THE GROMOV BOUNDARY OF THE RAY GRAPH

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ABSTRACT. The ray graph is a Gromov-hyperbolic graph on which the mapping class group of the plane minus a Cantor set acts by isometries. We give a description of the Gromov boundary of the ray graph in terms of cliques of long rays on the plane minus a Cantor set. As a consequence, we prove that the Gromov boundary of the ray graph is homeomorphic to a quotient of a subset of the circle.

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1. INTRODUCTION

1.1. **Background.** Let $S = \mathbb{R}^2$ – Cantor be the plane minus a Cantor set. The mapping class group Mod(S) arises naturally in dynamics problems of group actions on the plane and through *artinization* of groups acting by homeomorphisms on Cantor sets (see, for example, [5], [6]). The curve graph and curve complex are important tools in the study of mapping class groups of surfaces of finite type. However, the standard definition of the curve graph for S yields a graph with finite diameter since any two closed curves are disjoint from a sufficiently small loop. Therefore, we must find a more suitable object of study which usefully captures the action of Mod(S).

Such an object, the ray graph, was defined by Danny Calegari in [5]. The mapping class group Mod(S) naturally acts on the ray graph, and the first author proved in [1] that it has infinite diameter and is Gromov-hyperbolic. Hence, this graph can play a similar role to curve complexes for finite type surfaces.

Erica Klarreich gives in [9] a description of the Gromov boundary of curve complexes of finite topological type surfaces in terms of minimal filling laminations on the surface (see also [12] for a more recent proof using infinite unicorn paths). In the case of S, complications arise in generalizing this definition as the surface is of infinite type and in particular its lamination space is non-compact in the Hausdorff topology.

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Our goal in this paper is to give a simple characterization of the Gromov boundary of the ray graph.

1.2. **Results.** We give a brief overview of our results, leaving precise definitions for Section 2. A *short ray* is a simple arc embedded in S up to isotopy connecting infinity to a point on the Cantor set. The *ray graph* is the graph whose vertices are short rays and whose edges connect pairs of disjoint rays. The mapping class group Mod(S) acts by isometries on the ray graph, which, as proved in [1], is Gromov-hyperbolic and has infinite diameter.

A long ray is a simple arc embedded in S which limits to infinity on one end but has no limit point on the other (we define this condition precisely below). A long ray is more properly called a curve, since an arc implies its domain is closed, but inspired as we are by the arc graph for finite type surfaces, we think of long rays as arcs which accidentally fail to ever end. A *loop* is an embedded simple arc connecting infinity to itself. We define a *completed ray graph*, whose vertices are short and long rays and loops and whose edges are pairs of disjoint rays and loops. We prove that this completed ray graph has infinitely many connected components and describe them (Theorem 2.8.1). The main one contains the usual ray graph and is quasi-isometric to it. In particular, it has infinite diameter. The others are *cliques*, i.e., complete sub-graphs of diameter 0 or 1. We exhibit a natural bijection between the set of such cliques and the Gromov boundary of the main component, which is also the Gromov boundary of the ray graph. See Theorem 5.1.1. Thus, roughly speaking, the Gromov boundary of the ray graph is identified with a set of long rays.

Coming from another angle, we describe a specific cover of S, called the *conical* cover. This cover has a circle S^1 as a natural boundary, which plays a primary role in our description. Indeed, the main interest of this cover is the existence of a natural injection from the set of vertices of the completed ray graph to this circle boundary. This injection allows us to see the union of the vertices of the cliques as a subset \mathcal{H} of S^1 . In particular, \mathcal{H} inherits the induced topology of the circle S^1 . Quotienting \mathcal{H} by the equivalence relation "being in the same clique", we get a quotient space \mathcal{H}/\sim endowed with the quotient topology. We show that the Gromov boundary of the ray graph is homeomorphic to \mathcal{H}/\sim . See Theorem 5.4.3.

1.3. **Outline.** In Section 2, we define the completed ray graph and the different types of rays that we will encounter. In Section 3, we adapt the *infinite unicorn paths* defined by Witsarut Pho-On in [12] to our situation. Section 4 reviews Gromov boundaries of non-proper Gromov-hyperbolic spaces and provides some necessary technical background relating unicorn paths to quasi-geodesics. Finally, in Section 5, we describe a bijection between the set of cliques and the Gromov boundary, and we prove that this bijection is a homeomorphism.

2. The completed ray graph

2.1. The plane minus a Cantor set. As above, let $S = \mathbb{R}^2$ – Cantor be the plane minus a Cantor set. Note that it does not matter how the removed Cantor set is embedded in the plane, as all embeddings are homeomorphic; see, for example, [11]. Therefore, we will take the Cantor set to be K, the standard middle-thirds Cantor set on the horizontal axis. In addition, S is homeomorphic to a 2-sphere minus a Cantor set and a single isolated point ∞ . For this reason, we will often abuse

notation and think of S as $\mathbb{S}^2 - (K \cup \{\infty\})$, where \mathbb{S}^2 is the Riemann sphere, from which we have removed the above-defined Cantor set K and the point ∞ . From this viewpoint, we think of both K and ∞ as lying on the equator (real axis) of \mathbb{S}^2 . The fact that K and ∞ divide this equator into infinitely many segments will be important for our purposes.

It will also be useful to fix a complete hyperbolic metric (of the first kind) on S. Figure 1 shows the different ways in which we can think about S.

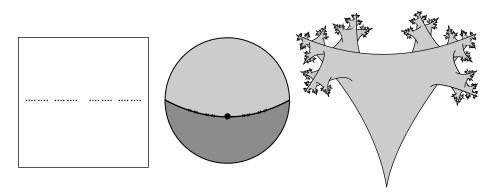


FIGURE 1. Three views of the plane minus a Cantor set: the usual viewpoint, as a sphere minus a Cantor set and a point, and as endowed with a hyperbolic metric

2.2. The conical cover. Since S is hyperbolic, we can think of the universal cover of S as the hyperbolic plane. We won't make particular use of the metric, but it is convenient to know that the boundary of the universal cover is a circle. We can work directly with the universal cover, but in fact there is a sub-cover which is even more amenable, the *conical cover*, defined as follows.

Think of S as $\mathbb{S}^2 - (K \cup \{\infty\})$, and recall that we have chosen an equator in S which contains $\{\infty\} \cup K$. A lift of each hemisphere is a half-fundamental domain (see Figure 2). Choose a lift of ∞ on the boundary of the universal cover, and quotient by parabolics around this lift in order to get a cover of S with one special lift $\tilde{\infty}$ of ∞ which bounds only two half-fundamental domains, as the cover on the right of Figure 2. Such a cover is said to be a *conical cover*. We fix a conical cover and denote it by \tilde{S} and its (circle) boundary by $\partial \tilde{S}$. Call the covering $\pi : \tilde{S} \to S$. See Figure 3 for two views of a conical cover.

2.3. Different types of rays.

2.3.1. Short rays and loops. An embedding γ of the segment]0,1[in S is said to be:

- a *loop* if it can be continuously extended in \mathbb{S}^2 by $\gamma(0) = \gamma(1) = \{\infty\}$,
- a short ray if it can be continuously extended in \mathbb{S}^2 by $\gamma(0) = \{\infty\}$ and $\gamma(1) \in K$.

We are interested in loops and short rays only up to isotopy. Hereafter, we conflate a short ray or loop with its isotopy equivalence class. We say that two short rays or loops are disjoint if there are disjoint rays or loops in their equivalence

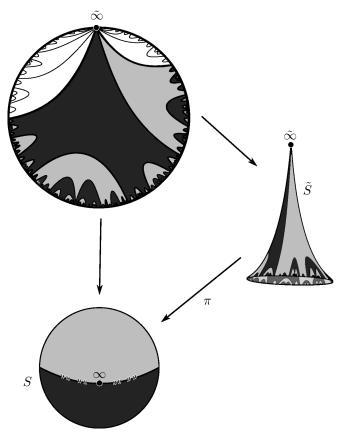


FIGURE 2. The universal cover and a *conical* cover \tilde{S} of $\mathbb{S}^2 - (\{\infty\} \cup K)$

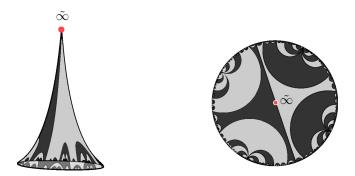


FIGURE 3. Two views of a conical cover: one as a cone, and the other one seen from the top - we will mainly use the second one

class (i.e., if they can be made disjoint by separate isotopies). There is a technical issue with loops, which is that we actually consider *oriented* loops. In practice, this is never an issue, except that we must keep it in mind when we define unicorn paths. We can take a loop and its reverse to be disjoint.

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The conical cover is very useful for understanding short rays and loops because it allows us to make a natural choice of equivalence class representatives, as follows. Given a representative of a short ray r, there is a distinguished lift \tilde{r} from the basepoint $\tilde{\infty}$ in the conical cover. The ray \tilde{r} limits to some point $p \in \partial \tilde{S}$ (by definition, a lift of the endpoint of r in the Cantor set). It follows from standard facts about the hyperbolic plane that the equivalence class of r is specified by p, and there is a unique geodesic in the equivalence class of r with limit point p which we can choose as the class representative. Similarly, given a loop, there is a welldefined limit point in $\partial \tilde{S}$ (which is a lift of ∞ , but not the distinguished lift $\tilde{\infty}$), and we can choose a class representative which is a geodesic in the conical cover. Note that we conflate this geodesic lift and its image in S. Figure 4 shows examples of a short ray and a loop.

The hyperbolic metric on \tilde{S} is not necessary to decide whether two rays or loops are disjoint. However, it is very useful because if we choose class representatives which are these geodesics, then these representatives are disjoint if and only if there are any class representatives which are disjoint. Also, a ray or short loop is simple iff its geodesic representative is simple. These are standard facts about hyperbolic geodesics.

2.3.2. Long rays. We have defined geodesic representatives of short rays and loops. We observed that short rays and loops are specified by points on the boundary $\partial \tilde{S}$ and that there is a subset of $\partial \tilde{S}$ corresponding to simple short rays and loops. We could instead start with $\partial \tilde{S}$ and ask: which points $p \in \partial \tilde{S}$ are the endpoints of geodesics whose image under the covering map $\pi : \tilde{S} \to S$ is simple? We would find the lifts of points in K which are endpoints of simple short rays and the lifts of ∞ which are the endpoints of simple loops. However, we would also find other points on the boundary. These are the long rays, defined as follows:

• Let $\tilde{\gamma}$ be a geodesic ray from $\tilde{\infty}$ to $\partial \tilde{S}$. If $\gamma = \pi(\tilde{\gamma})$ is simple and γ is not a short ray or loop, then it is a *long ray*.

Note that again the hyperbolic metric is not necessary for the definition of long rays – any ray from ∞ which lifts to a ray with a well-defined limit point on $\partial \tilde{S}$ is a long ray, but the metric is convenient to pick out a geodesic representative. Figure 4 gives an example of a long ray.

2.4. Cover-convergence and *k*-beginning. The conical cover will be helpful to us especially because it is compact. We record the idea that rays converge on the boundary.

Definition 2.4.1. We say that a sequence of rays or oriented loops (x_n) coverconverges to a geodesic (ray or loop) l on the surface if the sequence of endpoints of the $\tilde{x_n}$ on the boundary of the conical cover converges to a point p such that the image $\pi((\tilde{\infty}p))$ of the geodesic $(\tilde{\infty}p)$ by the quotient map of the covering is l.

Lemma 2.4.2. Let $E \subseteq \partial \tilde{S}$ be the set of endpoints of short rays, loops, and long rays (i.e., the set of all endpoints of geodesics whose projection to S is simple). Then E is compact. The endpoints of loops are isolated (are not accumulation points of E).

Proof. Since $\partial \tilde{S}$ is compact, to show that E is compact it suffices to show that E is closed. Let (γ_i) be a sequence of geodesic rays in \tilde{S} such that $\pi(\gamma_i)$ is simple

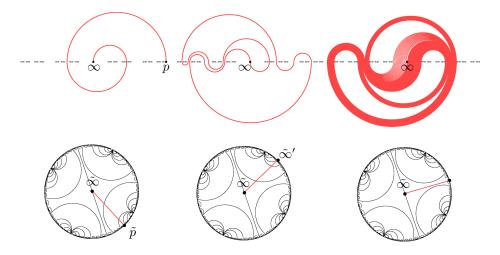


FIGURE 4. A short ray, a loop, and (the beginning of) a long ray. We show the lifts of the rays and loop in the conical cover. Note the short ray ends on a lift of a point in the Cantor set, the loop ends on another lift of ∞ , and the long ray ends on a boundary point which is neither. Note the pictures are drawn *on the sphere*, with ∞ highlighted in the middle.

(each is a short ray, loop, or long ray). Suppose that γ_i cover-converges to γ . To prove the lemma, it suffices to show that $\pi(\gamma)$ is simple. Suppose not. Then $\pi(\gamma)$ must have a self-intersection *not* at ∞ nor at a point in the Cantor set. Therefore, γ must intersect another lift of $\pi(\gamma)$, and this intersection must occur in the interior of \tilde{S} . Since γ_i cover-converges to γ , we must have γ_i converging to γ in the Hausdorff metric when γ_i and γ are thought of as subsets of the disk (analogous to the Poincaré disk). Therefore, for *i* sufficiently large, we must find a similar self-intersection between γ_i and another lift of $\pi(\gamma_i)$. This is a contradiction since all the γ_i are simple.

To show that loops are isolated, observe that no simple geodesic ray can cross the equator adjacent to ∞ more than twice (a simple geodesic cannot wrap around ∞). This produces a region about each lift of ∞ on $\partial \tilde{S}$ into which no simple geodesic can enter (except to terminate exactly on that lift of ∞).

Lemma 2.4.3. Let (x_i) and (y_i) be sequences of short rays, loops, or long rays such that the pair x_i , y_i is disjoint for all *i*. Suppose that (x_i) cover-converges to x and (y_i) cover-converges to y. Then x and y are disjoint.

Proof. By Lemma 2.4.2, we know that each of x, y is a simple short ray, loop, or long ray, so the only question is whether they are disjoint. Suppose they are not. Then thinking of x and y as their representatives in the conical cover \tilde{S} , we see they must have a point of intersection not at the boundary. But by taking i large enough, we can make x_i and y_i have geodesic representatives in \tilde{S} close enough to x and y that they must intersect. This is a contradiction.

Note that we have no information about whether x and y are short rays, loops, or long rays, but whatever they are, they are simple and disjoint.

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It is also helpful to have a topological or combinatorial notion of cover-convergence. This is the notion of k-beginning-like.

Definition 2.4.4. Let us think of S as $\mathbb{S}^2 - (\{\infty\} \cup K)$, and recall that K and ∞ divide the equator into infinitely many segments. We say that two oriented rays or loops k-begin like each other if they cross the same initial k equatorial segments in the same direction. Note that a loop or long ray is specified uniquely by the sequence of equatorial segments that it crosses, and note that it is important to record in which direction we cross the segments (or equivalently in which hemisphere the ray starts).

The definition applies even to short rays, but these might not be specified just by their equator crossings. Figure 5 gives an example of two rays which k-begin like each other.

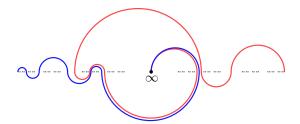


FIGURE 5. Two rays which 4-begin like each other.

Lemma 2.4.5. Let (x_i) be a sequence of rays or loops and let x be a long ray. Then (x_i) cover-converges to x if and only if for all k there is an I so that for all $i \ge I$, we have x_i k-begins like x.

Proof. Because x is a long ray, its endpoint in the boundary of the conical cover is not on the boundary of a fundamental domain of S. Hence cover convergence is equivalent to crossing the same fundamental domains as x; that is, k-beginning-like.

2.5. The ray graphs. The ray graph \mathcal{R}' is the graph whose vertices are short rays up to isotopy and whose edges are pairs of disjoint short rays. Assigning each edge a length of 1 makes \mathcal{R}' into a metric space. See Figure 6. Since homeomorphisms preserve the property of being disjoint, the mapping class group Mod(S) acts by isometries on \mathcal{R}' . In [1], the first author proved that \mathcal{R}' is infinite diameter and hyperbolic. The analogous *loop graph* defined for loops is, in fact, quasi-isometric to \mathcal{R}' . This relationship is critical in [1]. We can define yet another graph, the *short-ray-and-loop graph*, whose vertices are both short rays and loops and whose edges join pairs of objects (rays or loops) which are disjoint.

The completed ray graph \mathcal{R} is the graph whose vertices are short rays, long rays, and loops. The edges of \mathcal{R} are pairs of vertices which can be realized disjointly. It will be convenient to work in the context of \mathcal{R} because we will go back and forth between rays and loops: rays are easier to reason about, and they give a nice description of the Gromov boundary of \mathcal{R} , but loops are easier to work with than short rays. Note that \mathcal{R}' is a sub-graph of \mathcal{R} .

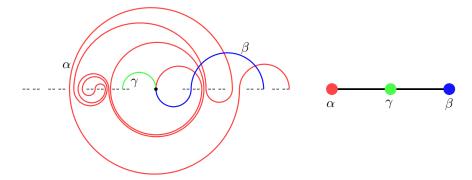


FIGURE 6. Example of three short rays giving a path of length 2 in the ray graph \mathcal{R}' , and the corresponding picture in \mathcal{R}' .

2.6. Actions. Although our overarching goal is to understand the mapping class group Mod(S) of the plane minus a Cantor set, we will not actually be acting by any homeomorphisms in this paper. Here we are concerned with understanding the boundary of the ray graph \mathcal{R}' so that later we may leverage this understanding to study Mod(S). However, it is important to note that Mod(S) does, in fact, act on \mathcal{R} .

Lemma 2.6.1. The mapping class group Mod(S) acts by homeomorphisms on the boundary $\partial \tilde{S}$ of the conical cover. It acts by isometries on \mathcal{R}' and \mathcal{R} .

Proof. First, we observe that $\operatorname{Mod}(S)$ acts by isometries on \mathcal{R}' . This is obvious from the definition of short rays and loops. Next, as we require that any mapping class $\phi \in \operatorname{Mod}(S)$ fixes ∞ , there is a well-defined lift $\tilde{\phi} : \tilde{S} \to \tilde{S}$ fixing $\tilde{\infty}$. Now, endpoints of short rays and loops (not necessarily simple) are dense in the circle $\partial \tilde{S}$. As any mapping class must preserve the cyclic order of geodesic rays from $\tilde{\infty}$, the mapping class must preserve the cyclic order of this dense set of ray and loop endpoints and thus extend to a homeomorphism on the boundary $\partial \tilde{S}$. To avoid issues of isotopy, our definition of a long ray is in terms of endpoints on $\partial \tilde{S}$. Hence only now, after establishing that $\tilde{\phi}$ induces a homeomorphism on $\partial \tilde{S}$, can we observe that this implies that ϕ acts on the set of long rays. Furthermore, this action preserves disjointness between long rays and short rays and loops (because this can be defined in terms of cyclic orders on the boundary $\partial \tilde{S}$). Thus ϕ induces an isometry on \mathcal{R} .

Lemma 2.6.2. The mapping class group Mod(S) acts transitively on short rays.

Proof. Let us be given two short rays a, b. We will think of them as being embedded in $S^2 - (\{\infty\} \cup K)$. The endpoints p_a, p_b of a, b are points in the Cantor set. By a theorem of Schoenflies (see for example [10, Theorem A.3]), there is a homeomorphism f of the sphere such that $f(\infty) = \infty$ and f(a) = b (and hence $f(p_a) = p_b$). If we had f(K) = K, we would be done, but this is probably not the case. We will find a homeomorphism g of the sphere which fixes ∞ , b, and p_b such that g(f(K)) = K. The composition $g \circ f$ will then be the desired element of Mod(S). We let $U = S^2 - \{\infty, p_b\} \cup b$. Note U is an open disk whose closure is the ray b. Every point of f(K) and K is contained in U except for the point which

is labeled $p_b \in K$. To alleviate confusion, we will denote by x the point which is labeled p_b in K (and is $f(p_a)$ in f(K)).

It is proved in [2, Proposition I.1], that if C, C' are two Cantor sets in the interior of a closed disk D, there is a homeomorphism of D which fixes the boundary and maps C to C'. We cannot quite apply this result to our situation because the Cantor sets K and f(K) have the point x on the boundary.

Because K and f(K) are both Cantor sets, the union $K \cup f(K)$ is also a Cantor set, and we can enclose it in a union of disks as follows. Let A_n be a sequence of closed disks such that:

- (1) $A_n \subseteq U$.
- (2) The A_n are pairwise disjoint.
- (3) As $n \to \infty$, the set A_n Hausdorff converges to the single point x.
- (4) For all n, A_n contains some point in K and some point in f(K) and ∂A_n does not intersect $K \cup f(K)$.
- (5) $(K \cup f(K)) \{x\} \subseteq \bigcup_n A_n.$

That is, A_0 contains the points in $K \cup f(K)$ which are far from x; then A_1 contains some points that are closer, and so on, limiting to the single point x. Note that both K and f(K) have x as an accumulation point, so there is no trouble containing points from both sets and also Hausdorff limiting to x.

Because ∂A_n does not intersect $K \cup f(K)$, we must have A_n containing two sub-Cantor-sets from K and f(K), and we can apply [2, Proposition I.1] to obtain a homeomorphism $g_n : A_n \to A_n$ taking $f(K) \cap A_n$ to $K \cap A_n$. Since the diameter of A_n is going to zero, the infinite composition $g = \ldots \circ g_1 \circ g_0$ is a well-defined homeomorphism of U which takes f(K) to K. Then $g \circ f$ is the desired homeomorphism in the lemma.

2.7. Filling rays. We will be interested in understanding how long rays interact with short rays and loops. The concept of *filling* rays is important. Recall rays and loops are defined up to isotopy; hence a long ray intersects a ray or loop if it intersects every representative of their equivalence class. Equivalently, the geodesic representatives intersect.

Definition 2.7.1. A long ray l is said to be:

- *loop-filling* if it intersects every loop.
- *ray-filling* if it intersects every short ray.
- *k*-filling if
 - (1) There exists a short ray l_0 and long rays $l_1, \ldots, l_k = l$ such that l_i is disjoint from l_{i+1} for all $i \ge 0$.
 - (2) k is minimal for this property.
- high-filling if it is ray-filling and not k-filling for any $k \in \mathbb{N}$.

That is, a long ray is k-filling if it is at distance exactly k from the set of short rays in the graph \mathcal{R} . A long ray is high-filling if it is not in the connected component of \mathcal{R} containing the short rays.

2.7.1. Examples of high-filling rays. In this section, we give some examples of high-filling rays. We will not need the fact that these rays are actually high-filling; they are provided as examples to help understand the ideas. Therefore, we just give a brief sketch of why they are high-filling.

Figure 7 gives an example of a high-filling ray. It is the limit ray of the sequence (α_k) defined in [1]. This sequence is such that for every $k \geq 1$, every ray disjoint from a ray which $\log(\alpha_{k+1})$ -begins like α_{k+1} has to $\log(\alpha_k)$ -begin like α_k , where for every ray x, $\log(x)$ denotes the number of segments of the equator that x intersects. Moreover, for every k, $\log(\alpha_{k+1})$ is greater than $\log(\alpha_k)$ (these properties are established in [1]). As points of the boundary of the universal cover, the endpoints of the α_k 's are cover-converging to a point, denoted by p. According to Lemma 2.4.2, the geodesic between $\tilde{\infty}$ and p is the lift of a simple ray, denoted by α_{∞} . We claim that this ray α_{∞} is high-filling. Indeed, according to the properties of (α_k) , there is no ray disjoint from α_{∞} (more precisely, every ray disjoint from α_{∞} has to ∞ -begin like α_{∞}). Thus, there are no edges in \mathcal{R} incident on α_{∞} ; it is a clique unto itself. As a consequence, we see that there exist infinitely many cliques in \mathcal{R} : any image of α_{∞} under an element ϕ of $Mod(\mathbb{R}^2 - Cantor)$ which does not preserve α_{∞} gives a clique with the one element $\phi(\alpha_{\infty})$.

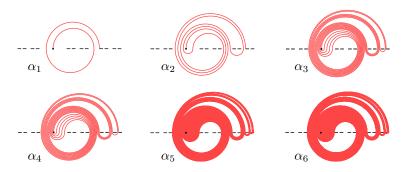


FIGURE 7. Example of (the beginning of) a high-filling ray – the high-filling ray itself is α_{∞} (which looks like α_6 if we do not zoom).

As the example above shows, it is non-trivial to prove that a long ray is highfilling. We can create more examples by finding rays fixed by pseudo-Anosov mapping classes of sub-surfaces. The details of this are not important for this paper; we mention it here as a comment that we have examples of cliques of high-filling rays with an arbitrary number of elements, (see, for example, Figure 8), but we don't know if there exist cliques of high-filling rays with infinitely many elements.

2.7.2. Some properties of filling rays.

Lemma 2.7.2. Every ray-filling ray is also loop-filling. Hence every high-filling ray is loop-filling.

Proof. Let r be a ray-filling ray. Suppose towards a contradiction that r is not loop-filling, so there is some non-trivial loop l which is disjoint from r. Now l divides S into two regions U_1, U_2 by the Jordan curve theorem, and since r and l are disjoint, it must be that r lies completely inside one of the two regions, say U_1 without loss of generality. Since l is a non-trivial loop, there are points in the Cantor set K in both U_1 and U_2 . Hence, it is simple to draw a short ray totally inside U_2 connecting ∞ to some point of the Cantor set. By construction, this ray is disjoint from r. This is a contradiction.

Conjecture 2.7.3. There are loop-filling rays which are not ray-filling.

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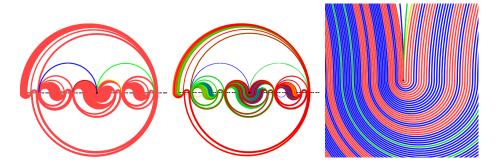


FIGURE 8. Example of (the beginning of) four mutually disjoint high-filling long rays. One of the rays is shown in red, left, with the beginnings of the others. We see them all (confusingly) drawn together, middle, and we show how they come together at ∞ , right.

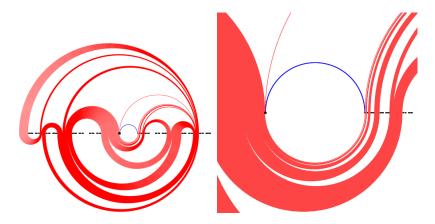


FIGURE 9. An apparently loop-filling ray (red) which is not ray-filling.

For an illustration of why we believe Conjecture 2.7.3, see Figure 9.

Lemma 2.7.4. Let L be a loop-filling ray. If l and l' are two rays (short or long) disjoint from L, then l and l' are disjoint.

Proof. Suppose that l and l' intersect. We must be careful, because a priori l' might accumulate on l in such a way that there is no "first" intersection as we follow l from ∞ . However, we can still let x be a point of intersection, and let l_x , l'_x be the initial segments of l and l' from ∞ to the point of intersection x. Let y be the first intersection point between l_x and l'_x as we follow l_x from ∞ . Let $l_{x,y}$ be the initial segment of l_x from ∞ to y, and let $l'_{x,y}$ be the similar initial segment of l'_x . Since l and l' are in minimal position, $l_{x,y} \cup l'_{x,y}$ is a non-trivial loop, and it is simple because y is that first intersection of l_x and l'_x . This non-trivial loop is disjoint from L, which is a loop-filling ray. This is a contradiction, so we conclude that l and l' are disjoint.

Corollary 2.7.5. Let *l* be a loop-filling ray which is not ray-filling; then every long ray disjoint from *l* is not ray-filling.

Proof. Since l is not ray-filling, there is a short ray x disjoint from l. Now suppose we are given a long ray l' disjoint from l. As l is loop-filling, according to Lemma 2.7.4, x and l' are disjoint. Hence l' is not ray-filling.

Lemma 2.7.6. There exists no k-filling ray for k > 2.

Proof. Suppose there is a k-filling ray l for k > 2. Then there is a short ray l_0 and long rays $l_1, \ldots, l_k = l$ such that (l_0, l_1, \ldots, l_k) is a path in the graph \mathcal{R} . Consider the portion of the sequence $l_0, l_1, l_2, l_3, \ldots$ Now, l_2 must be ray-filling (or else we could shorten the sequence by removing l_1) and hence loop-filling. Applying Lemma 2.7.4, we find that l_3 and l_1 must be disjoint. But in this case, we can remove l_2 from the sequence entirely, and in fact l is (k-1)-filling. This is a contradiction because l was assumed to be k-filling.

Note that Lemma 2.7.6 implies that any long ray which is in the connected component of \mathcal{R}' inside \mathcal{R} must be at distance at most 2 from \mathcal{R}' .

Question 2.7.7. Do there exist 2-filling rays?

Note that this question is equivalent to asking whether ray-filling rays are necessarily high-filling. We do not know the answer. It seems difficult to find a long ray which intersects every short ray but is disjoint from a long ray which is disjoint from a short ray.

Lemma 2.7.8. Any connected component of \mathcal{R} containing a high-filling ray is a clique of high-filling rays.

Proof. Let C be a connected component of \mathcal{R} containing a high-filling ray. By the definition of high-filling rays, C cannot contain any short rays. This implies it cannot contain any loops (which always have some disjoint ray). So any two elements of C are high-filling rays. Lemma 2.7.4 implies that if two long rays a,b are disjoint from l, then a,b are also disjoint from each other. Thus C is a clique. \Box

Lemma 2.7.9. As a subset of the boundary of the conical cover, every clique of high-filling rays is compact.

Proof. The boundary of the conical cover is compact, so it suffices to show that every clique is closed under cover-convergence. To see this, suppose we have a sequence (l_i) of high-filling rays, all pairwise disjoint, which cover-converges to a ray or loop l. By Lemma 2.4.2, l must be simple, and by Lemma 2.4.3, it must be disjoint from all the l_i . But since l is disjoint from all (any) of the l_i , it cannot be a short ray or loop because all the l_i are high-filling. Therefore, l is a high-filling ray, and as it is disjoint from all (any) of the l_i , it is in the same clique. So the clique is closed under cover-convergence, and we are done.

2.8. Connected components of \mathcal{R} . It was shown in [1] that the loop graph and the ray graph are quasi-isometric. More precisely, a quasi-isometry from the ray graph to the loop graph is given by any map which sends a short ray to a disjoint loop (see [1, section 3.2]). This implies that the natural inclusion from the loop graph to the short-ray-and-loop graph is a quasi-isometry. Thus the ray graph, the loop graph, and the short-ray-and-loop graph are quasi-isometric. We will now show that these graphs are also quasi-isometric to the connected component of \mathcal{R}' inside \mathcal{R} .

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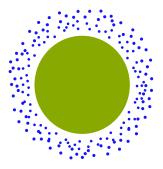


FIGURE 10. A representation of the connected components of the completed ray graph: the main one is infinite diameter – it contains the short-ray-and-loop graph and is quasi-isometric to it; others are cliques, of diameter 0 or 1. We will give a bijection between the Gromov boundary of the main one and the set of cliques (Theorem 5.1.1).

Theorem 2.8.1. There is a single connected component of \mathcal{R} containing all the short rays and loops. This component is quasi-isometric to the ray graph \mathcal{R}' , to the loop graph, a to the short-ray-and-loop graph. All other connected components (which are the high-filling rays) are cliques (complete sub-graphs).

We start with a technical lemma.

Lemma 2.8.2. Let a and b be two loops at distance 3 in the completed ray graph. Then they are at distance 3 in the short-ray-and-loop graph.

Proof. Let (a, λ, μ, b) be a path in the completed ray graph. We will show how to replace λ and μ with short rays, which will prove the lemma. This proof proceeds by chasing down various cases.

Recall the equator is divided into segments by the Cantor set. For each of λ and μ , there are two cases depending on whether they intersect finitely many or infinitely many equatorial segments. Suppose that λ intersects infinitely many segments. As a and b are loops, they can only intersect finitely many segments. Hence there is some segment I which λ intersects but a and b do not. We can therefore connect λ to a point in the Cantor set with a path disjoint from a and b, producing a short ray (this short ray is disjoint from a because λ is, but might not be disjoint from b). If μ intersects I as well, so much the better: we can connect it as well, producing two short rays. See Figure 11.

In the latter case, we are done: we have produced a path of length 3 between a and b in the short-ray-and-loop-graph. In the former case, we may now assume that λ is actually a short ray intersecting the equator finitely many times. If μ intersects infinitely many segments, we repeat the argument and we are done.

Hence we have two cases: λ and μ , long rays, both intersect finitely many equatorial segments, or λ is actually a short ray and μ intersects finitely many segments. Potentially by switching the roles of λ and μ , then, we may assume that λ is a long ray intersecting finitely many segments. Since λ intersects finitely many segments, there must be some segment I it intersects infinitely many times. Since a and b intersect the equator finitely many times, there must be some sub-segment $I' \subseteq I$ such that λ intersects I' at least twice, and a, b intersect it not at all. Let p_1, p_2 be

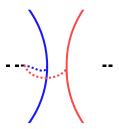


FIGURE 11. If λ (and possibly μ) intersect an equatorial segment which neither *a* nor *b* do, then we can connect their first intersections to the Cantor set with disjoint paths. Note λ and μ are allowed to crash through this picture in any way; as long as we take the first points of intersection, the resulting short rays are disjoint.

the first and second points of intersection of I' and λ as we travel along λ from ∞ . The union of the interval along λ between p_1 and p_2 and the interval along I' between p_1 and p_2 is a simple closed curve γ_{λ} . This curve divides the sphere into two connected components, one of which does not intersect a (because a is disjoint from both I' and λ). Call this component D_{λ} . Note that D_{λ} must contain some points in the Cantor set because λ is geodesic.

There are several cases depending on the status of μ . If μ is disjoint from D_{λ} , then we can easily connect λ to a Cantor set point in the interior with a path. This replaces λ with a short ray, and we can swap the roles of λ and μ and go back to the beginning of this case to handle μ . See Figure 12.

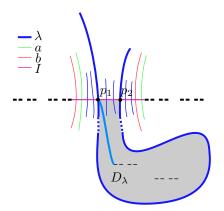


FIGURE 12. If μ is disjoint from D_{λ} , we just replace λ with a short ray which simply dead ends to any Cantor set point in D_{λ} (the change is shown in lighter blue).

If μ does intersect D_{λ} , then we must consider b. Note that b can intersect D_{λ} finitely many times. Hence, there is a subdisk $D'_{\lambda} \subseteq D_{\lambda}$ disjoint from b and containing the interval I'. Since μ is disjoint from λ and b, it must be that $\mu \cap D_{\lambda}$ is contained within D'_{λ} . As μ is geodesic, we must have some Cantor set points within D'_{λ} . Since a and b are both disjoint from D'_{λ} , we can connect both λ and μ to Cantor set points with paths in D'_{λ} . This replaces λ and μ with short rays λ' , μ' with the

same disjointness properties, producing the desired path in the short-ray-and-loop graph. See Figure 13. $\hfill \Box$

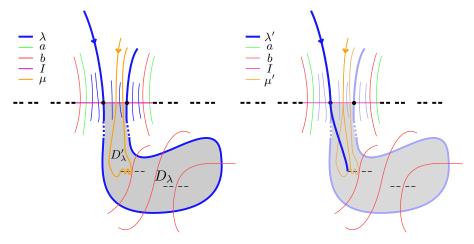


FIGURE 13. If μ is not disjoint from D_{λ} , then we create D'_{λ} . Here we can short-cut both λ and μ , as shown on right. Note it only matters that the new paths are disjoint from each other, not the original μ . For example, μ could even be ray-filling in D'_{λ} .

Proof of Theorem 2.8.1. We recall from [1] the definition of a unicorn path between two loops in the loop graph (or, see Section 3.2). The precise definition is not important: all we need is that this operation produces a path between any two loops. Thus the loop graph is connected.

By [1], we already know that the ray graph \mathcal{R}' is quasi-isometric to the loop graph and short-ray-and-loop graph. The loop graph is connected, so these other graphs are too. Because these graphs are quasi-isometric, we may conflate their definitions and show that the natural inclusion $\mathcal{R}' \to \mathcal{R}$ of the short-ray-and-loop graph into its connected component in \mathcal{R} is a quasi-isometry. Note all this means is that we must show that lengths in \mathcal{R}' are not distorted when we add the ability to take "detours" using long rays in \mathcal{R} . First we show it is a quasi-isometric embedding and then prove it is quasi-surjective.

Consider two loops or short rays a and b at distance $n \in \mathbb{N}$ in the completed ray graph. There exists two loops a' disjoint from a and b' disjoint from b. Consider a geodesic path (λ_i) between a' and b' in the completed ray graph. We claim that every element of this path is not loop-filling: by contradiction, if some λ_j is loopfilling, then because λ_j is different from a' and b', λ_{j-1} and λ_{j+1} are well-defined vertices in the path, and according to Lemma 2.7.4, they are disjoint. This gives a contradiction because (λ_i) is geodesic.

For every *i* such that λ_i is well defined, we choose a loop a_i disjoint from λ_i . According to Lemma 2.8.2, a'_i and a'_{i+1} at distance at most 3, hence the distance between a' and b' in the short-ray-and-loop graph is at most 3 times the distance between a' and b' in the completed ray graph. Hence we have a path of length 3n + 2 between a and b in the short-ray-and-loop graph. Thus we conclude that $\frac{1}{3}d_{\mathcal{R}'}(a,b) - 2 \leq d_{\mathcal{R}}(a,b) \leq d_{\mathcal{R}'}(a,b)$ so the inclusion is a quasi-isometry. Lemma 2.7.6 immediately implies that this inclusion is quasi-surjective.

Finally, Lemma 2.7.8 says that all other connected components of \mathcal{R} must be cliques because they are composed of high-filling rays.

Remark 2.8.3. After Theorem 2.8.1, we will conflate the definitions of the ray graph, the loop graph, and the short-ray-and-loop graph, using whichever is most convenient. We'll refer to all of these graphs, and also to the quasi-isometric connected component of \mathcal{R} , by \mathcal{R}' . We'll also call this component the *main component*, as distinguished from the *cliques*. We show in Section 5 that the cliques are in bijection with the Gromov boundary of \mathcal{R}' .

3. Infinite unicorn paths

3.1. **Outline.** We are interested in understanding the boundary of \mathcal{R}' . We will use the same general idea as [12] of *infinite unicorn paths*, although with a slightly different definition. That idea, in turn, was motivated by unicorn paths in the curve graph, as defined in [7]. See [1] for an introduction to unicorn paths in the loop graph. In this section, it is most convenient to work with the loop graph, although we recall that there is a bijection between its boundary and the boundary of the ray graph (and the short-ray-and-loop graph and the main component of \mathcal{R}).

3.2. Finite unicorn paths. There are two equivalent definitions of a finite unicorn path. We choose to use a slightly different definition from that which appears in [1,7] because it has some nice properties. However, it is therefore important for us to explain both definitions and prove their equivalence.

Definition 3.2.1 (Unicorn path, first definition, [1, 7]). Let a, b be loops representing vertices in the loop graph. Let them be oriented. We define the *unicorn path* $P_1(a, b)$ between a and b as follows. If a, b are disjoint, then the unicorn path is simply the pair (a, b). If a, b intersect, then we order their points of intersection x_1, x_2, \ldots, x_k in the reverse order they appear along a. We define $a_i = (\infty x_i)_b \cup (\infty x_i)_a$; that is, the loop which is the union of the interval along b from ∞ to x_i and the interval along a from ∞ to x_i (the beginnings of both a and b until they reach x_i). It is possible that this loop is not simple. The unicorn path $P_1(a, b)$ is the sequence $(a, a_{i_1}, a_{i_2}, \ldots, a_{i_n}, b)$, where i_1, i_2, \ldots are the indices of the simple loops defined above. That is, we define all the loops as above, then select just the simple ones, and that is the unicorn path. Note that because we order the points of intersection in the reverse order they appear along a, we will get a little less of a at each step along the path, and thus we interpolate between a and b

Definition 3.2.2 (Unicorn path, second definition). Let a, b be loops representing vertices in the loop graph. Let them be oriented. We define the unicorn path $P_2(a, b)$ as follows. Set $a_0 = a$, and proceed inductively. Let x_k be the first intersection between b and a_k as we travel along b from ∞ . If such an intersection does not exist, then $a_{k+1} = b$ and we are done. Otherwise, set $a_{k+1} = (\infty x_k)_b \cup (\infty x_k)_{a_k}$; that is, the union of the interval along b from ∞ to x_k followed by the reverse of the interval along a_k from ∞ to x_k . Note that a_k is always simple, and a_k, a_{k+1} are always disjoint. The unicorn path is the sequence $P_2(a, b) = (a = a_0, a_1, \ldots, b)$.

Lemma 3.2.3. Definitions 3.2.1 and 3.2.2 are equivalent. If $P_1(a,b) = (a = x_0, x_1, ..., x_n, b)$ and $P_2(a,b) = (a = y_0, y_1, ..., y_m, b)$, then n = m and $x_i = y_i$ for all *i*.

Proof. Although the definition of P_2 is inductive, we can still think of each element of the unicorn path (second definition) as the union of a beginning interval of b and a beginning interval of a (if we simply don't deform a_k , it is clearly the union of an interval on a and an interval on b; a_{k+1} is created by simply going farther along b before changing to a). Hence each y_i does appear in the sequence $P_1(a, b)$. We will prove that $x_i = y_i$ for all i inductively. Clearly $x_0 = a = y_0$ by definition. Now suppose that $x_i = y_i$. Write

$$y_i = (\infty p)_b \cup (\infty p)_a$$
 and $y_{i+1} = (\infty p')_b \cup (\infty p')_a$;

that is, y_i is constructed by following b until the intersection point p and then a backward, and y_{i+1} is constructed by following b until the point p' (which is farther along b, note) and then a backward. Note that in the interval (p, p') along a, we might find other points of intersection. If any of these points of intersection can produce a simple loop, then these loops will appear between $x_i = y_i$ and y_{i+1} in the first unicorn path $P_1(a, b)$. Thus if we can show that all these loops are, in fact, non-simple, then we will know that $y_{i+1} = x_{i+1}$.

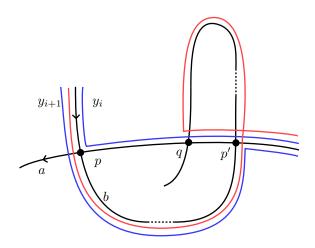


FIGURE 14. The proof of Lemma 3.2.3. Any intersection of b with the interval (p, p') along a must come after p'. Hence any loop produced by Definition 3.2.1 from any of these intersections will not be simple, and since $y_i = x_i$, we must have $y_{i+1} = x_{i+1}$.

Let q be a point between p and p' along a. Were q to appear before p along b, then y_i would not be simple (it would have a self-intersection at q); this is a contradiction to our assumption that it is simple, so we know that q must appear after p' along b. But now if we take the union

$$(\infty q)_b \cup (\infty q)_a,$$

then this loop cannot be simple, as it must have an essential self-intersection at p'. See Figure 14. Therefore, no intersection point in the interval (p, p') along b can contribute a loop to the sequence $P_1(a, b)$, since none of them are simple. We conclude that $x_{i+1} = y_{i+1}$ and the lemma is proved by induction.

Following Lemma 3.2.3, we may now freely use our preferred Definition 3.2.2 but cite facts about Definition 3.2.1. Hereafter, we define the unicorn path P(a, b) to be $P_2(a, b)$ and we will assume that construction. But we may use the following lemmas, which were adapted in [1] from lemmas in [7].

Lemma 3.2.4 (Lemma 3.3 of [1]). For every $0 \le i < j \le n$, either $P(a_i, a_j)$ is a subpath of $P(a, b) = \{a = a_0, a_1, ..., a_n = b\}$, or j = i + 2 and a_i and a_j represent adjacent vertices of the loop graph.

Lemma 3.2.5 (Proposition 3.5 and Corollary 3.6 of [1]). If g is a geodesic of the loop graph between a and b, then P(a, b) is included in the 6-neighborhood of g, and the Hausdorff distance between g and P(a, b) is at most 13.

An immediate consequence of the second lemma is that unicorn paths are close to quasi-geodesics.

Lemma 3.2.6. For any κ , ϵ , there exists $C \in \mathbb{N}$ such that if (x_n) is a (κ, ϵ) quasi-geodesic in the loop graph, then for every n, k with n > k, the unicorn path $P(x_0, x_k)$ is contained in the C-neighborhood of the unicorn path $P(x_0, x_n)$.

Proof. By Lemma 3.2.5, unicorn paths are uniformly close to geodesics, and for any κ , ϵ there is a constant C' such that (κ, ϵ) -quasi-geodesics are C'-close to geodesics. This constant is called the (κ, ϵ) -Morse constant of the loop graph; see Section 4. Putting these together yields the lemma.

One way to think about the unicorn path P(a, b) is that it slowly begins more and more like b. We can make this precise with the following lemma. Recall Definition 2.4.4, which defines the notion of k-beginning-like.

Lemma 3.2.7. Let a be a loop which intersects the equator p times. For any loop or ray b, enumerate the unicorn path $P(a,b) = (x_n)$. Then for $i > k\lceil p/2 \rceil$, the loop x_i must k-begin like b.

Proof. Divide $a = x_0$ into segments corresponding to which hemisphere of the sphere they are in. Since a intersects the equator p times, there are at most $\lceil p/2 \rceil$ in either hemisphere. In particular, b can intersect a at most $\lceil p/2 \rceil$ times before crossing the equator. By the definition of unicorn paths, this implies that after at most $k \lceil p/2 \rceil$ terms in the unicorn path, elements of the unicorn path will align with b for at least k equator crossings; that is, they will k-begin like b.

Note that it is certainly possible that the unicorn path does not have $\lceil p/2 \rceil$ terms in it, in which case this lemma is vacuous. The lemma just asserts that if the unicorn path is long enough compared to the number of intersections of a with the equator, then elements of the unicorn path must start beginning like b.

3.3. Infinite unicorn paths. We now adapt the idea of unicorn paths to build unicorn paths between a loop and a long ray. This idea is similar to the infinite unicorn paths defined by Witsarut Pho-On in [12, Section 3.1].

Definition 3.3.1. Let a be an oriented loop and let l be a long ray. Recall we have chosen geodesic class representatives for a and l, so in particular they are in minimal position. We think of l as oriented away from ∞ , and we define the unicorn path P(a, l) as follows (see Figure 15):

- (1) Set $a_0 = a$.
- (2) Let x_{k+1} be the first intersection between l and a_k as we follow l from ∞ . We define $a_{k+1} = (\infty x_{k+1})_l \cup (\infty x_{k+1})_{a_k}$.

Note as with finite unicorn paths, a_k and a_{k+1} are clearly disjoint. The unicorn path P(a, l) is the (possibly infinite) sequence $P(a, l) = (a, a_1, ...)$.

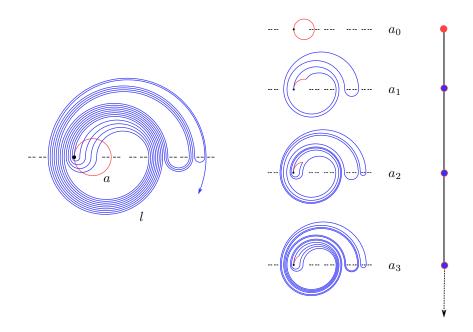


FIGURE 15. Example of the beginning of an infinite unicorn path P(a, l)

Since we always have a_{k+1} disjoint from a_k , P(a, l) is a path in the loop graph, which we can think of as a path in \mathcal{R}' if necessary. In this definition, we refer to the "first" intersection of a_k and l. To address possible confusion, we note that the set of intersection points $a_k \cap l$ is discrete as a subset of l, although it might not be discrete as a subset of a_k . That is, l is a long ray, so it is free to accumulate in complicated ways in the loop a_k , but a is a loop, so were it to accumulate anywhere in l, there would necessarily be points in the Cantor set arbitrarily close to a_k and thus on a_k , which is a contradiction.

We also remark that we can always remember that a_k is made up of one segment on l and one segment on a. As a consequence, if b is disjoint from a, then b is always disjoint from the segment $(\infty x_k)_{a_k}$, and if b is also disjoint from l, then it is disjoint from the entire unicorn path P(a, l).

Remark 3.3.2. It is important to note that an infinite unicorn path is so named because it is a unicorn path between a loop and an infinite object (a long ray). It is possible that an infinite unicorn path is actually a finite sequence of loops if the

long ray intersects the loop finitely many times. For our purposes, these cases will be degenerate, so there should generally be no confusion.

Lemma 3.3.3. Let l be a long ray, and let a and b be two oriented loops at distance n in the loop graph. Then the unicorn paths P(a, l) and P(b, l) have Hausdorff distance at most n.

Proof. Assume first that a and b are disjoint and let us prove that every element of P(a, l) is disjoint from some element of P(b, l). The general result for a and b at distance $n \in \mathbb{N}$ follows by considering a path $(a = c_0, c_1, ..., c_n = b)$ and apply this first result to each pair of disjoint loops (c_j, c_{j+1}) .

Consider an element $a_k := (\infty x_k)_l \cup (\infty x_k)_a$. Denote by l_k the sub-segment $(\infty x_k)_l \subset l$. If l_k is disjoint from b, then a_k is disjoint from b and we are done. If l_k is not disjoint from b, it intersects it finitely many times. Consider the first intersection point y_k between b and l_k when we follow b from ∞ . The segment $(\infty y_k)_b$ is disjoint from l_k (by definition) and from a (because it is included in b). Hence the loop $b' := (\infty y_k)_l \cup (\infty y_k)_b$ is disjoint from a_k . Moreover, because l_k intersects b only finitely many times, b' is an element of P(b, l) (by definition of unicorn paths).

Lemma 3.3.4. Infinite unicorn paths restrict to finite unicorn paths. More precisely, if $a_j \in P(a, l)$ with $j \ge 3$, then $P(a, a_j) \subset P(a, l)$.

Proof. This follows from the construction of P(a, l) because a_j contains the initial segment of l which generates the first j intersections with a. The restriction $j \ge 3$ follows from the special case of j = i + 2 in Lemma 3.2.4.

3.4. Infinite unicorn paths and cover-convergence.

Lemma 3.4.1. Let (x_n) be a sequence of oriented loops which has a subsequence which cover-converges to a loop-filling ray l. Then for every $k \in \mathbb{N}$, there exists n_k such that the unicorn paths $P(x_0, l)$ and $P(x_0, x_{n_k})$ have the same k-first terms.

Proof. Because l is loop-filling, $P(x_0, l)$ is infinite (by definition of unicorn path). The result follows then from Lemma 2.4.5 and the definition of unicorn paths. \Box

3.5. Infinite unicorn paths and filling rays. In this section, we relate whether or not a ray is filling to the behavior of the path P(a, l). It will turn out that the long rays whose unicorn paths give points on the boundary of \mathcal{R}' are exactly the high-filling rays. The definition of high-filling is in terms of rays, while the definition of unicorn paths is in terms of loops. Therefore, we must do some translating.

Note that all distances in this section are computed in the loop graph. This is for simplicity. However, by Theorem 2.8.1, all statements apply equally well to the loop, ray, or short-ray-and-loop graphs, possibly with a change of constants.

3.5.1. Bounded infinite unicorn paths.

Lemma 3.5.1. Let l be a long ray which is not loop-filling. For every oriented loop a, P(a, l) is bounded.

Proof. Consider a loop *a* disjoint from *l*. The only element of P(a, l) is *a*, hence P(a, l) is bounded. According to Lemma 3.3.3, for every oriented loop *b*, P(b, l) is included in the d(a, b)-neighborhood of P(a, l), hence P(b, l) is bounded.

Lemma 3.5.2. Let l be a long ray which is not ray-filling. Then there exists a loop a such that P(a, l) is included in the 2-neighborhood of a. Moreover, for every oriented loop b, P(b, l) is bounded.

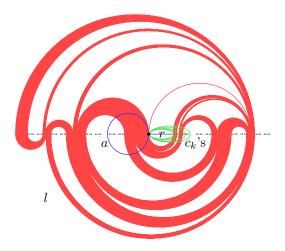


FIGURE 16. Picture for the proof of Lemma 3.5.2. Example of a nested sequence of loops.

Proof. Let r be a short ray disjoint from l. Consider a nested sequence of mutually disjoint loops (c_k) "around" r, smaller and smaller (as in Figure 16).

Let a be a loop disjoint from all of the c_k (this can be obtained by closely looping around a ray disjoint from r). By the construction of (c_k) , for every $x \in P(a, l)$, there exists k such that c_k is disjoint from x. Note that l might be loop-filling, and hence the c_k might not be disjoint from l. However, each $x \in P(a, l)$ contains only a sub-segment of l, followed by a sub-segment of a, both of which are disjoint from c_k for sufficiently large k. Hence for all $x \in P(a, l)$ there is a k so that the sequence $a - c_k - x$ is a path in the loop graph and therefore P(a, l) is always within distance 2 of a.

According to Lemma 3.3.3, for every oriented loop b, P(b, l) is included in the d(a, b)-neighborhood of P(a, l), hence P(b, l) is bounded.

Lemma 3.5.3. Let l be a 2-filling ray. For every oriented loop a, P(a, l) is bounded. Moreover, there exists a such that P(a, l) is included in the 2-neighborhood of a.

Proof. By definition of a 2-filling ray, there exists a ray l' disjoint from l and which is not ray-filling. By Corollary 2.7.5, l' cannot be loop-filling (this would imply that l is not loop-filling, a contradiction). Let a be a loop disjoint from l'. Consider any $x_k = (\infty p_k)_l \cup (\infty p_k)_a$ in P(a, l). We will find a loop y_k disjoint from a and x_k : hence x_k is at distance at most 2 from a. See Figure 17.

Denote by l_k the sub-segment $(\infty p_k)_l$ of l between ∞ and p_k . Consider $l_k \cup a$: this is a compact set of the sphere, whose complement is a finite union of open sets.

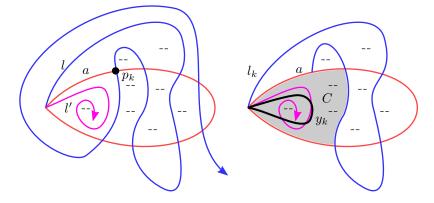


FIGURE 17. The picture for the proof of Lemma 3.5.3. Here y_k is drawn in black on the right; note it does not matter if it intersects l' as long as it is disjoint from a and l_k .

Each component of the complement of $l_k \cup a$ is arcwise connected and homeomorphic to a disk or to a disk minus a Cantor set. As l' is disjoint from both l and a, it is included in a connected component of the complement of $l_k \cup a$. Denote it by C. Note that C:

- is not simply connected, hence contains some points of the Cantor set (because l' is a geodesic included in C);
- contains ∞ in its boundary (because it contains the ray l').

It follows that we can draw a loop y_k included in C: this loop is disjoint from both a and x_k . Hence, P(a, l) is in the 2-neighborhood of a.

According to Lemma 3.3.3, for every oriented loop b, P(b, l) is included in the d(a, b)-neighborhood of P(a, l), hence P(b, l) is bounded.

3.5.2. Unbounded infinite unicorn paths.

Lemma 3.5.4. Let a be an oriented loop and let l be a long ray. The three following statements are equivalent:

- (1) The ray l is high-filling.
- (2) The infinite unicorn path P(a, l) is unbounded in \mathcal{R}' .
- (3) The infinite unicorn path P(a,l) goes to infinity in \mathcal{R}' ; that is, if $(x_i)_i = P(a,l)$, then $\lim_{i\to\infty} d(a,x_i) = \infty$.

Proof. First note that (2) is an immediate consequence of (3). We prove that (2) implies (1). Lemmas 3.5.2 and 3.5.3 show that any non-ray-filling and any 2-filling long ray l has P(a, l) bounded for all a. By Lemma 2.7.6, there is no k-filling ray for k > 2. Hence any long ray l in the main component of \mathcal{R}' (i.e., not high-filling) produces a bounded unicorn path P(a, l) for any loop a. So if P(a, l) is unbounded, then l must be high-filling.

Next, we must show that P(a, l) goes to infinity if l is high-filling (i.e., (1) implies (3)). Towards this end, let l be high-filling and let a loop a be given. Suppose towards a contradiction that P(a, l) does not go to infinity. We may assume that d(a, x) = N for infinitely many $x \in P(a, l)$, and by taking a further subsequence we can produce a subsequence (y_i) of P(a, l) so that $d(y_i, a) = N$ for all *i*. Now for each y_i , we can produce a sequence $a = y_i^0, y_i^1, \ldots, y_i^N = y_i$ of loops in the loop graph such that y_i^j and y_i^{j+1} are disjoint.

As we have discussed, each ray or loop gives a point on the boundary $\partial \tilde{S}$ of the conical cover, which is compact. Therefore, after possibly taking a subsequence of the y_i , we may assume that all of the sequences $(y_i^j)_i$ for $0 \leq j \leq N$ cover-converges in $\partial \tilde{S}$. These limits might be loops or rays, but whatever they are, we name them $a = y_{\infty}^0, y_{\infty}^1, \ldots, y_{\infty}^N$. We apply Lemma 2.4.3, which says that the limits of sequences of pairs of disjoint simple rays or loops are disjoint simple rays or loops (again, we don't actually know if they are rays or loops), we observe that this sequence is a path in the completed ray graph \mathcal{R} starting at a. Furthermore, by the construction of an infinite unicorn path, the limit object $y_{\infty}^N = \lim_{i \to \infty} y_i$ must be the original ray l. Hence we have produced a path in the completed ray graph \mathcal{R} from a to l. Since l is high-filling, this is a contradiction.

Lemma 3.5.5. If l and l' are two disjoint long rays, and l' is loop-filling, then for every loop a, P(a, l) is included in the 1-neighborhood of P(a, l').

Proof. Let x be any element of P(a, l). We will find $y \in P(a, l')$ such that x and y are disjoint. By definition, the loop x is the union of a sub-segment \overline{l} of l and a sub-segment \overline{a} of a. We first prove that l' intersects \overline{a} . If not, as l' is disjoint from l, l' is in particular disjoint from \overline{l} , hence l' is disjoint from x. This is a contradiction because l' is loop-filling. Hence l' intersects \overline{a} . Denote by p the first intersection point between l' and \overline{a} when we follow l' from ∞ . The union $y = (\infty p)_{l'} \cup (p\infty)_a$ is an element of P(a, l'). Moreover, by construction, it is disjoint from x.

Lemma 3.5.6. Let l, l' be two distinct high-filling rays, let a, a' be two loops, and let there be some $N \in \mathbb{N}$ such that P(a, l) is included in the N-neighborhood of P(a', l'). Then l and l' are disjoint.

Proof. This proof is similar to the proof of Lemma 3.5.4. We denote P(a,l) by (x_n) . According to the hypothesis, for every n there exists $y_n \in P(a', l')$ such that $d(x_n, y_n) \leq N$. For every n, we choose a path $(x_n = x_n^0, x_n^1, ..., x_n^N = y_n)$ of length N between x_n and y_n . Because $(x_n) = P(a, l)$, (x_n) cover-converges to l, and (y_n) cover-converges to l'. By taking subsequences, we can assume that each sequence x_n^i cover-converges to some loop or ray l_i . Applying the same idea as Lemma 3.5.4, we see that l_i is disjoint from l_{i+1} , and hence $l = l_0, \ldots, l_N = l'$ is a path in the completed ray graph. Since l and l' are high-filling rays, it follows from Theorem 2.8.1 that l and l' are disjoint.

4. Gromov boundaries and unicorn paths

We are building up to Theorem 5.1.1, for which we need to understand some background about what the Gromov boundary is and how unicorn paths are related to quasi-geodesics. In this section, we prove several key facts which will be assembled in Section 5.

4.1. Non-properness. Our aim is to study the boundary of the ray graph \mathcal{R}' , or equivalently the boundary of the loop graph or main component of \mathcal{R} . Recall that a metric space is proper if every closed ball is compact. In proper Gromov-hyperbolic spaces, we can readily understand the boundary using geodesic rays in the space. Unfortunately, \mathcal{R}' , like the standard curve complex for finite type surfaces, is not proper. So we must be more careful.

4.2. The Gromov boundary. We first recall the definition of the *Gromov-bound*ary of a δ -hyperbolic space X (see [4], or [8]). Choose a basepoint $p \in X$. A sequence (x_n) of elements of X is said to be *Gromov-converging at infinity* if:

$$\liminf_{i,j\to\infty} (x_i \cdot x_j)_p = \infty$$

where $(x_i \cdot x_j)_p$ is the Gromov product:

$$(x_i \cdot x_j)_p := \frac{1}{2}(d(x_i, p) + d(x_j, p) - d(x_i, x_j)).$$

Two such sequences (x_n) and (y_n) are said to be *equivalent* if:

$$\liminf_{i,j\to\infty} (x_i \cdot y_j)_p = \infty.$$

Definition 4.2.1. The *Gromov-boundary* of a δ -hyperbolic space X, denoted by ∂X , is the set of equivalent classes of sequences Gromov-converging to infinity.

This definition does not depend on the choice of the basepoint p. If X is proper and geodesic, it is known that every equivalence class of sequences Gromovconverging to infinity contains a sequence which is a geodesic of X. The graph \mathcal{R}' is not proper, and we don't know if there is a geodesic in every equivalence class in its boundary. However, according to Remark 2.16 of [8], there exists a $(1, 10\delta)$ quasi-geodesic path in every element of the boundary. This quasi-geodesic will be enough for our purposes.

Remark 4.2.2. A quasi-isometry between Gromov-hyperbolic spaces extends to a bijection of their boundaries (see, for example, Proposition 6.3 of [3]).

4.3. Quasi-geodesics and cover-convergence. In this section, we show that the only way to approach the boundary of the loop graph is with a sequence coverconverging to a loop-filling ray.

Lemma 4.3.1. Let (x_n) be a quasi-geodesic sequence of loops in the loop graph. For example, let (x_n) represent a point on the Gromov boundary of the loop graph. If α is a ray or loop such that some subsequence of (x_n) cover-converges to α , then α must be a loop-filling ray.

Proof. Recall from Lemma 2.4.2 that in the set of endpoints on the boundary of the conical cover of lifts of simple rays and loops, the set of endpoints of loops is isolated. Therefore, if a subsequence of x_n converges to a loop α , that means that α appears infinitely often in the sequence x_n , which contradicts the fact that it is quasi-geodesic. Thus α must be a ray, and we are left proving that α must be a loop-filling ray.

Suppose towards a contradiction that there exists some loop β disjoint from α . Up to extending (x_n) by a finite path between β and x_0 if necessary, we can assume that (x_n) is a quasi-geodesic which has a subsequence which cover-converges to α and whose first term x_0 is disjoint from α . Abusing notation, we denote again by (x_n) the subsequence which cover-converges to α .

Denote by $\tilde{\alpha}$ and \tilde{x}_n the lifts of α and x_n , respectively, which start at the preferred lift $\tilde{\infty}$ of ∞ . Consider also all the lifts of x_0 in the conical cover: none of them intersect $\tilde{\alpha}$. Because (x_n) is a quasi-geodesic, there exists N such that for any $n \geq N$, x_n intersects x_0 . For every such n, we denote by s_n the first lift of x_0 that x_n intersects.

Because x_0 is a loop, each preimage of each hemisphere only intersects finitely many preimages of x_0 . Consequently, for any ϵ , there are only finitely many lifts of x_0 whose endpoints on the circle are distance greater than ϵ apart. And because α is a ray, the endpoint of $\tilde{\alpha}$ on the boundary of the conical cover is not a lift of ∞ . Therefore, the sequence s_n must have infinitely many distinct elements, and the endpoints must both converge to the endpoint of $\tilde{\alpha}$. It follows that if a sequence of oriented loops (w_n) is such that for every n, the first lift of x_0 intersected by w_n is s_n , then w_n cover-converges to α .

Now we apply Lemma 3.2.6: since (x_n) is a subsequence of a quasi-geodesic, there is a C such that for all n, k with n > k, we have $P(x_0, x_k)$ contained in the C-neighborhood of $P(x_0, x_n)$. Choose x_k at distance at least C + 2 from α . For each n > k there is some loop y_n in $P(x_0, x_n)$ within distance C of x_k . For every n > k, we can make a path of length C in the loop graph between x_k and y_n . Denote it by $x_k, z_{1,n}, z_{2,n}, \ldots, y_n$. By taking a further subsequence, we can make $z_{j,n}$ cover-converge for all j. These limits give us a path $x_k, z_{1,\infty}, z_{2,\infty}, \ldots, y_\infty$ of length C in the completed ray graph. Now note that because $y_n \in P(x_0, x_n)$ for every n and $d(y_n, x_0) \ge 2$, by the definition of unicorn paths we have that the first lift of x_0 intersected by y_n is s_n (x_0 and x_n must intersect, and all entries in the unicorn path $P(x_0, x_n)$ after the first follow x_n until its first intersection with x_0). By the discussion we had in the previous paragraph, it follows that $y_\infty = \alpha$. Hence, we have constructed a path of length C between x_k and α , which are at distance at least C + 2: this is a contradiction.

4.4. **Infinite unicorn paths and quasi-geodesics.** In this section, we relate infinite unicorn paths to quasi-geodesics in the loop graph. This, in turn, relates unicorn paths to the boundary via our remark above that equivalence classes on the boundary contain quasi-geodesics. The lemma in this section is approximately saying that sequences cover-converging to a long ray are close to unicorn paths generated by that ray.

Recall that in every δ -hyperbolic geodesic space X, quasi-geodesics stay closed to geodesics (see, for example, Theorem 1.7, page 401 of [4]). More precisely, for all $\kappa \geq 1$ and $\epsilon \geq 0$, there exists a constant M, called the (κ, ϵ) -Morse constant of X, such that every (κ, ϵ) -quasi-geodesic between any two points $x, y \in X$ is included in the M-neighborhood of any geodesic between x and y.

Lemma 4.4.1. Let $\kappa, \epsilon \geq 0$. Denote by M the (κ, ϵ) -Morse constant of the loop graph. Let (x_n) be a (κ, ϵ) -quasi-geodesic path of the loop graph such that a subsequence of (x_n) cover-converges to a high-filling ray l. Then (x_n) is included in the (13 + M)-neighborhood of the infinite unicorn path $P(x_0, l)$.

Proof. It will be helpful during this proof to consult Figure 18, as there are many details to keep track of.

We denote by (y_n) the unicorn path $P(x_0, l)$, which is infinite and tends to infinity according to Lemma 3.5.4. Let $j \in \mathbb{N}$. Let us show that x_j is in the (13 + M)-neighborhood of $P(x_0, l)$. We denote by d_j the distance $d(x_0, x_j)$.

Let $A \in \mathbb{N}$ such that y_A is the first element of $P(x_0, l)$ at distance $d_j + 2M + 6$ from x_0 ; that is:

- $d(x_0, y_A) = d_i + 2M + 6;$
- for every 0 < k < A, $d(x_0, y_k) < d_i + 2M + 6$.

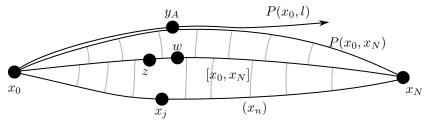


FIGURE 18. The picture for the proof of Lemma 4.4.1. The gray lines indicate the sequences are uniformly close. The sequences $P(x_0, l)$ and $P(x_0, x_N)$ are actually the same until at least y_A .

According to Lemma 3.4.1, there exists N such that $P(x_0, x_N)$ and $P(x_0, l)$ have the same A first terms. Consider a geodesic $[x_0, x_N]$. According to Lemma 3.2.5, the path $P(x_0, x_N)$ is included in the 6-neighborhood of $[x_0, x_N]$. Let $w \in [x_0, x_N]$ be the closest point from x_0 which satisfies $d(y_A, w) \leq 6$. For every $n \leq A$, y_n is 6-close to $[x_0, w]$. Hence $[x_0, w]$ is in the 13-neighborhood of $P(x_0, l)$. The triangle inequality gives us $d(x_0, y_A) \leq d(x_0, w) + d(y_A, w)$, hence $d(x_0, w) \geq d_i + 2M$.

Because (x_n) is a (κ, ϵ) -quasi-geodesic, x_j is in the *M*-neighborhood of $[x_0, x_N]$. Let $z \in [x_0, x_N]$ such that $d(x_j, z) \leq M$. The triangle inequality gives us $d(x_0, z) \leq M + d_j$. Hence $z \in [x_0, w]$. It follows that x_j is at distance at most M + 13 from $P(x_0, l)$.

4.5. Infinite unicorn paths and the Gromov boundary. Finally, we show that unicorn paths of high-filling rays produce points on the boundary of the ray graph \mathcal{R}' .

Lemma 4.5.1. If l is high-filling, then P(a, l) Gromov-converges to infinity for any loop a.

Proof. The proof follows the steps of the proof of Proposition 3.6 of [12]. Recall that a sequence (x_n) of loops Gromov-converges to infinity if

$$\liminf_{i,j\to\infty} (x_i, x_j)_x = \infty,$$

where $(x_i, x_j)_x$ is the Gromov product:

$$(x_i, x_j)_x := \frac{1}{2}(d(x_i, x) + d(x_j, x) - d(x_i, x_j)).$$

We will prove this holds for $(a_n) := P(a, l)$. According to Lemma 3.5.4, for every R > 0, there exists N > 0 such that $d(a, a_n) > R$ for every n > N. For all $m, n \ge N$, we have $(a_n, a_m)_a \ge d(a, [a_n, a_m]) - 2\delta$. Since $[a_n, a_m]$ and $P(a_n, a_m)$ have Hausdorff distance at most 13 by Lemma 3.2.5, this implies that:

$$(a_n, a_m)_a \ge d(a, P(a_n, a_m)) - 13 - 2\delta > R - 13 - 2\delta.$$

For |m-n| > 2, $P(a_n, a_m)$ is contained in P(a, l) by Lemma 3.3.4, so we have the desired property for the sequence P(a, l).

5. The boundary of the ray graph

In Sections 3 and 4, we proved many technical facts about how infinite unicorn paths behave and how the filling-ness of a ray determines the behavior of its unicorn path. In this section, we use these facts to provide a description of the boundary of the ray graph \mathcal{R}' .

5.1. Bijection.

Theorem 5.1.1. Given a point p on the Gromov boundary $\partial \mathcal{R}'$, there is a nonempty set of long rays R such that the point p is the equivalence class of P(a, l) for all $l \in R$ and any loop a. The set R is a clique (and an entire connected component) of high-filling rays in \mathcal{R} . Conversely, given a connected component R of high-filling rays in \mathcal{R} (which is necessarily a clique), there is a single boundary point $p \in \partial \mathcal{R}'$ such that p is the equivalence class of P(a, l) for all $l \in R$ and any loop a, and Ris exactly the set of rays with this property.

Hence the set of high-filling cliques is in bijection with the boundary of the ray graph $\partial \mathcal{R}'$.

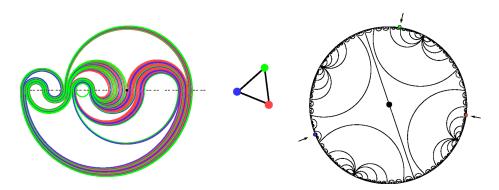


FIGURE 19. Example of the different versions of a point of the boundary: the corresponding clique seen as disjoint high-filling rays on the surface (left), as connected component of the completed ray graph (middle), and as subset of the boundary of the conical cover (right)

Proof. By Lemma 4.5.1, given any short ray a and and any high-filling ray l, P(a, l) Gromov-converges to infinity (is in the equivalence class of a point on the Gromov boundary), and this boundary point does not depend on a by Lemma 3.3.3.

Conversely, let p be a point on the boundary. According to Remark 2.16 of [8], there is a $(1, 10\delta)$ -geodesic path (x_i) in the equivalence class of p. Since $\partial \tilde{S}$ is compact, after passing to a subsequence we may assume that x_i cover-converges to an element $l \in \partial \tilde{S}$. By Lemma 4.3.1, l must be a loop-filling ray. We claim that l is a high-filling ray. To see this, we observe that by Lemma 3.2.5, $P(x_0, x_i)$ stays uniformly close to the geodesic $[x_0, x_i]$, whose length goes to infinity. By Lemma 3.4.1, the length of $P(x_0, l)$ must then be infinite. By Lemma 3.5.4, we must have l be high-filling. By Lemma 4.4.1, we have that $P(x_0, l)$ is a bounded distance from (x_i) . Hence the equivalence class of $P(x_0, l)$ is the same as (x_i) , which is p, as desired.

We have proved that every high-filling ray has a unicorn path which produces a point on the boundary, and we have shown that for every boundary point, there is a high-filling ray whose unicorn path produces it. To complete the proof, we must show that two high-filling rays produce the same boundary point if and only if they are in the same connected component (which is a clique by Theorem 2.8.1). The combination of Lemmas 3.5.5 and 3.5.6 shows that two high-filling rays are disjoint

if and only if their unicorn paths stay a bounded distance apart, i.e., if and only if they have the same equivalence class as boundary points. $\hfill\square$

Following Theorem 5.1.1, we can define a bijection between cliques in \mathcal{R} and the boundary of \mathcal{R}' . Let \mathcal{H} be the set of high-filling rays, and let \mathcal{E} be the set of cliques of high-filling rays in the graph \mathcal{R} . If l is a high-filling ray, we let [l] denote the clique in \mathcal{E} containing l. By Theorem 5.1.1, we can define the map $F : \mathcal{E} \to \partial \mathcal{R}'$ by F([l]) = [P(a, l)], and F is a bijection. That is, we take a clique in \mathcal{E} to the equivalence class of P(a, l) on the boundary, for any l in the clique and any loop a.

5.2. **Topologies.** In this section, we define topologies on both the high-filling rays and the Gromov boundary of the ray graph. The topology on the high-filling rays induces a topology on the set of cliques in \mathcal{R} , and in the space \mathcal{R} endowed with this topology which we will show is homeomorphic to the Gromov boundary. See Figure 20 for an outline of the situation.

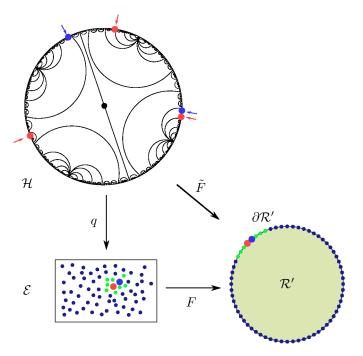


FIGURE 20. The space \mathcal{H} of high-filling rays can be thought of as a subset of the boundary of the conical cover. In the quotient \mathcal{E} , cliques are close if any of their constituent rays are close. We will show $F: \mathcal{E} \to \partial \mathcal{R}'$ is a homeomorphism.

5.2.1. Topology on the set of high-filling cliques. There is a natural topology on the set \mathcal{H} of high-filling rays given by the definition of k-beginning-like: for every $l \in \mathcal{H}$ and every $k \in \mathbb{N}$, we define the k-neighborhood N(l, k) of l as the set of high-filling rays which k-begin like l. The N(l, k)'s are a basis of open sets for \mathcal{H} . Note that they are also closed. Also note that this topology is the same as the topology induced by thinking about the high-filling rays as a subset of the boundary of the conical cover $\partial \tilde{S}$, i.e., as a subset of the circle.

There is a quotient map $q : \mathcal{H} \to \mathcal{E}$ defined by q(l) = [l] taking a ray l to its clique. We define the topology on \mathcal{E} to be the quotient topology of \mathcal{H} under the map q. Note that there is a lift of this map $\tilde{F} : \mathcal{H} \to \partial \mathcal{R}'$ defined by $\tilde{F}(l) = [P(a, l)]$. Clearly we have $\tilde{F} = F \circ q$. See Figure 20.

5.2.2. Topology on the Gromov boundary. There is a standard topology on the Gromov boundary, as follows. For every $s \in \partial \mathcal{R}'$ and for every $r \in \mathbb{N}$, we denote by U(s,r) the set of boundary points t such that there exist two sequences (x_n) and (y_n) such that $s = [(x_n)], t = [(y_n)]$ and $\liminf_{i,j\to\infty} (x_i \cdot y_j)_p \ge r$.

The collection $\{U(s,r)|r > 0\}$ defines a basis of open neighborhoods for s (see [8]).

Remark 5.2.1. Note that if $t \in U(s, r)$, then for any sequences (x_n) and (y_n) such that $s = [(x_n)], t = [(y_n)]$, we have $\liminf_{i,j\to\infty} (x_i \cdot y_j)_p \ge r - 2\delta$.

The remark follows from the fact that for any x, y, z, p in a δ -hyperbolic space, we have $(x \cdot y)_p \geq \min\{(x \cdot z)_p, (y \cdot z)_p\}$.

It is helpful to relate this topology to the usual topology on the boundary of the conical cover.

Lemma 5.2.2. Let (l_i) be a sequence of high-filling rays which cover-converges to a high-filling ray l. Then $[P(a, l_i)]$ converges to [P(a, l)] in the topology on the Gromov boundary $\partial \mathcal{R}'$.

Proof. By Lemma 2.4.5, as $i \to \infty$, l_i k-begins like l for $k \to \infty$. Hence $P(a, l_i)$ and P(a, l) agree for the first n terms, for $n \to \infty$. Thus the Gromov product of these two sequences must go to infinity.

5.3. High-filling rays and neighborhoods. It will be important for us to be able to relate the notion of "close" in terms of rays, i.e., *k*-beginning-like or coverconverging, to the notions of "close" in the Gromov boundary. This is the key step in relating the two topologies defined above. First, we show that being close in the ray graph translates to being close as rays.

Lemma 5.3.1. Let l be a high-filling ray. Let $m \in \mathbb{N}$. For every $k \in \mathbb{N}$, there exists $N \in \mathbb{N}$ such that every loop at distance at most m from a loop which N-begins like l k-begins like one of the rays of the clique of l.

Proof. Let k be given, and suppose towards a contradiction that for all N > 0, there are loops l_N, l'_N such that:

- (1) The distance between l_N and l'_N is at most m.
- (2) l_N N-begins like l.
- (3) l'_N does not k-begin like any ray in the clique of l.

We now apply the same idea as the proof of Lemma 3.5.6. For each N, we can write the path between l_N and l'_N as

$$l_N = l_N^0, l_N^1, \dots, l_N^m = l'_N.$$

By taking a subsequence, we can assume that each sequence $(l_N^i)_N$ cover-converges to some ray or loop l_{∞}^i . Note that $l_{\infty}^0 = l$ by the construction of l_N . By Lemma 2.4.3, this produces a path in the completed ray graph:

$$l = l_{\infty}^0, \dots, l_{\infty}^m$$

By the construction of l'_N , we have that l^m_∞ cannot k-begin like any ray in the clique of l. But by Theorem 2.8.1, since l is high-filling, we must actually have l^m_∞ in the clique of l. This is a contradiction.

Next, we generalize this result to the boundary, showing that being close on the boundary means being close as rays. The following lemma says that for any point P on the boundary, if Q is close to P in the boundary, then for any high-filling ray representative λ' of Q, there is a high-filling ray representative λ of P which begins like λ' .

Proposition 5.3.2. Let $P \in \partial \mathcal{R}'$. For every $k \in \mathbb{N}$, there exists N > 0 such that for every $Q \in U(P, N)$, for every $\lambda' \in \tilde{F}^{-1}(Q)$, there exists $\lambda \in \tilde{F}^{-1}(P)$ which k-begins like λ' .

Proof. Consider $P \in \partial \mathcal{R}'$ and choose $k \in \mathbb{N}$. Let $\lambda \in \tilde{F}^{-1}(P)$. Note that λ is high-filling, and $[P(a, \lambda)] = P$, and the clique of high-filling rays disjoint from λ is $\tilde{F}^{-1}(P)$. These are all consequences of Theorem 5.1.1.

According to Lemma 5.3.1 with $m = 2\delta + 20$, there exists $N' \in \mathbb{N}$ such that every loop at distance less than $2\delta + 20$ from a loop which N'-begins like λ must k-begin like one of the rays of $\tilde{F}^{-1}(P)$. Let us prove that $N := N' + 2\delta$ satisfies the proposition.

Pick a loop a which has only one intersection point with the equator. Denote by (x_n) the infinite unicorn path $P(a, \lambda)$. Because a has only one intersection with the equator, we have that for every $n \geq 1$, x_n at least (n-1)-begins like λ by Lemma 3.2.7. Hence $x_{N'+1}$ N'-begins like λ .

Choose $Q \in U(P, N' + 2\delta)$ and $\lambda' \in \tilde{F}^{-1}(Q)$. Again by Theorem 5.1.1, we have λ' high-filling and $[P(a, \lambda')] = Q$. Denote by (y_n) the unicorn path $P(a, \lambda')$. By definition of $U(P, N' + 2\delta)$, we have $\liminf_{i,j\to\infty} (x_i \cdot y_j)_a \geq N'$. Choose i > N' + 1 and $j \in \mathbb{N}$ such that $(x_i \cdot y_j)_a \geq N'$: any two geodesics $[a, x_i]$ and $[a, y_j]$ have their N' first terms (2δ) -Hausdorff close.

Because geodesics and unicorn paths are uniformly close (Lemma 3.2.5), it follows that some point of (y_n) , say y_p , is in the $2\delta + 20$ neighborhood of $x_{N'+1}$. According to Lemma 5.3.1, y_p k-begins like one of the rays in $\tilde{F}^{-1}(P)$. Thus λ' also k-begins like one of the rays in $\tilde{F}^{-1}(P)$.

5.4. Homeomorphism. In this section, we show that $F : \mathcal{E} \to \partial \mathcal{R}'$ is a homeomorphism, given the two topologies we have defined above.

Lemma 5.4.1. The map $\tilde{F} = F \circ q : \mathcal{H} \to \partial \mathcal{R}'$ is continuous. Consequently, F is continuous.

Proof. Consider a sequence of high-filling rays l_n which converges to l in the topology on \mathcal{H} . That is, for all k there is N such that for all n > N, we have $l_n k$ -begins like l. This is equivalent to l_n cover-converging to l, so applying Lemma 5.2.2, we conclude that $([P(a, l_n)])_n$ converges to [P(a, l)] in the Gromov topology, i.e., $(\tilde{F}(l_n))_n$ converges to $\tilde{F}(l)$.

The quotient map F is then continuous by the definition of the quotient topology.

Lemma 5.4.2. The map $F : \mathcal{E} \to \partial \mathcal{R}'$ is open.

Proof. Consider an open set $\mathcal{U} \subset \mathcal{E}$. Let us show that $F(\mathcal{U})$ is open in ∂X . Let $P \in F(\mathcal{U})$ be a point of the boundary. We denote by $R = F^{-1}(P) \in \mathcal{E}$ the preimage of P. We want to prove that some neighborhood of P is included in $F(\mathcal{U})$.

Because \mathcal{U} is open, by definition of the quotient topology, this means that $q^{-1}(\mathcal{U})$ is open in \mathcal{H} , where $q : \mathcal{H} \to \mathcal{E}$ is the quotient from the set of high-filling rays to the set of cliques, as previously defined. It follows that for every high-filling ray λ in $q^{-1}(R)$, there exists $k(\lambda) > 0$ such that $N(\lambda, k(\lambda)) \subset q^{-1}(\mathcal{U})$, i.e.:

$$\bigcup_{\lambda \in R} N(\lambda, k(\lambda)) \subset q^{-1}(\mathcal{U}).$$

According to Lemma 2.7.9, the clique R is compact as a subset of the boundary of the conical cover, and since each set $N(\lambda, k(\lambda))$, thought of as a subset of the (circle) boundary of the conical cover is open (the intersection of an open interval with the set of high-filling rays), it follows that there exists a finite family $(\lambda_i)_{1 \le i \le n}$ of rays in R such that:

$$R \subset \bigcup_{1 \le i \le n} N(\lambda_i, k(\lambda_i)).$$

Note for any ray $\mu \in R$, we must have $\mu \in N(\lambda_i, k(\lambda_i))$ for some *i*. Set $k := \max_{1 \leq i \leq n} k(\lambda_i)$. According to Proposition 5.3.2, there exists N such that for every $Q \in U(P, N)$, for every $\lambda' \in \tilde{F}^{-1}(Q)$, there exists $\mu \in R$ which k-begins like λ' . Now we have:

- (1) λ' must k-begin like μ for $\mu \in R$.
- (2) Since $\mu \in R$, we have $\mu \in N(\lambda_i, k(\lambda_i))$ for some *i*, so $\mu k(\lambda_i)$ -begins like λ_i .
- (3) $k > k(\lambda_i)$, so as λ' k-begins like μ , and $\mu k(\lambda_i)$ begins like λ_i , we have $\lambda' k(\lambda_i)$ -begins like λ_i .

That is, $\lambda' \in N(\lambda_i, k(\lambda_i))$. We specifically constructed this set so that $N(\lambda_i, k(\lambda_i)) \subset q^{-1}(\mathcal{U})$. Hence, $\lambda' \in q^{-1}(\mathcal{U})$.

We have shown that for all $Q \in U(P, N)$ and for all rays $\lambda' \in \tilde{F}^{-1}(Q)$, we have $\lambda' \in q^{-1}(\mathcal{U})$. That is, $Q \in F(\mathcal{U})$. We conclude that $U(P, N) \subset F(\mathcal{U})$. Since $P \in F(\mathcal{U})$ was arbitrary and U(P, N) is an open set, we conclude that $F(\mathcal{U})$ is open. Hence F is open.

Theorem 5.4.3. The map $F : \mathcal{E} \to \partial \mathcal{R}'$ is a homeomorphism.

Proof. We have proved that F is a bijection (Theorem 5.1.1) which is continuous (Lemma 5.4.1) and open (Lemma 5.4.2).

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