

Finite-Difference Methods and the Eigenvalue Problem for Nonselfadjoint Sturm-Liouville Operators*

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Abstract. In this paper we analyze the convergence of a centered finite-difference approximation to the nonselfadjoint Sturm-Liouville eigenvalue problem

$$(1) \quad \begin{aligned} \mathfrak{L}[u] &\equiv -[a(x)u']' - b(x)u' + c(x)u = \lambda u, \quad 0 < x < 1, \\ u(0) &= u(1) = 0 \end{aligned}$$

where \mathfrak{L} has smooth coefficients and $a(x) \geq a_0 > 0$ on $[0, 1]$. We show that the rate of convergence is $O(\Delta x^2)$ as in the selfadjoint case for a scheme of the same accuracy. We also establish discrete analogs of the *Sturm* oscillation and comparison theorems. As a corollary we obtain the result

$$(2) \quad \limsup_{M \rightarrow \infty; \Delta x \rightarrow 0; (M+1)\Delta x = 1} \left\{ \sum_{p=1}^M \frac{\|V^p\|_\infty}{\Lambda_p} \right\} < \infty$$

where $\Delta x = 1/(M + 1)$ is the mesh size and Λ_p, V^p are the characteristic pairs of L , the $M \times M$ matrix which approximates \mathfrak{L} , and V^p is normalized so that $\|V^p\|_2 = 1$.

1. Introduction. Many authors (e.g. [1], [6], [8], [9]) have studied the convergence of finite-difference methods for selfadjoint Sturm-Liouville eigenvalue problems. In this report we are concerned with the nonselfadjoint problem

$$(1.1) \quad \begin{aligned} \mathfrak{L}(u) &\equiv -[a(x)u']' - b(x)u' + c(x)u = \lambda u, \quad 0 < x < 1, \\ u(0) &= u(1) = 0 \end{aligned}$$

where $a(x) \geq a_0 > 0$, $c(x) \geq 0$, and $b(x)$ are all smooth functions. This problem has an infinite sequence of positive [12, p. 37] and distinct [13, p. 212] eigenvalues

$$0 < \lambda_1 < \lambda_2 < \lambda_3 < \dots$$

and a corresponding sequence of smooth eigenfunctions $u^1(x), u^2(x), u^3(x), \dots$ which we assume normalized so that

$$(1.2) \quad \int_0^1 |u^p|^2 dx = 1, \quad p = 1, 2, \dots$$

Of course, as is well known, the transformation

$$(1.3) \quad u(x) = \left[\exp \left(-\frac{1}{2} \int_0^x \frac{b(t)}{a(t)} dt \right) \right] v(x)$$

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puts (1.1) into the selfadjoint form

$$(1.4) \quad \begin{aligned} \hat{L}[v] &\equiv -(av')' + (c + \frac{1}{2} b' + \frac{1}{4} (b^2/a)v = \lambda v, \\ v(0) &= v(1) = 0. \end{aligned}$$

However, we consider the direct approximation of (1.1) by means of the finite-difference equations

$$(1.5) \quad \begin{aligned} - \frac{\{a_{k+1/2}(w_{k+1} - w_k) - a_{k-1/2}(w_k - w_{k-1})\}}{\Delta x^2} - \frac{b_k(w_{k+1} - w_{k-1})}{2\Delta x} \\ + c_k w_k = \Lambda w_k, \quad k = 1, 2, \dots, M, \\ w_0 = w_{M+1} = 0 \end{aligned}$$

where M is a large positive integer, $\Delta x = 1/(M + 1)$ is the mesh spacing and the notation g_k is used for $g(k \Delta x)$. Equivalently, we may write (1.5) as the finite-dimensional eigenvalue problem:

$$(1.6) \quad LW = \Lambda W$$

where W is the M component vector

$$W = \begin{bmatrix} w_1 \\ w_2 \\ \cdot \\ \cdot \\ \cdot \\ w_M \end{bmatrix}$$

and L the $M \times M$ tridiagonal matrix

$$(1.7) \quad L = \frac{1}{\Delta x^2} \begin{bmatrix} \alpha_1 & \beta_1 & & & 0 \\ \gamma_2 & \alpha_2 & \beta_2 & & \\ \cdot & \cdot & \cdot & \cdot & \\ \cdot & \cdot & \cdot & \cdot & \\ \cdot & \cdot & \cdot & \cdot & \\ 0 & & & & \beta_{M-1} \\ & & & \gamma_M & \alpha_M \end{bmatrix}$$

with

$$(1.8) \quad \begin{aligned} \alpha_k &= [a_{k+1/2} + a_{k-1/2}] + c_k \Delta x^2 \beta_k = -[a_{k+1/2} + b_k \Delta x/2] \quad \text{and} \\ \gamma_k &= [b_k \Delta x/2 - a_{k-1/2}] \quad k = 1, 2, \dots, M. \end{aligned}$$

We will show that the latter procedure preserves the rate of convergence, namely $O(\Delta x^2)$, which obtains in the selfadjoint case for a scheme of the same accuracy, (see [6]). This is Theorem 1.

The matrix L defined above will be shown to be similar to an *oscillation matrix*, by means of a diagonal transformation \tilde{D} . Using the basic theorem on oscillation matrices, (see [4], [5]) and the fact that the entries of \tilde{D} alternate in sign, one immediately has a discrete analog of the *Sturm Oscillation Theorem* [13, p. 212, Theorem

Hence,

$$\lim_{\Delta x \rightarrow 0, i \rightarrow \infty; i\Delta x = \bar{x}} [\log Q_i] = -\frac{1}{2} \int_0^{\bar{x}} \frac{b(t)}{a(t)} dt.$$

Similarly,

$$\lim_{\Delta x \rightarrow 0, i \rightarrow \infty; i\Delta x = x} [\log P_i] = \frac{1}{2} \int_0^{\bar{x}} \frac{b(t)}{a(t)} dt.$$

Consequently,

$$\lim d_i = \left[\exp \left(-\frac{1}{2} \int_0^{\bar{x}} \frac{b(t)}{a(t)} dt \right) \right] \leq K_0 < \infty$$

which shows both $\|D\|_2, \|D^{-1}\|_2$ remain bounded as $\Delta x \rightarrow 0, M \rightarrow \infty, (M + 1)\Delta x = 1$.

LEMMA 2. For Δx sufficiently small, the eigenvalues of L are strictly positive and they remain bounded away from zero as $M \rightarrow \infty, \Delta x \rightarrow 0, (M + 1)\Delta x = 1$.

Proof. For Δx sufficiently small, $\gamma_k, \beta_k < 0$. Hence if $L = (l_{ij})$ and $\Omega_i = \sum_{j \neq i} |l_{ij}|$, then

$$\Omega_i = (a_{i+1/2} + a_{i-1/2})/\Delta x^2$$

and $l_{ii} = (a_{i+1/2} + a_{i-1/2})/\Delta x^2 + c_i \geq \Omega_i$ since $c_i \geq 0$.

By Gershgorin's theorem, [7], the eigenvalues of L lie in the union of the discs $|z - l_{ii}| \leq \Omega_i$ in the complex plane. Hence if Λ is an eigenvalue of L , then $\Lambda \geq 0$ since Λ is real.

Now let l_h be the finite-difference operator corresponding to $-L$, i.e.

$$\begin{aligned} [l_h v]_k &\equiv - \left[\frac{(a_{k+1/2} + a_{k-1/2}) + c_k \Delta x^2}{\Delta x^2} \right] v_k + \left[\frac{a_{k+1/2} + b_k \Delta x/2}{\Delta x^2} \right] v_{k+1} \\ &\quad + \left[\frac{a_{k-1/2} - b_k \Delta x/2}{\Delta x^2} \right] v_{k-1}. \end{aligned}$$

Then, for sufficiently small Δx , l_h is of positive type [3, p. 181] and so satisfies the discrete maximum principle [16, p. 23, Lemma 2.3]. Consequently [16, p. 108, Theorem 7.1] if $w(k\Delta x), k = 0, 1, \dots, M + 1$ is an arbitrary real-valued mesh function, there exists positive constants K and δ such that if $0 < \Delta x < \delta$,

$$(2.2) \quad \|w\|_\infty \equiv \text{Max}_k |w_k| \leq \text{Max} \{|w_0|, |w_{M+1}|\} + K \|(l_h w)\|_\infty.$$

Now let $V = \{v_k\}_{k=1}^M$ be an eigenvector of L corresponding to Λ . We may assume V to be real. Defining $v_0 = v_{M+1} = 0, LV = \Lambda V$ is equivalent to

$$(2.3) \quad [l_h v]_k = -\Lambda v_k, \quad k = 1, \dots, M.$$

Hence, using (2.2) and the fact that $\Lambda \geq 0$,

$$\|v\|_\infty \leq K \|(l_h v)\|_\infty = \Lambda K \|v\|_\infty$$

i.e. $\Lambda \geq 1/K > 0$. Q.E.D.

COROLLARY. Let Γ be the $M \times M$ matrix given by

$$(3.6) \quad U^p = \sum_{j=1}^M \sigma_j DX^j$$

so that

$$LU^p = \sum_{j=1}^M \sigma_j LDX^j = \sum_{j=1}^M \sigma_j \Lambda_j DX^j$$

then

$$\tau = (\lambda_p - L)U^p = \sum_{j=1}^M \sigma_j (\lambda_p - \Lambda_j) DX^j$$

and

$$(3.7) \quad \sum_{j=1}^M \sigma_j^2 |\lambda_p - \Lambda_j|^2 = \|D^{-1}\tau\|_2^2 \leq \|D^{-1}\|_2^2 \|\tau\|_2^2 \leq K_1(p) \Delta x^4$$

where K_1 is a constant.

Now, the eigenvalues of L are distinct and converge to the corresponding distinct eigenvalues of \mathfrak{L} . It follows that

$$(3.8) \quad \inf_{j \neq p} \{|\lambda_p - \Lambda_j|\} \geq \omega_0 > 0$$

for all sufficiently small Δx . Hence, on using (3.7),

$$(3.9) \quad \sum_{j \neq p} \sigma_j^2 \leq K_1 \Delta x^4.$$

From (3.9), (3.6) we obtain

$$(3.10) \quad \sigma_p^2 = \|D^{-1}U^p\|_2^2 + O(\Delta x^4) \geq \omega_1 > 0$$

for all sufficiently small Δx .

Thus

$$(3.11) \quad |\lambda_p - \Lambda_p| \leq K_2(p) \Delta x^2.$$

Since $V^p = \beta DX^p$ for some β and $\|X^p\|_2 = 1$ we have

$$|\beta| = \|D^{-1}V^p\|_2.$$

On taking square roots in (3.10), we have

$$\sigma_p = \|D^{-1}U^p\|_2 + O(\Delta x^4)$$

and we may assume that σ_p and β have the same sign; hence using (3.1),

$$(3.12) \quad (\sigma_p - \beta) = O(\Delta x^4).$$

Writing $U^p - V^p = \sum_{j \neq p} \sigma_j DX^j + (\sigma_p - \beta) DX^p$ we have

$$(3.13) \quad \|D^{-1}(U^p - V^p)\|_2^2 = \sum_{j \neq p} \sigma_j^2 + (\sigma_p - \beta)^2 = O(\Delta x^4)$$

i.e.

$$(3.14) \quad \|U^p - V^p\|_2^2 \leq \|D\|_2^2 \|D^{-1}(U^p - V^p)\|_2^2 \leq K_3(p) \Delta x^4. \quad \text{Q.E.D.}$$

Notice that the above inequality also implies uniform convergence at the rate of $O(\Delta x)^{3/2}$.

4. Proof of Theorem 2.

LEMMA 3. Let $0 < \Lambda_1 < \dots < \Lambda_M$ be the eigenvalues of L . Then there exists a positive integer j_0 , independent of M , such that for $j_0 \leq j \leq M$ we have

$$(4.1) \quad K_1 j^2 \pi^2 \leq \Lambda_j \leq K_2 j^2 \pi^2, \quad K_1, K_2 \text{ positive constants.}$$

Proof. In the selfadjoint case this result may be found in Bückner [1]. In the present more general case we will need to estimate the off-diagonal elements of the matrix \hat{L} in Lemma 1.

With the notation of (1.8) let

$$(4.2) \quad q_k^2 = \gamma_{k+1} \beta_k = \left(a_{k+1/2} - \frac{b_{k+1} \Delta x}{2} \right) \left(a_{k+1/2} + \frac{b_k \Delta x}{2} \right), \quad k = 1, \dots, M - 1.$$

Since $b(x) \in C^1[0, 1]$, we have by the mean-value theorem,

$$(4.3) \quad q_k^2 = (a_{k+1/2})^2 [1 - 2\mu_k \Delta x^2 + O(\Delta x^3)]$$

where $2\mu_k = [b_k^2 + 2a_{k+1/2} b'(\xi_k)] / 4a_{k+1/2}$ for some ξ_k such that $k \Delta x < \xi_k < (k + 1) \Delta x$. Hence on taking square roots

$$(4.4) \quad q_k = a_{k+1/2} [1 - \mu_k \Delta x^2 + O(\Delta x^3)], \quad k = 1, \dots, M - 1.$$

We now proceed to estimate the quadratic form $\langle X, \hat{L}X \rangle$ where X is any complex M vector of norm 1. Defining $x_0 = x_{M+1} = 0$, and using (4.3), we may write

$$(4.5) \quad \begin{aligned} \langle X, \hat{L}X \rangle &= \Delta x \sum_{k=0}^M \frac{|x_k - x_{k+1}|^2}{\Delta x^2} + \Delta x \sum_{k=1}^M c_k |x_k|^2 \\ &\quad + 2\Delta x \sum_{k=0}^M \mu_k a_{k+1/2} x_k \bar{x}_{k+1} + O(\Delta x) \Delta x \sum_{k=0}^M x_k \bar{x}_{k+1}. \end{aligned}$$

Now let $0 < a_0 \leq a(x) \leq a_1$ on $[0, 1]$ and let

$$\|c\|_\infty = \text{Max}_k |c_k|, \quad \|\mu\|_\infty = \text{Max}_k |\mu_k|.$$

We have

$$(4.6) \quad \langle X, \hat{L}X \rangle \leq a_1 \Delta x \sum_{k=0}^M \frac{|x_{k+1} - x_k|^2}{\Delta x^2} + \|c\|_\infty + 2a_1 \|\mu\|_\infty + |O(\Delta x)|$$

and

$$(4.7) \quad \langle X, \hat{L}X \rangle \geq a_0 \Delta x \sum_{k=0}^M \frac{|x_{k+1} - x_k|^2}{\Delta x^2} - 2a_1 \|\mu\|_\infty - |O(\Delta x)|.$$

Let H be the tridiagonal $M \times M$ matrix defined by

$$(4.8) \quad H = \frac{1}{\Delta x^2} \begin{bmatrix} 2 & & -1 & & 0 \\ & \cdot & & \cdot & \\ -1 & & \cdot & & \cdot \\ & \cdot & & \cdot & \cdot \\ & & \cdot & & \cdot \\ & & & \cdot & -1 \\ 0 & & & -1 & 2 \end{bmatrix}.$$

It is easily verified that

$$(4.9) \quad \langle X, HX \rangle = \Delta x \sum_{k=0}^M \frac{|x_{k+1} - x_k|^2}{\Delta x^2}$$

and that the eigenvalues $\theta_j, j = 1, \dots, M$, of H , arranged in increasing order, are given by

$$(4.10) \quad \theta_j = \frac{4}{\Delta x^2} \sin^2 \frac{j\pi\Delta x}{2}, \quad j = 1, \dots, M.$$

Inserting (4.9) into (4.6), (4.7) and using the maximum principle for the eigenvalues of real symmetric matrices shows that

$$(4.11) \quad a_0\theta_j - 2\|\mu\|_\infty - |O(\Delta x)| \leq \Lambda_j \leq a_1\theta_j + \|c\|_\infty + 2a_1\|\mu\|_\infty + |O(\Delta x)|.$$

Using (4.10) and an elementary calculation, the proof follows from (4.11).

Proof of Theorem 2. Let

$$W^j = \begin{bmatrix} w_1^j \\ \cdot \\ \cdot \\ \cdot \\ w_M^j \end{bmatrix}$$

be an eigenvector of L corresponding to Λ_j . Then W^j satisfies the difference equations:

$$(4.12) \quad -\left[2 + \frac{(c_k - \Lambda_j)\Delta x^2}{\omega_k}\right]w_k^j + \left[\frac{a_{k+1/2} + b_k\Delta x/2}{\omega_k}\right]w_{k+1}^j + \left[\frac{a_{k-1/2} - b_k\Delta x/2}{\omega_k}\right]w_{k-1}^j = 0, \quad k = 1, \dots, M$$

where $w_0^j = w_{M+1}^j = 0$ and $\omega_k = \frac{1}{2} (a_{k+1/2} + a_{k-1/2})$.

Let

$$\begin{aligned} \tilde{\alpha}_k &= -\left[2 + \frac{(c_k - \Lambda_j)\Delta x^2}{\omega_k}\right], & \tilde{\beta}_k &= \left[\frac{a_{k+1/2} + \frac{1}{2} b_k\Delta x}{\omega_k}\right], \\ \tilde{\gamma}_k &= \left[\frac{a_{k-1/2} - \frac{1}{2} b_k\Delta x}{\omega_k}\right], \end{aligned}$$

and let A be the tridiagonal $M \times M$ matrix

$$(4.13) \quad A = \begin{bmatrix} \tilde{\alpha}_1 & \tilde{\beta}_1 & & & 0 \\ & \cdot & \cdot & & \\ & \tilde{\gamma}_2 & \cdot & \cdot & \\ & & \cdot & \cdot & \\ & & & \cdot & \cdot \\ & & & & \cdot \\ & & & & \cdot \\ & & & & \cdot \\ & & & & \cdot \\ 0 & & & \tilde{\gamma}_M & \tilde{\alpha}_M \end{bmatrix}.$$

Then we may write (4.12) as

$$(4.14) \quad AW^j = 0$$

or equivalently

$$(4.15) \quad (P^{-1}AP)P^{-1}W^j = 0$$

if P is any nonsingular matrix.

Choose P to be the diagonal matrix

$$(4.16) \quad P = \begin{bmatrix} p_1 & & & & 0 \\ & \cdot & & & \\ & & \cdot & & \\ & & & \cdot & \\ & & & & \cdot \\ 0 & & & & & p_M \end{bmatrix}$$

where $p_1 = 1$ and $p_i^2 = \prod_{k=1}^{i-1} (\tilde{\gamma}_{k+1}/\tilde{\beta}_k)$, $i = 2, \dots, M$.

For all sufficiently small Δx , $p_i^2 > 0$ and as in Lemma 1, P symmetrizes A . Let $\sigma_k = (\tilde{\gamma}_{k+1}\tilde{\beta}_k)^{1/2}$, then

$$(4.17) \quad P^{-1}AP = \begin{bmatrix} \tilde{\alpha}_1 & \sigma_1 & & & 0 \\ & \cdot & \cdot & & \\ \sigma_1 & & \cdot & \cdot & \\ & \cdot & \cdot & \cdot & \\ & & \cdot & \cdot & \sigma_{M-1} \\ 0 & & & \sigma_{M-1} & \tilde{\alpha}_M \end{bmatrix}.$$

Observe that by the mean-value theorem

$$(4.18) \quad \omega_k\omega_{k+1} = (a_{k+1/2})^2[1 + O(\Delta x^2)] \quad \text{as } \Delta x \rightarrow 0.$$

Also if $b(x) \in C^1[0, 1]$,

$$(4.19) \quad \begin{aligned} (\tilde{\gamma}_{k+1}\tilde{\beta}_k) &= \left[\frac{(a_{k+1/2})^2 + \frac{a_{k+1/2}(b_k - b_{k+1})\Delta x}{2} - \frac{b_k b_{k+1}\Delta x^2}{4}}{\omega_k\omega_{k+1}} \right] \\ &= \frac{(a_{k+1/2})^2[1 + O(\Delta x^2)]}{(a_{k+1/2})^2[1 + O(\Delta x^2)]} \quad \text{as } \Delta x \rightarrow 0. \end{aligned}$$

Hence,

$$(4.20) \quad \sigma_k = (\tilde{\gamma}_{k+1}\tilde{\beta}_k)^{1/2} = 1 + O(\Delta x^2) \quad \text{as } \Delta x \rightarrow 0.$$

Let $V = P^{-1}W^j$ and write the system (4.15) as

$$(4.21) \quad \begin{aligned} - \left[2 + \frac{(c_k - \Lambda_j)\Delta x^2}{\omega_k} \right] v_k + \sigma_k v_{k+1} + \sigma_k v_{k-1} &= 0, \\ v_0 = v_{M+1} &= 0, \end{aligned} \quad k = 1, \dots, M.$$

Let K_1 and K_2 be the constants in Lemma 3 and define

$$(4.22) \quad \beta_j^2 = \Lambda_j/K_2.$$

Let $y(x) = \sin \beta_j x$. Then $y_k = y(k \Delta x)$ satisfies the difference equations:

$$(4.23) \quad -[2 - \mu_j \Delta x^2]y_k + y_{k+1} + y_{k-1} = 0, \quad k = 1, 2, \dots$$

where

$$(4.24) \quad \mu_j = \frac{4}{\Delta x^2} \sin^2 \frac{\beta_j \Delta x}{2}.$$

The distance between successive zeros of $y(x)$ is $\pi/\beta_j = (K_2 \pi^2/\Lambda_j)^{1/2} \geq 1/j$ for j large enough by Lemma 3.

Let $v(x)$ be the piecewise-linear function corresponding to "graph" of vector $V = P^{-1}W^j$. Define the auxiliary function $z(x)$ by

$$z(x) = y(x)/v(x) \quad \text{whenever } v(x) \neq 0.$$

We proceed to estimate the distance between successive nodes of $v(x)$ by investigating the difference equation satisfied by $z(x)$.

We may assume that $\delta_{\max}(V) > 3\Delta x$; for if $\delta_{\max}(V) \leq 3\Delta x$, then in particular, $\delta_{\max}(V) \leq 3/(M+1) < 3/j \leq 3\pi(K_2/\Lambda_j)^{1/2}$ for all sufficiently large j . If $\delta_{\max} > 3\Delta x$, then there exists a set N of consecutive mesh points, containing at least three members on which $v(x)$ is strictly positive (or strictly negative). Let N' be N minus the two end points of N . Since $z_k = y_k/v_k$ for $k \in N'$,

$$(4.25) \quad [l_h z]_k \equiv - \left[\frac{(2 - \mu_j \Delta x^2) \sigma_k}{2 + (c_k - \Lambda_j) \Delta x^2 / \omega_k} (v_{k+1} + v_{k-1}) \right] z_k \\ + v_{k+1} z_{k+1} + v_{k-1} z_{k-1} = 0, \quad k \in N'.$$

We now show that for all sufficiently large j , the difference operator l_h (or $-l_h$ if v is strictly negative) occurring in (4.25) is of positive type, and hence satisfies the discrete maximum principle:

It is sufficient to show that if j is sufficiently large,

$$(4.26) \quad \frac{[2 - \mu_j \Delta x^2] \sigma_k}{2 + (c_k - \Lambda_j) \Delta x^2 / \omega_k} \geq 1, \quad \text{if } k \in N'.$$

From (4.24) we have $\mu_j \leq \Lambda_j/K_2 \leq \Lambda_j/2a_1$ if K_2 is chosen so that $K_2 \geq 2a_1$, where a_1 is an upper bound for $a(x)$ on $[0, 1]$. Hence,

$$(4.27) \quad (2 - \mu_j \Delta x^2) \sigma_k = 2 - \mu_j \Delta x^2 + O(\Delta x^2)$$

since $\mu_j \Delta x^2 \leq 4$ and $\sigma_k = 1 + O(\Delta x^2)$. Now,

$$2 - \mu_j \Delta x^2 + O(\Delta x^2) \geq 2 - \Lambda_j \Delta x^2 / K_2 + O(\Delta x^2) \\ \geq 2 - \Lambda_j \Delta x^2 / 2\omega_k + O(\Delta x^2) \\ = 2 + \frac{(c_k - \Lambda_j) \Delta x^2}{\omega_k} + \frac{(\Lambda_j - 2c_k) \Delta x^2}{2\omega_k} + O(\Delta x^2)$$

i.e.

$$(4.28) \quad (2 - \mu_j \Delta x^2) \sigma_k \geq 2 + (c_k - \Lambda_j) \Delta x^2 / \omega_k$$

if j is sufficiently large, since we assume $c(x)$ is bounded.

Furthermore, $2 + (c_k - \Lambda_j) \Delta x^2 / \omega_k$ is positive for $k \in N'$ since v_k, v_{k+1}, v_{k-1} have the same sign, on using (4.21). Thus (4.26) is satisfied.

Suppose now that $z(x)$ has two zeros in the interval spanned by N . At any mesh point lying between the two zeros we must have $z(x) = 0$ by the maximum principle. Since $z(x) = 0$ if and only if $y(x) = 0$, this means that the distance between successive zeros of $y(x)$ is $\leq \Delta x = 1/(M + 1)$. However, as already noted, this distance is $\geq 1/j$ and $j \leq M$.

Thus $y(x)$ has at most one zero in the interval spanned by N . Hence the maximum distance between successive nodes of $v(x)$ must be less than or equal to $\pi/\beta_j + 2\Delta x$. Since $\Lambda_j = O(1/\Delta x^2)$, we have

$$(4.29) \quad \delta_{\max}(V) \leq K(\Lambda_j)^{-1/2}.$$

A similar estimate is valid for the eigenvector W^j of L since $W^j = PV$ and P is a positive diagonal matrix. Q.E.D.

COROLLARY 1. *Let the eigenvectors $\{V^p\}$ of L be normalized so that $\|V^p\|_2 = 1$. Then there exists a constant K and an integer p_0 , both independent of M such that if $p_0 \leq p \leq M$*

$$(4.30) \quad \|V^p\|_\infty \equiv \text{Max}_{k=1 \dots M} |v_k^p| \leq Kp^{1/2}.$$

Proof. Let W^p be the normalized eigenvector of $\hat{L} = D^{-1}LD$ corresponding to Λ_p . Since $W^p = D^{-1}V^p/\|D^{-1}V^p\|_2$ and D^{-1} is a positive diagonal matrix, the distance between successive nodes of W^p satisfies an estimate similar to (4.29). Since W^p is normalized we have

$$(4.31) \quad \langle W^p, \hat{L}W^p \rangle = \Lambda_p.$$

Hence, using inequality (4.7) in the proof of Lemma 3, we get,

$$(4.32) \quad \Delta x \sum_{k=0}^M \frac{|w_{k+1}^p - w_k^p|^2}{\Delta x^2} \leq \frac{2\Lambda_p}{a_0}$$

for all sufficiently large p .

Let r, s be any two positive integers with $1 \leq s < r \leq M$. Then,

$$(4.33) \quad \begin{aligned} |w_r^p - w_s^p| &= \left| \Delta x \sum_{k=s}^{r-1} \frac{w_{k+1}^p - w_k^p}{\Delta x} \right| \\ &\leq [(r - s)\Delta x]^{1/2} \left(\Delta x \sum_{k=0}^M \frac{|w_{k+1}^p - w_k^p|^2}{\Delta x^2} \right)^{1/2} \\ &\leq [(r - s)\Delta x]^{1/2} (2\Lambda_p/a_0)^{1/2} \end{aligned}$$

on using Schwarz's inequality and (4.32). Now choose r so that $|w_r^p| = \|W^p\|_\infty > 0$ and let s be the integer nearest r with the property that $w_s^p w_r^p \leq 0$. (s need not necessarily be less than r .) We then have for sufficiently large p , by Theorem 2,

$$(4.34) \quad |(r - s)\Delta x| < 2\delta_{\max}(W^p) \leq K'(\Lambda_p)^{-1/2}.$$

Hence using (4.33), (4.34)

$$\begin{aligned} \|W^p\|_\infty &\leq |w_r^p - w_s^p| \leq [(r-s)\Delta x]^{1/2} (2\Lambda_p/a_0)^{1/2} \\ &\leq K''(\Lambda_p)^{1/4} \end{aligned}$$

for sufficiently large p and the proof follows from Lemma 3.

Remark. The estimate (4.30) was obtained by Bückner [1] in the selfadjoint case using an elementary device. It would be interesting to know whether or not the discrete eigenvectors display this growth as $M \rightarrow \infty$. In the case of the analytic problem (1.1) it is known (see [15, p. 334]) that the normalized eigenfunctions are uniformly bounded in the supremum norm.

COROLLARY 2. *Let $\{V^p\}_{p=1}^M$ be the eigenvectors of L normalized so that $\|V^p\|_2 = 1$, $p = 1, \dots, M$. Then,*

$$\limsup_{M \rightarrow \infty; \Delta x \rightarrow 0; (M+1)\Delta x = 1} \left\{ \sum_{p=1}^M \frac{\|V^p\|_\infty}{\Lambda_p} \right\} < \infty.$$

Proof. This follows immediately from Lemmas 2, 3, and Corollary 1.

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