

## GAUSSIAN MEASURE OF LARGE BALLS IN A HILBERT SPACE<sup>1</sup>

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**ABSTRACT.** Let  $P$  be a zero mean Gaussian measure in a Hilbert space. The asymptotic behavior of  $P\{\|x - b\|^2 > \varepsilon\}$  as  $\varepsilon \rightarrow \infty$  is studied in this note.

Let  $P$  be a Gaussian measure in a separable Hilbert space  $\mathcal{H}$  and  $b$  be an arbitrary fixed element in  $\mathcal{H}$ . We shall study how fast  $P\{\|x - b\|^2 > \varepsilon\} \rightarrow 0$  as  $\varepsilon \rightarrow \infty$ .

Without loss of generality,  $P$  is assumed to be of mean 0. Let  $B$  be the covariance operator of  $P$ , and let the eigenvalues (corresponding eigenvectors) of  $B$  be ordered by  $\lambda_1 \geq \lambda_2 \geq \dots$  ( $\{e_i\}$ ). Let  $k$  be the multiplicity of the largest eigenvalue,  $x_i = \langle x, e_i \rangle$ ,  $b_i = \langle b, e_i \rangle$  and  $a = (\sum_1^k b_i^2)^{1/2}$ . Finally, let  $F$  denote the distribution function of  $\|x - b\|^2 = \sum_1^\infty (x_i - b_i)^2$ . The Laplace transform of  $F$  is

$$\begin{aligned} \phi(c) &= \int_0^\infty c^{-ct} dF(t) = \prod_{i=1}^\infty [(1 + 2c\lambda_i)^{-1/2} \exp(-b_i^2 c(1 + 2c\lambda_i)^{-1})], \\ \psi(c) &= \int_0^\infty e^{-ct + t/2\lambda_1} (1 - F(t)) dt \\ &= \left(1 - \phi\left(c - \frac{1}{2\lambda_1}\right)\right) \left(c - \frac{1}{2\lambda_1}\right)^{-1} \\ &= \left(c - \frac{1}{2\lambda_1}\right)^{-1} (2c\lambda_1)^{-k/2} \exp\left(\frac{a^2}{4c\lambda_1^2} - \frac{a^2}{2\lambda_1}\right) \\ &\quad \times \left[ (2c\lambda_1)^{k/2} \exp\left(\frac{a^2}{2\lambda_1} - \frac{a^2}{4c\lambda_1^2}\right) - \prod_{k+1}^\infty \left(1 + 2c\lambda_i - \frac{\lambda_i}{\lambda_1}\right)^{-1/2} \right. \\ &\quad \left. \times \exp - \sum_{k+1}^\infty b_i^2 \left(c - \frac{1}{2\lambda_1}\right) \left(1 + 2c\lambda_i - \frac{\lambda_i}{\lambda_1}\right)^{-1} \right]. \quad (1) \end{aligned}$$

If  $a = 0$ , then

$$\psi(c) \sim (2\lambda_1)^{1-k/2} \prod_{k+1}^\infty \left(1 - \frac{\lambda_i}{\lambda_1}\right)^{-1/2} \exp\left(\frac{1}{2\lambda_1} \sum_{k+1}^\infty b_i^2 \left(1 - \frac{\lambda_i}{\lambda_1}\right)\right)^{-1} c^{-k/2}$$

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as  $c \rightarrow 0$ . By a Tauberian theorem [1], as  $\varepsilon \rightarrow \infty$ ,

$$\int_0^\varepsilon e^{t/2\lambda_1}(1 - F(t)) dt \sim \frac{1}{\Gamma(k/2 + 1)} (2\lambda_1)^{1-k/2} \prod_{k+1}^\infty \left(1 - \frac{\lambda_i}{\lambda_1}\right)^{-1/2} \\ \times \exp\left(\frac{1}{2\lambda_1} \sum_{k+1}^\infty b_i^2 \left(1 - \frac{\lambda_i}{\lambda_1}\right)^{-1}\right) \varepsilon^{k/2}.$$

Since  $e^{t/2\lambda_1}(1 - F(t))$  is continuous and positive, L'Hôpital's rule is applicable. Hence we have

**THEOREM 1.** *If  $a = (\sum_1^k b_i^2)^{1/2} = 0$ , then*

$$P\{\|x - b\|^2 > \varepsilon\} \sim K_1 e^{-\varepsilon/2\lambda_1} \varepsilon^{k/2-1},$$

where

$$K_1 = \frac{1}{\Gamma(k/2)} (2\lambda_1)^{1-k/2} \prod_{k+1}^\infty \left(1 - \frac{\lambda_i}{\lambda_1}\right)^{-1/2} \exp \frac{1}{2\lambda_1} \sum_{k+1}^\infty b_i^2 \left(1 - \frac{\lambda_i}{\lambda_1}\right)^{-1}. \quad (2)$$

If  $a > 0$ , then  $\psi(c) \rightarrow \infty$  exponentially fast as  $c \rightarrow 0$ . The ordinary Tauberian theorem is not applicable here. We shall try another approach.

Let  $\bar{b} \in \mathcal{K}$  such that  $\bar{b}_i = 0$  if  $i < k$  and  $\bar{b}_i = b_i$  otherwise. Then, for  $\varepsilon$  large enough

$$P\{\|x - b\|^2 > \varepsilon\} \\ \leq P\{\|x - \bar{b}\|^2 > (\sqrt{\varepsilon} - a)^2\} \sim K_1 \left(\exp \frac{-1}{2\lambda_1} (\sqrt{\varepsilon} - a)^2\right) \varepsilon^{k/2-1}.$$

On the other hand,  $P\{\|x - b\|^2 > \varepsilon\}$  may be regarded as  $P_b\{\|x\|^2 > \varepsilon\}$ , where  $P_b$  has mean  $-b$  and covariance  $B$ . It is easily seen that  $P_b$  is equivalent to  $P_{\bar{b}}$  and

$$\frac{dP_b}{dP_{\bar{b}}}(x) = \exp\left(\frac{-\sum_1^k x_i b_i}{\lambda_1} - \frac{a^2}{2\lambda_1}\right).$$

Therefore

$$P\{\|x - b\|^2 > \varepsilon\} = P_b\{\|x\|^2 > \varepsilon\} \\ = \int_{\|x\|^2 > \varepsilon} \exp\left(\frac{-\sum_1^k x_i b_i}{\lambda_1} - \frac{a^2}{2\lambda_1}\right) dP_{\bar{b}}(x) \\ > (2\pi\lambda_1)^{-k/2} \int_{x_1^2 + \dots + x_k^2 > \varepsilon} \exp\left(\frac{-1}{2\lambda_1} \sum_{i=1}^k (x_i + b_i)^2\right) dx_1 \cdots dx_k \\ = (2\pi\lambda_1)^{-k/2} \int_{x_1^2 + \dots + x_k^2 > \varepsilon} \\ \times \exp\left(\frac{-1}{2\lambda_1} [(x_1 - a)^2 + x_2^2 + \dots + x_k^2]\right) dx_1 \cdots dx_k. \quad (3)$$

(To get the last equality, just rotate the vector  $(-b_1, \dots, -b_k)$  to the first coordinate.)

For  $k > 3$ , change the last  $k - 1$  coordinates to polar coordinates; (3) becomes

$$\begin{aligned} & (2\pi\lambda_1)^{-k/2} \frac{2\pi^{(k-1)/2}}{\Gamma((k-1)/2)} \int_{\substack{x_1^2 + r^2 > \varepsilon \\ r > 0}} \left( \exp - \frac{(x_1 - a)^2 + r^2}{2\lambda_1} \right) r^{k-2} dr dx_1 \\ &= (2\pi\lambda_1)^{-k/2} \frac{2\pi^{(k-1)/2}}{\Gamma((k-1)/2)} \\ & \quad \times \int_0^\pi \int_{\sqrt{\varepsilon}}^\infty \left( \exp \frac{\rho^2 + a^2 - 2\rho a \cos \theta}{2\lambda_1} \right) \rho(\rho \sin \theta)^{k-2} d\rho d\theta. \end{aligned}$$

By L'Hôpital's rule

$$\begin{aligned} & \lim_{\varepsilon \rightarrow \infty} \frac{\int_0^\pi \int_{\sqrt{\varepsilon}}^\infty (\exp((2\rho a \cos \theta - \rho^2 - a^2)/2\lambda_1)) \rho(\rho \sin \theta)^{k-2} d\rho d\theta}{\varepsilon^{(k-3)/4} \exp(-(\sqrt{\varepsilon} - a)^2/2\lambda_1)} \\ &= \lim_{\varepsilon \rightarrow \infty} \frac{\varepsilon^{(k-2)/2} (\exp - (\varepsilon + a^2)/2\lambda_1) \int_0^\pi (\exp(\sqrt{\varepsilon} a \cos \theta)/\lambda_1) (\sin \theta)^{k-2} d\theta}{\lambda_1^{-1} \varepsilon^{(k-3)/4} \exp(-(\sqrt{\varepsilon} - a)^2/2\lambda_1)}. \end{aligned} \tag{4}$$

And, by Laplace's method [2]

$$\int_0^\pi \exp \frac{\sqrt{\varepsilon} a \cos \theta}{\lambda_1} (\sin \theta)^{k-2} d\theta \sim \frac{1}{2} \Gamma\left(\frac{k-1}{2}\right) \left(\frac{2\lambda_1}{a\sqrt{\varepsilon}}\right)^{(k-1)/2} \exp \frac{\sqrt{\varepsilon} a}{\lambda_1}.$$

Hence (4) is

$$\frac{1}{2} \lambda_1 \Gamma\left(\frac{k-1}{2}\right) \left(\frac{2\lambda_1}{a}\right)^{(k-1)/2},$$

and, consequently, (3) is asymptotic to

$$\sqrt{\frac{\lambda_1}{2\pi}} a^{(1-k)/2} \varepsilon^{(k-3)/4} \exp - \frac{1}{2\lambda_1} (\sqrt{\varepsilon} - a)^2.$$

(The special cases  $k = 1$  and  $k = 2$  easily lead to the same expression.)

To sum up, we have

**THEOREM 2.** *If  $a > 0$ , then*

$$\limsup_{\varepsilon \rightarrow \infty} P \{ \|x - b\|^2 > \varepsilon \} \varepsilon^{1-k/2} \exp \frac{1}{2\lambda_1} (\sqrt{\varepsilon} - a)^2 < K_1,$$

where  $K_1$  is defined by (2);

$$\liminf_{\varepsilon \rightarrow \infty} P \{ \|x - b\|^2 > \varepsilon \} \varepsilon^{(3-k)/4} \exp \frac{1}{2\lambda_1} (\sqrt{\varepsilon} - a)^2 > \sqrt{\frac{\lambda_1}{2\pi}} a^{(1-k)/2}.$$

REMARKS. 1. Zolotarev discussed the limiting behavior of  $P\{\sum_1^\infty x_i^2 > \varepsilon\}$  as  $\varepsilon \rightarrow \infty$  in [4], where  $\{x_i\}$  are independent  $\mathcal{U}(0, \sigma_i^2)$  with  $\sum_1^\infty \sigma_i^2 < \infty$ . This is a special case of Theorem 1. The proof here is much simpler.

2. It is tempting to try to apply Theorem 3 of [4], but this would require establishing that  $e^{e/(2\lambda)}(1 - F(\varepsilon))$  is nondecreasing, and would still give a weakened version of Theorem 2. Of course, our theorem also leaves open the question of the exact asymptotic behavior of  $1 - F(\varepsilon)$ , when  $a > 0$ .

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