

Some Integrals of the Arctangent Function

By M. L. Glasser

Integrals of the form $\int_0^\infty (\tan^{-1} cz)^{2n} R(z) dz$, where $R(z)$ is an even rational expression in z , occur in the theory of localized magnetic moments in metals. Since integrals of this general type do not appear to be tabulated, we present here a method for evaluation as well as some interesting related material.

By partial fraction decomposition all integrals of the above type may be reduced to the form

$$(1) \quad I_n(a) = \int_0^\infty (\tan^{-1} cz)^{2n} (z^2 + a^2)^{-1} dz,$$

where a is not required to be real. Also, by a simple change of variable, only the case $c = 1$ need be considered. By writing this as half the integral from $-\infty$ to ∞ and making the substitution $z = \tan \theta/2$ this may be brought into the form $2^{-(2n+1)}(a^2 + 1)^{-2} \int_{-\pi}^{\pi} \theta^{2n} (1 + \lambda \cos \theta)^{-1} d\theta$, where $\lambda = (a^2 - 1)/(a^2 + 1)$. In terms of $z = e^{i\theta}$, this may be written

$$(2) \quad I_n(a) = (-1)^{n+1} i 2^{-2n} (a^2 - 1)^{-1} \int_{\Gamma_1} (z - z_0)^{-1} (z - z_1)^{-1} \ln^{2n} z dz$$

where Γ_1 is the contour $z = e^{i\theta}$, $-\pi < \theta < \pi$ and $z_0 = 1/z_1 = (1 - a)/(1 + a)$. For the moment we assume $0 < a < 1$ so $z_0 > 0$ and lies inside the unit circle. Closing Γ_1 by the loop $\Gamma_2: z = \rho e^{\pm i\pi}$, $0 < \rho < 1$, we trap the pole at z_0 and thus

$$(3) \quad I_n(a) = \frac{(-1)^n 2\pi}{2^{2n}(a^2 - 1)} \left\{ \frac{\ln^{2n} z_0}{(z_0 - z_1)} + \frac{1}{2\pi i} \int_{\Gamma_2} \frac{\ln^{2n} z dz}{(z - z_0)(z - z_1)} \right\}.$$

The integral remaining in (3) is

$$(4) \quad J_n = \frac{1}{2\pi i} \int_0^1 \frac{(\ln x - i\pi)^{2n} - (\ln x + i\pi)^{2n}}{(x + z_0)(x + z_1)} dz.$$

This can be reduced to a sum of Kummer's Lambda functions [1]

$$(5) \quad \Lambda_{n+1}(x) = \int_0^x \frac{\ln^n |u|}{1 + u} du,$$

for example. However, since the resulting formula is somewhat unwieldy, the method will be illustrated by the cases $n = 1, 2$. We have

$$(6) \quad \begin{aligned} J_1 &= -2(z_1 - z_0)^{-1} \int_0^1 \ln x \{ (x - z_0)^{-1} - (x + z_1)^{-1} \} dx \\ &= \frac{2z_0}{z_0^2 - 1} \left\{ (1/2) \ln z_0 \ln \left[\frac{(1 + z_0)^2}{z_0} \right] + \text{Li}_2 \left(\frac{z_0}{z_0 + 1} \right) - \text{Li}_2 \left(\frac{1}{z_0 + 1} \right) \right\} \end{aligned}$$

where $\text{Li}_2(x)$ is the Euler dilogarithm [1]. Using the relation $\text{Li}_2(x) + \text{Li}_2(1-x) = \pi^2/6 - \log x \log(1-x)$, we find

$$(7) \quad I_1(a) = (\pi/4a) \left\{ \frac{\pi^2}{6} - \ln^2\left(\frac{1+a}{2}\right) - 2 \text{Li}_2\left(\frac{1-a}{2}\right) \right\}.$$

Since both sides of this equation are analytic functions of a for $\text{Re } a > 0$, the result is valid for all positive a . In addition, taking the limit $a \rightarrow 0$ leads to the known result $I_1(0) = \pi \ln 2$. The dilogarithm can be evaluated in closed form for a number of special cases [1], which leads to the apparently new results

$$\begin{aligned} I_1(\sqrt{5}) &= (\pi/4\sqrt{5}) \left\{ (3\pi^2/10) - 2 \ln^2\left(\frac{1+\sqrt{5}}{2}\right) \right\}, \\ I_1(3) &= (\pi/12) \{ (\pi^2/3) - \ln^2 2 \}, \\ (8) \quad I_1(\sqrt{5}-2) &= \frac{\pi}{4(\sqrt{5}-2)} \left\{ \frac{\pi^2}{30} + \ln^2\left(\frac{\sqrt{5}+1}{2}\right) \right\}, \\ I_1(\sqrt{5}+2) &= \frac{\pi}{4(\sqrt{5}+2)} \left\{ \frac{11\pi^2}{30} - 5 \ln^2\left(\frac{1+\sqrt{5}}{2}\right) \right\}. \end{aligned}$$

Other than for the trivial case $a = 1$, these are the only real values of a for which $I_1(a)$ may be expressed in elementary terms. The derivative of $I_1(a)$ is related to entry 3.813(5) of Gradshteyn and Ryzhik's tables [2] so (7) could also be obtained from that result by integration.

The case $n = 2$ leads to

$$(9) \quad \int_0^\infty \frac{(\tan^{-1} z)^4}{z^2 + a^2} dz = (\pi/4a) \left\{ \frac{\pi^4}{40} + \pi^2 \text{Li}_2\left(\frac{a-1}{a+1}\right) - 6 \text{Li}_4\left(\frac{a-1}{a+1}\right) \right\}$$

in terms of tabulated functions [1]. Unfortunately, the tetralogarithm cannot be evaluated in closed form for many special values. The case $a \rightarrow 0$ leads to

$$(10) \quad \int_0^{\pi/2} x^4 \csc^2 x dx = \frac{\pi^3}{2} \ln 2 - \frac{9\pi}{4} \zeta(3).$$

This integral can be obtained from [2, Eq. 3.748.2, p. 418], which leads to

$$(11) \quad \int_0^{\pi/2} x^{p+1} \csc^2 x dx = (p+1)(\pi/2)^p \left\{ p^{-1} - 2 \sum_{k=1}^{\infty} (p+2k)^{-1} 2^{-2k} \zeta(2k) \right\}$$

in terms of the Riemann Zeta function. Thus, for $p = 3$ we have the interesting relation

$$(12) \quad \zeta(3) = \frac{4\pi^2}{9} \left\{ \sum_{k=1}^{\infty} \zeta(2k)/(4^k(2k+3)) + \frac{1}{2} \ln 2 - \frac{1}{6} \right\}.$$

In a similar way we can sum the series $\sum_{k=1}^{\infty} \zeta(2k) 2^{-2k}/(2k+p)$ for any odd p .

The method used here can also be extended to arbitrary integrals of the form $\int_0^\infty (\tan^{-1} z)^n R(z) dz$ where n and R are not required to be even. When symmetry is not invoked, however, the Cauchy principal part rather than the residue is involved.

Finally, it is emphasized that a is not restricted to real values, so that cases such as $R(z) = (1+z^2)^{-1}$ may also be treated. The polylogarithms of complex

argument have been studied in detail [1] and we also obtain results as

$$(13) \quad \int_0^\infty \frac{(\tan^{-1} z)^2}{(z^2 - 3) - 4i} dz = \frac{\pi(1 + 2i)}{20} \left\{ \frac{7\pi^2}{48} - \frac{1}{4} \ln^2 2 + \frac{\pi \ln 2}{4} - 2i\beta(2) \right\}$$

where $\beta(2)$ is Catalan's constant 0.915965 \dots . Taking the real and imaginary parts of both sides of (13) gives

$$(14) \quad \int_0^\infty \frac{(\tan^{-1} z)^2}{z^4 - 6z^2 + 25} dz = \frac{\pi}{40} \left\{ \frac{7\pi^2}{48} - \frac{1}{4} \ln^2 2 + \frac{\pi \ln 2}{4} - \beta(2) \right\},$$

$$\int_0^\infty \frac{z^2 (\tan^{-1} z)^2}{z^4 - 6z^2 + 25} dz = \frac{\pi}{8} \left\{ \frac{7\pi^2}{48} - \frac{1}{4} \ln^2 2 + \frac{\pi \ln 2}{4} + \beta(2) \right\}.$$

These are only a few of the special cases which can be expressed in closed form.

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Battelle Memorial Institute
Columbus, Ohio 43201

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2. I. S. GRADSHTEYN & I. M. RYZHIK, *Table of Integrals, Series and Products*, translated from Russian, Academic Press, New York, 1965. MR 33 #5952.