

## More on the localization of a supersymmetric integral

The change of variables formula given in the display on Page 153 line 11 is incorrect, and in fact  $Z'$  does *not* vanish as incorrectly claimed in display (104). It turns out that the claimed method only works for supersymmetry transformations  $\delta$  satisfying  $\delta^2 = 0$ , which does not hold in this situation since, for example,

$$\delta^2(\psi_1) = \delta(\delta(\psi_1)) = \delta(h'(x)\epsilon_2) = h''(x)\delta(x)\epsilon_2 = h''(x)(\epsilon_1\psi_1 + \epsilon_2\psi_2)\epsilon_2 \neq 0.$$

Nevertheless, the conclusion (106) of the localization method is correct, an important point in the treatment. It requires a lengthier argument, which is supplied here.

The argument can be repaired by replacing  $h$  by  $\lambda h$  with  $\lambda$  a real parameter and letting  $\lambda$  approach infinity. This changes the action (100) to

$$S_\lambda(x, \psi_1, \psi_2) = \frac{\lambda^2 h'(x)^2}{2} - \lambda h''(x)\psi_1\psi_2,$$

which changes the partition function from  $Z$  to

$$Z(\lambda) = \int \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2) dx.$$

By the deformation invariance explained on page 155, it follows that  $Z(\lambda)$  is independent of  $\lambda$ , i.e., coincides with the original  $Z$ . In particular,

$$Z = \lim_{\lambda \rightarrow \infty} Z(\lambda).$$

We will reach (106) by calculating  $\lim_{\lambda \rightarrow \infty} Z(\lambda)$ .

Note also that the replacement of  $h$  by  $\lambda h$  does not alter the location of the critical points of  $h$ .

For any interval  $[a, b]$  not containing any critical points of  $h$ , including the possibility that  $a = -\infty$  or  $b = \infty$ , we have

$$\lim_{\lambda \rightarrow \infty} \int_a^b \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2) dx = \int_a^b \lim_{\lambda \rightarrow \infty} \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2) dx = 0,$$

due to the vanishing of the second limit which follows from the rapid decay of the exponential.

Denoting the critical points as  $x_{c_1} < x_{c_2} < \dots < x_{c_k}$ , choosing any positive  $\epsilon < \min\{(x_{c_{i+1}} - x_{c_i})/2\}$ , and denoting the integrand of  $Z(\lambda)$  by  $H_\lambda(x) = \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2)$  for ease of notation, we get for  $Z(\lambda)$  by splitting up the domain of integration the expression

$$\int_{-\infty}^{x_{c_1}-\epsilon} H_\lambda(x) dx + \sum_{i=1}^k \int_{x_{c_i}-\epsilon}^{x_{c_i}+\epsilon} H_\lambda(x) dx + \sum_{i=1}^{k-1} \int_{x_{c_i}+\epsilon}^{x_{c_{i+1}}-\epsilon} H_\lambda(x) dx + \int_{x_{c_k}+\epsilon}^{\infty} H_\lambda(x) dx.$$

It follows immediately that

$$Z = \lim_{\lambda \rightarrow \infty} Z(\lambda) = \lim_{\lambda \rightarrow \infty} \sum_{i=1}^k \int_{x_{c_i}-\epsilon}^{x_{c_i}+\epsilon} \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2) dx$$

because the limits of all of the other terms in the expansion of  $Z(\lambda)$  vanish as we have seen before, since the intervals that  $H_\lambda(x)$  is being integrated over contain no critical points.

By Taylor's Theorem with remainder, we now expand  $h$  near a critical point  $x_c$  as a truncated power series

$$h(x) = h(x_c) + \frac{h''(x_c)}{2} (x - x_c)^2 + O(x - x_c)^3,$$

where  $O(x - x_c)^3$  refers to a function of "order  $(x - x_c)^3$ ", which more precisely means a function  $f(x)$  such that  $|f(x)| \leq M|(x - x_c)^3|$  as  $x$  approaches  $x_c$ , for some constant  $M$ . We have

$$\begin{aligned} & \int_{x_c-\epsilon}^{x_c+\epsilon} \lambda h''(x) \exp(-\lambda^2 h'(x)^2/2) dx = \\ & \int_{x_c-\epsilon}^{x_c+\epsilon} \lambda (h''(x_c) + O(x - x_c)) \exp\left(-\lambda^2 \left(\frac{h''(x_c)(x - x_c)^2}{2} + O(x - x_c)^3\right)\right) dx. \end{aligned}$$

We next make the change of variables  $y = \lambda(x - x_c)$ , and we get

$$\int_{-\lambda\epsilon}^{\lambda\epsilon} \left(h''(x_c) + O\left(\frac{y}{\lambda}\right)\right) \exp\left(-\frac{h''(x_c)^2 y^2}{2} + \lambda^2 \left(O\left(\frac{y}{\lambda}\right)^3\right)\right) dy$$

Taking the limit as  $\lambda$  approaches infinity, the terms  $O(y/\lambda)$  and  $O((y/\lambda)^3)$  approach 0 and we are left with

$$\int_{-\infty}^{\infty} h''(x_c) \exp\left(-\frac{h''(x_c)^2 y^2}{2}\right) dy.$$

This is precisely the second line of (105) after another change of variables  $y = x - x_c$ .<sup>1</sup> From the above expression of  $Z$  as a sum of integrals in the neighborhood of critical points, we are led to (106), completing the argument.

Thanks to Kentaro Hori for assistance with the correction.

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<sup>1</sup>At the risk of introducing some confusion, we have inconsistently introduced the substitution variable  $x$  here which has already been used, in order to match (105).