

## NOTE ON THE HEAT EQUATION<sup>1</sup>

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If  $u(x, t)$  is the temperature of an infinite insulated rod, then

$$(1) \quad H(a, b; t) = \int_a^b u(x, t) dx$$

is the amount of heat at time  $t$  in the interval  $[a, b]$ . We prove an existence and uniqueness theorem for the case in which the heat, rather than the temperature, is prescribed initially. If the prescribed heat is not absolutely continuous, its derivative  $\partial H/\partial b$ , the initial temperature, may not determine  $H$ , so that we are dealing with a more general problem than the usual one.

The following lemma, which we did not find in the literature, will be useful.

LEMMA 1. *Let  $f(x)$  be measurable, and  $|f(x)| \leq Me^{cx^2}$ . Define*

$$(2) \quad W_r(x, t) = (4\pi t)^{-1/2} \frac{\partial^r}{\partial x^r} [e^{-x^2/4t}].$$

*Then, for every integer  $r \geq 0$ , the integral*

$$(3) \quad u_r(x, t) = \int_{-\infty}^{\infty} W_r(x - \xi, t) f(\xi) d\xi$$

*is absolutely convergent and satisfies*

$$(a) \quad \partial u_r / \partial x = u_{r+1}, \quad \partial u_r / \partial t = u_{r+2} \quad \text{on } 0 < t < 1/4c,$$

*and, for all  $t_0$  satisfying  $0 < t_0 < 1/4c$ ,*

$$(b) \quad |u_r(x, t)| \leq K t^{-r/2} e^{N x^2}, \quad \text{on } 0 < t \leq t_0,$$

*where  $K = K(r, t_0)$ , and  $N = N(t_0)$ .*

PROOF. It is well known that  $W_r(x, t)$  is analytic, satisfies the heat equation, and has the specific form

$$(4) \quad W_r(x, t) = \frac{1}{(2\pi)^{1/2}} (2t)^{-(r+1)/2} H_r \left( \frac{x}{(2t)^{1/2}} \right) e^{-x^2/4t},$$

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Received by the editors June 26, 1953.

<sup>1</sup> Work done at Harvard University under Contract N5ori-07634, with the Office of Naval Research.

where  $H_r(z)$  is the Hermite polynomial of degree  $r$  in  $z = x/(2t)^{1/2}$ . The absolute convergence of (3) on  $0 < t < 1/4c$  is obvious from (4) and  $|f(x)| \leq Me^{cx^2}$ .

To prove the first equation of (a), note

$$(5) \quad h^{-1}[u_r(x + h, t) - u_r(x, t)] = \int_{-\infty}^{\infty} W_{r+1}(x + \theta h - \xi, t)f(\xi)d\xi, \quad 0 \leq \theta \leq 1,$$

by the Law of the Mean. But, as  $|\xi| \rightarrow \infty$ ,

$$|W_{r+1}(x + \theta h - \xi, t)| = O(|x - \xi| + 1)^{r+1}e^{-(|x-\xi|+1)^2/4t},$$

and is bounded in finite intervals. Hence, since  $t < 1/4c$ , the integrand in (5) tends *dominatedly* [3, p. 168] to  $W_{r+1}(x - \xi, t)f(\xi)$  as  $h \rightarrow 0$ , which shows that

$$\lim_{h \rightarrow 0} h^{-1}[u_r(x + h, t) - u_r(x, t)] = \int_{-\infty}^{\infty} W_{r+1}(x - \xi, t)f(\xi)d\xi,$$

proving the first equation of (a). A similar proof applies to the second equation.

To prove (b), note that the right side of (3) is by (4) bounded by a sum of terms of the form

$$Mt^{-s} \int_{-\infty}^{\infty} |x - \xi|^s e^{-((x-\xi)^2/4t) + c\xi^2} d\xi \quad [s \leq r].$$

Completing the square, the exponent becomes

$$[(1 - 4ct)/4t][\xi - x/(1 - 4ct)]^2 + cx^2/(1 - 4ct).$$

Writing  $\eta = t^{-1/2}[\xi - x/(1 - rct)]$ , and noting that  $c/(1 - 4ct) < c/(1 - 4ct_0) = N'(t_0)$ , we bound the right side of (3) by a sum of terms

$$M'e^{N'x^2}t^{-s} \int_{-\infty}^{\infty} |\eta t + N'xt|^s e^{-(1-4ct)\eta^2} d\eta.$$

Since  $1 - 4ct \geq 1 - 4ct_0 = L > 0$ , another bound is

$$M'e^{N'x^2}t^{-s/2} \int_{-\infty}^{\infty} |\eta + N'xt^{1/2}|^s e^{-L\eta} d\eta.$$

For any  $N > N'(t_0)$  and suitable  $K = K(t_0, r)$ , this evidently implies (b).

**COROLLARY.** *In Lemma 1,  $u_r \in C^\infty$  and  $\partial^{m+n}u_r/\partial x^m \partial t^n = u_{r+m+2n}$ .*

We shall also use the following essentially known facts.

**THEOREM A.** *Let  $f(x)$  be measurable, and let  $|f(x)| \leq Me^{cx^2}$ ,  $c > 0$ . Then*

$$(6) \quad u(x, t) = \int_{-\infty}^{\infty} W_0(x - \xi, t) f(\xi) d\xi,$$

*is absolutely convergent for  $0 < t < 1/4c$ , and*

$$(\alpha) \quad u(x, t) \in C^\infty \quad \text{and} \quad u_{xx} = u_t, \quad 0 < t < 1/4c,$$

$$(\beta) \quad |u(x, t)| \leq Ke^{Nx^2}, \quad 0 < t \leq t_0 < 1/4c, \quad K = K(t_0), \quad N = N(t_0).$$

*If  $f(x)$  has bounded variation on every finite interval, and  $2f(x) = f(x^+) + f(x^-)$ , then*

$$(\gamma) \quad \lim_{t \rightarrow 0^+} u(x, t) = f(x).$$

*If  $f(x)$  is continuous, then*

$$(\gamma') \quad \lim_{t \rightarrow 0^+, x \rightarrow x_0} u(x, t) = f(x_0),$$

*so that if  $u(x, 0)$  is defined as  $f(x)$ , then  $u(x, t)$  is continuous for  $0 \leq t < 1/4c$ .*

**PROOF.** For  $(\alpha)$ ,  $(\gamma)$ , and  $(\gamma')$  see [1, p. 298];  $(\beta)$  is part of Lemma 1.

**THEOREM B.** *Conversely, if  $u(x, t)$  is defined for  $0 < t \leq t_0$  and satisfies*

$$(\alpha) \quad u(x, t) \in C^2 \quad \text{and} \quad u_{xx} = u_t, \quad 0 < t \leq t_0,$$

$$(\beta) \quad |u(x, t)| \leq Ke^{Nx^2}, \quad 0 < t \leq t_0, \quad \text{for some finite constants}$$

$$K = K(t_0), \quad N = N(t_0),$$

$$(\gamma) \quad \lim_{t \rightarrow 0^+} |u(x, t)| = 0, \quad \text{for all } x,$$

*then*

$$u(x, t) \equiv 0, \quad 0 < t \leq t_0.$$

The usual theorem, assuming simultaneous continuity in  $x$  and  $t$ ,  $t \geq 0$ , is in [2, p. 88]. To prove our slightly sharper result, note that by the usual theorem, if  $0 < \tau < t < 1/4N$ , then

$$u(x, t) = \lim_{\tau \rightarrow 0^+} \int_{-\infty}^{\infty} u(\xi, \tau) W_0(x - \xi, t - \tau) d\xi.$$

If  $\tau < t/2$ , then  $W_0(x - \xi, t - \tau) \leq (2\pi t)^{-1/2} e^{-(x-\xi)^2/2t}$ . Hence  $(\beta)$  gives dominated convergence [3, p. 168], and we can pass to the limit

under the integral sign. We conclude that  $u(x, t) = 0$  for all  $t$  such that  $t < 1/4N$ . Suppose  $t_1$  is the l.u.b. of the values of  $t$  for which  $\tau < t \rightarrow u(x, \tau) = 0$ , and suppose  $t_1 < t_0$ . Then, by continuity,  $u(x, t_1) = 0$ , and we may repeat the preceding argument to prove that  $u(x, t) \equiv 0$ ,  $t \leq t_1 + \eta$ , for some positive  $\eta$ . This contradiction shows that  $t_1 = t_0$ , completing the proof of Theorem 2.

We shall now prove our existence theorem.

**THEOREM 1.** *Let  $H(x)$  have bounded variation on every finite interval, let  $2H(x) = H(x^+) + H(x^-)$ ,  $H(0) = 0$ , and let  $|H(x)| \leq Me^{cx^2}$ . Then the improper integral*

$$(7) \quad u(x, t) = \int_{-\infty}^{\infty} W_0(x - \xi, t) dH(\xi) = \lim_{A \rightarrow \infty} \int_{-A}^A W_0(x - \xi, t) dH(\xi)$$

exists for  $0 < t < 1/4c$ , and satisfies

$$(a) \quad u(x, t) \in C^\infty \quad \text{and} \quad u_{xx} = u_t \quad \text{on} \quad 0 < t < 1/4c,$$

$$(b) \quad |u(x, t)| \leq Kt^{-1/2}e^{Nx^2} \quad \text{on} \quad 0 < t \leq t_0 < 1/4c,$$

$$(c) \quad \lim_{t \rightarrow 0^+} \int_0^x u(\xi, t) d\xi = H(x) \quad \text{for all } x,$$

$$(d) \quad \int_0^x u(\xi, t) d\xi \leq Le^{Nx^2} \quad \text{on} \quad 0 < t \leq t_0 < 1/4c,$$

where  $L = L(t_0)$  and  $N = N(t_0)$ .

**PROOF.** If  $0 < t < 1/4c$ , then by direct calculation,

$$\begin{aligned} u(x, t) &= \lim_{A \rightarrow \infty} \int_{-A}^A W_0(x - \xi, t) dH(\xi) \\ &= \lim_{B \rightarrow \infty} \left\{ [W_0(x - \xi, t) H(\xi)]_{-A}^A + \int_{-A}^A W_1(x - \xi, t) H(\xi) d\xi \right\}. \end{aligned}$$

By (4) and  $|H(x)| \leq Me^{cx^2}$ , the term in square brackets tends to zero. Hence

$$(8) \quad u(x, t) = \int_{-\infty}^{\infty} W_1(x - \xi, t) H(\xi) d\xi.$$

Lemma 1 assures us that (8) converges absolutely for  $0 < t < 1/4c$ , and has the properties (a), (b). Hence we can interpret  $u(x, t)$  as coming from an initial dipole distribution of density  $H(x)$ .

To prove (c), integrate (8). Defining

$$(9) \quad H(x, t) = \int_0^x u(\xi, t) d\xi,$$

we get

$$(9') \quad H(x, t) = \int_{-\infty}^{\infty} [W_0(x - \xi, t) - W_0(-\xi, t)] H(\xi) d\xi,$$

where the Fubini theorem is used to interchange the order of integration. Letting  $t$  tend to zero, and using  $(\gamma)$  of Theorem A, we get (c). To prove (d), we apply  $(\beta)$  of Theorem A to both terms of (9') and add.

To prove uniqueness, we shall want the following result.

**THEOREM C.** *If  $u(x, t) \in C^2$ ,  $u_{xx} = u_t$ , and  $|H(x, t)| \leq Ke^{N^2}$ , where  $H(x, t)$  is defined by (9) for all  $0 \leq t \leq T$ , then*

$$(10) \quad u(x, t) = \int_{-\infty}^{\infty} W_1(x - \xi, t) H(\xi, 0) d\xi, \quad \text{for } t < 1/4N.$$

**PROOF.** We first note that

$$H_t = \int_0^x u_t(\xi, t) d\xi = \int_0^x u_{xx}(\xi, t) d\xi = u_x(x, t) - u_x(0, t),$$

$$H_{xx} = u_x(x, t).$$

Hence, if we define  $\bar{H}(x, t) = H(x, t) + \int_0^t u_x(0, \tau) d\tau$ , then  $\bar{H}(x, t)$  satisfies the heat equation,  $\bar{H}_{xx} = \bar{H}_t$ . Also  $\bar{H}(x, t) \in C^2$  and  $|\bar{H}(x, t)| \leq K'e^{N^2}$ . Hence Theorem B applies, and

$$\bar{H}(x, t) = \int_{-\infty}^{\infty} W_0(x - \xi, t) \bar{H}(\xi, 0) d\xi = \int_{-\infty}^{\infty} W_0(x - \xi, t) H(\xi, 0) d\xi.$$

Using part (a) of Lemma 1, we differentiate with respect to  $x$  and obtain (10), q.e.d.

**THEOREM 2.** *Let  $u(x, t) \in C^2$ ,  $u_t = u_{xx}$ ,  $|H(x, t)| \leq Ke^{N^2}$ , where  $H(x, t)$  is defined by (9), for all  $0 < t \leq T$ , and  $H(x, t) \rightarrow 0$  as  $t \rightarrow 0^+$  for almost all  $x$ . Then  $u(x, t) \equiv 0$  on  $0 < t \leq T$ .*

**PROOF.** Let  $0 < \tau < t < \min [1/4N, T]$ , and  $u_1(x, t - \tau) = u(x, t)$ ,  $H_1(x, t - \tau) = H(x, t)$ . Then for  $t \geq \tau$ ,  $u_1(x, t - \tau)$ ,  $H_1(x, t - \tau)$  satisfy the conditions of Theorem C. Hence

$$u_1(x, t - \tau) = \int_{-\infty}^{\infty} W_x(x - \xi, t - \tau) H_1(\xi, 0) d\xi,$$

or

$$u(x, t) = \int_{-\infty}^{\infty} W_x(x - \xi, t - \tau)H(\xi, \tau)d\xi.$$

By hypothesis  $H(\xi, \tau) \rightarrow 0$  almost everywhere *dominatedly* as  $\tau \rightarrow 0$ . Hence, using the Lebesgue dominated convergence theorem again, we have  $u(x, t) \equiv 0$ ,  $0 < t < \min [1/4N, T]$ . If  $1/4N < T$ , we repeat the argument, considering  $\bar{u}(x, t) = h(x, t + 1/4N)$  instead of  $u(x, t)$ :  $\bar{u}(x, 0) = 0$  by the continuity of  $u(x, t)$  in  $0 < t < T$ . This shows  $u(x, t) \equiv 0$ ,  $0 < t < \min [1/4N, T]$ . The proof is complete after at most  $[4NT]$  steps.<sup>2</sup>

COROLLARY. *Under the hypotheses of Theorem 1,  $u(x, t)$  as defined by (6) is the only function satisfying conditions (a), (c), (d).*

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<sup>2</sup> J. L. B. Cooper (J. London Math. Soc. vol. 25 (1950) pp. 173-180) has proved a result related to Theorem 2. However, he requires that  $\int_a^b g(x)u(x, t)dx \rightarrow 0$  for all bounded integrable  $g(x)$ ; we require only  $\int_a^b u(x, t)dx \rightarrow 0$ , for all  $a, b$ . Cf. also J. Kampé de Fériet, C. R. Acad. Sci. Paris vol. 236 (1953) pp. 1527-1929.