

## ON FOURIER INTEGRAL OPERATORS

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**ABSTRACT.** We consider operators of the form:  $\int_{-\infty}^{\infty} F_t \varphi(t) dt$ , where  $F_t$  is a 1-parameter family of Fourier integral operators and  $\varphi(t) dt$  a tempered distribution on the real line and show that these operators are sums of pseudo-differential and Fourier integral operators. Here, we give the typical case where  $\varphi(t) dt = \text{p.v.}\{1/t\}$ . An application to singular integrals on variable curves is given.

**THEOREM.** Let  $F_t$  be a 1-parameter family of local Fourier integral operators such that the kernel of  $F_t$  has an integral representation of the form:  $\int_{\mathbf{R}^2} e^{i\phi(x,y,t,\theta)} a(x,y,t,\theta) d\theta$  with  $a$  and  $\phi$  depending smoothly on  $t$ ,  $a$  (resp.  $\phi$ ) a symbol in the class  $S_{1,0}^0$  i.e.:  $|\partial_x^\alpha \partial_y^\beta \partial_t^k \partial_\theta^\gamma a| \leq C_{\alpha,\beta,k,\gamma} (1+|\theta|)^{-|\gamma|}$  (resp. nondegenerate phase function) and  $a(x,y,0,\theta) = 1$  (resp.  $\phi(x,y,0,\theta) = (x-y) \cdot \theta$ ). Furthermore, we assume that  $\dot{\phi}_\theta \times \ddot{\phi}_\theta \neq 0$ ; then the operator  $T = \int_{-1}^1 F_t dt/t$  can be written in the form  $T = P + F$  where  $P$  is a  $\psi$ .d.o. with symbol in  $S_{1/2,1/2}^0$  and  $F$  a Fourier integral operator with amplitude in  $S_{1,0}^{-1/2}$ .

The proof follows from the asymptotic expansion of the function

$$\omega(x,y,\theta) = B(\theta) \int_{-1}^1 e^{i\phi(x,y,t,\theta)} a(x,y,t,\theta) \frac{dt}{t}$$

where  $B$  is a  $C^\infty$  function on  $\mathbf{R}^2$ , identically zero near the origin and homogeneous of degree zero for all  $|\theta| \geq 1$ . For simplicity, we will not, in the sequel, make mention of  $B$ .

We, now, use Taylor expansion to have:

$$\phi(x,y,t,\theta) = (x-y) \cdot \theta + tv(x,y,\theta) + t^2 w(x,y,t,\theta)/2$$

where  $v = v(x,y,\theta) = \dot{\phi}(x,y,0,\theta)$ . We will also write  $w = w(x,y,\theta) = \ddot{\phi}(x,y,0,\theta)$ .

From the hypothesis, we have  $v_\theta \times w_\theta \neq 0$ ; hence using the homogeneity of  $v$  and  $w$ , for any  $(x_0, y_0, \theta_0)$ , there exists a conic neighborhood  $K \times \Gamma$  of  $(x_0, y_0, \theta_0)$  on which either  $|v(x,y,\theta)| \geq a|\theta|$  or  $|w(x,y,\theta)| \geq b|\theta|$  where  $a, b$  are some positive constants depending only on  $K \times \Gamma$ . We consider separately the two types of cones. We let

$$\Gamma_1 = \{\theta \in \mathbf{R}^2 : v(x,y,\theta) \geq a|\theta| \text{ for all } (x,y) \in K\},$$

$$\Gamma_2 = \{\theta \in \mathbf{R}^2 : w(x,y,\theta) \geq b|\theta| \text{ for all } (x,y) \in K\}.$$

The two other cones can be considered in a similar fashion.

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Received by the editors September 14, 1981.

1980 *Mathematics Subject Classification.* Primary 35S99, 42B99.

*Key words and phrases.* Fourier transform, operator, singular integral.

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0002-9939/82/0000-0261/\$02.00

We first consider  $\Gamma_2$ ,

LEMMA 1. *The symbol*

$$\sigma_1 = \int_{|t| < 1/\sqrt{w}} \exp \left\{ itv + \frac{t^2}{2} w(x, y, t, \theta) \right\} a(x, y, t, \theta) \frac{dt}{t}$$

belongs to  $S_{1/2, 1/2}^0(K \times \Gamma_2)$ .

PROOF. We write

$$\sigma_1 = \int_{|t| \leq 1/\sqrt{w}} \exp \left\{ itv + \frac{t^2}{2} w \right\} \frac{dt}{t} + E_1$$

where

$$\begin{aligned} E_1 &= \int_{|t| \leq 1/\sqrt{w}} \exp \left\{ itv + \frac{t^2}{2} w \right\} \left( \exp \left\{ i \frac{t^2}{2} (w(x, y, t, \theta) - w) \right\} a(x, y, t, \theta) - 1 \right) \frac{dt}{t} \\ &= \int_{|t| \leq 1} \exp \left\{ it \frac{v}{\sqrt{w}} \right\} \frac{1}{t} \left( \exp \left\{ i \frac{t^2}{2} (w(t/\sqrt{w})/w - 1) \right\} a\left(\frac{t}{\sqrt{w}}\right) - 1 \right) \exp \left\{ i \frac{t^2}{2} \right\} dt \end{aligned}$$

where  $w(t) = w(x, y, t, \theta)$  and  $a(t) = a(x, y, t, \theta)$ . We then have

- (1)  $e^{itv/\sqrt{w}}$  belongs to  $S_{1/2, 1/2}^0(K \times \Gamma_2)$  uniformly in  $t$ ,  $|t| \leq 1$ ,
- (2)  $w(t/\sqrt{w})/w$  belongs to  $S_{1,0}^0(K \times \Gamma_2)$  uniformly in  $t$ ,  $|t| \leq 1$ . Also  $\underline{a} = a(t/\sqrt{w})$  verifies:

$$|\partial_x^\alpha \partial_y^\beta \partial_t^k \partial_\theta^\gamma \underline{a}| \leq c_{\alpha, \beta, k, \gamma} (1 + |\theta|)^{-(|\gamma| + k/2)}$$

and

$$\left| \frac{1}{t} \left( \exp \left\{ i \frac{t^2}{2} (w(t/\sqrt{w})/w - 1) \right\} a(t/\sqrt{w}) - 1 \right) \right| \leq (\text{const})(1 + |\theta|)^{-1/2}$$

on  $K \times \Gamma_2$ . Thus  $\frac{1}{t} (\exp \{ i \frac{t^2}{2} (w(t/\sqrt{w})/w - 1) \} a(t/\sqrt{w}) - 1)$  and therefore  $E_1$ , verifies the conclusion of Lemma 1.

For  $\sigma_1 - E_1$ , we have

$$\begin{aligned} \sigma_1 - E_1 &= \int_{|t| \leq 1} e^{it \frac{v}{\sqrt{w}}} \frac{e^{it^2/2}}{t} dt = \int_{|t| \leq 1} \exp \left\{ it \frac{v}{\sqrt{w}} \right\} \frac{dt}{t} \\ &\quad + \int_{|t| \leq 1} \exp \left\{ it \frac{v}{\sqrt{w}} \right\} \frac{\exp \left\{ i \frac{t^2}{2} \right\} - 1}{t} dt = \sigma_{11} + \sigma_{12}. \end{aligned}$$

$\sigma_{12}$  is easily seen to be in  $S_{1/2, 1/2}^0(K \times \Gamma_2)$  since  $v/\sqrt{w}$  is in  $S_{1/2, 1/2}^1(K \times \Gamma_2)$  and  $\frac{1}{t}(e^{it^2} - 1)$  is a smooth function. For  $\sigma_{11}$ , we write

$$\sigma_{11} = \int_0^{v/\sqrt{w}} \frac{\sin u}{u} du = F\left(\frac{v}{\sqrt{w}}\right)$$

from which one deduces that  $\sigma_{11}$  belongs to  $S_{1/2, 1/2}^0(K \times \Gamma_2)$ . This completes the proof of Lemma 1.

We now wish to estimate,

$$\sigma_2 = \int_{1/\sqrt{w} \leq |t| \leq 1} e^{i\phi(x,y,t,\theta)} a(x,y,t,\theta) \frac{dt}{t}.$$

We let  $t(x,y,\theta)$  be such that  $\dot{\phi}(x,y,t(x,y,\theta),\theta) = 0$ . From the hypothesis we have  $\phi(x,y,t(x,y,\theta),\theta) \neq 0$ . Also,  $t(x,y,\theta)$  is clearly homogeneous of degree 0 in  $\theta$ .

LEMMA 2. *The operator corresponding to  $\sigma_2$  is a Fourier integral operator with phase function  $\tilde{\phi}(x,y,\theta) = \phi(x,y,t(x,y,\theta),\theta)$  and amplitude in  $S_{1,0}^{-1/2}(K \times \Gamma_2)$ .*

PROOF. We use the stationary phase method to write the full expansion of  $\sigma_2$  whose first term is

$$\frac{e^{i\pi/4}}{t(x,y,\theta)} e^{i\phi(x,y,t(x,y,\theta),\theta)} \frac{a(x,y,t(x,y,\theta),\theta)}{\sqrt{\tilde{\phi}(x,y,t(x,y,\theta),\theta)}}.$$

Clearly  $\tilde{\phi} = \phi(x,y,t(x,y,\theta),\theta)$  is homogeneous of degree 1 in  $\theta$ . Also from  $\dot{\phi}(x,y,t(x,y,\theta),\theta) = 0$ , we see that

$$\begin{aligned}\tilde{\phi}_x &= \phi_x(x,y,t,\theta)|_{t=t(x,y,\theta)}, & \tilde{\phi}_\theta &= \phi_\theta(x,y,t,\theta)|_{t=t(x,y,\theta)}, \\ \tilde{\phi}_{x\theta} &= \phi_{x\theta}(x,y,t,\theta)|_{t=t(x,y,\theta)}\end{aligned}$$

which shows that  $\tilde{\phi}$  is nondegenerate since  $\phi$  is. Also from  $\dot{\phi}(x,y,t(x,y,\theta),\theta) = 0$  we have  $\dot{\phi}_x + \dot{\phi}_{t_x} = 0$ ; hence  $t_x = -\dot{\phi}/\dot{\phi}$ ; in particular  $t_x$  is bounded on  $K \times \Gamma_2$ . Similarly for  $t_y$  and all higher derivatives in  $x$  and  $y$  of  $t(x,y,\theta)$ . Now, it is clear that

$$\frac{1}{t(x,y,\theta)} \frac{a(x,y,t(x,y,\theta),\theta)}{\sqrt{\tilde{\phi}(x,y,t(x,y,\theta),\theta)}}$$

belongs to  $S_{1,0}^{-1/2}(K \times \Gamma_2)$ . The expansion given by the stationary phase method gives an asymptotic expansion of the amplitude (use formula 2.14, p. 431, [2] with  $\rho = 1$   $f(t) = \phi(x,y,t,\phi)$  and  $g(t) = a(t)/t$  and the fact that  $t(x,y,\theta)$  stays away from 0 on  $K \times \Gamma_2$ ).

We now consider  $\Gamma_1$ ,

LEMMA 3. *The symbol*

$$\nu_1 = \int_{|t| \leq 1/\sqrt{v}} \exp\left\{itv + \frac{t^2}{2}w(t)\right\} a(t) \frac{dt}{t}$$

belongs to  $S_{1/2,1/2}^0(K \times \Gamma_1)$ .

PROOF. We write

$$\nu_1 = \int_{|t| \leq 1/\sqrt{v}} \exp\left\{itv + \frac{t^2}{2}w\right\} \frac{dt}{t} + F_1$$

where

$$\begin{aligned}F_1 &= \int_{|t| \leq 1/\sqrt{v}} \exp\left\{itv + \frac{t^2}{2}w\right\} \left(\exp\left\{i\frac{t^2}{2}(w(x,y,t,\theta) - w)\right\} a(x,y,t,\theta) - 1\right) \frac{dt}{t} \\ &= \frac{1}{\sqrt{v}} \int_{|t| \leq 1} e^{it\sqrt{v}} \frac{1}{t/\sqrt{v}} \left(\exp\left\{i\frac{t^2}{2}\left(w\left(x,y,\frac{t}{\sqrt{v}},\theta\right) - w\right)/v\right\}\right. \\ &\quad \left.\cdot a\left(x,y,\frac{t}{\sqrt{v}},\theta\right) - 1\right) e^{it^2/2} dt\end{aligned}$$

which is easily seen to be in  $S_{1,0}^{-1/2}(K \times \Gamma_1)$  (see similar argument for  $E_1$ ).

For  $\nu_1 - F_1$ , we have

$$\begin{aligned}\nu_1 - F_1 &= \int_{|t| \leq 1} \exp \left\{ it\sqrt{v} + \frac{t^2 w}{2v} \right\} \frac{dt}{t} \\ &= \int_{|t| \leq 1} e^{it\sqrt{v}} \frac{dt}{t} + \frac{w}{v} \int_{|t| \leq 1} e^{it\sqrt{v}} \frac{e^{it^2 w/v} - 1}{t^2 \frac{w}{v}} t dt.\end{aligned}$$

Both terms are clearly in  $S_{1/2,1/2}^0(K \times \Gamma_1)$  which finishes the proof of Lemma 3.

We now let  $\psi_0(u)$  be a smooth function,  $\psi_0(-u) = \psi_0(u)$ ,  $\psi_0 \equiv 1$  for  $-\frac{1}{2} \leq u \leq \frac{1}{2}$  and  $\psi_0 \equiv 0$  for  $|u| \geq 1$  and put  $\psi_1 = 1 - \psi_0$ . The conclusion of Lemma 3 is clearly also valid for

$$\nu'_1 = \int_{-1}^1 \exp \left\{ itv + \frac{t^2}{2} w(x, y, t, \theta) \right\} a(x, y, t, \theta) \psi_0(t\sqrt{v}) \frac{dt}{t}.$$

LEMMA 4. *The operator corresponding to*

$$\nu'_2 = \int_{-1}^1 \exp \left\{ itv + \frac{t^2}{2} w(x, y, t, \theta) \right\} a(x, y, t, \theta) \psi_1(t\sqrt{v}) \frac{dt}{t}$$

*is a Fourier integral operator with phase function equal to  $\phi(x, y, \pm 1, \theta)$  and amplitude in  $S_{1,0}^{-1}(K \times \Gamma_1)$ .*

PROOF. We write

$$\begin{aligned}\nu'_2 &= \int_{-\sqrt{v}}^{\sqrt{v}} e^{it\sqrt{v}} \exp \left\{ i \frac{t}{2} w \left( x, y, \frac{t}{\sqrt{v}}, \theta \right) / v \right\} a \left( x, y, \frac{t}{\sqrt{v}}, \theta \right) \psi_1(t) \frac{dt}{t} \\ &= \int_{-\sqrt{v}}^{-1} + \int_1^{\sqrt{v}}.\end{aligned}$$

Now, an integration by parts gives

$$\begin{aligned}\int_1^{\sqrt{v}} &= \frac{1}{v} e^{iv + \frac{1}{2} w(x, y, 1, \theta)} a(x, y, 1, \theta) \psi_1(\sqrt{v}) \\ &\quad - \frac{1}{\sqrt{v}} \int_1^{\sqrt{v}} e^{it\sqrt{v}} \frac{d}{dt} \left( \exp \left\{ i \frac{t^2}{2} w \left( x, y, \frac{t}{\sqrt{v}}, \theta \right) / v \right\} a \left( x, y, \frac{t}{\sqrt{v}}, \theta \right) \psi_1(t) \frac{1}{t} \right) dt.\end{aligned}$$

Since  $e^{iv + w(x, y, 1, \theta)/2} = e^{i\phi(x, y, 1, \theta) - i(x-y)\theta}$ , the first term verifies clearly the conclusion of Lemma 4. Also, in the second term, a repeated integration by parts gives an asymptotic expansion of the amplitude. We have the same conclusion for the integral  $\int_{-\sqrt{v}}^{-1}$ . This finishes the proof of Lemma 4.

The proof of our result follows from the above lemmas and the use of a micro-local partition of unity associated with the different cones and a neighborhood of the origin. Notice that

$$\begin{aligned}E(x, y, \theta) &= (1 - B(\theta)) \int_{-1}^1 (e^{i\phi(x, y, t, \theta)} a(x, y, t, \theta)) \frac{dt}{t} \\ &= (1 - B(\theta)) \int_{-1}^1 (e^{i\phi(x, y, t, \theta)} a(x, y, t, \theta) - e^{i(x-y)\cdot\theta}) \frac{dt}{t}\end{aligned}$$

is a symbol of a smoothing operator.

We have the following:

COROLLARY. Let  $\gamma(x, t)$  be a variable curve of class  $C^\infty$  in  $\mathbf{R}^2$  such that  $\gamma(x, 0) = 0$  and  $\dot{\gamma} \times \ddot{\gamma} \neq 0$  then the operator

$$\mathcal{H}f(x) = \text{p.v.} \int_{-1}^1 f(x - \gamma(x, t)) \frac{dt}{t}$$

is a bounded operator from  $L^2_{\text{comp}}(\mathbf{R}^2)$  to  $L^2_{\text{loc}}(\mathbf{R}^2)$  (see [3]).

PROOF. We write

$$f(x - \gamma(x, t)) = \int_{\mathbf{R}^2} \int_{\mathbf{R}^2} e^{i[(x-y)\theta - \gamma(x, t) \cdot \theta]} f(y) dy d\theta$$

and apply our result with  $a \equiv 1$  and  $\phi(x, y, t, \theta) = (x - y)\theta - \gamma(x, t) \cdot \theta$ . In particular there exists a  $\psi$ .d.o.  $P$  with symbol in  $S^0_{1/2, 1/2}$  such that  $(\mathcal{H} - P)^*(\mathcal{H} - P)$  is a  $\psi$ .d.o. with symbol in  $S^{-1}_{1, 0}$ .

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