

WIENER-HOPF OPERATORS AND ABSOLUTELY CONTINUOUS SPECTRA¹

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If A is a selfadjoint operator on a Hilbert space \mathfrak{H} with spectral resolution $A = \int \lambda dE_\lambda$, it is known that the set of elements x in \mathfrak{H} for which $\|E_\lambda x\|^2$ is an absolutely continuous function of λ is a subspace, $\mathfrak{S}_a(A)$, reducing A ; cf. Halmos [1, p. 104]. In case $\mathfrak{S}_a(A) = \mathfrak{H}$, A is said to be absolutely continuous. The following was proved in Putnam [4]; see also [5] and will be stated as a

LEMMA. *Let T be a bounded operator on a Hilbert space \mathfrak{H} and let*

$$(1) \quad T^*T - TT^* = C, \quad C \geq 0.$$

If $A = T + T^$, then $\mathfrak{S}_a(A) \supset \mathfrak{M}_T$, where \mathfrak{M}_T is the least subspace of \mathfrak{H} reducing T (that is, invariant under T and T^*) and containing the range of C .*

The above will be used to give a short proof of the absolute continuity of certain bounded selfadjoint Wiener-Hopf operators on $L^2(0, \infty)$. For an extensive account of Wiener-Hopf operators on the half-line see Krein [2].

Let $k(t)$, for $-\infty < t < \infty$, satisfy

$$(2) \quad k \in L^1(-\infty, \infty) \cap L^2(-\infty, \infty) \quad \text{and} \quad k(-t) = \bar{k}(t).$$

Then the operator T on $\mathfrak{H} = L^2(0, \infty)$ defined by

$$(3) \quad (Tf)(t) = \int_0^t k(s-t)f(s)ds, \quad 0 \leq t < \infty,$$

is bounded. (In fact, the hypothesis $k \in L^1(-\infty, \infty)$ alone implies the boundedness of T , even $\|T\| \leq \int_{-\infty}^{\infty} |k(t)| dt$; cf. Krein [2, pp. 201-202].) The adjoint T^* , which is given by

$$(4) \quad (T^*f)(t) = \int_t^{\infty} k(s-t)f(s)ds,$$

and the selfadjoint operator $A = T + T^*$, where

$$(5) \quad (Af)(t) = \int_0^{\infty} k(s-t)f(s)ds,$$

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are of course also bounded on $\mathfrak{S} = L^2(0, \infty)$. There will be proved the following

THEOREM. *Let $k(t)$ satisfy (2) and suppose that*

$$(6) \quad K(\lambda) \equiv \int_{-\infty}^{\infty} k(t)e^{i\lambda t}dt \neq 0 \text{ a.e.,} \quad -\infty < \lambda < \infty.$$

Then the bounded selfadjoint operator A of (5) is absolutely continuous.

PROOF. A calculation similar to that in Putnam [3, p. 517], shows that, for $f \in \mathfrak{S}$, $\|Tf\|^2 - \|T^*f\|^2 = \|Bf\|^2$, where T is defined in (3) and

$$(7) \quad (Bf)(t) = \int_0^{\infty} k(t+s)f(s)ds,$$

so that (1) holds with $C = B^*B$. It will be shown that the set \mathfrak{M}_T of the Lemma is, in the present case, the entire space $\mathfrak{S} = L^2(0, \infty)$, and hence the absolute continuity of A will follow.

For $f \in L^2(0, \infty)$, define the Fourier transform $\hat{f}(\lambda)$ and the functions $F_+(\lambda)$ and $F_-(\lambda)$ by

$$(8) \quad \hat{f}(\lambda) = \int_0^{\infty} e^{-i\lambda t}f(t)dt \equiv F_-(\lambda)$$

and

$$(9) \quad F_+(\lambda) = \int_0^{\infty} e^{i\lambda t}f(t)dt.$$

The space of elements F_+ is a subspace of $L^2(-\infty, \infty)$ and will be denoted by R_+ ; similarly, the space of elements F_- will be denoted by $R_- (= \bar{R}_+)$. (It is clear that R_+ [R_-] can be regarded as the space of Fourier transforms of functions in $L^2(-\infty, \infty)$ which are 0 on the left [right] half-line. Since orthogonality is preserved under Fourier transforms it follows in particular that $R_+ \perp R_-$.)

If $f \in -L^2(0, \infty)$ then

$$(Tf)^\wedge(\lambda) = \int_0^{\infty} e^{-i\lambda t} \left[\int_0^t k(s-t)f(s)ds \right] dt$$

and hence, on inverting the order of integration, $(Tf)^\wedge(\lambda) = \bar{K}_+(\lambda)\hat{f}(\lambda)$, where $K_+(\lambda)$ is defined by

$$(10) \quad K_+(\lambda) = \int_0^{\infty} e^{i\lambda t}k(t)dt.$$

(It may be noted that $\int_0^t k(s-t)f(s)ds$ is the convolution of \bar{k} and f on $0 \leqq t < \infty$.) More generally, an iteration shows that

$$(11) \quad (T^n f)^\wedge(\lambda) = \bar{K}_+^n \hat{f}(\lambda), \quad n = 0, 1, 2, \dots$$

Let $g \perp \mathfrak{H}_{B^*B}$, so that $Bg = 0$, that is $\int_0^\infty k(t+s)g(s)ds$ is the 0 element of $\mathfrak{S} = L^2(0, \infty)$. Since, by (2), $k(t+s)$ belongs to $L^2(0, \infty)$ for any fixed t , it follows that the last integral is a continuous function of t on $0 \leqq t < \infty$ and that

$$(12) \quad m_c \in \mathfrak{H}_{B^*B}, \quad \text{where } m_c(t) = \bar{k}(t+c) \quad \text{and } c \geqq 0.$$

It is readily verified that

$$(13) \quad \hat{m}_c(\lambda) = e^{i\lambda c} \int_c^\infty e^{-i\lambda t} \bar{k}(t) dt.$$

Also, in view of the definition of $K(\lambda)$ in (6), one has

$$(14) \quad K(\lambda) = K_+(\lambda) + \bar{K}_+(\lambda),$$

where $K_+(\lambda)$ is defined in (10).

In order to prove that $\mathfrak{S}_a(A) = \mathfrak{S} (= L^2(0, \infty))$, it is sufficient, as noted above, to prove that $\mathfrak{M}_T = \mathfrak{S}$. Now, if $\mathfrak{M}_T \neq \mathfrak{S}$, then there exists a function $q \in \mathfrak{S}$ such that $q \neq 0$ and $q \perp \mathfrak{M}_T$. Let $Q = Q(\lambda)$ denote the Fourier transform of q , so that

$$(15) \quad Q(\lambda) = \int_0^\infty e^{-i\lambda t} q(t) dt \quad (\in R_-).$$

In view of (12), it follows from the relation $q \perp M_T$ and the fact that orthogonality is preserved under Fourier transforms that

$$(16) \quad Q \perp (T^n f)^\wedge(\lambda), \quad n = 0, 1, 2, \dots, \quad \text{where } f(t) = m_c(t), \quad c \geqq 0.$$

Thus, by (11) and (13), $Q \perp \bar{K}_+^n e^{i\lambda c} \int_c^\infty e^{-i\lambda t} \bar{k}(t) dt$, for $c \geqq 0$, that is,

$$(17) \quad Q \perp \bar{K}_+^{n+1} e^{i\lambda c} - \bar{K}_+^n e^{i\lambda c} \int_0^c e^{-i\lambda t} \bar{k}(t) dt \quad (n = 0, 1, 2, \dots).$$

Since Q and $e^{i\lambda c} \int_0^c e^{-i\lambda t} \bar{k}(t) dt$ belong to R_- and R_+ respectively, it follows from (17) for $n=0$ that $Q \perp \bar{K}_+ e^{i\lambda c}$ (therefore $Q \perp \bar{K}_+ R_+$) and hence, by induction, that

$$(18) \quad Q \perp e^{i\lambda c} \bar{K}_+^n, \quad n = 1, 2, \dots, \quad c \geqq 0.$$

Relations (14) and (18) and the fact that $Q \perp R_+$ imply that $Q \perp K^n(\lambda) R_+$ for $n = 0, 1, 2, \dots$, that is,

$$(19) \quad \int_{-\infty}^{\infty} K^n(\lambda) F_+(\lambda) \bar{Q}(\lambda) d\lambda = 0 \quad \text{for } n = 0, 1, 2, \dots,$$

where $K(\lambda)$ is given in (6) and $F_+(\lambda)$ is an arbitrary element of R_+ .

Since $k(t) \in L^1(-\infty, \infty)$, the function $K(\lambda)$ is continuous and satisfies

$$(20) \quad K(\lambda) \rightarrow 0 \quad \text{as } |\lambda| \rightarrow \infty.$$

Also, by (2), $K(\lambda)$ is real. Let $f(K)$ denote the characteristic function of the K -set: $|K| \geq 1/n$, where n is a positive integer. It follows from (19), Weierstrass' approximation theorem, and the fact that $F_+ \bar{Q}$ is in $L^1(-\infty, \infty)$, that $\int_{-\infty}^{\infty} f(K(\lambda)) F_+(\lambda) \bar{Q}(\lambda) d\lambda = 0$ and hence

$$(21) \quad \int_{E_n} F_+(\lambda) \bar{Q}(\lambda) d\lambda = 0,$$

where $E_n = \{\lambda: |K(\lambda)| \geq 1/n\}$. Since, for $c \geq 0$, $e^{i\lambda c} \bar{Q}(\lambda)$ is in R_+ , one can choose F_+ in (21) to be $e^{i\lambda c} \bar{Q}$ and so $\int_{E_n} e^{i\lambda c} \bar{Q}^2 d\lambda = 0$. In view of (20),

$$(22) \quad E_n \text{ is a bounded set.}$$

Since \bar{Q}^2 is in $L^1(-\infty, \infty)$, one can therefore differentiate under the last integral with respect to c and let $c \rightarrow 0+$ to obtain $\int_{E_n} \lambda^m \bar{Q}^2 d\lambda = 0$ ($m = 0, 1, 2, \dots; n = 1, 2, \dots$). Again, using (22) and Weierstrass' theorem, one concludes that $Q(\lambda) = 0$ a.e. on E_n . In virtue of (6), the set $\bigcup_{n=1}^{\infty} E_n$ differs from $(-\infty, \infty)$ by a set of measure zero and hence $Q(\lambda) = 0$ a.e. on $(-\infty, \infty)$. This implies that $q(t) = 0$ a.e. on $0 \leq t < \infty$, a contradiction. Thus $\mathfrak{M}_T = \mathfrak{S}$ and so $\mathfrak{S}_a(A) = \mathfrak{S}$ as was to be proved.

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