

# AN INEQUALITY FOR A STEKLOFF EIGENVALUE BY THE METHOD OF DEFECT<sup>1</sup>

J. R. KUTTLER<sup>2</sup> AND V. G. SIGILLITO

1. **Introduction.** In this paper we give a lower bound to the first nonzero eigenvalue  $p_2$  of the Stekloff problem [8] for plane regions having two axes of symmetry. Such bounds lead to a priori inequalities which are useful in giving error bounds for approximate solutions to the Neumann problem for Poisson's equation (see [5, §2]).

In [9], Weinstock has given an isoperimetric upper bound for  $p_2$ , but good lower bounds, which are more useful, are more elusive (see [1], [4], [6]).

We first prove a nodal-line theorem of some interest in itself. This nodal-line theorem does not seem to appear anywhere in the literature, although its proof is a straightforward application of familiar methods (see [7]). The results of the nodal-line theorem then permit us to use the method of defect (see, e.g., [2]) to obtain our bound by integrating an easily obtained one-dimensional version of the desired inequality.

2. **Preliminaries.** Let  $B$  be a bounded, connected domain of the  $x_1, x_2$ -plane with piecewise smooth boundary  $\partial B$ . We consider the Stekloff eigenvalue problem

$$(1) \quad \Delta u = 0 \text{ in } B, \quad \partial u / \partial n = pu \text{ on } \partial B,$$

where  $\Delta$  is the Laplacian and  $n$  the unit outer normal on  $\partial B$ . The problem has a discrete spectrum of eigenvalues  $0 = p_1 < p_2 \leq p_3 \leq \dots$  with corresponding eigenfunctions  $u_1 = \text{constant}$ ,  $u_2, u_3, \dots$ . The eigenvalues can be characterized by

$$(2) \quad p_n = \min \frac{\int_B |\text{grad } v|^2 dx}{\int_{\partial B} v^2 ds},$$

where the minimum is taken over all continuous and piecewise continuously differentiable functions  $v$  satisfying

$$(3) \quad \oint_{\partial B} v u_k ds = 0, \quad k = 1, 2, \dots, n-1.$$

---

Received by the editors October 30, 1967.

<sup>1</sup> Supported in part by Contract No. AS-7-283, Aerospace Research Laboratories, Office of Aerospace Research, United States Air Force.

<sup>2</sup> Post-doctoral Fellow at the Applied Physics Laboratory, The John Hopkins University.

The minimum of (2) is attained when and only when  $v$  is an eigenfunction of (1) associated with  $p_n$ .

**3. The nodal-line theorem.** A curve in  $B$  along which an eigenfunction  $u_n$  vanishes is called a nodal line of  $u_n$ . What we prove is

**THEOREM.** *The nodal lines of  $u_n$  divide  $B$  into no more than  $n$  subdomains, and no nodal line is a closed curve.*

**PROOF.** We prove the second part of the theorem first. If a nodal line were closed, we would have  $u_n \equiv 0$  in the interior of the curve since  $\Delta u_n = 0$  in  $B$ . But then  $u_n$  would vanish identically in  $B$  by the Unique Continuation Theorem for harmonic functions (see, e.g., [3, Chapter X, §5]). Thus no nodal line can be a closed curve.

Now suppose that the nodal lines of  $u_n$  divide  $B$  into more than  $n$  subdomains. Let  $D_1, D_2, \dots, D_n$  be  $n$  of these subdomains. Note that  $\partial B \cap \partial D_i$  is not empty. Define  $w_i$  to agree with  $u_n$  on  $D_i$  and vanish on  $B - D_i$ ,  $i = 1, 2, \dots, n$ . Notice that  $w_i \not\equiv 0$  on  $\partial B \cap \partial D_i$ , otherwise  $w_i \equiv 0$  in  $D_i$  (since  $\Delta w_i = 0$  in  $D_i$ ). Thus, we can find a linear combination, say  $v = \sum_{i=1}^n a_i w_i$ , such that

$$\oint_{\partial B} v^2 ds = 1$$

and  $v$  satisfies (3). Moreover, the  $w_i$ , hence  $v$ , are continuous and piecewise continuously differentiable (see [3, Chapter X, §9]). Since  $\partial v / \partial n = p_n v$  on  $\partial B$ , it follows from Green's first identity that

$$\frac{\int_B |\text{grad } v|^2 dx}{\oint_{\partial B} v^2 ds} = p_n.$$

Thus  $v$  minimizes (2) and is therefore an eigenfunction of (1). Hence  $v$  is harmonic and vanishes on the nonempty subdomain  $B - D_1 \cup \dots \cup D_n$ . Using the Unique Continuation Theorem we arrive at a contradiction. This completes the theorem.

The application we wish to make of this theorem is the following. Since an eigenfunction  $u_2$  associated with  $p_2$  satisfies  $\oint_{\partial B} u_2 ds = 0$ , we see that  $u_2$  must have a nodal line. By the theorem  $u_2$  cannot have more than one, hence has exactly one.

**4. The method of defect.** Suppose now our region  $B$  has two distinct axes of symmetry. They may be assumed perpendicular, and, with no loss of generality, we take them to be the  $x_1$  and  $x_2$  axes. We will call a function defined on  $B$  even-even, odd-odd, even-odd, or odd-even as  $u$  is respectively even in both  $x_1$  and  $x_2$ , odd in both  $x_1$

and  $x_2$ , even in  $x_1$  and odd in  $x_2$ , or odd in  $x_1$  and even in  $x_2$ . Every eigenfunction of (1) can be assumed to belong to one of these *symmetry classes*. From the previous section,  $u_2$  must be in the even-odd or odd-even symmetry classes. (Otherwise,  $u_2$  has an even number of nodal lines.) Hence, the nodal line is an axis of symmetry.

The one-dimensional inequality

$$(4) \quad [u(0)]^2 + [u(l)]^2 \leq \frac{l}{2} \int_0^l [u'(y)]^2 dy,$$

for continuous and piecewise continuously differentiable functions on the interval  $(0, l)$  which satisfy  $u(0) = -u(l)$ , is easily shown by solving the Euler equation.

Let us first consider the case when the eigenfunction  $u_2$  is odd across the  $x_1$ -axis. Suppose the boundary  $\partial B$  can be expressed by  $x_2 = \pm f_1(x_1)$ ,  $-a_1 \leq x_1 \leq a_1$ . Then, employing (4),

$$\begin{aligned} \oint_{\partial B} u_2^2 ds &= 2 \int_{-a_1}^{a_1} [u_2(x_1, f_1(x_1))]^2 (1 + [f_1'(x_1)]^2)^{1/2} dx_1 \\ &\leq \int_{-a_1}^{a_1} (1 + [f_1'(x_1)]^2)^{1/2} \left[ f_1(x_1) \int_{-f_1(x_1)}^{f_1(x_1)} \left[ \frac{\partial u_2(x_1, x_2)}{\partial x_2} \right]^2 dx_2 \right] dx_1 \\ &\leq \left[ \max_{-a_1 \leq x_1 \leq a_1} f_1(x_1) (1 + [f_1'(x_1)]^2)^{1/2} \right] \int_{-a_1}^{a_1} \int_{-f_1(x_1)}^{f_1(x_1)} \left( \frac{\partial u_2}{\partial x_2} \right)^2 dx_2 dx_1 \\ &\leq \left[ \max_{-a_1 \leq x_1 \leq a_1} f_1(x_1) (1 + [f_1'(x_1)]^2)^{1/2} \right] \int_B |\text{grad } u_2|^2 dx. \end{aligned}$$

By similarly treating the case when  $u_2$  is odd across the  $x_2$ -axis, we have our inequality:

$$(5) \quad p_2^{-1} \leq \max_{i=1,2} \left[ \max_{-a_i \leq x_i \leq a_i} f_i(x_i) (1 + [f_i'(x_i)]^2)^{1/2} \right],$$

where  $x_1 = \pm f_2(x_2)$ ,  $-a_2 \leq x_2 \leq a_2$ , is another representation of  $\partial B$ .

**5. Some examples.** We give a few examples of the application of (5) to particular regions. For an upper bound, we use the isoperimetric inequality of Weinstock [9], which says

$$(6) \quad p_2 \leq 2\pi/L,$$

where  $L$  is the length of  $\partial B$ .

First, consider the rhombus

$$|x_1/a_1| + |x_2/a_2| \leq 1.$$

We have  $L = 4(a_1^2 + a_2^2)^{1/2}$  and

$$f_i(x_i) = a_j(1 - |x_i/a_i|), \quad j \neq i.$$

Thus, from (5) and (6), we have

$$(7) \quad \min(a_2/a_1, a_1/a_2) \leq p_2(a_1^2 + a_2^2)^{1/2} \leq \pi/2.$$

When  $a_1 = a_2 = S/\sqrt{2}$ , the rhombus is a square of side  $S$ , and (7) becomes  $1 \leq p_2 S \leq \pi/2$ , whereas the exact value to five places is  $p_2 S = 1.3765 \dots$

Next, consider the ellipse  $(x_1/a_1)^2 + (x_2/a_2)^2 \leq 1$  for which

$$f_i(x_i) = a_j(1 - (x_i/a_i)^2)^{1/2}, \quad j \neq i.$$

For simplicity, we worsen (6) by combining it with the classical isoperimetric inequality

$$(8) \quad L^2 \geq 4\pi A,$$

where  $A$ , the area of the ellipse, is  $\pi a_1 a_2$ . Thus, we have

$$(9) \quad [\max(a_1, a_2)]^{-1} \leq p_2 \leq (a_1 a_2)^{-1/2}.$$

For the special case of a disc when  $a_1 = a_2 = R$ , we attain equality on both sides and  $p_2 R = 1$ .

#### REFERENCES

1. J. H. Bramble and L. E. Payne, *Bounds in the Neumann problem for second order uniformly elliptic operators*, Pacific J. Math. **12** (1962), 823-833.
2. J. Hersch, *Sur la fréquence fondamentale d'une membrane vibrante: évaluations par défaut et principe de maximum*, Z. Angew. Math. Phys. **11** (1960), 387-413.
3. O. D. Kellogg, *Foundations of potential theory*, Dover, New York, 1953.
4. J. R. Kuttler and V. G. Sigillito, *Inequalities for Stekloff and membrane eigenvalues*, J. Math. Anal. Appl. **23** (1968), 148-160.
5. L. E. Payne, *Isoperimetric inequalities and their applications*, SIAM Rev. **9** (1967), 453-488.
6. L. E. Payne and H. F. Weinberger, *New bounds in harmonic and biharmonic problems*, J. Math. Phys. **33** (1955), 291-307.
7. A. Pleijel, *Remarks on Courant's nodal line theorem*, Comm. Pure Appl. Math. **9** (1956), 543-550.
8. M. W. Stekloff, *Sur les problèmes fondamentaux de la physique mathématique*, Ann. Sci. École Norm. Sup. **19** (1902), 455-490.
9. R. Weinstock, *Inequalities for a classical eigenvalue problem*, J. Rational Mech. Anal. **3** (1954), 745-753.

THE JOHNS HOPKINS UNIVERSITY APPLIED PHYSICS LABORATORY