

**A VARIATIONAL PRINCIPLE FOR THE FIRST EIGENVALUE  
OF A SEMIFREE MEMBRANE\***

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It is well known that the eigenvalues of

$$-\Delta u + cu = \lambda u \text{ in } G, \quad u = 0 \text{ on } \partial G, \quad (1)$$

vary inversely with the domain  $G$ , in the sense that a decrease in the domain  $G$  will increase the eigenvalues of (1). Physically, we may think of (1) as determining the standing waves of an elastic membrane embedded in a medium which exerts a force  $-c(x)u$  on each point of the membrane. If the membrane is tied down on  $\partial G$ , then a decrease in  $G$  will increase the fundamental frequency of the membrane.

If the membrane is not tied down on  $\partial G$  but satisfies some other selfadjoint boundary condition then the above results need not carry over. Specifically, consider the case of a free membrane

$$-\Delta u + cu = \lambda u \text{ in } G, \quad \partial u / \partial \nu = 0 \text{ on } \partial G, \quad (1')$$

where  $\underline{c} = \inf_{x \in G} c(x)$  and  $\bar{c} = \sup_{x \in G} c(x)$ . Then the first eigenvalue of (1') satisfies  $\underline{c} \leq \lambda_1 \leq \bar{c}$ , with equality iff  $c(x)$  is constant. If in a certain neighborhood of  $\partial G$   $\lambda_1 > c(x)$ , then the elastic forces within the membrane dominate the external force  $-cu$ , and we can expect that removing such a piece of the free membrane will increase the fundamental frequency. If, on the other hand,  $\lambda_1 < c(x)$  in a certain neighborhood of  $\partial G$ , then the elastic forces within the membrane are dominated by the external force, and we can expect to decrease the fundamental frequency by removing a portion of the free membrane in which  $\lambda_1 < c(x)$ .

In order to verify these heuristic conclusions (in a slightly more general setting) we shall consider the dependence of the first eigenvalue of

$$\begin{aligned} -\Delta u + cu &= \lambda u \text{ in } G, \\ \partial u / \partial \nu &= 0 \text{ on } \Gamma_1 \subset \partial G, \end{aligned} \quad (2)$$

$$\partial u / \partial \nu + \sigma u = 0 \text{ on } \Gamma_2 \equiv \partial G - \bar{\Gamma}_1, \quad -\infty < \sigma(x) \leq +\infty,$$

with changes in  $G$  brought about by small variations in  $\Gamma_1$ . Specifically, we consider the smaller domain  $G^* \subset G$  such that  $\partial G^* \cap \partial G = \Gamma_2$  and  $\partial G^* \cap G = \Gamma_1^*$  and seek a relation between the first eigenvalue of (2) and

$$\begin{aligned} -\Delta u + cu &= \lambda^* u \text{ in } G^*, \\ \partial u / \partial \nu &= 0 \text{ on } \Gamma_1^*, \\ \partial u / \partial \nu + \sigma u &= 0 \text{ on } \Gamma_2. \end{aligned} \quad (2')$$

We shall show that if  $\lambda_1$  and  $\lambda_1^*$  are the first eigenvalues of (2) and (2') respectively, and if  $c(x) - \lambda_1^* < 0$  in  $G - \bar{G}^*$ , then  $\lambda_1 < \lambda_1^*$  (as in the case of Dirichlet boundary conditions). However, if  $c(x) - \lambda_1 \geq 0$  in  $G - G^*$ , then  $\lambda_1 > \lambda_1^*$ .

\*Received November 22, 1967; revised version received February 1, 1968.

It is assumed throughout that the coefficient  $c$  is continuous and that the domains  $G$  and  $G^*$  are of bounded curvature so that the classical variational theory for eigenvalues [1] is valid and the extremal functions so obtained are solutions of the related Euler-Lagrange equations in the classical sense. The positive normalized eigenfunctions corresponding to  $\lambda_1$  and  $\lambda_1^*$  will be denoted by  $v(x)$  and  $v^*(x)$ , respectively. The exterior normal derivative is denoted by  $\partial/\partial\nu$ .

**THEOREM 1.** *If  $c(x) - \lambda_1^* < 0$  in  $G - \overline{G^*}$ , then  $\lambda_1 < \lambda_1^*$ .*

*Proof.* By means of the usual variational formula we obtain

$$\lambda_1 - \lambda_1^* = \min_{u \in \Phi} \frac{\iint_G [|\nabla u|^2 + (c - \lambda_1^*)u^2] dx + \int_{\Gamma_2} \sigma u^2 ds}{\iint_G u^2 dx}$$

where the class  $\Phi$  consists of nontrivial functions which are continuous in  $G$ , have piecewise continuous first partial derivatives in  $G$ , and vanish whenever  $\sigma(x) = +\infty$  (such points being excluded from the boundary integral over  $\Gamma_2$ ). We shall extend  $v^*(x)$  into  $G$  as follows: consider the first order partial differential equation

$$|\nabla u|^2 + \frac{c - \lambda_1^*}{2} u^2 = 0, \quad u = v^* \quad \text{on} \quad \Gamma_1^*. \tag{3}$$

Since  $c - \lambda_1^* < 0$ , this equation makes sense. By choosing the proper square roots of the  $(\partial u/\partial x_i)^2$ , we can assume locally that  $\Gamma_1^*$  is not a characteristic for the differential equation. Finally, for sufficiently small exterior displacements  $\Gamma_1$  of  $\Gamma_1^*$ , classical theorems [1] guarantee the local existence of a continuously differentiable solution of (3) in  $G - G^*$ . Denoting such a solution by  $\tilde{v}(x)$ , we note that

$$\begin{aligned} w(x) &= v^*(x) \quad \text{in} \quad G^* \\ &= \tilde{v}(x) \quad \text{in} \quad G - G^* \end{aligned}$$

belongs to the class  $\Phi$ , and that

$$\begin{aligned} \lambda_1 - \lambda_1^* &\leq \frac{\iint_G [|\nabla w|^2 + (c - \lambda_1^*)w^2] dx + \int_{\Gamma_2} \sigma w^2 ds}{\iint_G w^2 dx} \\ \lambda_1 - \lambda_1^* &\leq \frac{\iint_{G^*} [|\nabla v^*|^2 + (c - \lambda_1^*)v^{*2}] dx + \int_{\Gamma_2} \sigma v^{*2} ds}{\iint_G w^2 dx} + \frac{\iint_{G-G^*} \frac{(c - \lambda_1^*)}{2} \tilde{v}^2 ds}{\iint_G w^2 dx}. \end{aligned} \tag{4}$$

By Green's Theorem and (2') the first term on the right side of (4) is zero. Since  $(c - \lambda_1^*) < 0$  in  $G - G^*$ , the second term is negative and  $\lambda_1 < \lambda_1^*$ .

In order to prove the converse relationship we shall make an additional assumption regarding the regularity of (2): for sufficiently small displacements along an interior direction normal to  $\Gamma_1$ ,  $\partial v/\partial\nu$  is to be a monotone function of distance from  $\Gamma_1$ . With this we can prove

**THEOREM 2.** *If  $c(x) - \lambda_1 \geq 0$  in  $G - \overline{G^*}$ , then  $\lambda_1 > \lambda_1^*$ .*

*Proof.* Let  $\Gamma_0$  be the subset of  $\Gamma_1$  for which  $\partial v/\partial \nu$  is a strictly decreasing function of distance from  $\Gamma_1$  (measured along an interior normal direction). We shall show that  $\Gamma_0$  is empty. For if  $x_0$  is an isolated point of  $\Gamma_0$ , then the level curve  $v = v(x_0)$  must form a cusp with vertex at  $x_0$  and with the interior normal to  $\Gamma_1$  at  $x_0$  lying inside the cusp. Since  $\partial v/\partial \nu \uparrow 0$  at boundary points near  $x_0$ , an examination of level curves near  $v = v(x_0)$  shows that  $v(x)$  has a positive maximum along the level curve  $v = v(x_0)$ . But since  $c(x) - \lambda_1 \geq 0$  near  $\Gamma_1$ , this contradicts the maximum principle and shows that  $\Gamma_0$  has no isolated points. If  $x_0$  is an interior point of  $\Gamma_0$  then we can construct a sphere  $S_0$  which lies in  $G - \overline{G^*}$ , is tangent to  $\Gamma_1$  at  $x_0$ , and is sufficiently small so that  $v(x_0) \geq v(x)$  for all  $x \in S_0$ . But by Hopf's second lemma [2], this implies  $\partial v/\partial \nu > 0$  at  $x_0$ , which contradicts (2) and shows that  $\Gamma_0$  is empty.

Therefore  $\partial v/\partial \nu \uparrow 0$  as  $x \rightarrow \Gamma_1$  along an interior normal direction, and for a sufficiently small interior displacement from  $\Gamma_1$  to  $\Gamma_1^*$ , we have  $\partial v/\partial \nu \leq 0$  on  $\Gamma_1^*$ ,  $\partial v/\partial \nu \neq 0$ . Therefore, in  $G^*$ ,  $v(x)$  will satisfy

$$\begin{aligned} -\Delta v + cv &= \lambda_1 v & \text{in } G^*, \\ \partial v/\partial \nu + \sigma^* v &= 0 & \text{on } \Gamma_1^*, \\ \partial v/\partial \nu + \sigma v &= 0 & \text{on } \Gamma_2, \end{aligned} \tag{5}$$

where  $\sigma^* \geq 0$  on  $\Gamma_1^*$ ,  $\sigma^* \neq 0$ . By classical variational principles, the first eigenvalue of (5) is larger than the first eigenvalue of (2'). Since  $v(x)$  is a nonnegative eigenfunction of (5) corresponding to  $\lambda_1$ , we have  $\lambda_1 > \lambda_1^*$ .

#### REFERENCES

- [1] R. Courant and D. Hilbert, *Methods of mathematical physics*, Vol. 1, Interscience, New York
- [2] E. Hopf, *A remark on linear elliptic differential equations of second order*, Proc. Amer. Math. Soc. **7**, 791-793 (1952)