

NOTE ON A THEOREM OF BÔCHER AND KOEBE*

BY J. J. GERGEN

1. *Introduction.* In this paper a generalization of the following theorem, discovered independently by Bôcher† and Koebe,‡ is established.

THEOREM 1. *If $u(x, y)$ is continuous with its first partial derivatives in a plane region R , and if, for every circle C contained in R ,*

$$\int_C \frac{\partial u}{\partial n} ds = 0,$$

where n is the exterior normal to C , then u is harmonic in R .

The generalization obtained is embodied in Theorem 2.

THEOREM 2. *If $v(x, y)$ is harmonic and positive in R , if $u(x, y)$ is continuous with its first partial derivatives in R , and if*

$$(1) \quad \int_C v \frac{\partial u}{\partial n} ds = \int_C u \frac{\partial v}{\partial n} ds$$

for every circle C contained in R , then u is harmonic in R .

Taking v as the constant one in Theorem 2, Theorem 1 is obtained.

Like Theorem 1, § Theorem 2 has an analog in space, but,

* Presented to the Society, April 3, 1931.

† Bôcher, M., *On harmonic functions in two dimensions*, Proceedings of the American Academy of Arts and Sciences, vol. 41 (1906), pp. 577-583.

‡ Koebe, P., *Herleitung der partiellen Differentialgleichung der Potentialfunktion aus der Intergraleigenschaft*, Sitzungsberichte der Berliner Mathematischen Gesellschaft, vol. 5 (1906), pp. 39-42.

§ Koebe, loc. cit. For generalizations of Bôcher's and Koebe's Theorem of another type, see G. C. Evans, *Fundamental points of potential theory*, Rice Institute Pamphlets, vol. 7 (1920), pp. 252-329, especially p. 286, and *Note on a theorem of Bôcher*, American Journal of Mathematics, vol. 50 (1928), pp. 123-126; and G. E. Raynor, *On the integro-differential equation of the Bôcher type in three space*, this Bulletin, vol. 52 (1926), pp. 654-658. Evans, using the notion of the potential function of a gradient vector, shows that the conclusion of Theorem 1 holds with much lighter hypotheses both on u and the character of the curves C .

since no new essentially different details present themselves in the proof for space, we simply state this analog, and consider in detail only the plane case.

THEOREM 3. *If $v(x, y, z)$ is harmonic and positive in a region R in space, if $u(x, y, z)$ is continuous with its first partial derivatives in R , and if, for every sphere C contained in R ,*

$$\iint_C v \frac{\partial u}{\partial n} ds = \iint_C u \frac{\partial v}{\partial n} ds,$$

where n is the exterior normal to C , then u is harmonic in R .

The proof of Theorem 2 is elementary in character. The idea is to express uv as a sum of integrals and deduce the character of u from the properties of these integrals.

2. Proof of Theorem 2. We first observe that it is enough to prove the theorem in the case that R is the interior of a circle C , and the hypotheses hold in the interior R' and on the boundary* of a circle C' concentric with C but of larger radius. The problem, then, is to show that u_{xx} and u_{yy} exist and are continuous in R , and that

$$(2) \quad \nabla^2 u = u_{xx} + u_{yy} = 0$$

there.

Let $P(x, y)$ be any point in R . Let α' be the radius of C' , α the radius of C , and

$$\rho = \frac{1}{2}(\alpha' - \alpha).$$

Then, by (1), the hypothesis on v , and a classical formula, we have, for $0 < t \leq \rho$,

$$(3) \quad \int_{C(P, t)} v \frac{\partial u}{\partial n} ds = \int_{C(P, t)} u \frac{\partial v}{\partial n} ds = \iint_{\sigma(P, t)} \phi d\sigma,$$

where $\sigma(P, t)$ is the interior of the circle $C(P, t)$ of radius t about P , and

$$\phi(x, y) = \nabla u \cdot \nabla v = u_x v_x + u_y v_y.$$

* That is to say, the hypotheses hold in a region containing R' and its boundary.

It is the third integral in (3) that enables us to express uv as a sum of integrals, whose properties lead to the conclusion of the theorem. Writing

$$\sigma = \sigma(P, \rho), \quad S' = R' + C', \quad r = \{(x - \xi)^2 + (y - \eta)^2\}^{1/2},$$

we find that

$$\begin{aligned} & \pi\rho^2 u(x, y)v(x, y) \\ &= \left\{ \iint_{\sigma} uv d\sigma + \rho^2 \left(\frac{1}{2} + \log \frac{1}{\rho} \right) \iint_{\sigma} \phi d\sigma \right. \\ (4) \quad & \left. - \frac{1}{2} \iint_{\sigma} \phi(\xi, \eta) r^2 d\xi d\eta - \rho^2 \iint_{S'-\sigma} \phi \log r d\sigma \right\} \\ & \quad + \rho^2 \iint_{S'} \phi \log r d\sigma \\ &= J'(P) + J''(P), \text{ say.} \end{aligned}$$

In fact, by (3),

$$\begin{aligned} \int_0^{\rho} \tau d\tau \int_0^{\tau} \frac{dt}{t} \int_{C(P,t)} v \frac{\partial u}{\partial n} ds &= \int_0^{\rho} \tau d\tau \int_0^{\tau} \frac{dt}{t} \int_{C(P,t)} u \frac{\partial v}{\partial n} ds \\ &= \int_0^{\rho} \tau d\tau \int_0^{\tau} \frac{dt}{t} \iint_{\sigma(P,t)} \phi d\sigma, \end{aligned}$$

or $K_1(P) = K_2(P) = K_3(P)$, say. Introducing, then, a system of polar coordinates (r, θ) with pole at P , we have

$$\begin{aligned} K_1(P) &= \int_0^{\rho} \tau d\tau \int_0^{2\pi} d\theta \int_0^{\tau} v \frac{\partial u}{\partial r} dr \\ &= \int_0^{\rho} \tau d\tau \int_0^{2\pi} \left\{ uv - u(x, y)v(x, y) - \int_0^{\tau} u \frac{\partial v}{\partial r} dr \right\} d\theta \\ &= \iint_{\sigma} uv d\sigma - \pi\rho^2 u(x, y)v(x, y) - K_2(P), \end{aligned}$$

so that

$$\pi\rho^2 u(x, y)v(x, y) = \iint_{\sigma} uv d\sigma - 2K_3(P).$$

But, we see that

$$\begin{aligned}
 K_3(P) &= \int_0^{2\pi} d\theta \int_0^\rho \tau d\tau \int_0^\tau \frac{dt}{t} \int_0^t r\phi dr \\
 &= \int_0^{2\pi} d\theta \int_0^\rho \left\{ \frac{1}{2}\rho^2(\log \rho - \frac{1}{2}) + \frac{1}{4}r^2 - \frac{1}{2}\rho^2 \log r \right\} r\phi dr,
 \end{aligned}$$

upon changing the order of integration twice. Hence (4) follows.

Consider, now, the derivatives of J' . We have

$$\begin{aligned}
 J'(P) &= \int_{y-\rho}^{y+\rho} d\eta \int_{x-\psi}^{x+\psi} w d\xi + \left\{ \int_{-\alpha'}^{y-\rho} d\eta \int_{-\psi'}^{\psi'} + \int_{y-\rho}^{y+\rho} d\eta \int_{-\psi'}^{x-\psi} \right. \\
 &\quad \left. + \int_{y-\rho}^{y+\rho} d\eta \int_{x+\psi}^{\psi'} + \int_{y+\rho}^{\alpha'} d\eta \int_{-\psi'}^{\psi'} \right\} w' d\xi,
 \end{aligned}$$

where

$$\begin{aligned}
 w &= w(x, y; \xi, \eta) = u(\xi, \eta)v(\xi, \eta) + \phi(\xi, \eta) \left\{ \frac{1}{2}\rho^2 - \rho^2 \log \rho - \frac{1}{2}r^2 \right\}, \\
 w' &= w'(x, y; \xi, \eta) = -\rho^2\phi(\xi, \eta) \log r, \\
 \psi &= \psi(\eta, y) = \{\rho^2 - (\eta - y)^2\}^{1/2}, \psi' = \psi'(\eta) = (\alpha'^2 - \eta^2)^{1/2}.
 \end{aligned}$$

We see, then, by using the fact that u, v and ϕ are continuous in S' , and the formula for differentiation under the integral sign, that, for P in R, J'_x exists and is given by

$$\begin{aligned}
 J'_x &= \int_{y-\rho}^{y+\rho} \{w(x, y; x + \psi, \eta) - w(x, y; x - \psi, \eta)\} d\eta \\
 &\quad + \int_{y-\rho}^{y+\rho} \{w'(x, y; x - \psi, \eta) - w'(x, y; x + \psi, \eta)\} d\eta \\
 &\quad + \iint_\sigma w_x d\sigma + \iint_{S'-\sigma} w'_x d\sigma \\
 &= \int_{y-\rho}^{y+\rho} \{u(x + \psi, \eta)v(x + \psi, \eta) - u(x - \psi, \eta)v(x - \psi, \eta)\} d\eta \\
 &\quad - \iint_\sigma (x - \xi)\phi d\sigma - \rho^2 \iint_{S'-\sigma} \phi \frac{(x - \xi)}{r^2} d\sigma;
 \end{aligned}$$

which reduces to

$$(5) \quad J_x' = \iint_{\sigma} \{uv_x + vu_x - (x - \xi)\phi\} d\sigma - \rho^2 \iint_{S'_{-\sigma}} \phi \frac{(x - \xi)}{r^2} d\sigma,$$

upon writing

$$\begin{aligned} u(x + \psi, \eta)v(x + \psi, \eta) - u(x - \psi, \eta)v(x - \psi, \eta) \\ = \int_{x-\psi}^{x+\psi} \{uv_x + vu_x\} d\xi. \end{aligned}$$

From (5) and the continuity of u and v and their first partial derivatives, we deduce immediately that J_{xx}' and, as it is worth while noting for future purposes, J_{yy}' exist and are continuous in R . By analogy, J_{yy}' exists and is continuous in R .

Next, consider J'' . The existence of J_x'' and J_y'' can be inferred from (4) and the existence of J_x' , J_y' and the first partial derivatives of u and v . We wish to know, further, that J_{xx}'' and J_{yy}'' exist and are continuous in R . To prove this, we first observe that $u, v^{-1}, J_x', J_y', J_x'', J_y''$ satisfy uniform Hölder conditions* in any closed domain S'' bounded by a circle C'' contained in R . The first four of these functions have this property because their first partial derivatives are continuous in R and R is convex† and contains S'' , the last two because of a theorem of Dini.‡ We next observe that from this property of u, v^{-1}, \dots, J_y'' , it follows that ϕ satisfies a uniform Hölder condition in S'' , for

$$\phi = \{\nabla v \cdot \nabla J' + \nabla v \cdot \nabla J'' - \pi \rho^2 u \nabla v \cdot \nabla v\} / (\pi \rho^2 v),$$

and thus ϕ is equal to a combination of sums and products of functions each of which satisfies a uniform Hölder condition in

* A function $f(P)$, defined on a set E , satisfies a uniform Hölder condition on E if, P and Q being any two points of E ,

$$|f(P) - f(Q)| < A |PQ|^\lambda,$$

where A and λ are independent of P and Q , and $\lambda > 0$. Evidently, if $f_1(P)$ and $f_2(P)$ satisfy uniform Hölder conditions on E , $f_1 + f_2$ and $f_1 f_2$ have the same property.

† A region R is convex if each segment, whose end points lie in R , lies in R .

‡ U. Dini, *Sur la méthode des approximations successives pour les équations aux dérivées partielles du deuxième ordre*, Acta Mathematica, vol. 25 (1901), pp. 185–230. The function J'' is, of course, the potential function due to a distribution of continuous density $-\rho^2\phi$ over S' .

S'' . The existence and continuity of J_{xx}'' and J_{yy}'' can now readily be deduced. We write

$$\begin{aligned}
 J'' &= \rho^2 \iint_{S'-S''} \phi \log r d\sigma + \rho^2 \iint_{S''} \phi \log r d\sigma \\
 &= L' + L'', \text{ say.}
 \end{aligned}$$

Now L_{xx}' and L_{yy}' evidently exist and are continuous for P in the interior of S'' , while L_{xx}'' and L_{yy}'' exist and are continuous in S'' by a theorem of Hölder.* Thus J_{xx}'' and J_{yy}'' exist and are continuous in the interior of S'' . But C'' was arbitrary in R ; and hence it follows that J_{xx}'' and J_{yy}'' exist and are continuous in R .

The proof is now almost complete. Since $J_{xx}', J_{yy}', J_{xx}'', J_{yy}''$ exist and are continuous in R , and since v is positive and harmonic there, u_{xx} and u_{yy} exist and are continuous in R . It remains, then, only to prove that (2) holds. To prove this, we show that a contrary assumption leads to a contradiction. Suppose there is a point P in R at which $|\nabla^2 u| = 2\beta$ is different from zero. Then we can choose t so small that $\sigma(P, t)$ lies in R and $|\nabla^2 u| > \beta$ in $\sigma(P, t)$. Thus, since v exceeds a positive constant ϵ in $\sigma(P, t)$ and $\nabla^2 u$ is continuous there,

$$(6) \quad \left| \iint_{\sigma(P,t)} v \nabla^2 u d\sigma \right| \geq \epsilon \beta \text{ area } \sigma(P, t) > 0.$$

But

$$(7) \quad \iint_{\sigma(P,t)} v \nabla^2 u d\sigma = \int_{C(P,t)} v \frac{\partial u}{\partial n} ds - \int_{C(P,t)} u \frac{\partial v}{\partial n} ds = 0,$$

by the continuity of $\nabla^2 u$, our hypotheses, and Green's formula. In (6) and (7) we reach the desired contradiction. The proof is now complete.

HARVARD UNIVERSITY

* Otto Hölder, *Beiträge zur Potentialtheorie*, Inaugural Dissertation, Stuttgart, 1882, p. 17. Hölder considers the second partial derivatives of a volume distribution. The same type of analysis, however, holds for plane distributions.