leq Ce(a + b - \cos e) < Ce^3.$$ 

If $\cos e < b - a$, a similar estimate works.

Case (3), $a > b$ and $3\pi < \phi < \pi/2 - 3\pi$. The following relations now hold:

$$a > \sqrt{3}/4, \quad a - b < \cos e < c, \quad \mu[\phi - e, \phi + e] \sim eb^{\beta+1/2}, \quad 1 + B \sim (c - (a - b))/2b.$$ 

It follows again by substitution of (3.8) in (3.10) that

$$J \leq Ceb^{\beta+1/2} \int_{\cos e}^{a+b} b^{-\beta-1/2} \log \left( \frac{2b}{c - (a - b)} \right) \, dc$$

and, since $\int \log x = x(\log x - 1)$,

$$J \leq Ce \left[ \log \left( \frac{2b}{c - (a - b)} \right) + 1 \right]$$

$$= Ce \left[ \log 2b - \log(\cos e - (a - b)) \right]$$

$$\leq Ce[a + b - \cos e] < Ce^3$$

by the mean value theorem.

Case (4), $a > b$ and $\pi/2 - 3\pi < \phi < \pi/2$. The following relations hold:

$$a > \sqrt{3}/4, \quad b < Ce^2, \quad \mu[\phi - e, \phi + e] \sim e^{2e+2},$$

$$a - b - c < e^2 \quad \text{if} \quad \cos e < c < a - b,$$

$$1 + B \sim (c - (a - b))/2b,$$
and if \( c < a - b, \) \( B < -1 \) so
\[
1 - B^{-2} \sim 1 + B^{-1} \sim (a - b - c)/(a^2 - b^2 - c^2).
\]
Now if \( \cos \varepsilon < a - b, \) then
\[
J \leq C e^{2\beta + 2} \int_{a - b}^{a + b} \frac{1}{1 - B^{-2}} \log\left(\frac{1}{1 - B^{-2}}\right) dc.
\]
In \( I, \) \( c < a - b \) so substitution of (3.9) in (3.10) yields
\[
I \leq C e^{2\beta + 2} \int_{a - b}^{a + b} \frac{1}{1 + B^{-1}} \log\left(\frac{1}{1 + B^{-1}}\right) dc.
\]
Since \( B < -1, \) \( \log(1 + B^{-1})^{-1} = -\log(1 + B^{-1}) \) can be expanded in powers of \( B^{-1} \) to obtain
\[
I \leq C e^{2\beta + 2} \int_{a - b}^{a + b} \sum_{k=1}^{\infty} k^{-1}(2bc)^k (a^2 - b^2 - c^2)^{-k-\beta-1/2} 2c dc.
\]
Since \( c < a - b \) in the range of integration, Fubini's theorem yields
\[
I \leq C e^{2\beta + 2} (a^2 - b^2 - \cos^2 \varepsilon)^{1/2-\beta} \sum_{k=1}^{\infty} k^{-2} \left(\frac{2b(a - b)}{a^2 - b^2 - c^2}\right)^{k} \int_{a - b}^{a + b} \frac{1}{1 - \frac{1}{2}(1 - B)} \log\left(\frac{2b(a - b)}{c - (a - b)}\right) \frac{1}{1 - \frac{1}{2}(1 - B)} dc.
\]
since \( \frac{1}{2} - \beta > 0. \) But \( c < a - b \) implies that \( 2b(a - b)/(a^2 - b^2 - c^2) < 1 \) so that the series has an absolute bound independent of \( \varepsilon. \) Now it is immediate that
\[
I \leq C e^{2\beta + 2} (e^2)^{1/2-\beta} = C e^{3}.
\]
Finally, evaluate \( II \) by substituting (3.8) in (3.10):
\[
II \leq C e^{2\beta + 2} \int_{a - b}^{a + b} b^{-\beta} \log\left(\frac{1}{1 - \frac{1}{2}(1 - B)}\right) dc.
\]
If \( \cos \varepsilon > a - b, \) the same methods work with the simplification that only \( II \) must be evaluated. This concludes the proof of the theorem.

It remains now to relate the maximal functions \( P_{\ast}f \) and \( S_{\ast}^{\delta} \) to \( \mathcal{M}f. \) Since the arguments in these cases are somewhat similar, the development will be simplified by proving the following general theorem, which is an easy adaption of theorems given by Hardy and Littlewood in [13].
Let \( w_p(\theta) \) be an integrable function that depends on a parameter \( p \) and define \( h(p, \theta) = f \ast w_p(\theta) \).

**Theorem 5.** Assume either

(a) there is a \( B > 0 \), independent of \( p \) such that

\[
\left| \int_0^{\pi/2} w_p(\theta) \, d\mu(\theta) \right| < B
\]

and

\[
\int_0^{\pi/2} \left| \frac{d}{d\theta} w_p(\theta) \right| \, d\theta < B,
\]

or

(b) there are functions \( W_p(\theta) > |w_p(\theta)| \) and \( C \) independent of \( p \) such that \( W_p(\pi/2) < C \) and

\[
\int_0^{\pi/2} \left| \mu[0, \theta] \frac{d}{d\theta} W_p(\theta) \right| \, d\theta < C;
\]

then there is \( K \) independent of \( p \) such that

\[
|h(\theta, p)| < K \Re f(\theta).
\]

The proof is based on integration by parts. Let \( f_1(\psi) = f \ast \delta(\psi) \, d\mu(\phi) \) so that

\[
h(p, \theta) = \int_0^{\pi/2} f \ast \delta(\psi) \, w_p(\psi) \, d\mu(\psi)
\]

\[
= f_1(\psi) w_p(\psi)|_{\psi=0}^{\psi=\pi/2} - \int_0^{\pi/2} f_1(\psi) \mu[0, \psi] \frac{d}{d\psi} w_p(\psi) \, d\psi,
\]

and proceed as in [13, Theorems 15 and 19].

**Theorem 6.** Suppose \( f \) is nonnegative; then there exist positive constants \( A \) and \( B \) such that

\[
A \Re f(\phi) < P_\epsilon f(\phi) \quad \text{if} \quad \alpha + \beta > 0 \quad \text{or} \quad \beta > -\frac{1}{2},
\]

\[
P_\epsilon f(\phi) < B \Re f(\phi) \quad \text{if} \quad \alpha + \beta > -1.
\]

**Proof.** Observe that Stirling’s formula implies

\[
\left( \frac{a}{n} \right)_n \frac{(b)_n}{(c)_n} \to \frac{\Gamma(a + b - c)\Gamma(c)}{\Gamma(a)\Gamma(b)},
\]

so that Appell’s theorem yields

\[
(1 - s)^{a+b-c}F(a, b; c; s) \to \frac{\Gamma(a + b - c)\Gamma(c)}{\Gamma(a)\Gamma(b)} \quad (a + b - c > 0).
\]
Since the formula of Bailey (2.2) above reduces to a hypergeometric series when \( \theta = 0 \),

\[
P_r(\phi) = P_r(\phi, 0) \sim C \frac{(1 - r)}{(1 + r)^{a + \beta + 2}} \left( 1 - \frac{\cos^2 \phi}{k^2} \right)^{-a - \frac{3}{2}}
\]

Now \( 1 - \cos^2 \phi/k^2 \sim [(1 - r)^2 + \sin^2 \phi] \) so finally there must exist positive constants \( C_1 \) and \( C_2 \) such that

\[
C_1 P_r(\phi) \leq (1 - r) \left[ (1 - r)^2 + \sin^2 \phi \right]^{-a - \frac{3}{2}} \leq C_2 P_r(\phi). \tag{3.11}
\]

In particular, if \( \phi < (1 - r) \),

\[
C_2 P_r(\phi) \geq 2^{-a - \frac{3}{2}} (1 - r)^{-2a - 2},
\]

hence, for some positive constant \( A \), \( P_r(\phi) > Ah_{1-r}(\phi) \). The restrictions on \( \alpha \) and \( \beta \) are chosen to guarantee a positive convolution so that

\[
f(r, \phi) = f \ast P_r(\phi) > Af \ast h_{1-r}(\phi) = A^\mathbb{M}_{1-r} f(\phi).
\]

To obtain the reverse inequality let \( W_r(\theta) \) be the middle expression in (3.11) and use Theorem 5.

The following result generalizes a theorem in Bonami and Clerc [4].

**Theorem 7.** Suppose \( \alpha + \beta > -1 \) and \( \alpha + \frac{1}{2} < \delta < [\alpha + \frac{3}{2}] \) (the greatest integer in \( \alpha + \frac{3}{2} \)). Then

\[
S^\delta f(\phi) \leq A \left[ \mathbb{M} f(\phi) + \int \int |f(\theta)| \left| \frac{1}{2} \pi - \psi \right|^{\beta - \frac{1}{2}} d|m_{\phi, \psi}(\psi)| d\mu(\theta) \right].
\]

**Proof.** The following inequalities hold:

\[
|K^\delta_L(\theta)| < CL^{2a + 2}, \quad 0 < \theta < \frac{1}{4} \pi,
\]

\[
|K^\delta_L(\theta)| < CL^{a + 1/2 - \delta} |\theta|^{-a - \delta - \frac{3}{2}}, \quad 2L^{-1} < \theta < \frac{1}{4} \pi,
\]

\[
|K^\delta_L(\theta)| < C \left| \frac{1}{2} \pi - \theta \right|^{-\beta - \frac{1}{2}}, \quad \frac{1}{4} \pi < \theta < \frac{1}{2} \pi.
\]

The first two relations are found in [4, Theorem 2.1] where the notation \( s^\delta_L(\cos 2\theta) = K^\delta_L(\theta) \) is used. The remaining inequality is obtained by the techniques used in the cited work. Now let \( w_L(\theta) = K^\delta_L(\theta)X_{(0, \pi/4)}(\theta) \). The first two relations above then imply that \( W_L(\theta) = CL^{a + 1/2 - \delta}(L^{-1} + \theta)^{-(\alpha + \beta + \frac{3}{2})} \) is a majorant for \( w_L(\theta) \), so the result follows by Theorem 5.

It is now possible to state the main result of this section.

**Theorem 8.** Let \( 1 < p < \infty \).

(a) For \( \alpha, \beta > -1 \), \( M \) is of strong type \( p - p \) and weak type \( 1 - 1 \).
(b) For \( \alpha + \beta > -1 \), \( \mathbb{M} \) is of strong type \( p - p \) and weak type \( 1 - 1 \).
(c) For \( \alpha + \beta > -1 \), \( P_\phi \) is of strong type \( p - p \) and weak type \( 1 - 1 \).
(d) For \( \alpha + \beta > -1 \) and \( \alpha + \frac{1}{2} < \delta \), \( S^\delta_\phi \) is of strong type \( p - p \) and weak type \( 1 - 1 \).
Proof. (a) Since the interval \([0, \pi/2]\) with the measure \(d\mu\) is of "homogeneous type" the argument in \([5, p. 71]\) yields the result.
(b) By Theorem 4 and the above case, \(\mathcal{M}\) has the desired properties.
(c) In the region \(\alpha + \beta > -1\), Theorem 6 and the above implies the result since for \(f\) of arbitrary sign \(P_\ast f < P_\ast |f| < \mathcal{M} f\) because \(P_\ast (\theta, \phi) > 0\).
(d) The argument is similar to (c).

Remark. An alternate proof that \(\mathcal{M}\) is of strong type \(p - p\) can easily be constructed from Theorem 6 and the observation that the semigroup \(T^f(\phi) = f(e^{-t}, \phi)\) satisfies properties I and II of the axioms for a symmetric diffusion semigroup, and by Stein \([18, p. 73]\) \(P_\ast\) must be of strong type \(p - p\).

Finally, it is now possible to complete the proof of Theorem 1(c). The argument is based on the same ideas as in \([8, p. 62]\), so that it will suffice to establish that

\[\|QF\|_q < A\|F\|_q \quad (1 < q < \infty)\]

where

\[QF(\phi) = \sup_{0 < r < 1} \int (1-r)^{-2a-2}k_{\eta,r}(\psi)F \ast \delta_\phi(\psi)\ d\mu(\psi).\]

It suffices to show \(QF(\phi) < C\mathcal{M} F(\phi)\) whenever the right-hand side is finite. Let \(w_r(\theta) = (1-r)^{-2a-2}k_{\eta,r}(\theta)\) and apply Theorem 5 and Theorem 8(b).

4. Multipliers for Jacobi series. An immediate application of the above ideas is the following multiplier theorem.

Theorem 9. The Banach space \(S(2, \eta)\) is a space of \(L^p\) multipliers for \(1 < p < \infty\), if \(\eta > \alpha + 1\), \(\alpha > 0\), \(\beta > -1\).

The spaces \(S(q, \eta)\) are not defined here (see \([8, p. 32]\) for complete definitions and a discussion of the properties of the \(S[q, \eta]\) spaces). Suffice it to say that when \(k\) is a positive integer, \(1 < q < \infty\), \(m \in S[q, k]\), if and only if

(a) \(\sup_n |m(n)| = C_0 < \infty\),
(b) \(\sup_M \left\{ M^{-1} \sum_{k=0}^M |n^k \Delta^k m(n)|^q \right\}^{1/q} = C_1 < \infty\)

and the norm of \(m\) in \(S[q, k]\) is equivalent to \(C_0 + C_1 (\Delta^k m(n)\) is the \(k\)th difference of the sequence \(\{m(n)\}\)). Thus the \(S(q, \eta)\) spaces contain as special cases the classical multiplier sequences.

What is needed is the following

Lemma. If \(\eta > 1\), \(m \in S(2, \eta)\), \(m(j) = 0\) for \(j = 0, 1, \ldots, [\eta] + 1\). Define \(M(r, \theta) = \sum_{n=0}^{\infty} m(n)n^r h_n R_n(\theta)\), then

(a) \(|d_\theta M(r, \theta)| < A(1-r)^{-\eta-2a-2}\|m\|_\infty\),
(b) \(\|(1 - \cos \theta)\|^{1/2}d_\theta M(r, \theta)\|_2 < Ar(1-r)^{-\alpha-1}\|m\|_{S(2,\eta)'}\)

The proof of (a) is an easy consequence of the estimate \(h_n \sim n^{2a+1}\) and the
fact that \( m(n) \) is bounded. The proof of (b) is an easy adaptation of the argument in [8]. The only change is that the following lemma must be used in place of [8, p. 69, equation (2)].

**Lemma.** Suppose \( \eta \) is an integer and \( f \sim \sum c_n h_n R_n \); then

\[
\int (1 - \cos \theta)^\eta \left( f(\theta) \right)^2 d\mu(\theta) < K \sum_{n=0}^{\infty} \sum_{j=1}^{\eta} (n+1)^{j-\eta} \Delta c_n^2 h_n.
\]

The lemma is similar to [6, Lemma 4.3] but it makes use of the difference computations in [7, §5].

The proof of the theorem now follows the well-understood path given in [8, p. 73] so it is omitted.

It is now possible to interpolate between the above theorem and the following results:

A. If \( m \in S(1, \alpha + \frac{3}{2}) \), then \( m \) defines an \( L^p \) multiplier for \( 1 < p < \infty \), \( \alpha > \beta > -\frac{1}{2} \) (Gasper and Trebels [11], [12]).

B. If \( m \in S(\gamma, \frac{1}{q} + \epsilon) \), then \( m \) is a \( L^2 \) multiplier for \( 1 < q < \infty \).

C. If \( m \in S(1, 1) \), then \( m \) is an \( L^p \) multiplier for all \( p \) such that \( |1/p - \frac{1}{2}| \leq 1/2(2\alpha + 2) \), \( \alpha > \beta > -\frac{1}{2} \) (Askey [1]).

If \( S_\alpha^{\alpha, \beta} \) is defined as the set of parameters \( (1/p, 1/q, \gamma) \) for which the operator \( \mathcal{T}: L^p \otimes S(q, \gamma) \to L^p \) defined by \( \mathcal{T}(m, f) = \sum m(n) f(n) h_n R_n \) is continuous, then \( S_\alpha^{\alpha, \beta} \) has a convex interior and contains the above mentioned results as “edge” results. The proof is an immediate application of the ideas of weak-stability and the fact that the \( S(q, \gamma) \) spaces are Sobolev modules (see [8, §10] for details).

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