## THE DE'TERMINATION OF THE CONSTANTS IN THE PROBLEM OF THE BRACHISTOCHRONE.*

The general solution of Euler's differential equation for the problem of the brachistochrone in a vertical plane is the doubly infinite system of cycloids $\dagger$

$$
\begin{align*}
& x-x_{0}+h= \pm r(\omega-\sin \omega)  \tag{1}\\
& y-y_{0}+k=r(1-\cos \omega)
\end{align*}
$$

referred to a rectangular system of coördinates whose $(x, y)$ plane is the given vertical plane and whose positive $y$-axis is directed vertically downward ; $x_{0}, y_{0}$ are the coördinates of the starting point $A, k$ is a given constant, viz.,

$$
k=v_{0}^{2} / 2 g
$$

where $v_{0}$ is the initial velocity and $g$ the constant of gravity; finally $h$ and $r$ are the two constants of integration, and $r$ is essentially positive.

We suppose that the endpoints $A$ and $B$ are fixed and propose to determine the constants of integration so that the cycloid

[^0]shall pass through the two given points $A$ and $B$, with the additional condition that no cusp * of the cycloid shall lie between $A$ and $B$ and that the angle $\omega$ increases as the point $x, y$ describes the cycloid from $A$ to $B$.

We shall prove that this problem has always one and but one solution provided that

$$
\begin{equation*}
x_{1} \neq x_{0} \quad \text { and } \quad y_{1}-y_{0}+k \geqq 0 \tag{2}
\end{equation*}
$$

If we denote by $x_{1}, y_{1}$ the coördinates of the endpoint $B$ and by $\omega_{0}$ and $\omega_{1}$ the values of the angle $\omega$ at the two points $A$ and $B$ respectively, we have to determine the four quantities $h$, $r, \omega_{0}, \omega_{1}$, from the four equations

$$
\begin{align*}
& h= \pm r\left(\omega_{0}-\sin \omega_{0}\right), x_{1}-x_{0}+h= \pm r\left(\omega_{1}-\sin \omega_{1}\right), \\
& k=r\left(1-\cos \omega_{0}\right), \quad y_{1}-y_{0}+k=r\left(1-\cos \omega_{1}\right), \tag{3}
\end{align*}
$$

with the additional inequality $\dagger$

$$
\begin{equation*}
0 \leqq \omega_{0}<\omega_{1} \leqq 2 \pi \tag{4}
\end{equation*}
$$

Since the function $\omega-\sin \omega$ always increases as $\omega$ increases, the upper sign must be taken when $x_{1}>x_{0}$, the lower sign when $x_{1}<x_{0}$.

Leaving aside the two special cases

$$
v_{0}=0, \quad \text { and } \quad v_{0} \neq 0, \quad y_{1}=y_{0}
$$

in which the solution is rather trivial, we assume, in addition to the inequalities (2), that $v_{0} \neq 0$ and therefore
and that

$$
\begin{equation*}
k>0 \tag{5}
\end{equation*}
$$

$$
\begin{equation*}
y_{1} \neq y_{0} . \tag{6}
\end{equation*}
$$

If we eliminate $h$ and $r$ between the four equations (3), we obtain

$$
\begin{gather*}
m\left[\omega_{1}-\omega_{0}-\left(\sin \omega_{1}-\sin \omega_{0}\right)\right]=\cos \omega_{0}-\cos \omega_{1}  \tag{7}\\
a \sin \frac{\omega_{0}}{2}=\sin \frac{\omega_{1}}{2} \tag{8}
\end{gather*}
$$

[^1]where
$$
m= \pm \frac{y_{1}-y_{0}}{x_{1}-x_{0}}
$$
and
$$
a=\sqrt{\frac{y_{1}-y_{0}+k}{k}}
$$
the square root being taken positive.
According as $y_{1} \gtrless y_{0}$, we have $m \gtrless 0$ and $a \gtrless 1$.
If we set
\[

$$
\begin{equation*}
u=\frac{\omega_{0}+\omega_{1}}{2}, \quad v=\frac{\omega_{1}-\omega_{0}}{2} \tag{9}
\end{equation*}
$$

\]

we obtain from (8)

$$
\begin{equation*}
\operatorname{tg} \frac{u}{2}=\frac{a+1}{a-1} \operatorname{tg} \frac{v}{2} \tag{10}
\end{equation*}
$$

Hence if we introduce $u$ and $v$ in (7), express $u$ in terms of $v$ by means of (10) and put for brevity

$$
\begin{equation*}
z=\operatorname{tg} \frac{v}{2}, \quad c=\frac{a-1}{a+1}, \quad b=\frac{c}{m} \tag{11}
\end{equation*}
$$

the problem reduces to the solution of the transcendental equation

$$
\begin{equation*}
F(z) \equiv \operatorname{arctg} z-\frac{z\left(c^{2}+2 b z-z^{2}\right)}{\left(1+z^{2}\right)\left(c^{2}+z^{2}\right)}=0 \tag{12}
\end{equation*}
$$

On account of (4), the angle

$$
\frac{v}{2}=\operatorname{arctg} z
$$

lies between 0 and $\pi / 2$ and therefore $z$ must be positive.
For the derivative of $F(z)$ we obtain

$$
F^{\prime}(z)=\frac{z\left[4 b z^{4}+2\left(1+3 c^{2}\right) z^{3}+2 c^{2}\left(3+c^{2}\right) z-4 b c^{2}\right]}{\left(z^{2}+1\right)^{2}\left(z^{2}+c^{2}\right)^{2}}
$$

Since we suppose $x_{0} \neq x_{1}$ and $y_{0} \neq y_{1}, m$ and $c$ are both finite and different from zero and have the same sign, viz., the same sign as $y_{1}-y_{0}$. Hence $b$ is always positive. It follows therefore from Descartes's rule that the numerator of $F^{\prime}(z)$ has one and but one real positive root, say $z=z_{0}$. Since moreover $F^{\prime}(z)$
is negative for sufficiently small positive values of $z$, it follows that $F(z)$ decreases from the value zero to a certain negative minimum value, as $z$ increases from 0 to $z_{0}$; then as $z$ increases

from $z_{0}$ to $+\infty, F(z)$ increases continually and approaches for $z=+\infty$ the limit $+\pi / 2$. Hence the equation (12) has one and but one real positive root.

This root being found, the equations (11), (10), and (9) yield a unique solution $\omega_{0}, \omega_{1}$, of (7) and (8), satisfying the inequality (4). Finally the values of $h$ are $z$ follow unambiguously from (3).

The existence and uniqueness of the solution of the proposed problem are therefore proved.

University of Chicago,
November 2, 1903.

## ON THREE TYPES OF SURFACES OF THE THIRD ORDER REGARDED AS DOUBLE SURFACES OF TRANSLATION.

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(Read before the American Mathematical Society, October 31, 1903.)
This note serves the double purpose of making a slight addition to the theory of three types of surfaces of the third order and of exhibiting the double surfaces of translation of lowest order. The latter surfaces enjoy all the properties of the double minimum surfaces * except those immediately dependent on the

[^2]
[^0]:    * Very little attention seems to have been paid to the question of the determination of the constants in the problem of the brachistochrone. I have been able to find only the following few references on the subject in the literature of the calculus of variations:

    Johann Bernoulli who first proposed the problem of the brachistochrone in 1696, and Jacob Bernoulli give a geometrical construction for the special case where the initial velocity is zero, in which case the starting point is a cusp of the cycloid (compare Ostwald's Klassiker der exacten Wissenschaften, no. 46, pp. 12 and 18).

    Dienger, Grundriss der Variationsrechnung (1867), p. 38, reduces for the general case the determination of the constants to the two transcendental equations (7) and (8) of the text without entering into a further discussion of these equations.

    Weierstrass in his lectures (1882) states without proof that it is always possible to construct a cycloid upon a given base passing through two given points $A$ and $B$, and only one cycloid which contains no cusp between $A$ and $B$.
    $\dagger$ Compare for instance Lindelöf-Moigno, Calcul des variations, p. 228, and Pascal, Die Variationsrechnung, \& 31, where numerous historical references are given.

[^1]:    * An arc of a cycloid containing a cusp in its interior can never furnish a minimum for the time (Schwarz, Lectures, 1898-99, unpublished).
    $\dagger$ If $\omega_{0}, \omega_{1}, h, r$ is a solution of (4) and if $\mu$ is an integer, then $\omega_{0}+2 \mu \pi$, $\omega_{1}+2 \mu \pi, h \pm 2 \mu \pi r, r$ is also a solution, furnishing however the same cycloid. Hence we may restrict $\omega_{0}$ and $\omega_{1}$, to the above interval without loss of generality.

[^2]:    * Lie, Math. Annalen, vol. 14 (1879), p. 346 et seq. ; Darboux, Théorie des surfaces, vol. 1, p. 348 et seq.

