

## AN OSCILLATION THEOREM FOR SECOND ORDER SUBLINEAR DIFFERENTIAL EQUATIONS

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**ABSTRACT.** An oscillation criterion is given for the second order sublinear differential equation  $x'' + a(t)|x|^\gamma \operatorname{sgn} x = 0$ ,  $0 < \gamma < 1$ , where the coefficient  $a(t)$  is not assumed to be nonnegative for all large values of  $t$ . The result extends a condition recently discovered by Butler, Erbe, and Mingarelli for the linear equation.

Consider the second order sublinear differential equation

$$(1) \quad x'' + a(t)|x|^\gamma \operatorname{sgn} x = 0, \quad 0 < \gamma < 1,$$

where  $a(t) \in C[0, \infty)$ . We restrict our attention to solutions of (1) which exist on some ray  $[t_0, \infty)$ , where  $t_0 \geq 0$  may depend on the particular solution. Such a solution is said to be oscillatory if it has arbitrarily large zeros. Equation (1) is called oscillatory if all such solutions are oscillatory. For a general discussion on sublinear oscillation problems, we refer the reader to [11]. We are here concerned with sufficient conditions on  $a(t)$  for the oscillation of (1) when  $a(t)$  is allowed to assume negative values for arbitrarily large values of  $t$ . In the linear case, the well-known Wintner's oscillation theorem states that if  $a(t)$  satisfies

$$(2) \quad \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \int_0^t a(s) ds dt = +\infty,$$

then equation (1) is oscillatory for  $\gamma = 1$ , see [10]. Butler [1] proved that Wintner's theorem remains valid for equation (1) where  $\gamma > 1$ . In the sublinear case, i.e.,  $0 < \gamma < 1$ , condition (2) can be relaxed to

$$(3) \quad \limsup_{T \rightarrow \infty} \frac{1}{T} \int_0^T \int_0^t a(s) ds dt = +\infty,$$

an earlier result due to Kamenev [6]. We note that Condition (3) alone is not sufficient for oscillation for  $\gamma \geq 1$ ; in the linear case, see Willett [9].

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Wintner's condition (2) was later improved by Hartman [3], see also [4], who proved that the following conditions, i.e.,

$$(4) \quad \liminf_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(t) dt > -\infty,$$

where  $A(t) = \int_0^t a(s) ds$  and that the limit in (2) does not exist, namely,

$$(5) \quad \underline{\lim}_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(t) dt < \overline{\lim}_{T \rightarrow \infty} \frac{1}{T} \int_0^T A(t) dt,$$

are sufficient for oscillation of the linear equation. In particular, conditions (3) and (4) form an oscillation criterion for equation (1) with  $\gamma = 1$ . Butler [1] also proved that in the sublinear case, i.e., equation (1), condition (5) alone will suffice, (see also Kwong and Wong [8, Theorem 6], and Wong [14] for alternative proofs of this result). Recently in connection with the study of oscillation theory for linear systems, Butler, Erbe, and Mingarelli [2] showed that condition (4), together with

$$(6) \quad \limsup_{T \rightarrow \infty} \frac{1}{T} \int_0^T A^2(t) dt = +\infty,$$

are sufficient for oscillation of equation (1) in the linear case. Since (3) implies (6) by an application of Schwarz's inequality, their result extends part of Hartman's theorem, i.e., condition (3) plus (4), and is hence a new oscillation criterion thus far undiscovered. It is therefore natural to ask whether this new oscillation criterion remains valid for the sublinear equation (1). The purpose of this note is to answer this question in the affirmative by proving

**Theorem.** *Suppose that conditions (4) and (6) hold. Then equation (1) is oscillatory for  $0 < \gamma < 1$ .*

*Proof.* Assume to the contrary that there exists a nonoscillatory solution  $x(t)$ , which may be assumed to be positive on  $[t_0, \infty)$ . For  $0 < \gamma < 1$ , define  $y(t) = x^{1-\gamma}(t)$ . It is easy to verify from (1) that  $y(t)$  satisfies the second order nonlinear differential equation

$$(7) \quad y'' + (1-\gamma)a + \gamma(1-\gamma)^{-1}y^{-1}y'^2 = 0,$$

on  $[t_0, \infty)$ . For convenience, denote the positive numbers  $\alpha = 1 - \gamma$ ,  $\beta = \gamma(1 - \gamma)^{-1}$ . Integrating (7) once obtains

$$(8) \quad y'(t) - y'(t_0) + \alpha \int_{t_0}^t a(s) ds + \beta \int_{t_0}^t y^{-1}(s)y'^2(s) ds = 0.$$

A further integration yields

$$(9) \quad y(t) - y(t_0) + \beta \int_{t_0}^t \int_{t_0}^s y^{-1}y'^2 = C_0(t - t_0) - \alpha \int_{t_0}^t A(s) ds,$$

where  $C_0 = y'(t_0) + \alpha A(t_0)$ .

We distinguish two mutually exclusive cases, namely (i)  $y^{-1}y'^2 \notin L^1(T_0, \infty)$  and (ii)  $y^{-1}y'^2 \in L^1(t_0, \infty)$ , and verify in each case that the existence of a positive solution leads to a contradiction. If  $y^{-1}y'^2 \notin L^1(t_0, \infty)$ , then the weighted average also diverges, i.e.,

$$(10) \quad \lim_{t \rightarrow \infty} \frac{1}{t} \int_{t_0}^t \int_{t_0}^s y^{-1}y'^2 = +\infty,$$

Dividing (9) by  $t$ , noting that  $y(t)$  is positive and using (10) and condition (4), we easily arrive at a desired contradiction.

We now turn to the case when  $y^{-1}y'^2 \in L^1(t_0, \infty)$ , and first show that

$$(11) \quad \lim_{t \rightarrow \infty} \frac{y(t)}{t} = 0.$$

Let  $\varepsilon > 0$ , choose  $t_1 \geq t_0$  such that  $\int_{t_1}^\infty y^{-1}y'^2 < \frac{\varepsilon}{4}$ . By Schwarz's inequality,

$$(12) \quad y(t) - y(t_1) = \int_{t_1}^t y' \leq \left( \int_{t_1}^t y^{-1}y'^2 \right)^{1/2} y \left( \int_{t_1}^t y \right)^{1/2} \leq \frac{\sqrt{\varepsilon}}{2} \left( \int_{t_1}^t y \right)^{1/2} y.$$

Suppose that  $y \in L^1(t_1, \infty)$ . Then (12) implies  $y(t)$  is bounded, hence (11) follows. Thus we assume that  $y \notin L^1(t_1, \infty)$ , so one can choose  $t_2 \geq t_1$  so that  $y^2(t_1) \leq \frac{\varepsilon}{4} \int_{t_1}^t y$ , for  $t \geq t_2$ . Using this fact in (12),

$$(13) \quad y(t) \leq y(t_1) + \frac{\sqrt{\varepsilon}}{2} \left( \int_{t_1}^t y \right)^{1/2} y \leq \left( \frac{\sqrt{\varepsilon}}{2} + \frac{\sqrt{\varepsilon}}{2} \right) \left( \int_{t_1}^t y \right)^{1/2} y.$$

Dividing (13) through by  $\left( \int_{t_1}^t y \right)^{1/2} y$  and integrating from  $t_2$  to  $t$  obtains

$$(14) \quad \left( \int_{t_1}^t y \right)^{1/2} y - \left( \int_{t_1}^{t_2} y \right)^{1/2} y \leq \frac{\sqrt{\varepsilon}}{2}(t - t_2) \leq \frac{\sqrt{\varepsilon}}{2}t.$$

Now choose  $t_3 \geq t_2$  such that  $\int_{t_1}^{t_3} y \leq \frac{\varepsilon}{4}t_3^2$ , hence, for  $t \geq t_3$ , we have from (14),

$$(15) \quad \left( \int_{t_1}^t y \right)^{1/2} y \leq \frac{\sqrt{\varepsilon}}{2}t + \frac{\sqrt{\varepsilon}}{2}t_3 \leq \left( \frac{\sqrt{\varepsilon}}{2} + \frac{\sqrt{\varepsilon}}{2} \right) t.$$

Combining (13) and (15), we conclude that  $y(t) \leq \varepsilon t$  for  $t \geq t_3$ , proving (11).

Returning to (8), one can express  $A(t)$  in the form

$$(16) \quad \alpha A(t) = y'(t_0) - y'(t) + \alpha A(t_0) - \beta \int_{t_0}^t y^{-1}y'^2.$$

Recalling that  $C_0 = y'(t_0) + \alpha A(t_0)$ , one can use (16) to estimate  $A^2(t)$  by

$$(17) \quad \alpha^2 A^2(t) \leq 3C_0^2 + 3y'^2(t) + 3\beta^2 \int_{t_0}^t y^{-1}y'^2 \leq 3C_1^2 + 3y'^2(t),$$

where  $C_1^2 = C_0^2 + \beta^2 \int_{t_0}^{\infty} y^{-1} y'^2$ . Integrating (17) from  $t_0$  to  $T$  and dividing through by  $T$ , we obtain

$$(18) \quad \frac{\alpha^2}{T} \int_{t_0}^T A^2(t) dt \leq 3C_1^2 \left(1 - \frac{t_0}{T}\right) + \frac{3}{T} \int_{t_0}^T y'^2.$$

Finally, we note

$$(19) \quad \frac{1}{T} \int_{t_0}^T y'^2 y^{-1} y \leq \frac{1}{T} \max_{t_0 \leq t \leq T} |y(t)| \int_{t_0}^T y^{-1} y'^2.$$

By (11) we can choose  $T_0 \geq t_0$  such that  $|y(t)| \leq t$  for  $t \geq T_0$ . Using this,

$$(20) \quad \max_{t_0 \leq t \leq T} |y(t)| \leq \max_{t_0 \leq t \leq T_0} |y(t)| + T \leq M_0 + T,$$

where  $M_0$  is a constant of  $T$ . Substituting (19) and (20) in (18), we find the right hand side is bounded as  $T$  tends to infinity, which is incompatible with condition (6).  $\square$

*Remark 1.* It is easy to give an example of  $a(t)$  which satisfies conditions (4) and (6) but fails to satisfy conditions (3) and (4). Take  $A(t) = t^\lambda \sin t$ ,  $1/2 < \lambda < 1$ . We note in passing that condition (5) alone suffices for the oscillation of equation (1) in the sublinear case, so does condition (3). Nevertheless, the example

$$a(t) = t^\lambda \cos t + \lambda t^{\lambda-1} \sin t, \quad \frac{1}{2} < \lambda < 1$$

does not satisfy either (3) or (5), but our theorem applies.

*Remark 2.* We note that the weighted average used here is the simple arithmetic mean used by Wintner [10] in condition (2). Hartman [5] showed that some of his results, which include that of Wintner, remain valid for a much wider class of "general means". It would be interesting to investigate to what extent oscillation criteria with more general weighted means for the linear equation can be extended to the sublinear case. Some of the oscillation theorems using weighted means by Willett [9] have been successfully extended by Butler [1].

*Remark 3.* Condition (3) has been extended by Kamenev [7] to the weaker condition for the linear equation, i.e., for some  $\mu > 1$ ,

$$(21) \quad \limsup_{T \rightarrow \infty} \frac{1}{T^\mu} \int_0^T (T-t)^\mu a(t) dt = +\infty.$$

It is known that (21) is also sufficient for the sublinear equation, see [13].

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