

the analogue of the procedure. We consider the $(n - 1)$ -dimensional spaces V_{k_i} ($i = 1, 2, \dots, n$) through O , whose equations are $\sum a_{ik}x_k = 0$. When we take arbitrary values for the unknown quantities and correct them in the order x_{k_1}, x_{k_2}, \dots , the construction runs as follows. An arbitrary point P_0 is chosen in n -space; P_1 is the orthogonal projection of P_0 on V_{k_1} ; P_2 is the projection of P_1 on V_{k_2} , etc. And so the procedure is illustrated in a simple geometrical way. We add some supplementary remarks. The point P_m ($m > 0$) lies always in one of the spaces V_{k_i} . These spaces are linearly independent if $|a_{ik}| \neq 0$. At each step (after the first) one of the V_{k_i} is projected by parallel projection onto the following one. In this way, an affine correspondence is established between the two successive V_{k_i} , the *modulus* of the affinity being $\cos \alpha$, where α is the angle between them. Thus the convergence of the procedure can easily be proved, provided that the corrections take place in a fixed cyclic order $x_{k_1}, x_{k_2}, \dots, x_{k_n}$ where k_1, k_2, \dots, k_n is a permutation of $1, 2, \dots, n$.

If two successive V_{k_i} are perpendicular to each other, the projection of the first onto the second coincides with their $(n - 2)$ -dimensional space of intersection. It follows, therefore, that if all the V_{k_i} are mutually perpendicular (that is, if the matrix $\|a_{ik}\|$ is orthogonal), the point P_1 lies on V_{k_1} , P_2 on the intersection of V_{k_1} and V_{k_2} , P_3 on $(V_{k_1}, V_{k_2}, V_{k_3})$ and so on; hence P_n coincides with O . In this case the procedure ends automatically after n steps.

A SIMPLIFIED METHOD OF DIFFERENTIATING AND EVALUATING FUNCTIONS REPRESENTED BY FOURIER SERIES*

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1. Introduction. This paper shows how to eliminate the difficulties caused by discontinuities of Fourier sine series at the ends of the interval of periodicity.

Applications of Fourier series to exact solutions of problems in mathematical physics involve the following essential considerations. In an interval $-a \leq x \leq a$, it is assumed that a function $f(x)$ and its successive derivatives up to some finite order $f^{(m)}(x)$ all comply with sufficient conditions of continuity, bounded variation, differentiability and integrability. They thus permit representation by Fourier series, which can be differentiated to give the derivative of next higher order, and integrated to give an expression for the derivative of next lower order.

When $f(x)$ is an x -odd function, and $f(a) \neq f(-a) \neq 0$, particular difficulties are encountered. The corresponding Fourier sine series is discontinuous at $x = \pm a$ and does not conveniently represent the values of $f(a)$ and $f(-a)$. In addition, the derivative $f'(x)$ is represented by a complicated Fourier series which is not readily evaluated at $x = \pm a$. Thus

$$f(x) = \sum_1^{\infty} b'_n \sin \beta_n x, \quad (1)$$

in which $\beta_n = n\pi/a$. The expression for the derivative is

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$$f'(x) = \frac{1}{2a} [f(a) - f(-a)] + \sum_1^{\infty} \left\{ \frac{(-1)^n}{a} [f(a) - f(-a)] + b'_n \beta_n \right\} \cos \beta_n x. \quad (2)$$

2. Modification of Fourier sine series. These difficulties of evaluation and differentiation are alleviated by rewriting Eq. (1) in the form

$$f(x) = cx + \left(\sum_1^{\infty} b'_n \sin \beta_n x - cx \right), \quad (3)$$

where $c = f(a)/a$. The last term of Eq. (3) may be written

$$cx = \sum_1^{\infty} b''_n \sin \beta_n x;$$

then

$$\sum_1^{\infty} b'_n \sin \beta_n x - cx = \sum_1^{\infty} b_n \sin \beta_n x, \quad (4)$$

in which $b_n = b'_n - b''_n$.

In Eq. (4), $\sum_1^{\infty} b_n \sin \beta_n x$ represents a function having the value zero at $x = \pm a$. It consequently permits termwise differentiation.

Thus the original equation (1) is now

$$f(x) = cx + \sum_1^{\infty} b_n \sin \beta_n x. \quad (1a)$$

At $x = \pm a$, $f(x) = \pm ca$, respectively. Also the derivative $f'(x)$, correctly obtained by termwise differentiation, is

$$f'(x) = c + \sum_1^{\infty} a_n \cos \beta_n x, \quad (2a)$$

where $a_n = b_n \beta_n$.

In this manner, with complete generality, Eqs. (1) and (2) can always be rewritten in the form of Eqs. (1a) and (2a).

Furthermore, in the case of integration of an x -even function $f'(x)$ expressed by Eq. (2a), termwise integration is a legitimate process. If it is known that the indefinite integral $f(x)$ is an x -odd function, there is no additional constant of integration, and $f(x)$ is expressed by Eq. (1a), in which $b_n = a_n/\beta_n$. Also at $x = \pm a$, $f(x) = \pm ca$, respectively. This follows from the fact that termwise differentiation of Eq. (1a) gives Eq. (2a). Therefore, $\sum_1^{\infty} b_n \sin \beta_n x$, complying with the necessary condition of termwise differentiation, must represent a function having the value zero at $x = \pm a$.

3. Derivatives of any finite order. More general results for a mixed function are obtained by successive integrations of $f^{(m)}(x)$ with introduction of constants of integration. Thus

$$f^{(m)}(x) = c' + \sum_1^{\infty} (a'_n \cos \beta_n x + b'_n \sin \beta_n x) \quad (5)$$

after m successive integrations gives an expression of the form

$$f(x) = \sum_0^m c_n x^n + \sum_1^{\infty} (a_n \cos \beta_n x + b_n \sin \beta_n x). \quad (6)$$

Conversely, derivatives up to the order $f^{(m)}(x)$ may be obtained from Eq. (6) by term-wise differentiation. In the expressions for $f(x)$ and any derivative up to $f^{(m-1)}(x)$, any sine series represents a function having the value zero at $x = \pm a$, thus facilitating evaluation at these points. In some solutions Eq. (5) is a convenient expression for $f^{(m)}(x)$. If desired, Eq. (5) can be modified by adding a term $c''x$ combined with $\sum_1^\infty b'_n \sin \beta_n x$ so that the sine series in this equation also represents a function having the value zero at $x = \pm a$. Then Eq. (6), modified by this additional term, is

$$f(x) = \sum_0^{m+1} c_n x^n + \sum_1^\infty (a_n \cos \beta_n x + b_n \sin \beta_n x). \quad (6a)$$

Equations (6) and (6a) are special cases of Borel's theorem¹ which apply to derivatives of any finite order up to $f^{(m)}(x)$.

The preceding methods of derivation, stated in terms of a single variable x , also apply completely to a three-dimensional member bounded by parallel planes $x = \pm a$ with assigned boundary conditions. In this case $f(x)$ is replaced by $f(x, y, z)$ and the derivatives are written as partial derivatives with respect to x , $f^{(m)}(x)$ becoming $\partial^m f / \partial x^m$. Also in Eqs. (1a), (2a), (6), (6a), every coefficient a_n , b_n , c_n , c , while independent of x , is a function of y and z . For example, in Eq. (1a), b_n is $b_n(y, z)$, and c is $c(y, z) = f(a, y, z)/a$. Thus $f(x, y, z)$ and each of its x -derivatives of finite order can be expressed as the sum of a finite power series and a Fourier series without discontinuity at $x = \pm a$; the expressions can, therefore, be evaluated definitely at these boundaries.

¹E. Borel, *Leçons sur les fonctions de variables réelles*, Gauthier-Villars, Paris, 1905, p. 68.

A TAPERED LINE TERMINATION AT MICROWAVES*

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1. Introduction. In the field of ultra-high frequency, one method of power transmission is by means of concentric metallic conductors called a co-axial transmission line. The maximum amount of power can be delivered over these lines if the impedance of the load is equal to the characteristic impedance¹ of the line. In addition, in various measurements at ultra-high frequency it is essential to have a matched termination over a broad band of frequencies. This problem was approached experimentally and led to the tapered line termination as shown in Fig. 1.

To the left of $x = 0$, the co-axial line has an inner metallic conductor of radius a . From $x = 0$ to $x = L$, the inner conductor is a glass tube coated with a thin metallic film of resistive material.² To the left of $x = 0$, the outer conductor has a radius b and from $x = 0$ to $x = L$, the outer conductor has a linear taper down to a radius c . Also at $x = L$, the inner and outer conductors are short circuited.

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¹J. C. Slater, *Microwave transmission*, McGraw Hill, 1942, pp. 71-74.

²Publication P. B. 6588 (U. S. Commerce Dept.), 1945.