

THE RECIPROCAL OF A FOURIER SERIES¹

R. J. DUFFIN

Edrei and Szegö [1] have posed the following problem: *Given the Fourier coefficients of a function $G(x)$ find the Fourier coefficients of the reciprocal of the function without actually evaluating $G(x)$.* They were able to solve this problem in the case that $G(x) \geq 0$. Unfortunately this restriction eliminates the interesting case of complex-valued functions such as those which arise in Laurent series.

In this note it is shown possible to obtain a solution without the restriction, $G(x) \geq 0$. This is achieved by first treating the more general problem of finding the coefficients of $F(x)/G(x)$, given the coefficients of $F(x)$ and $G(x)$.

Edrei and Szegö confine attention to classical Fourier series. Their problem is treated here for arbitrary orthogonal series expansions.²

Let $\theta_j(x)$, $j=1, 2, \dots$, be a sequence of bounded orthonormal functions in some region R of a space S . Thus

$$(1) \quad \int_R \theta_j(x) \theta_k^*(x) dx = \delta_{jk}$$

where the asterisk denotes the complex conjugate. The Fourier coefficients of a function $F(x) \in L$ are denoted by f_j and are given by the integral

$$(2) \quad f_j = \int_R F(x) \theta_j^*(x) dx, \quad j = 1, 2, \dots$$

If $f_j=0$ for all j it is assumed that $F(x)=0$ almost everywhere. In other words the orthonormal sequence is closed in L .

PROBLEM 1. *Given the Fourier coefficients f_n and g_n of two functions $F(x)$ and $G(x)$ find the Fourier coefficients h_n of the function $H(x) = F(x)/G(x)$.*

A related question concerns a direct determination of $H(x)$ without first finding $F(x)$ and $G(x)$. This may be stated as

PROBLEM 2. *In terms of f_n and g_n find expansion coefficient sets $(\rho_1^{(m)}, \rho_2^{(m)}, \dots, \rho_m^{(m)})$ for $m=1, 2, \dots$ such that*

$$(3) \quad \lim_{m \rightarrow \infty} \int_R \left| H(x) - \sum_{k=1}^m \rho_k^{(m)} \theta_k(x) \right| dx = 0.$$

Received by the editors April 28, 1961 and, in revised form, December 2, 1961.

¹ Prepared under Contract DA-36-061 ORD-490, U. S. Army Research Office.

² This generalization was suggested to the writer by David Moskovitz.

Given a solution of Problem 2 then

$$(4) \quad \lim_{m \rightarrow \infty} \int_R \left[H(x) - \sum_{k=1}^m \rho_k^{(m)} \theta_k(x) \right] \theta_j^*(x) dx = 0.$$

This is seen to give

$$(5) \quad h_j = \lim_{m \rightarrow \infty} \rho_j^{(m)}$$

and consequently the solution of Problem 2 also furnishes a solution of Problem 1. (In the following $\rho_j^{(m)}$ is denoted by ρ_j .)

PROBLEM 3. *The solution of Problem 2 is sought but given the "Fourier matrix"*

$$(6) \quad g_{jk} = \int_R G(x) \theta_j^*(x) \theta_k(x) dx \quad ,$$

instead of the Fourier coefficients g_j . In other words find h_n in terms of f_n and g_{jk} . The following formal considerations indicate the close relationship between Problem 2 and Problem 3. Let

$$(7) \quad \beta_{jkr} = \int_R \theta_j^*(x) \theta_k(x) \theta_r(x) dx.$$

Then

$$(8) \quad \theta_j^*(x) \theta_k(x) = \sum_{r=1}^{\infty} \beta_{jkr} \theta_r^*(x).$$

This relation is multiplied by $G(x)$ and integrated termwise to yield

$$(9) \quad g_{jk} = \sum_{r=1}^{\infty} \beta_{jkr} g_r.$$

Thus if (9) is valid then a solution of Problem 3 will yield a solution of Problem 2.

THEOREM 1. *Let $G(x) \geq 0$, $G(x) \in L$, $1/G(x) \in L$, and $F^2(x)/G(x) \in L$. Then for each integer m the equations*

$$(10) \quad f_j = \sum_{k=1}^m g_{jk} \rho_k, \quad j = 1, \dots, m,$$

may be solved for $\rho_1, \rho_2, \dots, \rho_m$ and these are expansion coefficient sets for $H(x) = F(x)/G(x)$ and this solves Problem 3.

PROOF. Since $F = HG^{1/2}G^{1/2}$ the Schwarz inequality gives

$$\left(\int |F| dx \right)^2 \leq \int |H^2G| dx \int |G| dx,$$

and

$$\left(\int |H| dx \right)^2 \leq \int |H^2G| dx \int |G^{-1}| dx.$$

This proves that $F \in L$ and $H \in L$. Then it is seen that the following integral exists for any choice of the constants ρ_j :

$$(11) \quad E = \int \left| H(x) - \sum_1^m \rho_k \theta_k(x) \right|^2 G(x) dx.$$

Since $G(x) \geq 0$ this may be written in the form

$$(12) \quad E = \int \left| HG^{1/2} - \sum_1^m \rho_k \theta_k G^{1/2} \right|^2 dx.$$

This leads to consideration of the relation

$$(13) \quad \int U(x) \theta_j^*(x) G^{1/2} dx = 0, \quad j = 1, 2, \dots,$$

where $U(x) \in L^2$. Then $UG^{1/2} \in L$ and since the sequence θ_j is closed in L relation (13) implies $UG^{1/2} = 0$ almost everywhere. Thus $U = UG^{1/2}G^{-1/2} = 0$ almost everywhere. This proves that the sequence $\theta_j G^{1/2}$ is closed in L^2 .

Closure in L^2 implies completeness in L^2 so it is possible to make E arbitrarily small. This may be accomplished by choosing, for each m , the set (ρ_1, \dots, ρ_m) which minimizes E . As is well known this optimal choice is determined by the orthogonality equations

$$(14) \quad 0 = \int \left(H - \sum_1^m \rho_k \theta_k \right) \theta_j^* G dx, \quad j = 1, 2, \dots, m.$$

Since $HG = F$ it is apparent that these are precisely equations (10) of Theorem 1.

By the Schwarz inequality

$$(15) \quad \left(\int \left| H - \sum_1^m \rho_k \theta_k \right| dx \right)^2 \leq E \int |G^{-1}| dx.$$

The left side approaches zero if E approaches zero. This shows that the optimal choice of ρ_j leads to the satisfaction of condition (3) of Problem 2 and completes the proof.

THEOREM 2. Suppose $F(x) \in L^2$ and $1/F(x) \in L^2$. Let

$$(16) \quad \beta_{jkr\theta} = \int_R \theta_j^*(x)\theta_k(x)\theta_r^*(x)\theta_s(x)dx,$$

$$(17) \quad g_{jk} = \sum_1^\infty \sum_1^\infty \beta_{jkr\theta} f_r^* f_s.$$

Then the equations

$$(18) \quad f_j = \sum_1^m g_{jk}\rho_k, \quad j = 1, 2, \dots, m,$$

may be solved for $\rho_1, \rho_2, \dots, \rho_m$. For $m = 1, 2, \dots$ these solutions form a set of expansion coefficients of the function $H(x) = 1/F^*(x)$.

PROOF. Define $G(x) = F(x)F^*(x)$ then $H(x) = F(x)/G(x)$. Then $G(x) \geq 0, G(x) \in L, (G(x))^{-1} \in L$, and $H^2(x)G(x) \in L$. Thus the conditions of Theorem 1 are satisfied. Substituting $G = FF^*$ in relation (6) for the Fourier matrix gives

$$(19) \quad g_{ik} = \int \theta_j^* \theta_k F^* F dx.$$

The series $\sum_1^\infty f_r^* \theta_r^*$ and $\sum_1^\infty f_s \theta_s$ converge in L^2 mean to F^* and F . Therefore it is permissible to substitute these series in (19) and to integrate termwise. This justifies (17); the proof then follows from Theorem 1.

THEOREM 3. Let $G(x) \geq 0, G(x) \in L, 1/G(x) \in L$, and $F^2(x)/G(x) \in L$ in the interval $0 \leq x \leq 1$. The Fourier coefficients are now defined as

$$(20) \quad f_j = \int_0^1 F(x) \exp(-2\pi i j x) dx, \quad j = 0, \pm 1, \dots$$

Then the equations

$$(21) \quad f_j = \sum_{-m}^m g_{j-k}\rho_k, \quad j = -m, \dots, m,$$

may be solved for ρ_{-m}, \dots, ρ_m and

$$(22) \quad \int_0^1 \left| F(x)/G(x) - \sum_{-m}^m \rho_k \exp(2\pi i k x) \right| dx \rightarrow 0 \text{ as } m \rightarrow \infty.$$

PROOF. Theorem 1 is specialized with the sequence $\theta_j(x)$ being the sequence $\exp(2\pi i j x)$ suitably reordered. It is then seen that

$$(23) \quad g_{jk} = g_{j-k}$$

and thus equation (10) becomes equation (21). The proof of relation (22) then follows from Theorem 1.

THEOREM 4. *Suppose $F(x) \in L^2$ and $1/F(x) \in L^2$ on the interval $0 \leq x \leq 1$. Let f_j denote the Fourier coefficients of $F(x)$ relative to the orthonormal sequence $\exp(2\pi i j x)$. Let*

$$(24) \quad g_j = \sum_{-\infty}^{\infty} f_{j+r} f_r^*.$$

Then the equations

$$(25) \quad f_{-j}^* = \sum_{-m}^m g_{j-k} \alpha_k$$

may be solved for $\alpha_{-m}, \dots, \alpha_m$ and

$$(26) \quad \int_0^1 \left| (F(x))^{-1} - \sum_{-m}^m \alpha_k \exp(2\pi i k x) \right| dx \rightarrow 0 \text{ as } m \rightarrow \infty.$$

PROOF. First Theorem 2 is applied to obtain $H = 1/F^*$. It is seen from (16) that $\beta_{jkr s} = 1$ if $s = j - k + r$ and $\beta_{jkr s} = 0$ otherwise. Thus (17) gives

$$g_{jk} = g_{j-k} = \sum_{-\infty}^{\infty} f_{j-k+r} f_r^*.$$

Then (18) is

$$f_j = \sum_{-m}^m g_{j-k} \rho_k.$$

Taking the complex conjugate of this relation gives

$$f_{-j}^* = \sum_{-m}^m g_{k-j}^* \rho_{-k}^*.$$

But $g_{-k}^* = g_k$ and $\rho_{-k}^* = \alpha_k$ so this proves (25).

The questions studied in this note originated from the need for an algorithm which would give the reciprocal of a Laurent series. (In this connection it is worth noting that there is a simple algorithm for the reciprocal of a Taylor series.) An algorithm for the reciprocal of a Laurent series is stated in the following Corollary of Theorem 4.

COROLLARY. *Consider the Laurent series*

$$(27) \quad S(z) = \sum_{-\infty}^{\infty} a_n z^n$$

which is assumed to converge when the complex variable z is in the region D defined by the relation $r_1 < |z| < r_2$. If $S(z)$ does not vanish in D it follows, of course, that the reciprocal $T(z) = 1/S(z)$ is a Laurent series in D . Thus

$$(28) \quad T(z) = \sum_{-\infty}^{\infty} c_n z^n.$$

Let coefficients b_j be defined as

$$(29) \quad b_j = \sum_{-\infty}^{\infty} a_{j+r} a_r^*.$$

Then the equations

$$(30) \quad a_{-j}^* = \sum_{-m}^m b_{j-k} C_k^{(m)}, \quad j = -m, \dots, m,$$

may be solved for $C_{-m}^{(m)}, \dots, C_m^{(m)}$. Let

$$(31) \quad T_m(z) = \sum_{-m}^m C_k^{(m)} z^k.$$

Then as $m \rightarrow \infty$ the sequence $T_m(z)$ converges uniformly to $T(z)$ in any closed region contained in D . Moreover, $C_k^{(m)} \rightarrow c_k$ as $m \rightarrow \infty$.

PROOF. The function $F(x) = S(re^{2\pi iz})$ satisfies the condition of Theorem 4 for any value of r in the range $r_1 < r < r_2$. The Fourier coefficients of F are given by $f_j = r^j a_j$. Defining b_j and $C_j^{(m)}$ by the relations $g_j = r^j b_j$ and $\alpha_j = r^j C_j^{(m)}$ it results that relations (24) and (25) yield relations (29) and (30). It follows that T_m converges in mean to T for $|z| = r$. Likewise T_m converges in mean for $|z| = r'$ where $r_1 < r < r' < r_2$. Then by a standard theorem from complex function theory it follows that T_m converges uniformly to T for $r < |z| < r'$. This is seen to complete the proof of the Corollary.

REFERENCE

1. A. Edrei and G. Szegő, *A note on the reciprocal of a Fourier series*, Proc. Amer. Math. Soc. **4** (1953), 323-329.

CARNEGIE INSTITUTE OF TECHNOLOGY