SPHERICAL HARMONIC EXPANSION OF THE
POISSON-SZEGÖ KERNEL FOR THE BALL

G. B. FOLLAND

ABSTRACT. By applying the theory of unitary spherical harmonics, we obtain an expansion of the Poisson-Szego kernel for the unit ball in complex $n$-space in terms of Jacobi polynomials and hypergeometric functions.

1. Dictionary. We list here the various constants and special functions we will be using. The reader should proceed to §2 and refer back here as necessary. For the relevant properties of these special functions we refer to the Bateman Manuscript Project [1].

(A) For $n = 1, 2, 3, \cdots$, $\omega_n$ is the area of the unit sphere in $\mathbb{R}^n$:

$$\omega_n = 2\pi^{n/2}\Gamma(n/2)^{-1}.$$ 

(B) For $k = 0, 1, 2, \cdots$ and $n = 2, 3, 4, \cdots$, $d(k; n)$ is the dimension of the space of homogeneous harmonic polynomials of degree $k$ on $\mathbb{R}^n$:

$$d(k; n) = (2k + n - 2)(k + n - 3)!(n - 2)!k!^{-1}.$$ 

(C) For $p, q = 0, 1, 2, \cdots$ and $n = 2, 3, 4, \cdots$, we set

$$D(p, q; n) = \frac{(p + q + n - 1)(p + n - 2)!(q + n - 2)!}{p!q!(n - 1)!(n - 2)!}.$$ 

(D) For $\alpha, \beta > -1$ and $m = 0, 1, 2, \cdots$, $P_{m}^{(\alpha, \beta)}$ is the Jacobi polynomial of degree $m$ associated to $(\alpha, \beta)$:

$$P_{m}^{(\alpha, \beta)}(t) = \frac{(1 - t)^{\alpha}(1 + t)^{\beta}}{m!2^{m}} \int \frac{d^m}{dt^m} [(1 - t)^{\alpha+m}(1 + t)^{\beta+m}].$$

We have

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\[
p_m^{(\alpha, \beta)}(1) = \Gamma(\alpha + m + 1)m!\Gamma(\alpha + 1)^{-1}.
\]

(E) For \( p, q = 0, 1, 2, \cdots; n = 2, 3, 4, \cdots, \) and \( z = re^{i\theta} \in \mathbb{C}, \) we set
\[
H_n^{p, q}(z) = D(p, q; n)r^{|p-q|}e^{i(p-q)\theta} \frac{p(n-2, |p-q|)(2r^2 - 1)}{\omega_{2n} p(n-2, |p-q|)}.
\]

(F) For \( \lambda > -\frac{1}{2} \) and \( m = 0, 1, 2, \cdots, \) \( C^\lambda_m \) is the Gegenbauer polynomial of degree \( m \) associated to \( \lambda: \)
\[
C^\lambda_m(t) = \frac{\Gamma(2\lambda + m)\Gamma(\lambda + \frac{1}{2})}{\Gamma(2\lambda)\Gamma(\lambda + m + \frac{1}{2})} p^{(\lambda-\frac{1}{2}, \lambda-\frac{1}{2})}(t).
\]
We have \( C^\lambda_m(1) = \Gamma(2\lambda + m)m!\Gamma(2\lambda)^{-1}. \)

(G) For \( k = 0, 1, 2, \cdots; n = 3, 4, 5, \cdots, \) and \( t \in \mathbb{R} \) we set
\[
G_n^k(t) = \frac{\omega_{n-k} C_{n/2}^{(n-2)/2}(t)}{\omega_n C_{n/2}^{(n-2)/2}(1)}.
\]

(H) For \( a, b \in \mathbb{R}, c > 0, \) and \( |t| < 1, \) \( F(a, b, c; t) \) is the hypergeometric function of \( t \) associated to \( (a, b, c): \)
\[
F(a, b, c; t) = \sum_{n=0}^{\infty} \frac{\Gamma(a+n)\Gamma(b+n)\Gamma(c)}{\Gamma(a)\Gamma(b)\Gamma(c+n)} t^n/n!.
\]
If \( c > a + b, \) \( \lim_{t \to 1-} F(a, b, c; t) \) exists and equals
\[
F(a, b, c; 1) = \frac{\Gamma(c)\Gamma(c-a-b)\Gamma(c-a)\Gamma(c-b)}{\Gamma(c-a)\Gamma(c-b)\Gamma(c)} - 1.
\]

(I) For \( p, q = 0, 1, 2, \cdots; n = 2, 3, 4, \cdots, \) and \( 0 \leq r \leq 1, \) we set
\[
S_n^{p, q}(r) = r^{p+q} \frac{F(p, q, p + q + n; r^2)}{F(p, q, p + q + 1)}.
\]

2. The Poisson kernel. We recall some facts from classical potential theory (cf. Stein-Weiss [8]). Let \( B_n \) be the open unit ball in \( \mathbb{R}^n, \) \( \Omega_n = \partial B_n \) the unit sphere, and \( \Delta \) the Laplace operator. Given \( f \in C(\Omega_n), \) the Dirichlet problem of finding \( u \in C(B_n) \) such that \( \Delta u = 0 \) in \( B_n \) and \( u|\Omega_n = f \) is solved by the formula
\[ u(x) = \int_{\Omega_n} P_n(x, \xi)f(\xi) \, d\xi, \]

where \( P_n \) is the Poisson kernel defined for \((x, \xi) \in B_n \times \Omega_n\) by

\[ P_n(x, \xi) = \frac{(1 - |x|^2)/\omega_n}{|x - \xi|^n}. \]

If \( n \geq 3 \), the orthogonal projection \( \pi_k \) of \( L^2(\Omega_n) \) onto the space of spherical harmonics of degree \( k \) is given by

\[ \pi_k f(\eta) = \int_{\Omega_n} G^k_n(\eta \cdot \xi)f(\xi) \, d\xi, \]

where \( \eta \cdot \xi \) denotes the standard scalar product of \( \eta \) and \( \xi \) in \( \mathbb{R}^n \). It is well known, and can be shown by the methods used below, that \( P_n \) \((n \geq 3)\) has the following expansion in spherical harmonics:

\[ P_n(\eta, \xi) = \sum_{k=0}^{\infty} r^k G^k_n(\eta \cdot \xi) \quad (0 \leq r < 1; \eta, \xi \in \Omega_n). \]

If \( n = 2 \), the corresponding formula is of course

\[ P_2(r e^{i\theta}, e^{i\phi}) = \sum_{k=-\infty}^{\infty} r^k e^{ik(\theta - \phi)} = 1 + 2 \sum_{k=1}^{\infty} r^k \cos k(\theta - \phi), \]

where we have identified \( \mathbb{R}^2 \) with \( \mathbb{C} \).

3. The Poisson-Szego kernel. We now consider \( B_{2n} \) and \( \Omega_{2n} \) as the unit ball and unit sphere in \( \mathbb{C}^n \), and we denote by \( \Delta_B \) the Laplace-Beltrami operator associated to the Bergman metric on \( B_{2n} \). We have

\[ \Delta_B = \frac{4}{n + 1} (1 - |z|^2) \sum_{i,j=1}^{n} (\delta_{ij} - z_i \overline{z_j}) \frac{\partial^2}{\partial z_i \partial \overline{z_j}}. \]

\( \Delta_B \) is the basic invariant differential operator on the symmetric space \( B_{2n} \simeq SU(n, 1)/U(n) \). Although \( \Delta_B \) degenerates at the boundary, the Dirichlet problem of finding \( u \in C(\overline{B_{2n}}) \) such that \( \Delta_B u = 0 \) in \( B_{2n} \) and \( u|\Omega_{2n} = f \) is well posed for \( f \in C(\Omega_{2n}) \), and the solution is given by

\[ u(z) = \int_{\Omega_{2n}} \mathcal{P}_n(z, \xi)f(\xi) \, d\xi, \]

where \( \mathcal{P}_n \) is the Poisson-Szego kernel defined for \((z, \xi) \in B_{2n} \times \Omega_{2n} \) by

\[ \mathcal{P}_n(z, \xi) = \frac{(1 - |z|^2)^n/\omega_{2n}}{1 - z \cdot \xi}^{2n}. \]

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(Here, and in the sequel, $z \cdot \zeta$ denotes the Hermitian scalar product of $z$ and $\zeta$ in $\mathbb{C}^n$.) The uniqueness of the solution $u$ follows from the mean value property of $\Delta_B$.

For the proofs of the foregoing facts, see Stein [7], Hua [4], and Helgason [3].

The object of this note is to derive the analogue of the formula (1) for $\mathcal{P}_n$ ($n \geq 2$). (If $n = 1$, $\Delta_B = \frac{1}{2}(1 - |z|^2)^2\Delta$ and $\mathcal{P}_1 = P_2$, so the formula in this case is simply (2).) Henceforth we assume $n \geq 2$.

Since $\Delta_B$ is invariant under $U(n)$ but not under $O(2n)$, the appropriate analogue of spherical harmonics is the decomposition of $L^2(\Omega_{2n})$ into its $U(n)$-irreducible components. This theory is developed in Ikeda [5], Vilenkin-Sapiro [9], Folland [2], and Koornwinder [6], and the principal result is as follows.

Proposition 1. $L^2(\Omega_{2n}) = \bigoplus_{p,q=0}^{\infty} \mathcal{H}^{p,q}$, where $\mathcal{H}^{p,q}$ is the space of restrictions to $\Omega_{2n}$ of harmonic polynomials $f(z, \zeta)$ on $\mathbb{C}^n$ which are homogeneous of degree $p$ in $z$ and degree $q$ in $\zeta$. $\mathcal{H}^{p,q}$ is $U(n)$-irreducible, and its dimension is $D(p, q; n)$.

To identify the orthogonal projection onto $\mathcal{H}^{p,q}$, we use the formula of Koornwinder [6]:

Proposition 2. If $f_1, f_2, \ldots, f_{D(p,q;n)}$ is any orthonormal basis for $\mathcal{H}^{p,q}$, then

$$
D(p,q;n) \sum_{j=1}^{D(p,q;n)} f_j(\zeta)f_j(\eta) = H^p,q(\zeta \cdot \eta).
$$

Consequently, the orthogonal projection $\pi_{p,q}: L^2(\Omega_{2n}) \rightarrow \mathcal{H}^{p,q}$ is given by

$$
\pi_{p,q}f(\zeta) = \int_{\Omega_{2n}} H^p,q(\zeta \cdot \eta)f(\eta) d\eta.
$$

We now identify the $\Delta_B$-harmonic extensions of elements of $\mathcal{H}^{p,q}$.

Proposition 3. If $f \in \mathcal{H}^{p,q}$, the (unique) solution of the Dirichlet problem $\Delta_B u = 0$ in $B_{2n'}$, $u|\Omega_{2n} = f$ is

$$
u(r \zeta) = S^{p,q}_n(r)f(\zeta) \quad (0 \leq r \leq 1, \ z \in \Omega_{2n}).
$$

Proof. One element of $\mathcal{H}^{p,q}$ is the restriction $f_0$ to $\Omega_{2n}$ of $F_0(z) = z_1^{p-2}z_2$. Since $\mathcal{H}^{p,q}$ is irreducible, every element of $\mathcal{H}^{p,q}$ is a linear combination of translates of $f_0$ by elements of $U(n)$. Since moreover $\Delta_B$
commutes with the action of $U(n)$, it will suffice to prove the proposition for $f = f_0$. Thus we seek a $\Delta_B$-harmonic function on $B_{2n}$ of the form $u(z) = g(r^2)z_1^p z_2^q$ where $r = |z|$. (We choose $g$ to depend on $r^2$ instead of $r$ because $u(z)$ must be smooth at $z = 0$.) A straightforward computation shows that

$$
\frac{\partial^2 u}{\partial z_i \partial \overline{z}_j} = \left[ g''(r^2)z_1^{p-1}z_2^q + \delta_{ij}g'(r^2) \right] z_1^p z_2^q + \delta_{i1}g'(r^2)z_1^{p-1}z_2^q - \delta_{j2}pg(r^2)z_1^p z_2^{p-1}.
$$

Hence

$$
\sum_{i=1}^n \frac{\partial^2 u}{\partial z_i \partial \overline{z}_i} = [r^2 g''(r^2) + (n + p + q)g'(r^2)] z_1^p z_2^q,
$$

and

$$
\sum_{i,j=1}^n z_i z_j \frac{\partial^2 u}{\partial z_i \partial \overline{z}_j} = [r^4 g''(r^2) + (p + q + 1)r^2 g'(r^2) + pqg(r^2)] z_1^p z_2^q,
$$

so that

$$
\Delta_B u = \frac{4}{n+1} (1 - r^2)z_1^p z_2^q r^2 (1 - r^2)g''(r^2) + [(p + q + n) - (p + q + 1)t]g'(t) - pqg(t) = 0.
$$

Therefore if $\Delta_B u = 0$, $g$ must satisfy

$$
t(1 - t)g''(t) + [(p + q + n) - (p + q + 1)t]g'(t) - pqg(t) = 0.
$$

But this is the hypergeometric equation with parameters $a = p$, $b = q$, $c = p + q + n$, and the only solutions which are smooth at $t = 0$ are multiples of $F(p, q, p + q + n; t)$. Thus, writing $z = r\zeta$ with $\zeta \in \Omega_{2n}$,

$$
u(z) = cF(p, q, p + q + n; r^2)z_1^p z_2^q = cF(p, q, p + q + n; r^2) r^{p+q}f_0(\zeta),
$$

and the requirement that $u|_{\Omega_{2n}} = f$ shows that $c = |F(p, q, p + q + n; 1)|^{-1}$. This completes the proof.

Remark. Except in the trivial cases $p = 0$ or $q = 0$ (when $F(p, q, p + q + n; r^2) \equiv 1$), $S_n^{p,q}(r)$ has a branch point at $r = 1$. Thus, in contradistinction to the real case, the $\Delta_B$-harmonic extension of $f \in S_n^{p,q}$ to $B_{2n}$ cannot be analytically continued outside $B_{2n}$. This is not surpris-
ing since the boundary is "at infinity" from the point of view of the Bergman metric.

We now state the main result.

Theorem. For $0 < r < 1$ and $\eta, \zeta \in \Omega_{2n}$.

\begin{equation}
\mathcal{P}_n(r \eta, \zeta) = \sum_{p, q=0}^{\infty} S_n^{p, q}(r) H_n^{p, q}(\eta \cdot \zeta).
\end{equation}

The series on the right converges absolutely and uniformly for $\eta, \zeta \in \Omega_{2n}$ and $0 < r < \rho$ for each $\rho < 1$.

Proof. We first draw some consequences from Proposition 2. Let $f_1, \ldots, f_{D(p, q; n)}$ be any orthonormal basis for $\mathcal{S}^{p, q}$; then, first,

$$H_n^{p, q}(1) = H_n^{p, q}(\eta \cdot \eta) = \sum_{j=1}^{D(p, q; n)} |f_j(\eta)|^2$$

for any $\eta \in \Omega_{2n}$. Second,

$$H_n^{p, q}(1) = \omega_{2n}^{-1} \int_{\Omega_{2n}} H_n^{p, q}(1) \, d\eta = \omega_{2n}^{-1} \sum_{j=1}^{D(p, q; n)} \int_{\Omega_{2n}} |f_j(\eta)|^2 \, d\eta = \omega_{2n}^{-1} D(p, q; n).$$

Third, by orthogonality of the $f_j$'s,

$$\int_{\Omega_{2n}} |H_n^{p, q}(\eta \cdot \zeta)|^2 \, d\eta = \sum_{j=1}^{D(p, q; n)} \int_{\Omega_{2n}} |f_j(\eta)|^2 |f_j(\zeta)|^2 \, d\eta = \sum_{j=1}^{D(p, q; n)} |f_j(\zeta)|^2 = H_n^{p, q}(1) = \omega_{2n}^{-1} D(p, q; n).$$

Finally, $H_n^{p, q}(\eta \cdot \zeta)$ is Hermitian-symmetric in $\eta$ and $\zeta$ and is in $\mathcal{S}^{p, q}$ as a function of $\eta$ for each fixed $\zeta$; therefore by the Schwarz inequality,

$$|H_n^{p, q}(\eta \cdot \zeta)| \leq \int_{\Omega_{2n}} H_n^{p, q}(\eta \cdot \xi) H_n^{p, q}(\xi \cdot \zeta) \, d\xi \leq \omega_{2n}^{-1} D(p, q; n).$$

Moreover, since $\bigoplus_{p+q=k} \mathcal{S}^{p, q}$ is the space of all spherical harmonics of degree $k$,
Now since \( F(p, q; p + q + n; t) \) is an increasing function for \( t \in [0, 1] \), we have \( |S^p,q_n(r)| \leq r^{p+q} \). Thus if \( r \leq p < 1 \),

\[
\sum_{p, q=0}^\infty |S^p,q_n(r)H^p,q_n(\eta \cdot \zeta)| \leq \text{const} \sum_{p, q=0}^\infty r^{p+q}(p + q + 1)^{2n}
\]

\[
= \text{const} \sum_{k=0}^\infty r^k(k + 1)^{2n+1} < \infty.
\]

This establishes the second assertion.

In particular, the right-hand side of (3) converges uniformly to a continuous function of \( \zeta \) for each fixed \( r \in [0, 1) \) and \( \eta \in \Omega_{2n} \). Thus to prove (3) it will suffice to show that for each \( r \in [0, 1) \), \( \eta \in \Omega_{2n} \), and \( f \in C(\Omega_{2n}) \),

\[
\int_{\Omega_{2n}} S^p,q_n(r, \Omega)(\eta, \zeta) f(\zeta) d\zeta = \sum_{p, q=0}^\infty \int_{\Omega_{2n}} S^p,q_n(r)H^p,q_n(\eta \cdot \zeta) f(\zeta) d\zeta.
\]

But on interchanging the summation and integration, we see from Propositions 2 and 3 that this is true when \( f \) is a finite sum of spherical harmonics, and since such functions are dense in \( C(\Omega_{2n}) \), it is true in general. The proof is complete.

REFERENCES


DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WASHINGTON, SEATTLE, WASHINGTON 98195