# ORDER RELATION IN QUADRATIC JORDAN RINGS AND A STRUCTURE THEOREM 

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#### Abstract

It is shown that the relation defined by $x \leq y$ if and only if $V_{x} x=V_{x} y$ and $U_{x} x=U_{x} y=U_{y} x$ is an order relation for quadratic Jordan algebras without nilpotent elements, which extends our previous one for linear Jordan algebras, and reduces to the usual Abian order for associative algebras. We prove that a quadratic Jordan algebra is isomorphic to a direct product of division algebras if and only if the algebra has no nilpotent elements and is hyperatomic and orthogonally complete.


Preliminaries. Throughout this paper, $\Phi$ will denote a commutative associative ring with identity element and $J=(J, U, 1)$ a quadratic unital Jordan algebra over $\Phi$, that is, $J$ is a $\Phi$-module, 1 an element of $J$, and $U_{x} y$ is a product quadratic in $x$ and linear in $y$ which satisfies certain axioms analogous to properties of the product $x y x$ in associative algebras. We use the notation $V_{x, y} z=U_{x, z} y=\{x y z\}$ for the polarization $U_{x, z}=U_{x+z}-U_{x}-U_{z}, V_{x} y=U_{x, 1} y=U_{x, y} 1=x \circ y$.

If $A$ is a unital associative algebra, then $A^{+}$becomes a quadratic Jordan algebra via

$$
U_{x} y=x y x, \quad V_{x, y} z=x y z+z y x, \quad V_{x} y=x y+y x
$$

Jordan powers $x^{n}$ can be defined recursively (corresponding to the usual associative powers in $A^{+}$); we will need only powers $x^{2}=U_{x} 1$ and $x^{3}=U_{x} x . J$ has no nilpotent elements ( $x^{n}=0$ for some $n$ forces $x=0$ ) if and only if $x^{2}=0$ or $x^{3}=0$ forces $x=0$. MacDonald's principle asserts that any Jordan relation $f(x, y, z)=0$ in three variables which is of degree at most 1 in $z$, and which holds in all associative algebras $A^{+}$, necessarily holds in all Jordan algebras. Any Jordan identity that we use will be a consequence of this principle.

## 1. Orthogonality.

DEFINITION. $x_{1}$ and $x_{2}$ are orthogonal, $x_{1} \perp x_{2}$, if $x_{1} \circ x_{2}=U_{x_{1}} x_{2}=U_{x_{2}} x_{1}=0$.
This relation is symmetric in $x_{1}$ and $x_{2}$. Zero is orthogonal to everything; if $J$ has no nilpotent elements, then $x \perp x$ forces $x=0$. If $J \subset A^{+}$is special, then associative orthogonality $x_{1} x_{2}=x_{2} x_{1}=0$ is sufficient (and, if $A$ has no nilpotent elements, necessary: $\left(x_{i} x_{j}\right)^{2}=x_{i}\left(U_{x_{j}} x_{i}\right)=0$ forces $\left.x_{i} x_{j}=0\right)$ for Jordan orthogonality.

[^0]Lemma. If $x_{1} \perp x_{2}$ and $J$ has no nilpotent elements, then for $i \neq j$
(i) $x_{j}^{2} \perp x_{i}$,
(ii) $U_{x_{i}} U_{x_{j}}=V_{x_{i}, x_{j}}=0$ on $J$,
(iii) $x_{i}^{\perp}=\left\{y \in J \mid y \perp x_{i}\right\}$ is a linear subspace closed under squares,
(iv) $y \perp x_{1}+x_{2} \Rightarrow y \perp x_{1}, y \perp x_{2}$.

Proof. (i)

$$
\begin{aligned}
& x_{i} \circ x_{j}^{2}=\left(x_{i} \circ x_{j}\right) \circ x_{j}-2 U_{x_{j}} x_{i} \quad(\text { by MacDonald's principle })=0, \\
& U_{x_{j}}^{2} x_{i}=U_{x_{j}}\left(U_{x_{j}} x_{i}\right) \quad(\text { MacDonald })=0, \\
&\left(U_{x_{i}} x_{j}^{2}\right)^{2}= U_{x_{i}}\left\{U_{x_{j}}\left[\left(x_{i} \circ x_{j}\right)^{2}-x_{i} \circ U_{x_{j}} x_{i}\right]-\left(U_{x_{j}} x_{i}\right)^{2}\right\} \quad \text { (MacDonald) } \\
&=
\end{aligned}
$$

forces $U_{x_{i}} x_{j}^{2}=0$ by absence of nilpotents.

$$
\begin{equation*}
\left(U_{x_{i}} U_{x_{j}} a\right)^{2}=U_{x_{i}} U_{x_{j}} U_{a}\left(U_{x_{j}} x_{i}^{2}\right)=0 \quad \text { by (i) } \tag{ii}
\end{equation*}
$$

(from $\left(U_{x} y\right)^{2}=U_{x} U_{y} x^{2}$ by MacDonald),

$$
\begin{aligned}
& \left(V_{x_{i}, x_{j}} a\right)^{2}=U_{x_{i}} U_{x_{j}} a^{2}+U_{a} U_{x_{j}} x_{i}^{2}+U_{x_{i}, a} U_{x_{j}}\left(x_{i} \circ a\right)-U_{x_{i}} x_{j} \circ U_{a} x_{j} \\
& \text { (follows from linearizing }\left(U_{x} x_{j}\right)^{2}=\ldots \text { ) } \\
& =0+0+U_{x_{i}, a}\left[\left\{x_{j} \circ x_{i} a x_{j}\right\}-x_{i} \circ U_{x_{j}} a\right]-0 \quad \text { (MacDonald) } \\
& =-\left\{x_{i} x_{i} \circ U_{x_{j}} a a\right\}=-\left\{x_{i}^{2}\left(U_{x_{j}} a\right) a\right\}-U_{x_{i}}\left(U_{x_{j}} a\right) \circ a \quad \text { (MacDonald) } \\
& =-\left\{x_{i}^{2} x_{j} U_{a} x_{j}\right\}-0 \quad \text { (MacDonald) } \\
& =-\left\{x_{i} x_{i} \circ x_{j}\left(U_{a} x_{j}\right)\right\}+U_{x_{i}} x_{j} \circ\left(U_{a} x_{j}\right) \quad \text { (MacDonald) } \\
& =0 \text {. }
\end{aligned}
$$

(iii) If $y_{1}, y_{2} \perp x$ then clearly $\Phi y_{1} \perp x$, and $y_{1}+y_{2} \perp x$ since

$$
\begin{aligned}
\left(U_{y_{1}+y_{2}} x\right)^{2} & =\left(U_{y_{1}, y_{2}} x\right)^{2} \\
& =U_{y_{1}}\left(U_{x} y_{2}^{2}\right)+U_{y_{2}}\left(U_{x} y_{1}^{2}\right)+U_{y_{1}, y_{2}} U_{x}\left(y_{1} \circ y_{2}\right)-\left(U_{y_{1}} x\right) \circ\left(U_{y_{2}} x\right) \\
& \left.\quad \text { (linearizing }\left(U_{y} x\right)^{2}=\ldots\right) \\
& =U_{y_{1}, y_{2}}\left[\left\{x \circ y_{1} y_{2} x\right\}-y_{1} \circ U_{x} y_{2}\right] \quad \text { (MacDonald) } \\
& =0 .
\end{aligned}
$$

(iv) We have $U_{x_{i}} y=0$ since

$$
U_{x_{i}} y=U_{\left(x_{i}+x_{j}\right)-x_{j}} y=U_{x_{j}} y \quad \text { (by } x_{i}+x_{j} \perp y \text { and (ii) }
$$

and

$$
\left(U_{x_{i}} y\right)^{3}=\left(U_{x_{i}} U_{y} U_{x_{i}}\right)\left(U_{x_{i}} y\right) \quad(\text { MacDonald })=U_{x_{i}} U_{y}\left(U_{x_{i}} U_{x_{j}} y\right)=0 \quad \text { by }(\mathrm{ii}) ;
$$

then $U_{x_{i}} y^{2}=0$ since by (iii) $y^{2} \perp x_{i}+x_{j}$, hence $U_{y} x_{i}=U_{y} x_{i}^{2}=0$ since $\left(U_{y} x_{i}\right)^{2}=$ $U_{y}\left(U_{x_{i}} y^{2}\right)=0$ and $\left(U_{y} x_{i}^{2}\right)^{2}=U_{y} U_{x_{i}}\left(U_{x_{i}} y^{2}\right)=0$ (as we just saw), and finally $x_{i} \circ y=0$ since

$$
\left(x_{i} \circ y\right)^{2}=y \circ U_{x_{i}} y+U_{x_{i}} y^{2}+U_{y} x_{i}^{2} \quad(\text { MacDonald })=0
$$

by the previous.

## 2. Ordering.

DEFINITION. $x \leq y$ if $x-y \perp x$, i.e. if $y=x+x^{\prime}$ for $x^{\prime} \perp x$. By definition of orthogonality this means
(0i) $x \circ(x-y)=0$,
(0ii) $U_{x}(x-y)=0$,
(0iii) $U_{x-y} x=0$ and in the presence of ( 0 i ), (0ii) the latter is equivalent to $U_{x} x=U_{y} x$ (since $U_{x-y} x+\left(U_{x} y-U_{y} x\right)=\left(V_{x}^{2}-2 U_{x}\right)(x-y)($ MacDonald $\left.)=0\right)$, so $x-y \perp x$ iff
(0i) $V_{x} x=V_{x} y$,
(0ii) $U_{x} x=U_{x} y$,
(0iii) $U_{x} x=U_{y} x$. This agrees with our definition in (5) for linear Jordan algebras where $\frac{1}{2} \in \Phi$,
(Li) $x \cdot x=x \cdot y$,
(Lii) $x^{2} \cdot x=x^{2} \cdot y$,
(Liii) $x \cdot x^{2}=x \cdot y^{2}$, since $a \circ b=2 a \cdot b, U_{a} b=2 a \cdot(a \cdot b)-a^{2} \cdot b$. For $J=A^{+}$ with $A$ an associative algebra without nilpotent elements this ordering agrees with the Abian ordering in the associative algebra $A$

$$
x \leq y \Leftrightarrow x y=y x=x^{2}
$$

( $\Leftarrow$ is clear, for $\Rightarrow$ note $x-y \perp x \Rightarrow(x-y) x=x(x-y)=0$ (by our remark after the definition of orthogonality)). Note

$$
x \leq y \Rightarrow y^{\perp} \subset x^{\perp} \text { (since } z \perp y=x+x^{\prime} \Rightarrow z \perp x \text { by Lemma (iv)). }
$$

THEOREM 1. The relation $\leq$ is a partial ordering for a quadratic Jordan algebra $J$ if and only if $J$ has no nilpotent elements.

Proof. If $J$ has nilpotent elements it has $x \neq 0$ with $x^{2}=0$, so $0 \leq x \leq 0$ and $\leq$ is not antisymmetric. Now assume that $J$ has no nilpotent elements (so the Lemma applies). Clearly $\leq$ is reflexive ( $0 \perp x$ ); it is antisymmetric since $x \leq y \leq x \Rightarrow x-y \perp x, y \Rightarrow x-y \perp x-y$ (by Lemma (iii)) $\Rightarrow x-y=0$; it is transitive since $x \leq y \leq z \Rightarrow x-y \in x^{\perp}, y-z \in y^{\perp} \subset x^{\perp}$ (by (^)), so $x-z=(x-y)+(y-z) \in x^{\perp}$ (by Lemma (iii)) and $x \leq z$. Note in particular that our hypothesis (P), $(x, x, y)=0 \Rightarrow(x y, x, y)=0$, of (5) is superfluous, answering the question raised on [5, p. 383].

We remark that if $x \leq y$ then $p(x, y)=p(x, x)$ for all Jordan polynomials in $x$ and $y$ all of whose monomials have degree at least one in both $x$ and $y$. We also remark that if $J$ has no nilpotents then $x_{1} \perp x_{2}$ iff $U_{x_{1}} U_{x_{2}}=0$ on $J(\Rightarrow$ from Lemma (ii), $\Leftarrow$ since $U_{x_{i}} x_{j}=0$ (since $\left(U_{x_{i}} x_{j}\right)^{2}=U_{x_{i}} U_{x_{j}} U_{x_{i}} 1=0$ ), similarly $U_{x_{i}} x_{j}^{2}=0$, and so $x_{1} \circ x_{2}=0$ (since $\left(x_{1} \circ x_{2}\right)^{2}=x_{1} \circ U_{x_{2}} x_{1}+U_{x_{2}} x_{1}^{2}+U_{x_{1}} x_{2}^{2}=0$ ).

## 3. A structure theorem.

DEFINITION. $J$ is orthogonally complete if every family $\left\{x_{\alpha}\right\}$ of pairwise orthogonal elements has a supremum relative to the ordering $\leq$.

Definition. An element $e \in J$ is a hyperatom if it is a central division idempotent (so has Peirce decomposition $J=J+1 \oplus J_{0}$ for ideal $J_{i}$ with $J_{1}=U_{e} J$ a Jordan division algebra (in the sense that all operators $U_{x_{1}}$ are surjective on $J_{1}$, hence bijective there)).

Definition. $J$ is hyperatomic if for every $a \neq 0$ in $J$ there is a hyperatom $e$ with $U_{e} a \neq 0$.

STRUCTURE ThEOREM. $J$ is isomorphic to a direct product of Jordan division algebras if and only if $J$ has no nilpotent elements, is hyperatomic and orthogonally complete with respect to $\leq$.

Proof. Suppose first that $J \cong \pi J_{i}$ for division algebras $J_{i}$ with units $e_{i}$. Clearly $J$ has no nilpotent elements, and is hyperatomic since any $x \neq 0$ has some $x_{i}=U_{e_{i}} \neq 0$ where the $e_{i}$ are hyperatoms. To see that $J$ is orthogonally complete, if $x=\pi x_{i}, y=\pi y_{i}$ are orthogonal then $x_{i} \neq 0$ implies $y_{i}=0$ (since $U_{x} y=0$ by orthogonality forces $U_{x_{i}} y_{i}=0$ in the division algebra $J_{i}$ ), so $x$ and $y$ have disjoint support (support being the set of indices $i$ for which $x_{i} \neq 0$ ). If $\left\{x_{\alpha}\right\}$ is an orthogonal family then the support sets $S_{\alpha}=\operatorname{support}\left(x_{\alpha}\right)$ are disjoint, and $x=\pi x_{i}$ is the supremum of the $x_{\alpha}$, where $x_{i}=x_{\alpha i}$ if the index $i$ lies in a (unique) $S_{\alpha}$ and $x_{i}=0$ if $i$ lies in no $S_{\alpha}$.

Conversely, suppose $J$ is hyperatomic with hyperatoms $\left\{e_{i}\right\}$. Then the Peirce decomposition relative to $e_{i}$ is $J=J_{1}\left(e_{i}\right) \oplus J_{0}\left(e_{i}\right)$, and $U_{e_{i}}$ is an algebra homomorphism of $J$ onto the division algebra $J_{i}=U_{e_{i}} J$, and $\theta(x)=\pi\left(U_{e_{i}} x\right)$ is an algebra homomorphism of $J$ into $\pi J_{i}$. This is injective since $\theta(x)=0$ implies $U_{e_{i}} x=0$ for each $i$, hence $x=0$ by definition of hyperatomic. If $J$ has no nilpotent elements $\leq$ is a partial order, and if $J$ is orthogonally complete $\theta$ is surjective: given any family $x_{i} \in J_{i}$ their supremum $x$ must be the element such that $\theta(x)=\pi x_{i}$, since $U_{e_{j}} x=x_{j}$ for each $j: U_{e_{j}} x=\sup U_{e_{j}} x_{i}=U_{e_{j}} x_{j}=x_{j}\left(U_{e_{j}} x_{i}=U_{e_{j}} U_{e_{i}} x_{i}=0\right.$ for $i \neq j$ by Lemma (ii) since $e_{i} \perp e_{j}$ ).

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