OSCILLATION THEOREMS FOR SECOND ORDER NONLINEAR DIFFERENTIAL EQUATIONS

LYNN ERBE1

ABSTRACT. The oscillatory and nonoscillatory behavior of the nonlinear second order differential equation (1) x'' + p(t)f(x) = 0 is related to that of $(2)_{\lambda} x'' + \lambda p(t)x = 0$, $\lambda > 0$. Under certain conditions on p(t) and f(x) it is shown that all solutions of (1) are oscillatory if $(2)_{\lambda}$ is oscillatory for all $\lambda > 0$. In contrast to most of the literature on this subject, no sign or integrability conditions on p(t) are explicitly assumed.

Consider the second order nonlinear differential equation

$$(1) x'' + p(t)f(x) = 0$$

where $p(t) \in C[0, +\infty)$ and $f(x) \in C^{(1)}(-\infty, +\infty)$, with

(2)
$$f'(x) \ge \frac{f(x)}{x} > 0 \quad \text{for } x \ne 0.$$

As a special case we have

(3)
$$x'' + p(t)x^{2n+1} = 0.$$

In case p(t) is eventually positive, oscillation and nonoscillation criteria for (1) and (3) have been extensively developed. (See [1] and the bibliography therein for the nonlinear case. Willett in [2] has an extensive bibliography for the case when (1) is linear.) However, much less is known for the nonlinear case when p(t) is allowed to be negative for arbitrarily large values of t. It is the purpose of this paper to relate the oscillatory behavior of (1) with that of the linear equation

$$(4)_{\lambda} \qquad x'' + \lambda p(t)x = 0, \quad \lambda > 0,$$

which, presumably, is easier to handle. We shall restrict attention to solutions of (1) which exist on some ray $[T, +\infty)$ where T may depend on the particular solution. We shall at various times assume that the following condition holds:

(5)
$$\lim_{t\to\infty} \inf_{T} \int_{T}^{t} p(s)ds > 0 \quad \text{for all large } T.$$

Received by the editors May 19, 1969.

AMS Subject Classifications. Primary 3442, 3445.

Key Words and Phrases. Second order nonlinear oscillation, boundedness, linear oscillation.

¹ This research was supported by a University of Alberta Post-Doctoral Fellowship.

For the case when $f(x) = x^{2n+1}$, our main result generalizes a theorem due to Utz [4] who assumes $p(t) \ge 0$ and a theorem due to Waltman [5] who has shown that all solutions of (3) oscillate provided the following condition holds:

$$\int_{-\infty}^{\infty} p(s)ds = + \infty.$$

For n = 0 we have by the well-known Fite-Wintner Theorem (see [6]) that condition (6) implies all solutions of (3) oscillate. In fact, we see that $(4)_{\lambda}$ is oscillatory for all $\lambda > 0$.

LEMMA 1. Let u(t) be a nonoscillatory solution of (1) on $[T, +\infty)$ and let condition (5) hold. Then for all large t we have u(t)u'(t) > 0.

PROOF. Assume, to be specific, that u(t) > 0 for $t \ge T_1$, $T_1 \ge T$. Obvious modifications are valid when u(t) < 0. If the lemma is not true, then either u'(t) < 0 for all large t or u'(t) oscillates. In the former case we may suppose that T_1 is sufficiently large so that

$$\int_{T_1}^t p(s)ds \ge 0 \quad \text{for } t \ge T_1$$

and u'(t) < 0 for $t \ge T_1$. Hence, we have

(7)
$$\int_{T_1}^t p(s)f(u(s))ds = f(u(t)) \int_{T_1}^t p(s)ds - \int_{T_1}^t f'(u(s))u'(s) \int_{T_1}^s p(\sigma)d\sigma ds \ge 0, \quad t \ge T_1.$$

Now integrating (1) we have by (7) that $u'(t) \le u'(T_1) < 0$, $t \ge T_1$, which contradicts the fact that u(t) is nonoscillatory.

If u'(t) oscillates, let $T_n \to +\infty$ be such that $u'(T_n) = 0$. For $t \ge T_1$ we define

(8)
$$v(t) = -u'(t)/f(u(t)),$$

and differentiating, we get

$$(9) v'(t) = p(t) + w(t),$$

where

$$w(t) = (v(t))^2 f'(u(t)) \ge 0, \quad t \ge T_1.$$

Since $v(T_n) = 0$ we integrate (9) between T_n and T_{n+1} and sum on n to get an immediate contradiction to (5).

THEOREM 2. Let $(4)_{\lambda}$ be oscillatory and let u(t) be a nonoscillatory solution of (1) with u(t)u'(t) > 0 for all $t \ge T$. Then

(10)
$$\lim_{t \to \infty} \frac{f(u(t))}{u(t)} \le \lambda.$$

PROOF. Let g(x) = f(x)/x. We note that condition (2) implies that the limit in (9) exists (possibly infinite). If the theorem is not true, we may assume $g(u(t)) \ge \lambda$ for all $t \ge T$. Let z(t) be the solution of $(4)_{\lambda}$ satisfying z(T) = 0, z'(T) = 1, and let $T_1 > T$ be the first zero of z'(t) so that z'(t) > 0 on $[T, T_1)$. Then

(11)
$$\int_{T}^{T_1} (g(u(t)) - \lambda)(z')^2 dt \ge 0$$

so integrating by parts we get

$$\begin{split} \int_{T}^{T_{1}} (g(u(t)) - \lambda)(z')^{2} dt \\ &= \lambda \int_{T}^{T_{1}} pz^{2} (g(u(t)) - \lambda) dt - \int_{T}^{T_{1}} zz'g'(u(t))u' dt \\ &\leq \lambda \int_{T}^{T_{1}} pz^{2} (g(u(t)) - \lambda) dt \end{split}$$

since the integrand zz'g'(u(t))u' is nonnegative on $[T, T_1]$ by condition (2). But

$$\begin{split} \int_{T}^{T_{1}} pz^{2}(g(u(t)) - \lambda) dt \\ &= \int_{T}^{T_{1}} \frac{z}{u} (pzf(u(t)) - \lambda pzu) dt \\ &= \int_{T}^{T_{1}} \frac{z}{u} (z'u - u'z)' dt \\ &= - u'(T_{1})(z(T_{1}))^{2}/u(T_{1}) - \int_{T}^{T_{1}} ((z'u - u'z)/u)^{2} dt < 0, \end{split}$$

and this is a contradiction.

Lemma 1 along with Theorem 2 imply the following:

COROLLARY 3. Let $(4)_{\lambda}$ be oscillatory and assume condition (5) holds. Then all nonoscillatory solutions of (3) are bounded. In fact, if u(t) is a nonoscillatory solution of (3), then

$$\lim_{t\to\infty} |u(t)| = \gamma \le (\lambda)^{1/2n}.$$

THEOREM 4. Assume (4)_{λ} is oscillatory for all $\lambda > 0$ and assume condition (5) holds. Then all solutions of (1) oscillate.

PROOF. If not, let u(t) be a nonoscillatory solution of (1). Lemma 1 and Theorem 2 imply u(t) satisfies u(t)u'(t) > 0 for all large t and $\lim_{t\to\infty} g(u(t)) = 0$. But this is a contradiction since $d(g(u(t)))/dt \ge 0$ by (2). This proves the theorem.

Consider the following somewhat weaker condition than (5): There exists a sequence $T_n \to +\infty$ such that

(5*)
$$\int_{T_n}^t p(s)ds \ge 0, \qquad t \ge T_n.$$

The proof of Lemma 1 and Theorem 2 imply

COROLLARY 5. If p(t) satisfies condition (5*), and if (4)_{λ} is oscillatory for all $\lambda > 0$, then u'(t) oscillates for all solutions u(t) of (1).

EXAMPLES. Willett [3] has shown that $(4)_{\lambda}$ is oscillatory for all $\lambda > 0$ where $p(t) = t^{\eta} \sin t$ and $\eta > -1$. Thus, Corollary 5 implies that u'(t) oscillates for all solutions u(t) of (1) if $-1 < \eta \le 0$.

For the equation

(12)
$$x'' + (\rho t^{-2} + \mu t^{-1} \sin \nu t)x = 0$$

results in [3] imply oscillation if $\rho > \frac{1}{4} - \frac{1}{2}(\mu/\nu)^2$ and nonoscillation if $\rho < \frac{1}{4} - \frac{1}{2}(\mu/\nu)^2$. Moreover, p(t) satisfies condition (5) if $\rho > \mu/\nu \ge 0$. Letting $\mu = \nu$ and $\rho > 1$ it follows that $x'' + \lambda p(t)x = 0$ is oscillatory if $\lambda > \lambda_0 \equiv (\rho^2 + \frac{1}{2})^{1/2} - \rho$ so that all nonoscillatory solutions of (3) satisfy $|u(t)| \le (\lambda_0)^{1/2n}$ for all large t by Corollary 3.

REFERENCES

- 1. James S. W. Wong, On second order nonlinear oscillation, Funkcial. Ekvac. 11 (1969), 207-234.
- 2. D. Willett, Classification of second order linear differential equations with respect to oscillation, Advances in Math. 3 (1969), 594-623.
- 3. ——, On the oscillatory behavior of the solutions of second order linear differential equations, Ann. Polon. Math. 21 (1969), 175-194.
- **4.** W. R. Utz, Properties of solutions of $u'' + g(t)u^{2n-1} = 0$, Monatsh. Math. **66** (1962), 55-60. MR **25** #2275.
- 5. Paul Waltman, An oscillation criterion for a nonlinear second order equation, J. Math. Anal. Appl. 10 (1965), 439-441. MR 30 #3265.
- 6. A. Wintner, A criterion of oscillatory stability, Quart. Appl. Math. 7 (1949), 115-117. MR 10, 456.

University of Alberta