

On comparing the second and fourth rows of Table 3 we see that for the first four characteristic values the results obtained from the modified Method II are better than those given by the unmodified one for $N = 9$. Other computations have shown that this result does not hold for $m \neq 2$ if the class of admissible function is defined so as to have continuous derivatives at the points $x = x_i$ and if the matrix D is evaluated as in Method II.

In the third row of Table 3 the matrix D is evaluated as in Method II where the admissible functions are defined as for the first two rows.

The numerical results, although obtained only for a particular simple example, would seem to justify the use of the methods discussed above for the solution of more complicated Sturm-Liouville problems. They should be supplemented by a test for the accuracy of the results obtained for the characteristic values and characteristic functions applicable to problems for which the solution is not known. Various such tests are under investigation and will be reported on subsequently.

Digital Computer Laboratory
University of Illinois
Urbana, Illinois

University of California
Goleta, California

Digital Computer Laboratory
University of Illinois
Urbana, Illinois

C. C. FARRINGTON

R. T. GREGORY

A. H. TAUB

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On the Numerical Evaluation of the Stokes' Stream Function

1. Introduction. In the study of axially symmetric problems in fluid dynamics, the Stokes' stream function, that is, the function which satisfies

$$(1.1) \quad u_{zz} - \frac{1}{y} u_y + u_{yy} = 0, \quad y \neq 0,$$

is of considerable interest. This function is constant on the streamlines.

The problem to be considered is a Dirichlet type problem. Let G be a closed, bounded, simply connected plane region whose interior is denoted by R and whose boundary curve is denoted by S . Let G not contain any point where $y = 0$.

tion Procedure, matrix inversion, relaxation, iteration, gradient methods, and other techniques (Geiringer [4], Greenspan [6], Milne [7], Shortley et al [8], NBS [9], Young [13, 14]).

The author is indebted to D. M. Young for valuable criticisms and suggestions.

2. General Method. Suppose G_h consists of n points. Number these points in a one-to-one fashion with the integers $1, 2, 3, \dots, n$. Denote the coordinates of the point numbered k by (x_k, y_k) and the unknown stream function at (x_k, y_k) by $u(x_k, y_k) \equiv u_k$, for $k = 1, 2, \dots, n$.

Let (x_i, y_i) be an arbitrary point of S_h , the lattice boundary. Approximate u , by $g(x', y')$, where (x', y') is the nearest point of S to (x_i, y_i) . If (x', y') is not unique, choose any one of the set of nearest points and use it. The problem of

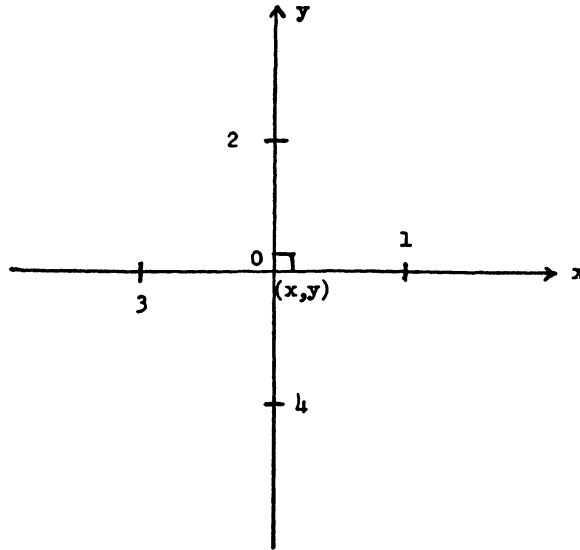


DIAGRAM 2

finding numerical approximations to $u(x, y)$ on the lattice boundary is, though crudely done, adequate for present purposes. For diagram 1 this means that at the points 1, 2, 3, 4, 5, 7, 9, 10, 14, 15, 19, 20, 21, 22, 23 the values of the stream function have been approximated by values at nearby points of S .

We then require that at each point (x_i, y_i) of R_h , the function u satisfy

$$(2.1) \quad 4u(x_i, y_i) - u(x_i + h, y_i) - \left(1 - \frac{h}{2y_i}\right)u(x_i, y_i + h) \\ - u(x_i - h, y_i) - \left(1 + \frac{h}{2y_i}\right)u(x_i, y_i - h) = 0.$$

It is convenient, in practice, to use the subscript notation. For example, for diagram 1, if $i = 12$, (2.1) becomes

$$4u_{12} - u_{13} - \left(1 - \frac{h}{2y_{12}}\right)u_7 - u_{11} - \left(1 + \frac{h}{2y_{12}}\right)u_{17} = 0.$$

Note that $u_{12}, u_{13}, u_{11}, u_{17}$ are unknowns while u_7 is a constant determined by the method for points of S_h . Application of (2.1) to each point of R_h thereby results in a system of linear equations which, when solved, yields the remaining numerical approximations.

Equation (2.1) is only one of a variety of possible difference equations which can be developed, as described in the next section.

3. Development of a General Difference Equation. We shall seek a solution $u(x, y)$ which is of class C^4 and use the notation $\partial^{i+j}u/\partial x^i\partial y^j \equiv u_{i,j}, i + j \leq 4$, and $A_{ij} \equiv u_{i,j}/i!j!$.

Let (x, y) be an arbitrary point of R_h . The aim will be to develop a "5-point formula," so consider here only (x, y) and its four adjacent points. Let $(x, y), (x + h, y), (x, y + h), (x - h, y),$ and $(x, y - h)$ be denoted, respectively, by 0, 1, 2, 3, 4, as in diagram 2. Hence $(x_0, y_0) = (x, y), (x_1, y_1) = (x + h, y), (x_2, y_2) = (x, y + h), (x_3, y_3) = (x - h, y), (x_4, y_4) = (x, y - h)$. Also let $O(h^n)$ represent any and all functions of order at least n in h .

In order to deduce a difference equation which approximates (1.1), let

$$(3.1) \quad L(u) = \sum_0^4 \alpha_i u_i.$$

By use of the finite Taylor series expansions about (x_0, y_0) for the various u_i and by the definition of A_{ij} , (3.1) becomes

$$(3.2) \quad L(u) = A_{00}(\alpha_0 + \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4) + h[A_{10}(\alpha_1 - \alpha_3) + A_{01}(\alpha_2 - \alpha_4)] + h^2[A_{20}(\alpha_1 + \alpha_3) + A_{02}(\alpha_2 + \alpha_4)] + h^3[A_{30}(\alpha_1 - \alpha_3) + A_{03}(\alpha_2 - \alpha_4)] + O(h^4)$$

Equation (1.1) may be rewritten as $yu_{2,0} - u_{0,1} + yu_{0,2} = 0$. Since u is assumed to be of class C^4 , further differentiation of (1.1) yields

$$yu_{3,0} - u_{1,1} + yu_{1,2} = 0$$

$$u_{2,0} + yu_{2,1} + yu_{0,3} = 0.$$

Rewriting these in terms of the A_{ij} 's and solving, one finds

$$(3.3) \quad A_{20} = \frac{A_{01}}{2y} - A_{02}; \quad A_{30} = \frac{A_{11}}{6y} - \frac{A_{12}}{3}; \quad A_{21} = -3A_{03} + \frac{A_{02}}{y} - \frac{A_{01}}{2y^2}.$$

Substitution of (3.3) into (3.2) yields

$$(3.4) \quad L(u) = A_{00}(\alpha_0 + \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4) + h \left[A_{01} \left(\frac{\alpha_1 h}{2y} + \frac{\alpha_3 h}{2y} + \alpha_2 - \alpha_4 \right) + A_{10}(\alpha_1 - \alpha_3) \right] + h^2[A_{02}(\alpha_2 + \alpha_4 - \alpha_1 - \alpha_3)] + h^3[A_{30}(\alpha_1 - \alpha_3) + A_{03}(\alpha_2 - \alpha_4)] + O(h^4).$$

Let

$$\begin{aligned}
 (3.5) \quad & \alpha_0 + \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4 = \epsilon_1 \\
 & \frac{\alpha_1 h}{2y} + \frac{\alpha_3 h}{2y} + \alpha_2 - \alpha_4 = \epsilon_2 \\
 & \alpha_1 - \alpha_3 = \epsilon_3 \\
 & \alpha_2 + \alpha_4 - \alpha_1 - \alpha_3 = \epsilon_4 \\
 & \alpha_2 - \alpha_4 = \epsilon_5,
 \end{aligned}$$

where

$$\epsilon_1 = O(h^4), \quad \epsilon_2 = O(h^3), \quad \epsilon_3 = O(h^3), \quad \epsilon_4 = O(h^2), \quad \epsilon_5 = O(h).$$

The solution to system (3.5) is

$$\begin{aligned}
 (3.6) \quad & \alpha_1 = \frac{\epsilon_2 y}{h} - \frac{\epsilon_5 y}{h} + \frac{\epsilon_3}{2} \quad \alpha_2 = \frac{\epsilon_2 y}{h} + \frac{\epsilon_4}{2} + \frac{\epsilon_5}{2} - \frac{\epsilon_6 y}{h} \\
 & \alpha_3 = \frac{\epsilon_2 y}{h} - \frac{\epsilon_5 y}{h} - \frac{\epsilon_3}{2} \quad \alpha_4 = \frac{\epsilon_2 y}{h} + \frac{\epsilon_4}{2} - \frac{\epsilon_5}{2} - \frac{\epsilon_6 y}{h} \\
 & \alpha_0 = \epsilon_1 - \epsilon_4 + 4 \frac{\epsilon_5 y}{h} - 4 \frac{\epsilon_2 y}{h}.
 \end{aligned}$$

Finally, setting (3.1) equal to zero and substituting (3.6) in (3.1) yield the general difference equation (3.7) which approximates the differential equation (1.1)

$$\begin{aligned}
 (3.7) \quad & \left(\epsilon_1 - \epsilon_4 + \frac{4\epsilon_5 y}{h} - \frac{4\epsilon_2 y}{h} \right) u_0 + \left(\frac{\epsilon_2 y}{h} - \frac{\epsilon_5 y}{h} + \frac{\epsilon_3}{2} \right) u_1 \\
 & + \left(\frac{\epsilon_2 y}{h} + \frac{\epsilon_4}{2} + \frac{\epsilon_5}{2} - \frac{\epsilon_6 y}{h} \right) u_2 + \left(\frac{\epsilon_2 y}{h} - \frac{\epsilon_5 y}{h} - \frac{\epsilon_3}{2} \right) u_3 \\
 & + \left(\frac{\epsilon_2 y}{h} + \frac{\epsilon_4}{2} - \frac{\epsilon_5}{2} - \frac{\epsilon_6 y}{h} \right) u_4 = 0.
 \end{aligned}$$

The process, then, was to find α_i such that: $L(u) + O(h^4) \equiv 0$, where u satisfies (1.1), and to then use $L(u) = 0$ for the general approximation.

Equation (2.1), which will be used throughout, is deduced from (3.7) by letting $\epsilon_1 = \epsilon_2 = \epsilon_3 = \epsilon_4 = 0$, $\epsilon_5 = h$, by dividing through by y , and by noting $u_0 = u(x, y)$, $u_1 = u(x + h, y)$, $u_2 = u(x, y + h)$, $u_3 = u(x - h, y)$, $u_4 = u(x, y - h)$.

Other difference equations may be constructed by selecting the ϵ_i 's differently. One of these of particular interest is one which would yield a linear system which has a symmetric coefficient matrix. This may be done in the following manner. Number the points as in diagram 1 so that the numbers increase from left to right and also increase if one reads down any column of points. Apply difference equation

$$\begin{aligned}
 (3.8) \quad & u_0 \left(+ \frac{2}{y} + \frac{2}{2y + h} + \frac{2}{2y - h} \right) - \frac{u_1}{y} - \frac{u_3}{y} \\
 & - u_2 \left(\frac{2}{2y + h} \right) - u_4 \left(\frac{2}{2y - h} \right) = 0, \quad |2y| \neq h
 \end{aligned}$$

to the points of R_h . Start with the first, or highest, row of points; traverse the row from left to right; proceed to the next row; traverse the row from left to right; proceed to the next row, etc.; continue the process until all the points of R_h have been used. Equation (3.8) may be constructed by a numerical procedure for certain self-adjoint elliptic differential equations (Young [15]), or directly from (3.7) by letting

$$\epsilon_1 = 0, \quad \epsilon_2 = \frac{-h}{y^2} + \frac{4h}{(2y+h)(2y-h)}, \quad \epsilon_3 = 0,$$

$$\epsilon_4 = \frac{-8y}{(2y+h)(2y-h)} + \frac{2}{y}, \quad \epsilon_5 = \frac{+4h}{(2y+h)(2y-h)}.$$

4. A Theorem on an Error Bound. In this section, let $u(x, y)$ be the solution of the Dirichlet problem being considered and $U(x, y)$ the solution of the numerical method described in section 2. This means that

$$(4.1) \quad -4U_0 + U_1 + U_3 + \left(1 - \frac{h}{2y}\right)U_2 + \left(1 + \frac{h}{2y}\right)U_4 = 0$$

for any arrangement of points of G_h like that of diagram 2.

LEMMA 1. *The solution of the system of linear equations which results by application of the process of section 2 with $y \neq 0$ and $\left|\frac{h}{y}\right| < 1$ is unique.*

Proof. It is sufficient to show that the determinant of the system of linear equations is not zero and this is done by demonstrating that the only solution of the homogeneous system which results by considering $g(x, y) \equiv 0$ on S is the zero vector. Suppose then there exists a nontrivial solution for the homogeneous system. For some point of R_h , $U \neq 0$. Suppose $U > 0$. Let the largest value M occur at (x_0, y_0) . Then

$$(4.2) \quad 4U_0 = U_1 + U_3 + \left(1 - \frac{h}{2y}\right)U_2 + \left(1 + \frac{h}{2y}\right)U_4,$$

and

$$(4.3) \quad M = U_0 \geq U_i, \quad i = 1, 2, 3, 4.$$

Case 1. $U_2 = U_4$. Then (4.1) becomes

$$(4.4) \quad 4U_0 = U_1 + U_2 + U_3 + U_4.$$

But (4.3) and (4.4) imply $U_0 = U_1 = U_2 = U_3 = U_4 = M$. Hence one can apply the same argument at (x_1, y_1) and continue in a finite number of steps to show that U at a boundary point is equal to M . This is a contradiction since $g(x, y) \equiv 0$ on S and the method described for points of S_h implies U at all lattice boundary points is zero.

Case 2. $U_2 > U_4$. Let $U_4 = U_2 - k$, $k > 0$. Then

$$4U_0 = U_1 + 2U_2 + U_3 - k \left(1 + \frac{h}{2y}\right),$$

and

$$4U_0 < U_1 + 2U_2 + U_3, \quad \text{since } 1 + \frac{h}{2y} > 0.$$

But this is impossible since $U_1 \leq U_0$, $U_2 \leq U_0$, $U_3 \leq U_0$.

Case 3. $U_4 > U_2$. The discussion is analogous to that for Case 2 and leads to a contradiction.

Hence, in all three cases, the assumption that $U > 0$ yields a contradiction. Contradictions are similarly reached if $U < 0$ is assumed. Hence the only solution is the trivial solution. This completes the proof.

Using the same ideas and techniques first set forth by Gershgorin [5], we apply these by assuming in all that follows that $\sigma = h/2y$, $|\sigma| < 1$ and by defining

$$(4.5) \quad \lambda[u(x, y)] \equiv h^{-2} \left\{ -4u(x, y) + u(x + h, y) + u(x - h, y) \right. \\ \left. + \left(1 - \frac{h}{2y}\right)u(x, y + h) + \left(1 + \frac{h}{2y}\right)u(x, y - h) \right\}.$$

LEMMA 2. If $\lambda[v] \leq 0$ on R_h and $v \geq 0$ on S_h , then $v \geq 0$ on R_h .

Proof. Suppose $v < 0$ for some point of R_h . Then there exists a point (x_0, y_0) of R_h where v is a minimum. Hence

$$(4.6) \quad v(x_0, y_0) < 0, \text{ and } v(x_0, y_0) \leq v(x, y), \text{ for all } (x, y) \text{ of } R_h.$$

Now, $\lambda[v] \leq 0$, by assumption, so that, equivalently

$$(4.7) \quad v_0 \geq \frac{1}{4} \left[v_1 + \left(1 - \frac{h}{2y}\right)v_2 + v_3 + \left(1 + \frac{h}{2y}\right)v_4 \right].$$

But (4.6) and (4.7) imply, in a manner similar to that used in Lemma 1, that $v_0 = v_1 = v_2 = v_3 = v_4$. In a finite number of steps this leads to the result that $v < 0$ for some point of S_h , which is a contradiction. This completes the proof.

LEMMA 3. If $-\lambda[v_1] \geq \lambda[v_2]$ in R_h and $|v_1| \leq v_2$ on S_h , then $|v_1| \leq v_2$ on R_h .

Proof. $v_2 - v_1 \geq 0$ on S_h and $\lambda[v_2 - v_1] = \lambda[v_2] - \lambda[v_1]$. Since $\lambda[v_2] + |\lambda[v_1]| \leq 0$, then $\lambda[v_2] - \lambda[v_1] \leq \lambda[v_2] + |\lambda[v_1]| \leq 0$. Hence $\lambda[v_2 - v_1] \leq 0$. By Lemma 2, $v_2 - v_1 \geq 0$ on R_h .

Also, $v_2 + v_1 \geq v_2 - |v_1| \geq 0$ on S_h . Then, $\lambda[v_2 + v_1] = \lambda[v_2] + \lambda[v_1] \leq 0$. Hence $v_2 + v_1 \geq 0$ on R_h . But $v_2 + v_1 \geq 0$ on R_h and $v_2 - v_1 \geq 0$ on R_h imply $v_2 \geq |v_1|$ on R_h . This completes the proof.

LEMMA 4. If $|\lambda[v]| \leq A$ on R_h and if $|v| \leq B$ on S_h , and r is the radius of a properly selected circle, as described below, which contains G , then $|v| \leq \frac{A}{4}r^2 + B$.

Proof. Let

$$w(x, y) = \left[\frac{Ar^2}{4} \left\{ 1 - \frac{(x - x_0)^2 + (y - y_0)^2}{r^2} \right\} + B \right],$$

where $(x - x_0)^2 + (y - y_0)^2 = r^2$ is the equation of any circle containing G for which $y_0/y \geq 1$, for all y in G . At least one such circle exists since G is simply connected, bounded, closed and $y \neq 0$, which imply that y is of constant sign.

Direct calculation yields $\lambda[w] = -\frac{A}{2}(1 + y_0/y) \leq -A$. Also, $w \geq B$ on S_h .

Now, since $|\lambda[v]| \leq A$ on R_h , by assumption, and it has been shown that $\lambda[w] \leq -A$, or equivalently, $-\lambda[w] \geq A$, it follows that $-\lambda[w] \geq -A \geq \lambda[v]$ on R_h . Since $|v| \leq B$ on S_h , and $w \geq B$ on S_h , $w \geq |v|$. Hence, $-\lambda[v] \geq \lambda[w]$ on R_h and $w \geq |v|$ on S_h . By Lemma 3, $w \geq |v|$ on R_h , or $|v| \leq w \leq \frac{Ar^2}{4} + B$.

This completes the proof.

THEOREM. *If $u(x, y)$ is of class C^5 in a closed region G , $y \neq 0$ in G , u denotes the solution of the stream equation, U denotes the solution of the linear system which results from application of (2.1) in the method of section 2, $|\sigma| = \left| \frac{h}{y} \right| < 1$, then*

$$|U - u| \leq \frac{r^2}{4} \left[\frac{h^2 M_3}{3! \bar{y}} + \frac{h^2 M_4}{3!} + \frac{h^4 M_5}{2 \cdot 4! \bar{y}} \right] + 2hM_1,$$

where

$$M_1 = \text{Max}_{(x,y) \in G} [\text{Max}_{(x,y) \in G} |u_{1,0}|, \text{Max}_{(x,y) \in G} |u_{0,1}|], \quad M_3 = \text{Max}_{(x,y) \in G} |u_{0,3}|,$$

$$M_4 = \text{Max}_{(x,y) \in G} [\text{Max}_{(x,y) \in G} |u_{0,4}|, \text{Max}_{(x,y) \in G} |u_{4,0}|], \quad M_5 = \text{Max}_{(x,y) \in G} |u_{0,5}|,$$

$\bar{y} = GLB|y|$ in G , $r =$ radius of any circle of type described in Lemma 4.

Proof. Let $R = \lambda[u] - \left(u_{xx} - \frac{1}{y} u_y + u_{yy} \right)$. Substitution of the finite Taylor series expansions for u_1, u_2, u_3, u_4 in $\lambda[u]$ and use of the Mean Value Theorem yields

$$R = \frac{h^2}{4!} [u_{4,0}(\rho_1, y) + u_{0,4}(x, \rho_2) + u_{4,0}(\rho_3, y) + u_{0,4}(x, \rho_4)] - \frac{h^2}{3!y} u_{0,3}(x, y) + \frac{h^4}{2 \cdot 4!y} [u_{0,5}(x, \rho_5)],$$

where $\rho_2 < \rho_5 < \rho_4$. Hence

$$|R| \leq \frac{M_4 h^2}{3!} + \frac{M_3 h^2}{3! \bar{y}} + \frac{M_5 h^4}{2 \cdot 4! \bar{y}}.$$

Since $u_{xx} - \frac{1}{y} u_y + u_{yy} = 0$,

$$(4.8) \quad \left| \lambda[u] - \left(u_{xx} - \frac{1}{y} u_y + u_{yy} \right) \right| = |\lambda[u]| = |R| \leq \frac{h^2 M_3}{3! \bar{y}} + \frac{h^2 M_4}{3!} + \frac{h^4 M_5}{2 \cdot 4! \bar{y}}.$$

Also, for any point of S_h , U was taken as the value of $g(x', y')$ at the nearest point (x', y') on the boundary S , and $g(x', y') = u(x', y')$ on S . Thereby, for any point (x, y) of S_h ,

$$(4.9) \quad |U(x, y) - u(x, y)| = |g(x', y') - u(x, y)| = |u(x', y') - u(x, y)|,$$

where (x', y') is a point of S and $(x - x')^2 + (y - y')^2 \leq 2h^2$. Therefore

$$(4.10) \quad \begin{aligned} |U(x, y) - u(x, y)| &= |u(x', y') - u(x, y)| \\ &\leq |u_x(\rho_6, y')(x - x')| + |u_y(x, \rho_7)(y - y')| \\ &\leq M_1[|y' - y| + |x' - x|] \leq 2hM_1. \end{aligned}$$

It must also be noted that $\lambda[U] = 0$, by (4.1). Hence

$$(4.11) \quad |\lambda[u - U]| = |\lambda[u] - \lambda[U]| = |\lambda[u]|.$$

Applying Lemma 4 to (4.8), (4.10), and (4.11), one finds that on R_h

$$(4.12) \quad |U - u| \leq \frac{r^2}{4} \left\{ \frac{h^2 M_3}{3! \bar{y}} + \frac{h^2 M_4}{3!} + \frac{h^4 M_6}{2 \cdot 4! \bar{y}} \right\} + 2hM_1.$$

This completes the proof.

COROLLARY. *Under the assumptions of the theorem just proved, $U \rightarrow u$, as $h \rightarrow 0$.*

Under general conditions set forth by Collatz [2], the term $2hM_1$ in the error bound for $|U - u|$ may be replaced by a term of type $O(h^2)$ by replacing the unesthetic method for deciding U on S_h as follows. Let (x_1, y_1) be a point of S , (x_2, y_2) of R_h , and (x_0, y_0) of S_h . Let PQ (diagram 3) be parallel to a coordinate axis and let δ be the absolute value of the distance between (x_0, y_0) and (x_1, y_1) . Choose U_0 by

$$U_0 = \frac{hg(x_1, y_1) + \delta U_2}{h + \delta}.$$

5. Use of Diagonal Points. After a numerical solution has been found on a given set G_h , solutions at other points may be approximated rapidly as follows. Let 5, 6, 7, 8 represent four points of G_h which are vertices of a square. (See

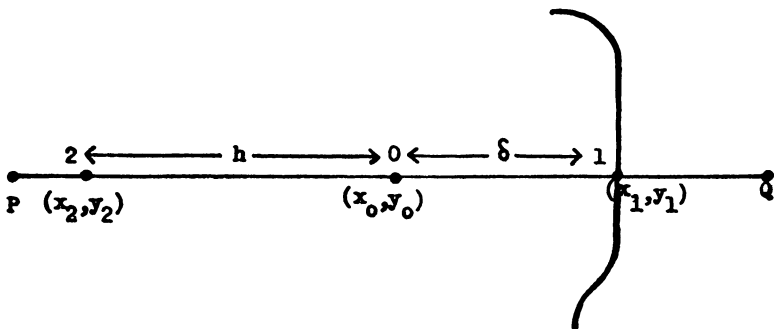


DIAGRAM 3

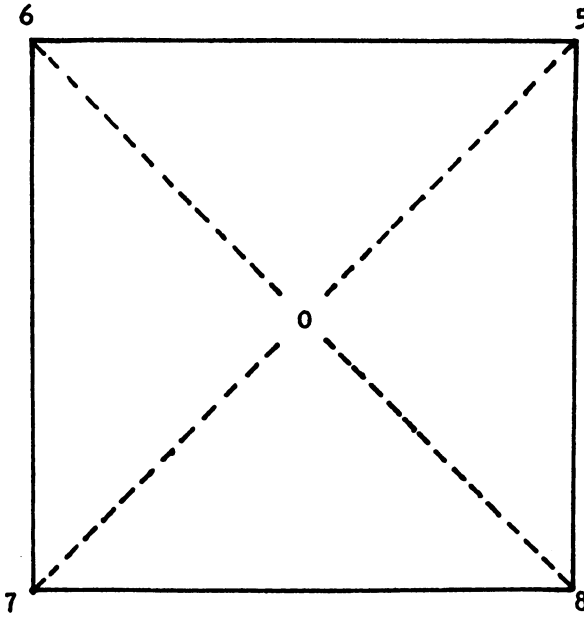


DIAGRAM 4

diagram 4.) Let 0 represent the diagonal point of the square. Use

$$(5.1) \quad U_0 = \frac{(2 - \sigma - \sigma^2)(U_5 + U_6) + (2 + \sigma - \sigma^2)(U_7 + U_8)}{4(2 - \sigma^2)}, \quad \sigma = \frac{h}{y_0}.$$

The unexciting calculations will here be replaced by a general outline of how to establish (5.1). Let $L(u) = \alpha_0 u_0 + \alpha_5 u_5 + \alpha_6 u_6 + \alpha_7 u_7 + \alpha_8 u_8$. Expand u_5, u_6, u_7, u_8 in Taylor series about (x_0, y_0) and substitute these expressions into $L(u)$. Simplify by replacing $u_{i,j}$ by $i!j!A_{ij}$. Apply identities (4.3) and combine as in (4.4). Set the coefficients equal to ϵ_i 's. This system of equations is

$$\begin{aligned} \alpha_0 + \alpha_5 + \alpha_6 + \alpha_7 + \alpha_8 &= \epsilon_1, & \alpha_5 - \alpha_6 - \alpha_7 + \alpha_8 &= \epsilon_2, & \alpha_5 + \alpha_6 - \alpha_7 - \alpha_8 &= \epsilon_5, \\ \alpha_5 \left(1 + \frac{\sigma}{2} - \frac{\sigma^2}{2}\right) + \alpha_6 \left(1 + \frac{\sigma}{2} - \frac{\sigma^2}{2}\right) & & & & & \\ & + \alpha_7 \left(-1 + \frac{\sigma}{2} + \frac{\sigma^2}{2}\right) + \alpha_8 \left(-1 + \frac{\sigma}{2} + \frac{\sigma^2}{2}\right) &= \epsilon_3, \\ \alpha_5 \left(1 + \frac{\sigma}{6}\right) - \alpha_6 \left(1 + \frac{\sigma}{6}\right) + \alpha_7 \left(1 - \frac{\sigma}{6}\right) - \alpha_8 \left(1 - \frac{\sigma}{6}\right) &= \epsilon_4, \end{aligned}$$

where $\epsilon_1 = O(h^4)$, $\epsilon_2 = O(h^3)$, $\epsilon_3 = O(h^3)$, $\epsilon_4 = O(h^2)$, $\epsilon_5 = O(h)$. A general formula can be established, but by letting $\epsilon_1 = \epsilon_2 = \epsilon_3 = \epsilon_4 = 0$, $\epsilon_5 = h$, solving for the α_i 's and using the method of section 3, one deduces (5.1).

6. Concluding Remarks. The method and ideas developed can be used in an analogous fashion to effect the numerical solution of problems involving any one of the class of equations

$$(6.1) \quad u_{xx} + \frac{k}{y} u_y + u_{yy} = 0,$$

where k is a fixed constant. Such equations are confronted frequently in physics (Shortley et al [8]).

Finally, two problems remain which are worthy of note. Many physical situations allow y to be zero on the boundary S and a solution which satisfies (6.1) in R is desired. An effective technique in the sense that $U \rightarrow u$ as $h \rightarrow 0$ is needed. Also, a reasonable method for approximating the error in the numerical solution is wanting. The error bound (4.12) involves unknown quantities and the only work which has been done toward establishing error bounds which can be calculated has been for harmonic functions (Walsh and Young [11, 12]).

DONALD GREENSPAN

Hughes Aircraft Company
Culver City, California

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Tables of Values of 16 Integrals of Algebraic-Hyperbolic Type

This paper gives tables of values of the following 16 integrals of algebraic-hyperbolic type.

$$(1) \quad \begin{aligned} I_k &= \frac{2^k}{k!} \int_0^\infty \frac{x^k dx}{\sinh 2x \pm 2x}, & (k \geq 1) \\ I_{k^*} &= \frac{2^k}{k!} \int_0^\infty \frac{x^k dx}{\sinh 2x \pm 2x}, & (k \geq 3) \end{aligned}$$

$$(2) \quad \begin{aligned} II_k &= \frac{2^k}{k!} \int_0^\infty \frac{x^k e^{-2x} dx}{\sinh 2x \pm 2x}, & (k \geq 1) \\ II_{k^*} &= \frac{2^k}{k!} \int_0^\infty \frac{x^k e^{-2x} dx}{\sinh 2x \pm 2x}, & (k \geq 3) \end{aligned}$$