

—NOTES—

VARIATIONAL SOLUTIONS OF THE  
THOMAS-FERMI EQUATION\*

BY N. ANDERSON AND A. M. ARTHURS (*University of York*)

**Abstract.** Variational solutions of the Thomas-Fermi equation are examined in the context of complementary extremum principles. A new one-parameter trial function is found to provide an accurate representation of the solution.

**1. Introduction.** The Thomas-Fermi screening function  $\phi(x)$  for neutral atoms with spherical symmetry satisfies the nonlinear second-order differential equation (see [6])

$$d^2\phi/dx^2 = \phi^{3/2}/x^{1/2}, \quad 0 \leq x < \infty, \quad (1)$$

subject to the boundary conditions

$$\phi(0) = 1; \quad \phi \rightarrow 0, \quad x\phi' \rightarrow 0 \quad \text{as } x \rightarrow \infty. \quad (2)$$

Near  $x = 0$  the function  $\phi(x)$  can be expanded as

$$\phi(x) = 1 + cx + \frac{4}{3}x^{3/2} + \frac{2}{5}cx^{5/2} + \frac{1}{3}x^3 + \dots, \quad (3)$$

where  $c$  is the unspecified value of  $\phi'(0)$ , while for large  $x$  the solution behaves like

$$\phi(x) \sim 144/x^3 \quad \text{as } x \rightarrow \infty. \quad (4)$$

An analytical solution of the problem in (1) and (2) is not known and so recourse must be taken in approximate solutions. Such solutions are either purely numerical [3, 5] or variational [4, 7]. Here we shall be concerned with solutions of the latter kind.

From the theory of complementary variational principles (see [2]) we find that the solution  $\phi$  of (1) and (2) is the function which minimizes the integral

$$J(\Phi) = \int_0^\infty \left\{ \frac{1}{2} (\Phi')^2 + \frac{2}{5} \frac{\Phi^{5/2}}{x^{1/2}} \right\} dx, \quad (5)$$

and which maximizes the integral

$$G(\Psi) = - \int_0^\infty \left\{ \frac{1}{2} (\Psi')^2 + \frac{3}{5} \frac{(x^{1/2}\Psi'')^{5/3}}{x^{1/2}} \right\} dx - \Psi'(0). \quad (6)$$

The extreme values of  $J$  and  $G$  agree and we have the global upper and lower bounds

$$G(\Psi) \leq G(\phi) = J(\phi) \leq J(\Phi). \quad (7)$$

---

\* Received July 1, 1980.

Here the trial function  $\Phi$  is subject to conditions (2) while the trial function  $\Psi$  is free of essential conditions. The extremum principles in (7) provide a basis for variational approximations to the exact function  $\phi$ , the difference  $J-G$  being a measure of the accuracy of  $\Phi$  and  $\Psi$ .

**2. Variational solutions.** The trial function

$$\Phi_1 = \{ae^{-\alpha x} + be^{-\beta x}\}^2, \quad a + b = 1, \quad (8)$$

has been proposed by Csavinsky [4] who determined the three free parameters by a method equivalent to minimizing  $J(\Phi_1)$ , with the results

$$J = 0.6816 \quad (9)$$

at  $a = 0.7111$ ,  $\alpha = 0.175$ ,  $\beta = 9.5\alpha$ . For the same class of trial functions the optimum value of the complementary integral  $G$  is [1]

$$G = 0.6010. \quad (10)$$

An improvement on these results is provided by the one-parameter trial function

$$\Phi_2 = (1 + \gamma x^{1/2})e^{-\gamma x^{1/2}} \quad (11)$$

of Roberts [7], who found the minimum value of  $J(\Phi_2)$  to be

$$J = 0.6810 \quad \text{at} \quad \gamma = 1.905. \quad (12)$$

The optimum value of  $G$  for the same class of trial functions is [1]

$$G = 0.6699 \quad \text{at} \quad \gamma = 1.750. \quad (13)$$

Neither function  $\Phi_1$  nor  $\Phi_2$  has the correct behavior for large  $x$ , while for small  $x$ , where the function is most significant,  $\Phi_1$  has an expansion involving only integral powers of  $x$  and cannot look like (3), and  $\Phi_2$  has the expansion

$$\Phi_2 = 1 - \frac{1}{2}\gamma^2 x + \frac{1}{3}\gamma^3 x^{3/2} - \frac{1}{8}\gamma^4 x^2 + \left(\frac{1}{4!} - \frac{1}{5!}\right)\gamma^5 x^{5/2} + O(x^3). \quad (14)$$

Comparison of (14) with (3) shows that these expressions agree on the absence of a term in  $x^{1/2}$ , but that  $\Phi_2$  contains a term in  $x^2$  which is not present in (3). To incorporate this feature of the exact function  $\phi$  near  $x = 0$ , and to retain expansion in powers of  $x^{1/2}$ , we shall therefore modify  $\Phi_2$  in (11) in the simplest possible way and introduce the one-parameter class of trial functions

$$\Phi_3 = (1 + \alpha x^{1/2} + \frac{1}{4}\alpha^2 x)e^{-\alpha^2 x} e^{-\alpha x^{1/2}}. \quad (15)$$

Near  $x = 0$  such a trial function has the expansion

$$\Phi_3 = 1 - \frac{1}{4}\alpha^2 x + \frac{1}{12}\alpha^3 x^{3/2} - \frac{1}{5!}\alpha^5 x^{5/2} + \frac{1}{4!12}\alpha^6 x^3 + O(x^{7/2}), \quad (16)$$

the terms in  $x^{1/2}$  and  $x^2$  being absent as required.

With function  $\Phi_3$  the optimum values of  $J$  and  $G$  and the variational parameter  $\alpha$  are

$$J = 0.6811 \quad \text{at} \quad \alpha = 2.472, \quad (17)$$

$$G = 0.6803 \quad \text{at} \quad \alpha = 2.528. \quad (18)$$

In terms of the agreement of these bounds the variational solution  $\Phi_3$  is the best of the three trial functions. Comparison with the numerical solution [3] shows a discrepancy of less than one percent in the region  $0 \leq x \leq 1$  where the function drops from 1 to 0.4. For  $x > 1$ ,  $\Phi_3$  falls off more rapidly than the numerical solution, but this behavior is to be expected through the built-in exponential decrease in the trial function.

## REFERENCES

- [1] N. Anderson, A. M. Arthurs and P. D. Robinson, *Nuovo Cimento* **57B**, 523 (1968)
- [2] A. M. Arthurs, *Complementary variational principles*, Clarendon Press, Oxford, second edition, 1980
- [3] V. Bush and S. H. Caldwell, *Phys. Rev.* **38**, 1898 (1931)
- [4] P. Csavinszky, *Phys. Rev.* **166**, 53 (1968)
- [5] S. Kobayashi, T. Matsukuma, S. Nagai and K. Umeda, *J. Phys. Soc. Japan* **10**, 759 (1955)
- [6] L. D. Landau and E. M. Lifshitz, *Quantum mechanics*, Pergamon Press, Oxford, 1958
- [7] R. E. Roberts, *Phys. Rev.* **170**, 8 (1968)