

# AN ABSTRACT OSCILLATION THEOREM

KURT KREITH

Questions of oscillatory behavior are generally associated with Sturm-Liouville equations of the form

$$(1) \quad -(d/dt)(p(t)dx/dt) + q(t)x = \mu x$$

where (1) is singular at one or both ends of an interval  $I = (a, b)$ ,  $-\infty \leq a < b \leq \infty$ . Numerous criteria exist, depending on the behavior of the coefficients  $p(t)$  and  $q(t)$ , which assure that solutions of (1) are or are not oscillatory near a singular end point of  $I$ . However there also exist criteria for oscillatory behavior involving only the spectrum of the Sturm-Liouville operator

$$(2) \quad \tau = -(d/dt)(p(t)d/dt) + q(t)$$

in the Hilbert space  $L_2(I)$  (see for example Dunford-Schwartz [1, Theorem XIII.7.40]). For lower semibounded operators these criteria state that if (1) is oscillatory then the essential spectrum of  $\tau$  intersects  $(-\infty, \mu]$ , whereas if (1) is nonoscillatory, then the essential spectrum of  $\tau$  does not intersect  $(-\infty, \mu)$ . These spectral criteria suggest that it may be of interest to generalize the notion of "oscillatory behavior" to solutions of certain operator equations of the form

$$(3) \quad Ax = \mu x$$

where  $A$  is an appropriate operator in a Hilbert space  $\mathfrak{H}$  and  $\mu$  is a real constant. The purpose of this paper is to examine one such generalization.

We shall assume throughout that  $A$  is a symmetric operator which is bounded below in a Hilbert space  $\mathfrak{H}$ . By adding sufficiently large positive multiples of  $x$  to both sides of (3), one may assume, without loss of generality, that  $A$  satisfies

$$(4) \quad (Ax, x) \geq \|x\|^2$$

for all  $x$  in  $\mathfrak{D}_A$  (the domain of  $A$ ). It will be useful to recall that a symmetric operator  $A$  satisfying (4) has a selfadjoint extension  $\bar{A}$ , the Friedrichs extension, obtained as follows: complete  $\mathfrak{D}_A$  under the norm  $\|x\|^2 = (Ax, x)$  to construct a Hilbert space  $\mathfrak{N}$  with inner product  $((\cdot, \cdot))$ ; define  $\mathfrak{D}_{\bar{A}}$  to consist of those elements  $y \in \mathfrak{H}$  for which there exists a sequence  $x_n$  in  $\mathfrak{D}_A$  such that

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$$\lim_{n \rightarrow \infty} \|y - x_n\| = 0 \quad \text{and} \quad \lim_{m, n \rightarrow \infty} \|x_m - x_n\| = 0.$$

Then, according to Friedrichs' theorem, there exists a unique self-adjoint extension  $\bar{A}$  of  $A$  for which  $\mathcal{D}_{\bar{A}} \subset \mathcal{M}$  and  $((x, y)) = (x, Ay)$  for all  $x \in \mathcal{M}$  and all  $y \in \mathcal{D}_{\bar{A}}$ . If  $\tau$  is nonsingular on  $I = (a, b)$  and  $\mathcal{D}_\tau$  is taken to be the class of infinitely differentiable functions with compact support on  $I$ , then  $\bar{\tau}$  will be the selfadjoint extension of  $\tau$  corresponding to the boundary conditions  $x(a) = x(b) = 0$ .

Let  $\mathfrak{S}_k$  be a closed linear subspace of  $\mathfrak{S}$  for which  $\mathcal{D}_A \cap \mathfrak{S}_k$  is dense in  $\mathfrak{S}_k$  and let  $P_k$  denote the projection operator with range  $\mathfrak{S}_k$ . In  $\mathfrak{S}_k$  we define a symmetric operator  $A_k$  as follows:

- (i)  $x \in \mathcal{D}_{A_k}$  if  $x \in \mathcal{D}_A$  and  $P_k x = x$ ,
- (ii) if  $x \in \mathcal{D}_{A_k}$ , then  $A_k x = P_k A x$ .

Since  $A_k x \in \mathfrak{S}_k$  for all  $x \in \mathcal{D}_{A_k}$ , it is clear that  $A_k$  is a symmetric operator in  $\mathfrak{S}_k$ , satisfies (4), and therefore has a Friedrichs extension, to be denoted by  $\bar{A}_k$ . Norms and inner products in  $\mathfrak{S}_k$  will be denoted by a subscript  $k$  outside the norm or inner product symbol.

Consider now the operator equation

$$(5) \quad A^* x = \mu x,$$

where  $A^*$  is the adjoint of  $A$ .

DEFINITION 1. We say that  $\mathfrak{S}_k$  is a nodal domain for (5) if  $Ax = P_k Ax$  for all  $x \in \mathfrak{S}_k \cap \mathcal{D}_A$  and

$$\inf_{x \in \mathcal{D}_{\bar{A}_k}} \frac{(\bar{A}_k x, x)_k}{\|x\|_k^2} = \mu,$$

where this infimum is achieved by an eigenfunction  $u_k \in \mathcal{D}_{\bar{A}_k}$ .

DEFINITION 2. We say that a solution  $x_0$  of (5) is oscillatory if there exists a decomposition of  $\mathfrak{S}$  into orthogonal closed subspaces.

$$(6) \quad \mathfrak{S} = \sum_{k=0}^{\infty} \oplus \mathfrak{S}_k$$

such that for  $k \geq 1$  each  $\mathfrak{S}_k$  is a nodal domain for (5) and  $P_k x_0$  is the required eigenfunction satisfying  $(\bar{A}_k P_k x_0, P_k x_0)_k = \mu \|P_k x_0\|_k^2$ . (The fact that we do not impose any conditions on  $\mathfrak{S}_0$  corresponds to a lack of boundary conditions in (1).)

In one direction these definitions lead to spectral criteria for oscillatory behavior of (5) very similar to those which exist for (1).

**THEOREM.** *If (5) has an infinite number of orthogonal nodal domains, then the essential spectrum of  $\bar{A}$  intersects  $[1, \mu]$ .<sup>1</sup>*

<sup>1</sup> Since  $(\bar{A}x, x) \geq \|x\|^2$ , we know that the spectrum of  $\bar{A}$  cannot intersect  $(-\infty, 1)$ .

PROOF. We shall show that given any  $\epsilon > 0$  and any positive integer  $N$ , there exists a set of  $N$  linearly independent vectors  $v_1, \dots, v_N$  in  $\mathfrak{D}_A$  such that for any nontrivial linear combination  $v = c_1 v_1 + \dots + c_N v_N$ , we have  $(Av, v) < (\mu + \epsilon) \|v\|^2$ . According to a principle adapted by Friedrichs [2], this implies either that  $\bar{A}$  has at least  $N$  eigenvalues in  $[1, \mu + \epsilon)$  or else that the spectrum of  $\bar{A}$  is not discrete below  $\mu + \epsilon$ ; since  $\epsilon$  is arbitrarily small,  $N$  is arbitrarily large and  $\bar{A}$  is bounded below, such a construction implies that the spectrum of  $\bar{A}$  is not discrete in  $[1, \mu]$ .

To construct the set  $v_1, \dots, v_N$ , let  $u_1, \dots, u_N$  denote the normalized eigenfunctions of  $\bar{A}_1, \dots, \bar{A}_N$  corresponding to the eigenvalue  $\mu$ . By the definition of  $\bar{A}_k$ , for every  $u_k$  there exists a sequence  $v_{kl}$  satisfying

- (i)  $v_{kl} \in \mathfrak{D}_A$ ;  $k = 1, \dots, N$ ;  $l = 1, 2, \dots$ ;
- (ii)  $P_k v_{kl} = v_{kl}$ ;  $k = 1, \dots, N$ ;  $l = 1, 2, \dots$ ;
- (iii)  $\lim_{l \rightarrow \infty} \|P_k v_{kl} - u_k\|_k = 0$ ;
- (iv)  $\lim_{l \rightarrow \infty} \|P_k v_{kl} - u_k\|_k = 0$ .

For  $k = 1, \dots, N$  we have

$$\| \|v_{kl}\| \| = \| \|P_k v_{kl}\| \|_k \leq \| \|P_k v_{kl} - u_k\| \|_k + \| \|u_k\| \|_k$$

and

$$\| \|u_k\| \|_k = \mu^{1/2} \| \|u_k\| \|_k \leq \mu^{1/2} (\| \|u_k - P_k v_{kl}\| \|_k + \| \|v_{kl}\| \|).$$

Combining these inequalities

$$\| \|v_{kl}\| \| \leq \| \|P_k v_{kl} - u_k\| \|_k + \mu^{1/2} (\| \|u_k - P_k v_{kl}\| \|_k + \| \|v_{kl}\| \|).$$

In light of (iii) and (iv) and the fact that  $\lim_{l \rightarrow \infty} \| \|v_{kl}\| \| = 1$ , we can choose  $l_0$  sufficiently large so that

$$\| \|v_{kl_0}\| \| < (\mu + \epsilon)^{1/2} \| \|v_{kl_0}\| \|; \quad k = 1, 2, \dots, N.$$

Defining  $v_k = v_{kl_0}$ , and noting that

$$\begin{aligned} (v_j, v_k) &= 0 \quad \text{for } j \neq k; & (Av_j, v_k) &= 0 \quad \text{for } j \neq k; \\ (Av_k, v_k) &< (\mu + \epsilon) \| \|v_k\| \|^2, \end{aligned}$$

it follows that for any  $v = c_1 v_1 + \dots + c_N v_N \neq 0$

$$(Av, v) = \sum_{k=1}^N |c_k|^2 (Av_k, v_k) < (\mu + \epsilon) \sum_{k=1}^N |c_k|^2 \| \|v_k\| \|^2 = (\mu + \epsilon) \| \|v\| \|^2,$$

which completes the proof.

As an immediate consequence we have the following.

COROLLARY. *If  $A^*x = \mu x$  has an oscillatory solution, then the spectrum of  $\bar{A}$  is not discrete in  $[1, \mu]$ .*

If the operator  $A$  has finite deficiency indices, then the essential spectrum of all selfadjoint extensions is the same, and this fact enables one to formulate criteria in terms of the essential spectrum of  $A$  instead of  $\bar{A}$ . This remark applies in particular to the operator  $\tau$  whose deficiency indices are at most 2.

For the operator  $\tau$  we also have the following converse of Theorem 1. If (1) has a nonoscillatory solution, then the essential spectrum of  $\tau$  does not intersect  $[1, \mu)$ . That such a converse does not hold in our more general setting is indicated by an elementary example.

Consider  $Ax \equiv -\partial^2 x / \partial s^2 - \partial^2 x / \partial t^2 + q(t)x$ , where  $q(t)$  is chosen such that the singular selfadjoint Sturm-Liouville system

$$\begin{aligned} -d^2y/dt^2 + qy &= \lambda y; & 0 < t \leq 1, \\ y(1) &= 0, \end{aligned}$$

has a spectrum composed of an eigenvalue at  $\lambda = -1$ , some continuous spectrum in the interval  $(-3, -2)$ , a finite number of eigenvalues for  $\lambda < -3$ , and an arbitrary mixture of discrete and continuous spectrum for  $\lambda > -1$ . Defining  $R = \{(s, t) \mid 0 < s < \pi; 0 < t < 1\}$  and setting  $\mathfrak{S} = L_2(R)$  we can solve the boundary value problem

$$\begin{aligned} Ax &= 0 \quad \text{in } R, \\ x &= 0 \quad \text{on } \partial R \cap \{(s, t) \mid t > 0\} \end{aligned}$$

by separation of variables. Denoting this boundary value problem by  $A^*x = 0$  and setting  $x(s, t) = S(s)T(t)$ , we get

$$S(s) = \sin ns, \quad -T'' + qT = -n^2T$$

for  $n = 1, 2, \dots$ . Setting  $n = 1$ , it follows that  $A^*x = 0$  has a square integrable oscillatory solution; however, for  $n > 1$  no solution of  $A^*x = 0$  is oscillatory.

This example suggests that a satisfactory converse to Corollary 1 is likely to be elusive in any abstract setting general enough to include singular elliptic operators.

#### BIBLIOGRAPHY

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UNIVERSITY OF CALIFORNIA, DAVIS