

been definitely excluded from this result by previous assumption that $x < n$.

For other values of d the expressions in (5) do not simplify, and we have to determine x from the general conditions

$$n = ax + c, c < x, (a \neq 0, a - c > (d - 1)(x + 1) \text{ and } \leq d(x + 1)).$$

It may be verified by actual trial that the solutions of (1) and (2) may in this way be effected with a degree of readiness which will make the method serviceable in a large number of cases.

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THE INNER FORCE OF A MOVING ELECTRON.

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1. *Introduction.* — Sommerfeld* has given a general method of determining the force which an electron exerts on itself when its motion is known. Schott† has applied the retarded potential to the same problem. In the present paper the known vector expression‡ for the force with which a moving point charge acts on another point charge is used to determine the action between any two elements of the electron. A double integration over the volume of the electron gives the inner force of the electron. The Abraham-Sommerfeld expressions for the longitudinal mass and Abraham's value for the transverse mass are thus very simply determined and the limitations on the solutions are made manifest.

2. *The Force Between Two Moving Point Charges.* — The electron is assumed to be a uniformly charged solid sphere which is moving in a straight line without rotation. Take the x -axis in the direction of motion. Measure the time from the instant at which the force is to be determined and choose the velocity of light as the unit of velocity. Let (x_0, y, z) be the coordinates of the point charge de_2 relative to the point charge de_1 . The distance moved by the electron in time t is

* *Göttinger Nachrichten*, 1904.

† *Ann. der. Physik*, 1908, No. 1.

‡ Abraham : *Theorie der Elektrizität*, vol. II, p. 98.

$$v_0 t + \frac{1}{2} \dot{v}_0 t^2 + \frac{1}{6} \ddot{v}_0 t^3 + \dots,$$

where v_0 is its velocity at time 0. The point charge de_1 will therefore act on de_2 at the time $-\tau$ which is determined so as to satisfy

$$(1) \quad \tau^2 = y^2 + z^2 + (x_0 + v_0 \tau - \frac{1}{2} \dot{v}_0 \tau^2 + \dots)^2.$$

The velocity of de_1 at this time is

$$(2) \quad v = v_0 - \dot{v}_0 \tau + \frac{1}{2} \ddot{v}_0 \tau^2 - \dots,$$

and its acceleration is

$$(3) \quad \dot{v} = \dot{v}_0 - \ddot{v}_0 \tau + \dots.$$

If the velocity of de_1 had remained constant and equal to v from the time $-\tau$, its coordinates at time 0 would have been (x, y, z) , where

$$x = x_0 + v_0 \tau - \frac{1}{2} \ddot{v}_0 \tau^2 + \dots - \tau(v_0 - \dot{v}_0 \tau + \dots)$$

or

$$(4) \quad x = x_0 + \frac{1}{2} \dot{v}_0 \tau^2 + \dots.$$

The x component of the force of de_1 on de_2 can be reduced from its vector form to

$$(5) \quad \frac{x(1 - v^2) - \dot{v}(y^2 + z^2)}{\{x^2 + (1 - v^2)(y^2 + z^2)\}^{\frac{3}{2}}} de_1 de_2.$$

The effect of de_2 on de_1 is obtained by changing the sign of x_0 .

3. *Uniform Motion.*—In this case v is v_0 and x is x_0 , so that (5) becomes

$$\frac{x_0(1 - v_0^2)}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{3}{2}}} de_1 de_2.$$

Hence the mutual actions of de_1 and de_2 are equal and opposite and the inner force is zero.

4. *Slow Uniform Acceleration.*—Here \dot{v} is \dot{v}_0 , and the unbalanced part of the force of de_1 on de_2 is

$$-\frac{\dot{v}_0(y^2 + z^2)de_1 \cdot de_2}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{3}{2}}},$$

which may be written

$$-\frac{\dot{v}_0 \partial}{v_0 \partial v_0} \left[\frac{de_1 de_2}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{1}{2}}} \right].$$

The inner force of the electron is twice the sum of this expression for every pair of elements of the electron. Let the coordinates of de_1 and de_2 referred to the center of the electron be (x_1, y_1, z_1) and (x_2, y_2, z_2) . Then

$$x_0 = x_2 - x_1, \quad y = y_2 - y_1, \quad \text{and} \quad z = z_2 - z_1.$$

If we now make the transformation

$$x = \sqrt{1 - v_0^2} \cdot x', \quad \rho = \frac{\rho'}{\sqrt{1 - v_0^2}},$$

where ρ is the volume density of the charge e of the electron, the charge is unaltered, but the spherical electron of radius a is transformed into the prolate spheroid

$$x'^2(1 - v_0^2) + y^2 + z^2 = a^2.$$

The expression which we have to integrate is therefore $2/\sqrt{1 - v_0^2}$ of the self potential of this spheroid and has the value*

$$\frac{6}{5} \cdot \frac{e^2}{2av_0} \log \left(\frac{1 + v_0}{1 - v_0} \right).$$

Hence the inner force of the electron is

$$-\frac{6}{5} \cdot \frac{e^2 \dot{v}_0}{2a} \cdot \frac{\partial}{v_0 \partial v_0} \left\{ \frac{1}{v_0} \log \left(\frac{1 + v_0}{1 - v_0} \right) \right\},$$

that is,

$$-\dot{v}_0 \cdot \frac{6}{5} \cdot \frac{e^2}{2av_0^2} \left(-\frac{1}{v_0} \log \frac{1 + v_0}{1 - v_0} + \frac{2}{1 - v_0^2} \right).$$

The coefficient of $-\dot{v}_0$ is Abraham's expression for the longitudinal mass in the case of quasi-stationary motion.

5. *Variable Acceleration.* — In the previous case the changes in the velocity and acceleration were neglected. If we retain the first power of τ , we have from (2) and (3)

$$v = v_0 - \dot{v}_0 \tau, \quad \dot{v} = \dot{v}_0 - \ddot{v}_0 \tau,$$

* Abraham, *loc. cit.*, p. 179.

and from (1)

$$\tau = \frac{v_0 x + \sqrt{x_0^2 + (1 - v_0^2)(y^2 + z^2)}}{1 - v_0^2}.$$

Substituting in (5), we find the additional unbalanced terms

$$\frac{\ddot{v}_0(y^2 + z^2)de_1de_2}{(1 - v_0^2)\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}} + \frac{3\dot{v}_0^2v_0(y^2 + z^2)^2de_1de_2}{(1 - v_0^2)\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^2}.$$

Denoting by θ the angle which the line joining de_1 to de_2 makes with the x axis, we have

$$\frac{\ddot{v}_0}{1 - v_0^2} \cdot \frac{\sin^2 \theta de_1de_2}{1 - v_0^2 \sin^2 \theta} + \frac{3\dot{v}_0^2}{2(1 - v_0^2)} \frac{\partial}{\partial v_0} \left(\frac{\sin^2 \theta de_1de_2}{1 - v_0^2 \sin^2 \theta} \right).$$

Consider

$$\int \int \frac{\sin^2 \theta de_1de_2}{1 - v_0^2 \sin^2 \theta};$$

since it is independent of the length of the line joining de_1 to de_2 , its value is

$$\frac{e^2}{4\pi} \int \frac{\sin^2 \theta}{1 - v_0^2 \sin^2 \theta} dw,$$

where dw is the element $2\pi \sin \theta d\theta$ of the complete solid angle 4π . Integrating from 0 to π , we find the value of the additional terms to be

$$\begin{aligned} \frac{\ddot{v}_0 e^2}{1 - v_0^2} \left(-\frac{1}{v_0^2} + \frac{1}{v_0^3 \sqrt{1 - v_0^2}} \tan^{-1} \frac{v_0}{\sqrt{1 - v_0^2}} \right) \\ + \frac{3\dot{v}_0^2 e^2}{2(1 - v_0^2)} \frac{\partial}{\partial v_0} \left(-\frac{1}{v_0^2} + \frac{1}{v_0^3 \sqrt{1 - v_0^2}} \tan^{-1} \frac{v_0}{\sqrt{1 - v_0^2}} \right). \end{aligned}$$

This result differs somewhat from the expression found by Schott and Abraham.

6. *Transverse Acceleration.*—When the velocity is constant but there is a small transverse acceleration f , we may choose the y axis in the direction of f . The vector expression for the y component of the unbalanced force may be reduced to

$$-f \frac{x_0^2 + (1 - v_0^2)z^2}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{3}{2}}} de_1de_2,$$

that is,

$$-f de_1 de_2 \left[\frac{1}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{3}{2}}} - \frac{(1 - v_0^2)y^2}{\{x_0^2 + (1 - v_0^2)(y^2 + z^2)\}^{\frac{3}{2}}} \right].$$

Integrating as in the previous cases, we have

$$-f \cdot \frac{6}{5} \cdot \frac{e^2}{2a} \left\{ \frac{1}{v_0} \log \frac{1 + v_0}{1 - v_0} - \frac{(1 - v_0^2)\partial}{2v_0\partial v_0} \left(\frac{1}{v_0} \log \frac{1 + v_0}{1 - v_0} \right) \right\}$$

or

$$-f \cdot \frac{6}{5} \cdot \frac{e^2}{2av_0^2} \left\{ \frac{1 + v_0^2}{2v_0} \log \frac{1 + v_0}{1 - v_0} - 1 \right\},$$

which gives Abraham's expression for the transverse mass.

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THE RECENTLY DISCOVERED MANUSCRIPT OF ARCHIMEDES.

PROFESSOR J. L. Heiberg has published two important accounts of his recent discovery of a new manuscript of Archimedes, both of which are of great interest to mathematicians. The first of these accounts is printed in volume 42 of *Hermes*. It contains the Greek text of a lost treatise of Archimedes, which is recovered nearly complete in the newly found manuscript. A German translation of the Greek text, and an interesting commentary by Zeuthen, is printed by Heiberg in the *Bibliotheca Mathematica*, volume 7, page 321.

Professor Heiberg's critical study of Archimedes has extended over a period of more than thirty years. His dissertation, "Quæstiones Archimedææ" (Copenhagen, 1879), is constantly referred to by students of Archimedes, both on account of its scholarly critique of Archimedes's work, and on account of the innumerable references to Archimedean literature which are there brought together for the first time. In 1880-1881 Heiberg published the definitive edition of Archimedes, with which his name is usually associated. It was while at work upon a second edition of this book that Heiberg's attention was directed to a palimpsest manuscript of mathematical content recently catalogued in a cloister at Constantinople. Failing in an attempt to have the manuscript sent to him at Copen-