

WILLMORE INEQUALITY ON HYPERSURFACES IN HYPERBOLIC SPACE

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(Communicated by Lei Ni)

ABSTRACT. In this article, we prove a geometric inequality for star-shaped and mean-convex hypersurfaces in hyperbolic space by inverse mean curvature flow. This inequality can be considered as a generalization of Willmore inequality for a closed surface in hyperbolic 3-space.

1. INTRODUCTION

The classical isoperimetric inequality and its generalization, the Alexandrov-Fenchel inequalities, play an important role in different branches of geometry. Let $\Omega \subset \mathbb{R}^n$ be a smooth bounded domain with boundary Σ . Then the classical isoperimetric inequality is

$$(1.1) \quad |\Sigma| \geq n \frac{n-1}{n} \omega_{n-1}^{\frac{1}{n}} |\Omega|^{\frac{n-1}{n}},$$

and equality in (1.1) holds if and only if Ω is a geodesic ball.

For $k \in \{1, \dots, n-1\}$, we denote by p_k the normalized k -th order mean curvature of Σ , and set $p_0 = 1$ by convention. The celebrated Alexandrov-Fenchel inequalities [1, 2, 15] for convex hypersurface $\Sigma^{n-1} \subset \mathbb{R}^n$ are

$$(1.2) \quad \frac{1}{\omega_{n-1}} \int_{\Sigma} p_k d\mu \geq \left(\frac{1}{\omega_{n-1}} \int_{\Sigma} p_j d\mu \right)^{\frac{n-1-k}{n-1-j}}, \quad 0 \leq j < k \leq n-1,$$

and equality in (1.2) holds if and only if Σ is a geodesic sphere.

Since the isoperimetric inequality holds for non-convex domains, it is natural to extend the original Alexandrov-Fenchel inequality to non-convex domains; see [7–9, 20, 21, 28, 31], etc. We should also mention that the Willmore inequality, which is a weaker form of Alexandrov-Fenchel inequality, has been established for closed surfaces in \mathbb{R}^3 ; see e.g. [10, 26, 30]. More precisely, for any closed surface $\Sigma \subset \mathbb{R}^3$, the Willmore inequality is

$$(1.3) \quad \int_{\Sigma} p_1^2 d\mu \geq \omega_2 = 4\pi,$$

and equality in (1.3) holds if and only if Σ is a geodesic sphere.

It is interesting to establish the Alexandrov-Fenchel inequalities for hypersurfaces in hyperbolic space; see [4, 16]. Recently, the following hyperbolic Alexandrov-Fenchel inequalities were obtained by Ge-Wang-Wu [17, 18] and Wang-Xia [32].

Received by the editors November 27, 2016, and, in revised form, September 13, 2017.

2010 *Mathematics Subject Classification*. Primary 53C42, 53C44.

Key words and phrases. Willmore inequality, inverse curvature flow, hyperbolic space.

Theorem A ([17, 18, 32]). *Let $k \in \{1, \dots, n - 1\}$. Any horospherical convex hypersurface $\Sigma \subset \mathbb{H}^n$ satisfies*

$$(1.4) \quad \int_{\Sigma} p_k d\mu \geq \omega_{n-1} \left[\left(\frac{|\Sigma|}{\omega_{n-1}} \right)^{\frac{2}{k}} + \left(\frac{|\Sigma|}{\omega_{n-1}} \right)^{\frac{2(n-1-k)}{k(n-1)}} \right]^{\frac{k}{2}}.$$

Equality in (1.4) holds if and only if Σ is a geodesic sphere.

Inequality (1.4) was proved in [17] for $k = 4$ and in [18] for general even k . For $k = 1$, (1.4) was proved in [18] with a help of a result of Cheng and Zhou [12]. For general integer k , (1.4) was proved in [32].

For $k = 2$, inequality (1.4) was proved by Li-Wei-Xiong [25] under a weaker condition that Σ is star-shaped and 2-convex (i.e., $p_2 > 0$). More precisely,

Theorem B ([25]). *Any star-shaped and 2-convex hypersurface $\Sigma \subset \mathbb{H}^n$ ($n \geq 3$) satisfies*

$$(1.5) \quad \int_{\Sigma} p_2 d\mu \geq \omega_{\frac{n-1}{2}} |\Sigma|^{\frac{n-3}{n-1}} + |\Sigma|.$$

Equality in (1.5) holds if and only if Σ is a geodesic sphere.

The Willmore inequality (1.3) has also been generalized to closed surface $\Sigma \subset \mathbb{H}^3$; see e.g. [11, 27, 29, 33].

Theorem C ([11, 27, 29, 33]). *Any closed surface $\Sigma \subset \mathbb{H}^3$ satisfies*

$$(1.6) \quad \int_{\Sigma} (p_1^2 - 1) d\mu \geq \omega_2 = 4\pi.$$

Equality in (1.6) holds if and only if Σ is a geodesic sphere.

To generalize the hyperbolic Willmore inequality to higher dimension, the positivity of the functional $\int_{\Sigma} (p_1^2 - 1) d\mu$ has already been known. This follows from the optimal Reilly inequality for submanifolds of hyperbolic space, which was achieved by El Soufi and Ilias [14].

Theorem D ([14]). *Let (M^m, g) ($m \geq 2$) be a compact and connected Riemannian manifold isometrically immersed in \mathbb{H}^n by ϕ . Then*

$$(1.7) \quad \lambda_1(M) \leq \frac{m}{|M|} \int_M (|H|^2 - 1) d\mu,$$

where $\lambda_1(M)$ is the first non-zero eigenvalue of Laplacian of (M, g) and H is the mean curvature vector of M . Furthermore, equality in (1.7) holds if and only if $\phi(M)$ is minimally immersed in a geodesic sphere of radius $\sinh^{-1}(\sqrt{\frac{m}{\lambda_1(M)}})$.

Inspired by these previous results, we prove the Willmore inequality for star-shaped and mean-convex (i.e., $p_1 > 0$) hypersurfaces in hyperbolic space.

Theorem 1. *Any star-shaped and mean-convex hypersurface $\Sigma \subset \mathbb{H}^n$ ($n \geq 3$) satisfies*

$$(1.8) \quad \int_{\Sigma} (p_1^2 - 1) d\mu \geq \omega_{\frac{n-1}{2}} |\Sigma|^{\frac{n-3}{n-1}}.$$

Equality in (1.8) holds if and only if Σ is a geodesic sphere.

We expect that the inequality (1.8) will be useful in defining the Hawking mass for hypersurfaces in \mathbb{H}^n . In [24], the *Hawking mass* for a closed embedded surface Σ in \mathbb{H}^3 is defined as

$$m_H(\Sigma) = \sqrt{\frac{|\Sigma|}{16\pi}} \left[1 - \frac{1}{4\pi} \int_{\Sigma} (p_1^2 - 1) d\mu \right].$$

We now give the outline of the proof of Theorem 1. In fact, we prove it in a more general setting. Motivated by [6, 13, 25], we adopt the inverse curvature flow (ICF)

$$\partial_t X = \frac{1}{(n-1)p_k^{1/k}} \nu$$

in our proof. When $k = 1$, this flow is just inverse mean curvature flow, which has been used by Huisken and Ilmanen [22, 23] to prove the Riemannian Penrose inequality in general relativity. We start from a given star-shaped and k -convex hypersurface Σ in hyperbolic space and evolve it by ICF. By the convergence results of Gerhardts [19], this flow exists for all time, and the evolving hypersurface Σ_t with $\Sigma_0 = \Sigma$ remains star-shaped and k -convex for all $t \geq 0$.

We next consider the quantity

$$Q_k(t) := |\Sigma_t|^{-\frac{n-3}{n-1}} \int_{\Sigma_t} (p_k^{2/k} - 1) d\mu.$$

We study the limit of $Q_k(t)$ as $t \rightarrow \infty$. Notice that the roundness estimate for Σ_t is not strong enough to calculate the limit of $Q_k(t)$. However, similar to [6, 25], we are able to give a positive lower bound for the limit of $Q_k(t)$, which will be used to establish the monotonicity of $Q_k(t)$. Finally, we prove that for $k = 1, 2$, $Q_k(t)$ is monotone decreasing under ICF. From this, Theorem 1 follows immediately.

2. PRELIMINARIES

In this article, we consider the hyperbolic space $\mathbb{H}^n = \mathbb{R}^+ \times \mathbb{S}^{n-1}$ equipped with the metric

$$\bar{g} = dr^2 + \sinh^2 r g_{\mathbb{S}^{n-1}},$$

where $g_{\mathbb{S}^{n-1}}$ is the standard round metric on the unit sphere \mathbb{S}^{n-1} . Let $\Sigma \subset \mathbb{H}^n$ be a closed hypersurface with its unit outward normal vector ν . The second fundamental form h of Σ is defined by

$$h(X, Y) = \langle \bar{\nabla}_X \nu, Y \rangle$$

for any $X, Y \in T\Sigma$. The principal curvature $\kappa = (\kappa_1, \dots, \kappa_{n-1})$ comprises the eigenvalues of h with respect to the induced metric g on Σ . For $k \in \{1, \dots, n-1\}$, the normalized k -th elementary symmetric polynomial of κ is defined as

$$p_k(\kappa) := \frac{1}{\binom{n-1}{k}} \sum_{i_1 < i_2 < \dots < i_k} \kappa_{i_1} \dots \kappa_{i_k},$$

which can also be viewed as a function of the second fundamental form $h_i^j = g^{jk} h_{ki}$. For simplicity, we write p_k for $p_k(\kappa)$. The following lemma can be regarded as a normalized version of Lemma 2 in [25].

Lemma 2. Let $(T_{k-1})_j^i := \frac{\partial p_k}{\partial h_j^i}$ and $(h^2)_i^j := h_i^\ell h_\ell^j$. Then we have

$$\begin{aligned} \sum_{i,j} (T_{k-1})_i^j h_j^i &= k p_k, & \sum_{i,j} (T_{k-1})_i^j \delta_j^i &= k p_{k-1}, \\ \sum_{i,j} (T_{k-1})_i^j (h^2)_j^i &= (n-1)p_1 p_{k-1} - (n-1-k)p_{k+1}. \end{aligned}$$

Moreover, if $\kappa \in \Gamma_k^+$, we have the following Newton-MacLaurin inequalities:

$$(2.1) \quad \frac{p_{k-1} p_{k+1}}{p_k^2} \leq 1, \quad \frac{p_1 p_{k-1}}{p_k} \geq 1,$$

and equality holds in (2.1) at a given point if and only if Σ is umbilical at this point.

We now consider the inverse curvature flow (ICF)

$$(2.2) \quad \partial_t X = \frac{1}{(n-1)p_k^{1/k}} \nu,$$

where $\Sigma_t = X(t, \cdot)$ is a family of hypersurfaces in \mathbb{H}^n , ν is the unit outward normal to $\Sigma_t = X(t, \cdot)$. Let $d\mu_t$ be its area element on Σ_t . By the divergence free property of T_{k-1} , we list the following evolution equations.

Lemma 3. Under ICF (2.2), we have

$$(2.3) \quad \begin{aligned} \partial_t p_k &= -\frac{1}{n-1} \nabla^i \left[(T_{k-1})_i^j \nabla_j \left(\frac{1}{p_k^{1/k}} \right) \right] \\ &\quad - \frac{1}{(n-1)p_k^{1/k}} [(n-1)p_1 p_k - (n-1-k)p_{k+1} - k p_{k-1}], \end{aligned}$$

$$(2.4) \quad \partial_t d\mu = \frac{p_1}{p_k^{1/k}} d\mu.$$

In [19], Gerhardt investigated the inverse curvature flow of star-shaped hypersurfaces in hyperbolic space and proved the following long-time existence and convergence result.

Theorem 4 ([19]). *If the initial hypersurface is star-shaped and k -convex, then the solution for inverse curvature flow (2.2) exists for all time t and preserves the condition of star-shapedness and k -convexity. Moreover, the hypersurfaces become strictly convex exponentially fast and more and more totally umbilical in the sense of*

$$|h_i^j - \delta_i^j| \leq C e^{-\frac{t}{n-1}}, \quad t > 0;$$

i.e., the principal curvatures are uniformly bounded and converge exponentially fast to one.

3. THE ASYMPTOTIC BEHAVIOR OF MONOTONE QUANTITY

We define the quantity

$$Q_k(t) := |\Sigma_t|^{-\frac{n-3}{n-1}} \int_{\Sigma_t} (p_k^{2/k} - 1) d\mu,$$

where $|\Sigma_t|$ is the area of Σ_t . In this section, we estimate the lower bound of the limit of $Q_k(t)$. First of all, we recall Lemma 7 of [25], which is an application of the sharp Sobolev inequality on \mathbb{S}^{n-1} due to Beckner [3].

Lemma 5. *For every positive function f on \mathbb{S}^{n-1} , we have*

$$(3.1) \quad \int_{\mathbb{S}^{n-1}} f^{n-3} dvol_{\mathbb{S}^{n-1}} + \frac{n-3}{n-1} \int_{\mathbb{S}^{n-1}} f^{n-5} |\nabla f|^2 dvol_{\mathbb{S}^{n-1}} \geq \omega_{\frac{n-1}{2}}^{\frac{2}{n-1}} \left(\int_{\mathbb{S}^{n-1}} f^{n-1} dvol_{\mathbb{S}^{n-1}} \right)^{\frac{n-3}{n-1}}.$$

Moreover, equality in (3.1) holds if and only if f is a constant.

Proposition 6. *Under ICF (2.2), we have*

$$(3.2) \quad \liminf_{t \rightarrow \infty} Q_k(t) \geq \omega_{\frac{n-1}{2}}^{\frac{2}{n-1}}.$$

Proof. Recall that star-shaped hypersurfaces can be written as graphs of function $r = r(t, \theta)$, $\theta \in \mathbb{S}^{n-1}$. Denote $\lambda(r) = \sinh(r)$; then $\lambda'(r) = \cosh(r)$. We next define a function $\varphi(\theta) = \Phi(r(\theta))$, where $\Phi(r)$ is a positive function satisfying $\Phi' = \frac{1}{\lambda}$. Let $\theta = \{\theta^j\}$, $j = 1, \dots, n-1$, be a coordinate system on \mathbb{S}^{n-1} and φ_i, φ_{ij} be the covariant derivatives of φ with respect to the metric $g_{\mathbb{S}^{n-1}}$. Define

$$v = \sqrt{1 + |\nabla\varphi|_{\mathbb{S}^{n-1}}^2}.$$

From [19], we know that

$$(3.3) \quad \lambda = O(e^{-\frac{t}{n-1}}), \quad |\nabla\varphi|_{\mathbb{S}^{n-1}} + |\nabla^2\varphi|_{\mathbb{S}^{n-1}} = O(e^{-\frac{t}{n-1}}).$$

Since $\lambda' = \sqrt{1 + \lambda^2}$, we have

$$(3.4) \quad \lambda' = \lambda \left(1 + \frac{1}{2}\lambda^{-2} + O(e^{-\frac{4t}{n-1}}) \right).$$

From (3.3), we also have

$$(3.5) \quad \frac{1}{v} = 1 - \frac{1}{2}|\nabla\varphi|_{\mathbb{S}^{n-1}}^2 + O(e^{-\frac{4t}{n-1}}).$$

In terms of φ , we can express the metric and the second fundamental form of Σ as

$$g_{ij} = \lambda^2(\sigma_{ij} + \varphi_i\varphi_j),$$

$$h_{ij} = \frac{\lambda'}{v\lambda}g_{ij} - \frac{\lambda}{v}\varphi_{ij},$$

where $\sigma_{ij} = g_{\mathbb{S}^{n-1}}(\partial_{\theta^i}, \partial_{\theta^j})$. Denote $a_i = \sum_k \sigma^{ik} \varphi_{ki}$ and note that $\sum_i a_i = \Delta_{\mathbb{S}^{n-1}}\varphi$. By (3.3), the principal curvatures of Σ_t take the following form

$$\kappa_i = \frac{\lambda'}{v\lambda} - \frac{a_i}{v\lambda} + O(e^{-\frac{4t}{n-1}}), \quad i = 1, \dots, n-1.$$

Then we have

$$p_k = \left(\frac{\lambda'}{v\lambda} \right)^k - \frac{k}{n-1} \left(\frac{\lambda'}{v\lambda} \right)^{k-1} \frac{\Delta_{\mathbb{S}^{n-1}}\varphi}{v\lambda} + O(e^{-\frac{4t}{n-1}}).$$

By using (3.4) and (3.5), we get

$$p_k = \left(1 + \frac{k}{2\lambda^2} - \frac{k|\nabla\varphi|_{\mathbb{S}^{n-1}}^2}{2} \right) - \frac{k}{n-1} \frac{\Delta_{\mathbb{S}^{n-1}}\varphi}{\lambda} + O(e^{-\frac{4t}{n-1}}).$$

Hence, we have

$$p_k^{2/k} - 1 = \frac{1}{\lambda^2} - |\nabla\varphi|_{\mathbb{S}^{n-1}}^2 - \frac{2}{n-1} \frac{\Delta_{\mathbb{S}^{n-1}}\varphi}{\lambda} + O(e^{-\frac{4t}{n-1}}).$$

On the other hand,

$$\sqrt{\det g} = \left[\lambda^{n-1} + O(e^{\frac{(n-3)t}{n-1}}) \right] \sqrt{\det g_{\mathbb{S}^{n-1}}}.$$

So we have

$$\begin{aligned} \int_{\Sigma_t} (p_k^{2/k} - 1) d\mu &= \int_{\mathbb{S}^{n-1}} \lambda^{n-1} (p_k^{2/k} - 1) dvol_{\mathbb{S}^{n-1}} + O(e^{\frac{(n-5)t}{n-1}}) \\ &= \int_{\mathbb{S}^{n-1}} (\lambda^{n-3} - \lambda^{n-1} |\nabla\varphi|_{\mathbb{S}^{n-1}}^2) dvol_{\mathbb{S}^{n-1}} \\ &\quad - \frac{2}{n-1} \int_{\mathbb{S}^{n-1}} \lambda^{n-2} \Delta_{\mathbb{S}^{n-1}}\varphi dvol_{\mathbb{S}^{n-1}} + O(e^{\frac{(n-5)t}{n-1}}) \\ &= \int_{\mathbb{S}^{n-1}} (\lambda^{n-3} - \lambda^{n-1} |\nabla\varphi|_{\mathbb{S}^{n-1}}^2) dvol_{\mathbb{S}^{n-1}} \\ &\quad + \frac{2(n-2)}{n-1} \int_{\mathbb{S}^{n-1}} \lambda^{n-3} \langle \nabla\lambda, \nabla\varphi \rangle_{\mathbb{S}^{n-1}} dvol_{\mathbb{S}^{n-1}} + O(e^{\frac{(n-5)t}{n-1}}). \end{aligned}$$

Since $\nabla\lambda = \lambda\lambda'\nabla\varphi$, it follows that $|\nabla\lambda - \lambda^2\nabla\varphi|_{g_{\mathbb{S}^{n-1}}} \leq O(e^{-\frac{t}{n-1}})$. We deduce that

$$(3.6) \quad \int_{\Sigma_t} (p_k^{2/k} - 1) d\mu = \int_{\mathbb{S}^{n-1}} \left(\lambda^{n-3} + \frac{n-3}{n-1} \lambda^{n-5} |\nabla\lambda|^2 \right) dvol_{\mathbb{S}^{n-1}} + O(e^{\frac{(n-5)t}{n-1}}).$$

Moreover,

$$|\Sigma_t|^{\frac{n-3}{n-1}} = \left(\int_{\mathbb{S}^{n-1}} \lambda^{n-1} dvol_{\mathbb{S}^{n-1}} \right)^{\frac{n-3}{n-1}} + O(e^{\frac{(n-5)t}{n-1}}).$$

Using Lemma 5, we achieve

$$\liminf_{t \rightarrow \infty} |\Sigma_t|^{-\frac{n-3}{n-1}} \int_{\Sigma_t} (p_k^{2/k} - 1) d\mu \geq \omega_{\frac{2}{n-1}}.$$

□

4. MONOTONICITY

In this section, we show that for $k = 1, 2$, the quantity $Q_k(t)$ is monotone decreasing under ICF (2.2).

Proposition 7. *Under ICF (2.2), the quantity $Q_k(t)$ is monotone decreasing for $k = 1, 2$. Moreover, $\frac{d}{dt}Q_k(t) = 0$ at some time t if and only if Σ_t is totally umbilical.*

Proof. Under ICF (2.2), by (2.3) we have

$$\begin{aligned}
 & \frac{d}{dt} \int_{\Sigma_t} (p_k^{\frac{2}{k}} - 1) d\mu \\
 &= \int_{\Sigma_t} \frac{2}{k} p_k^{\frac{2}{k}-1} \partial_t p_k + (p_k^{\frac{2}{k}} - 1) p_1 p_k^{-\frac{1}{k}} d\mu \\
 &= -\frac{2}{k(n-1)} \int_{\Sigma_t} p_k^{\frac{2}{k}-1} \nabla^i \left[(T_{k-1})^j_i \nabla_j (p_k^{-\frac{1}{k}}) \right] d\mu \\
 &\quad - \int_{\Sigma_t} \frac{2p_k^{\frac{1}{k}-1}}{k(n-1)} [(n-1)p_1 p_k - (n-1-k)p_{k+1} - kp_{k-1}] d\mu \\
 &\quad + \int_{\Sigma_t} (p_k^{\frac{2}{k}} - 1) p_1 p_k^{-\frac{1}{k}} d\mu \\
 &= -\frac{2}{k^2(n-1)} \left(\frac{2}{k} - 1 \right) \int_{\Sigma_t} p_k^{\frac{1}{k}-3} (T_{k-1})^j_i \nabla^i p_k \nabla_j p_k d\mu \\
 &\quad + \int_{\Sigma_t} \left[p_1 p_k^{-\frac{1}{k}} (p_k^{\frac{2}{k}} - 1) + \frac{2}{k} p_k^{\frac{1}{k}} \left(\frac{p_{k+1}}{p_k} - p_1 \right) - \frac{2}{n-1} p_k^{\frac{1}{k}-1} (p_{k+1} - p_{k-1}) \right] d\mu \\
 &=: I + II.
 \end{aligned}$$

Since $(T_{k-1})^j_i$ is positive definite if $p_k > 0$, we get $I \leq 0$ for $k = 1, 2$. To handle the second term, we analyze it for $k = 1, 2$ separately. By the Newton-MacLaurin inequality, if $p_k > 0$, then $p_1 \geq p_2^{\frac{1}{2}} \geq \dots \geq p_k^{\frac{1}{k}} > 0$.

(i) If $k = 1$, then

$$\begin{aligned}
 II &= \int_{\Sigma_t} \left[(p_1^2 - 1) + 2(p_2 - p_1^2) - \frac{2}{n-1}(p_2 - 1) \right] d\mu \\
 &= \int_{\Sigma_t} \left[\frac{2n-4}{n-1} p_2 - p_1^2 - \frac{n-3}{n-1} \right] d\mu \\
 &\leq \frac{n-3}{n-1} \int_{\Sigma_t} (p_1^2 - 1) d\mu.
 \end{aligned}$$

(ii) If $k = 2$, then

$$\begin{aligned}
 II &= \int_{\Sigma_t} \left[p_1 p_2^{-\frac{1}{2}} (p_2 - 1) + p_2^{\frac{1}{2}} \left(\frac{p_3}{p_2} - p_1 \right) - \frac{2}{n-1} p_2^{-\frac{1}{2}} (p_3 - p_1) \right] d\mu \\
 &= \frac{n-3}{n-1} \int_{\Sigma_t} p_2^{-\frac{1}{2}} (p_3 - p_1) d\mu \\
 &\leq \frac{n-3}{n-1} \int_{\Sigma_t} (p_2 - 1) d\mu.
 \end{aligned}$$

Combining with Proposition 6, we know that the quantity

$$\int_{\Sigma_t} (p_k^{2/k} - 1) d\mu$$

is positive under ICF (2.2). By (2.4) we get

$$\frac{d}{dt} |\Sigma_t| = \int_{\Sigma_t} \frac{p_1}{p_k^{1/k}} d\mu \geq |\Sigma_t|.$$

Therefore, for $k = 1, 2$, we have

$$\frac{d}{dt}Q_k(t) \leq 0.$$

If the equality holds, then the Newton-MacLaurin inequalities assure equalities everywhere on Σ_t . Therefore Σ_t is totally umbilical. \square

Now we complete the proof of Theorem 1.

Proof of Theorem 1. For $k = 1, 2$, since $Q_k(t)$ is monotone decreasing, we have

$$Q_k(0) \geq \liminf_{t \rightarrow \infty} Q_k(t) \geq \omega_{n-1}^{\frac{2}{n-1}}.$$

This implies that $\Sigma_0 = \Sigma$ satisfies

$$\int_{\Sigma} (p_k^{2/k} - 1) d\mu \geq \omega_{n-1}^{\frac{2}{n-1}} |\Sigma|^{\frac{n-3}{n-1}}.$$

Now if we assume that equality in (1.8) is attained, then $Q_k(t)$ is a constant. Then Proposition 7 indicates that Σ_t is totally umbilical and therefore a geodesic sphere. If Σ is a geodesic sphere of radius r , then $|\Sigma| = \omega_{n-1} \sinh^{n-1} r$ and $p_1 = \coth r$. Hence, we have

$$\int_{\Sigma} (p_k^{2/k} - 1) d\mu = \omega_{n-1} \sinh^{n-1} r (\coth^2 r - 1) = \omega_{n-1}^{\frac{2}{n-1}} |\Sigma|^{\frac{n-3}{n-1}}.$$

Therefore, equality in (1.8) holds on a geodesic sphere. This completes the proof of Theorem 1. \square

Remark 8. In fact, for $k = 1$ we prove Theorem 1, and for $k = 2$ we recover Theorem B proved by Li-Wei-Xiong [25].

It is natural to put forward the following question.

Question. For $k \in \{3, \dots, n - 1\}$, let $\Sigma \subset \mathbb{H}^n$ ($n \geq k + 1$) be a star-shaped and k -convex hypersurface. Then

$$(4.1) \quad \int_{\Sigma} (p_k^{2/k} - 1) d\mu \geq \omega_{n-1}^{\frac{2}{n-1}} |\Sigma|^{\frac{n-3}{n-1}}.$$

Equality in (4.1) holds if and only if Σ is a geodesic sphere.

ACKNOWLEDGMENTS

The author would like to thank the referee for valuable comments and suggestions. The author is very grateful to Professor Hongwei Xu for his support, encouragement, and stimulating discussions over the years. He would also like to thank Professor Haizhong Li for his interest and discussions.

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