

## THE ISOPERIMETRIC INEQUALITY FOR NONSIMPLE CLOSED CURVES

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**ABSTRACT.** The main purpose of this paper is the generalization to the hyperbolic and elliptic spaces of the isoperimetric inequality of Banchoff and Pohl (J. Differential Geom. 6 (1971), 175–192).

1

Banchoff and Pohl [1] showed that the classical isoperimetric inequality still holds for nonsimple closed plane curves of length  $L$ , if the area is replaced by the sum of the areas into which the curve divides the plane, each weighted with the square of the winding number, i.e.,

$$L^2 - 4\pi \int_{\mathbb{R}^2} w^2 dP \geq 0.$$

Equality holds only for one or several coincident circles, each traversed in the same direction a number of times. The authors also generalized the result to arbitrary dimension and codimension.

In this paper we show that a similar result is valid for compact oriented manifolds immersed in either the hyperbolic or the elliptic space. In §2 we obtain a generalization of formulae (4.7) and (4.8) of [1] in any  $n$ -dimensional Riemann-space of constant curvature by using some known formulae of integral geometry [3]. In §3 we show that the results of Banchoff and Pohl [1, §3] are valid without change in the hyperbolic and elliptic  $n$ -dimensional spaces. In §4 we obtain a general expression for the main result of Banchoff and Pohl [1], which can be stated as:

Let  $\mathbb{W}^n$  denote the hyperbolic, elliptic, or euclidean  $n$ -dimensional space with  $H(r)$  being  $\text{sh } r$ ,  $\text{sen } r$ , or  $r$ , respectively. Let  $M$  be a compact oriented manifold of dimension  $m$  and  $f: M \rightarrow \mathbb{W}^n$  be an immersion of class  $C^2$ . For  $(x, y) \in M \times M$  let  $r(x, y)$  denote the chord length from  $f(x)$  to  $f(y)$  and let  $dV_1, dV_2$  denote the volume elements on  $M$  at  $x$  and  $y$ , respectively. Let  $H_{n-m-1, n}$  denote the Grassmann manifold of all (unoriented)  $(n-m-1)$ -planes in  $\mathbb{W}^n$  (parallel planes are not identified), and let  $|dH_{n-m-1, n}|$  denote its invariant measure. For  $h \in H_{n-m-1, n}$  let  $\pm\lambda$  denote the linking number of

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$h$  with  $f$ . Then:

$$(1.1) \quad \int_{M \times M} H(r)^{-m+1} dV_1 dV_2 - (1+m)\Sigma_m K_{m,n} \int_{H_{n-m-1,n}} \lambda^2 |dH_{n-m-1,n}| \\ + \frac{(-1)^{m(m+1)/2}}{m!} \int_{G \cap M \times M \neq \emptyset} [H(r) - r] dG^m \geq 0,$$

where  $\Sigma_m$  denotes the surface volume of the unit  $m$ -sphere and  $K_{m,n}$  is a constant, which depends only on  $m$  and  $n$ . Equality holds only for one or several coincident spheres with coincident orientations or ( $m = 1$ ) one or several circles each traversed in the same direction a number of times.

For closed nonsimple plane curves we obtain

$$L^2 - 4\Pi \int_{W^2} w^2 dP - \int_{G \cap (M \times M)} [H(r) - r] d\vec{G} \geq 0,$$

where  $d\vec{G}$  denotes the density for oriented geodesics.

This formula, when applied to simple closed curves, taking into account (18.10) and (18.11) of [3], gives the classical isoperimetric inequality:

$$L^2 - 4\Pi F + \varepsilon F^2 \geq 0,$$

where  $\varepsilon$  equals  $-1$ ,  $1$ , or  $0$  for the hyperbolic, elliptic, or euclidean plane, respectively.

2

Let  $W^n$  be a Riemann manifold of dimension  $n$  endowed with the fundamental quadratic form (for a local coordinate system) (see [3, p. 331])

$$ds^2 = \sum_{i,j=1}^n g_{ij} dx_i dx_j.$$

Since the 2-form

$$dG = \sum_{i=1}^n dp_i \wedge dx_i$$

is an invariant form (any geodesic is defined by a point  $(x_1, \dots, x_n)$  and a direction  $(p_1, \dots, p_n)$ ), the exterior powers

$$dG^m = (dG)^m = m! \sum_{i_1 < \dots < i_m} dp_{i_1} \wedge dx_{i_1} \wedge \dots \wedge dp_{i_m} \wedge dx_{i_m}$$

are also invariant forms ( $m = 1, \dots, n - 1$ ) and they define a measure for  $2m$ -dimensional sets of geodesics. Consider the expression:

$$(2.1) \quad dG^m = m!(-1)^{m(m-1)/2} \sum_{i_1 < \dots < i_m} dp_{i_1} \wedge \dots \wedge dp_{i_m} \wedge dx_{i_1} \wedge \dots \wedge dx_{i_m}$$

for a family of geodesics parametrized in a special way; in particular, for geodesics that join the points of one  $m$ -dimensional submanifold of  $W^n$  to another. We restrict ourselves to  $W^n$  the  $n$ -dimensional hyperbolic, elliptic, or euclidean space, noting that they include any complete, simply-connected Riemann-space with constant curvature.

Let us consider as Pohl [2, p.1325] the following slightly more general situation:

Let  $M_1, M_2$  be oriented  $m$ -dimensional differentiable manifolds of class  $C^1$ . Let  $f: M_1 \rightarrow \mathbb{W}^n$  be an immersion, and let  $F: M_1 \times M_2 \rightarrow \mathbb{W}^n$  be a continuously differentiable map such that for each fixed  $x \in M_1$  the map  $F_x(y) = F(x, y)$  describes an immersion of  $M_2$ . We assume as Pohl that  $F(x, y) \neq f(x) \forall x, y$ . Let  $G$  denote the geodesic that joins  $f(x)$  to  $F(x, y)$  oriented from  $f(x)$  to  $F(x, y)$ . We are now going to prove equation (2.5) of [2] under these premises. For this purpose, let us define orthonormal frames  $Xe_1, \dots, e_n$  and  $Xa_1, \dots, a_n$  following Pohl [2, p.1325] with only the following modifications:

- $e_1$  is the tangent unit vector to  $G$  in  $f(x)$ ;
- $S_1$  contains  $T_1$  and  $e_1$  (but not  $G$ ); and
- $U_1$  is the  $(m+1)$ -dimensional manifold containing  $M_1$  and  $G$ , so that  $S_1$  is the tangent space to  $U_1$  in  $f(x)$ .

Since the orientations  $e_1 e_2$  and  $a_1 a_2$  agree on  $Q_1$ , we may write

$$(2.2) \quad de_2 = \text{sen } \sigma_1 da_2.$$

We now determine  $G$  in  $U_1$  by its intersection with  $S' (e_1 = 0) \subset U_1$ , i.e.,  $X$ . Following Santaló [3, p. 334] we have:

$$(2.3) \quad ds^2 = \sum_{i=1}^n g_{ii} dx_i^2, \quad p_i = (g_{ii})^{1/2} \cos \alpha_i, \quad \alpha_i = \angle(G, e_i).$$

Using expression (2.1)

$$(2.4) \quad dG^m = m!(-1)^{(m-1)m/2} dp_2 \wedge \dots \wedge dp_{m+1} \wedge de_2 \wedge \dots \wedge de_{m+1}.$$

Also

$$(2.5) \quad du_m = \text{sen } \alpha_2 \dots \text{sen } \alpha_{m+1} d\alpha_2 \wedge \dots \wedge d\alpha_{m+1},$$

where  $du_m$  denotes the area element on the unit  $m$ -dimensional sphere of center  $X$  corresponding to the direction of  $e_1$  [3, p. 337] (notice that  $\alpha = \angle(e_1, e_1)$ ). Since the volume element  $dV_1$  on  $M$  at  $X$  has the value:

$$(2.6) \quad dV_1 = (g_{22} \dots g_{m+1, m+1})^{1/2} da_2 \wedge \dots \wedge da_{m+1},$$

replacing (2.2), (2.3), (2.5), and (2.6) in (2.4) we obtain

$$(2.7) \quad dG^m = m!(-1)^{m(m-1)/2} \text{sen } \sigma_1 du_m \wedge dV_1.$$

Next we compute  $du_m$  with fixed  $x$ . To do this we translate the frame  $Xe_1, \dots, e_n$  parallel along  $G$  to  $F_x(y)$  obtaining  $X'e'_1, \dots, e'_n$ . Let  $S_2$  be the  $(m+1)$ -dimensional linear space through  $X'$  containing  $e'_1$  and  $T_2$ , the tangent space to  $F_x$  at  $y$ . We orient  $S_2$  by  $e'_1 T_2$  unless  $e'_1$  lies in  $T_2$ , in which case we take either orientation on  $S_2$ . Let  $\sigma_2 = \angle(T_2, e'_1)$ ,  $X'b_1, \dots, b_n$  be an orthonormal frame such that  $b_1 = e'_1$ ;  $b_2, \dots, b_{m+1}$  are in  $S_2$  and so chosen that  $b_1 \dots b_{m+1}$  agrees with the given orientation of  $S_2$ .

We use polar coordinates in  $S_2$ , with respect to the frame  $Xb'_1, \dots, b'_n$  (where  $b'_i$  is the parallel translated vector of  $b_i$  from  $X'$  to  $X$  along  $G$ ). Therefore, we have

$$ds^2 = dr^2 + H(r) d\lambda_i^2$$

and for fixed  $r$ ,

$$db_i = H(r) d\lambda_i, \quad d\sigma_p = H(r)^m du'_m,$$

where  $d\sigma_p$  denotes the  $m$ -dimensional volume element for fixed  $r$ , and  $du'_m$  is the area element on the unit  $m$ -dimensional euclidean sphere of center  $X$  corresponding to the direction of  $b'_1 = e_1$  (defined for the frame  $b'_1, \dots, b'_n$ ).

On the other hand if  $dV_2$  denotes the volume element on  $M_2$  at  $X'$ ,

$$d\sigma_p = dV_2 \text{sen } \sigma_2,$$

and if we define  $\cos \tau$  as Pohl [2, p. 1326] so that  $\tau$  is interpreted as the angle between the oriented  $(m+1)$ -planes  $S_1$  and  $S_2$  (or its parallel translation along  $G$ ) we have

$$\cos \tau = du_m/du'_m,$$

so that

$$du_m = H(r)^{-m} \cos \tau \text{sen } \sigma_2 dV_2.$$

Replacing in (2.7) we obtain

$$(2.8) \quad dG^m = m!(-1)^{m(m-1)/2} \text{sen } \sigma_1 \text{sen } \sigma_2 \cos \tau H(r)^{-m} dV_2 \wedge dV_1,$$

which is one of the results of Pohl [2, p. 1326].

In order to obtain the general result we have to compute

$$(2.9) \quad \int_{M \times M} r dG^m = \int_{M \times M} (r dp'_2) \wedge da_2 \wedge dG^{m-1},$$

where  $Xa_1, \dots, a_n$  is the frame defined above so that  $p'_2 = (g_{22})^{1/2} \cos \sigma_1$ . To do this we define as in [1] a certain 'secant space' and apply Stoke's theorem:

$$\int_{M \times M} r dp'_2 = r \cdot p'_2|_{r=0} - \int_{M \times M} p'_2 dr = - \int (g_{22})^{1/2} \cos \sigma_1 dr.$$

Since  $\cos \sigma_2 = dr/db_2$  (for  $Xb_1, \dots, b_n$  the frame defined above) we get:

$$(2.10) \quad \int_{M \times M} r dG^m = - \int_{M \times M} (g_{22})^{1/2} \cos \sigma_1 \cos \sigma_2 db_2 \wedge da_2 \wedge dG^{m-1}.$$

Defining frames as above and interpreting the angle  $\nu$  as in [1] we have

$$du_{m-1} = \cos \nu du'_{m-1}$$

and therefore,

$$dG^{m-1} = (m-1)!(-1)^{m(m-1)/2} \cos \nu (g_{33} \cdots g_{m+1 m+1})^{1/2} du'_{m-1} \wedge da_3 \wedge \cdots \wedge da_{m+1}.$$

Since

$$dV_1 = (g_{22} \cdots g_{m+1 m+1})^{1/2} da_2 \wedge \cdots \wedge da_{m+1}$$

and

$$dV_2 = H(r)^{-m+1} db_2 \wedge du'_{m-1},$$

replacing in (2.10)

$$(2.11) \quad \int_{M \times M} r dG^m = -(m-1)!(-1)^{m(m-1)/2} \times \int_{M \times M} H(r)^{-m+1} \cos \sigma_1 \cos \sigma_2 \cos \nu dV_2 \wedge dV_1,$$

which is the desired expression.

3

We now prove that the formulae of Banchoff and Pohl [1, §3] are valid without change in the elliptic and hyperbolic  $n$ -dimensional spaces. For this purpose we work in an analogous way to Banchoff and Pohl [1] but we make use of some known results of integral geometry [3].

We first consider the second integral in our result (see (1.1)) and write

$$(3.1) \quad \mathcal{A}(M) = K_{m,n} \int_{H_{n-m-1,n}} \lambda^2 |dH_{n-m-1,n}|,$$

where

$$K_{m,n} = \frac{\Sigma_{n-m-1} \cdots \Sigma_1}{\Sigma_n \cdots \Sigma_{m+2}} \quad (n \geq m+2), \quad K_{m,m+1} = 1,$$

is defined in Banchoff and Pohl [1, p. 180]. We now prove that  $\mathcal{A}(M)$  is stable under rising of the codimension, i.e., if  $f(M) \subseteq \mathbb{W}^n \subseteq \mathbb{W}^N$  then

$$(3.2) \quad K_{m,N} \int_{H_{N-m-1,N}} \lambda^2 |dH_{N-m-1,N}| = K_{m,n} \int_{H_{n-m-1,n}} \lambda^2 |dH_{n-m-1,n}|.$$

Comparing [3, (14.69)] firstly for  $q \rightarrow n$ ,  $r \rightarrow n-m$ ,  $n \rightarrow n+1$ ,  $M_q \rightarrow U_n$ , and secondly for  $q \rightarrow n$ ,  $r \rightarrow n-m-1$ ,  $n \rightarrow n$ ,  $M_q \rightarrow U_n$ , we obtain

$$(3.3) \quad \int_{H_{n-m,m+1} \cap U_n \neq \emptyset} dH_{n-m,m+1} = \frac{\Sigma_{n+1}}{\Sigma_{n-m}} \int_{H_{n-m-1,n} \cap U_n \neq \emptyset} dH_{n-m-1,n}.$$

On the other hand (see [2] or [3]),

$$dH_{n-m,n+1} = \bigwedge_{j,\beta} \omega_{j\beta} \wedge \omega_\beta$$

$$(j = 1, \dots, n-m; \beta = n-m+1, \dots, n+1),$$

$$dH_{n-m-1,n} = \bigwedge_{i,\alpha} \pi_{i\alpha} \wedge \pi_\alpha$$

$$(i = 1, \dots, n-m-1; \alpha = n-m, \dots, n)$$

with  $\omega_{j\beta} = de_\beta \cdot e_j$ ,  $\omega_\beta = dp \cdot e_\beta$ ,  $\omega_{j\beta} + \omega_{\beta j} = 0$ , for an orthonormal frame  $pe_1, \dots, e_{n+1}$  in  $p$ , and  $\pi_{i\alpha} = da_\alpha \cdot a_i$ ,  $\pi_\alpha = dp \cdot a_\alpha$ ,  $\pi_{i\alpha} + \pi_{\alpha i} = 0$  for  $pa_1, \dots, a_{n+1}$  the frame defined as follows:

$$a_i = e_i, \quad i = 1, \dots, n-m-1; \quad a_{n-m} = e_{n+1};$$

$$a_{n-m+1} = e_{n-m+1}, \quad i = 1, \dots, m; \quad a_{n+1} = e_{n-m};$$

therefore,

$$|dH_{n-m,n+1}| = |dH_{n-m-1,n}| \wedge \bigwedge_{\alpha} \pi_{n+1,\alpha}, \quad \alpha = n-m, \dots, n.$$

Since the linking number of an  $(n-m)$ -plane  $h \subseteq \mathbb{W}^{n+1}$  with  $M$  does not change regarding  $h$  as lying in  $\mathbb{W}^n$ , using Fubini's theorem and integrating  $\lambda^2$  we get

$$\int_{H_{n-m,n+1}} \lambda^2 |dH_{n-m,n+1}| = c_{n,m} \int_{H_{n-m-1,n}} \lambda^2 |dH_{n-m-1,n}|.$$

Applying (3.3) we obtain  $c_{n,m} = \Sigma_{n+1}/\Sigma_{n-m}$ . This proves the assertion for  $N = n+1$ , and so (3.2) is valid by induction.

Also since both the Cauchy formula and the remaining arguments used by Banchoff and Pohl [1] to demonstrate:

$$(3.4) \quad \mathcal{A}(M^0) = K_{0,n} \int_{H_{n-1,n}} \lambda^2 |dH_{n-1,n}| = -1/2 \sum_{i,j} r(x_i, x_j) i_i j_j$$

are also valid in the hyperbolic and elliptic spaces, equation (3.4) stays valid without change.

In order to prove that for  $q > n - m - 1$

$$(3.5) \quad \mathcal{A}(M^m) = 1_{m,n,q} \int_{H_{q,n}} \mathcal{A}(M^m \cap H_q) |dH_{q,n}|$$

with

$$1_{m,n,q} = \frac{K_{m,n} \Sigma_{q-n+m} \cdots \Sigma_0 \Sigma_{n-q-1} \cdots \Sigma_0}{K_{q-n+m,q} \Sigma_m \cdots \Sigma_0},$$

we use the fact that equations (12.36) and (12.52) of [3] are also valid in hyperbolic and elliptic spaces, so that

$$(3.6) \quad \int_{\text{total}} dH_{r,n^{[q]}} = \Sigma_{n-q-1} \cdots \Sigma_{n-r} / \Sigma_{r-q-1} \cdots \Sigma_0,$$

$$(3.7) \quad dH_{i+1,r} \wedge dH_{r,n}^* = dH_{r,n^{[i+1]}} \wedge dH_{i+1,n},$$

where  $dH_{r,n^{[i+1]}}$  is the density of  $H_{r,n}$  about  $H_{i+1,r}$  and  $H_{r,n}^*$  denote the oriented Grassmannians. Using that, for

$$(3.8) \quad H_{n-m-1,q} \subseteq H_{q,n},$$

it is

$$\lambda(M^m, h_{n-m-1}) = \lambda(M^m \cap h_q, h_{n-m-1}),$$

we now integrate  $\lambda^2$  over the set of pairs of linear spaces  $(H_{q,n}; H_{n-m-1,q})$  so that (3.8) is valid, and apply (3.7) to obtain

$$\int \lambda^2 dH_{n-m-1,q} \wedge dH_{q,n}^* = \int \lambda^2 dH_{q,n^{[n-m-1]}} \wedge dH_{n-m-1,n}.$$

Evaluating the right-hand side using equation (3.6) for  $r \rightarrow q, q \rightarrow n - m - 1$ ; and taking into account the definition (3.1) in both sides, we compare and get equation (3.5) after few calculations.

Comparing as above, the densities  $dH_{n-m,n}$  and  $dH_{n-m-1,n-1}$ , using orthonormal frames  $pe_1, \dots, e_n$  and  $pa_1, \dots, a_n$  with  $a_1 = e_1$ ;  $a_i = e_{n-m-1+i}$   $i = 2, \dots, m + 1$ ;  $a_{m+j} = e_j$   $j = 2, \dots, n - m$ ; and using:

$$|dG^m|/m! = |\pi_{12} \wedge \cdots \wedge \pi_{1m+1} \wedge \pi_2 \wedge \cdots \wedge \pi_{m+1}|$$

(which is easy to prove from expression (2.1)), we obtain

$$|dH_{n-m,n}| = |dH_{n-m-1,n-1^{[0]}}| \wedge (1/m!) |dG^m|.$$

We now integrate the function  $f(r) = r$  using Fubini's theorem to get

$$(3.9) \quad \int r |dH_{n-m,n}| = c_{m,n}/m! \int r |dG^m|.$$

Using a general result of differential geometry (see [2, p. 1328]) and a particular case of (3.5) combined with (3.4), and evaluating the constant  $c_{m,n}$  as Banchoff and Pohl [1, p. 186], equation (3.9) yields

$$(3.10) \quad \mathcal{A}(M^m) = \frac{(-1)^{1+m(m+1)/2}}{m! \Sigma_m} \int_{M \times M} r(dG^m),$$

which is Theorem 4 of [1].

#### 4

In order to obtain the general result we are now able to use the same arguments as Banchoff and Pohl [1] with  $H(r)$  instead of  $r$ . Using the results of §§2 and 3, we obtain

$$\int H(r)^{-m+1} dV_1 dV_2 - (1+m) \Sigma_m \mathcal{A}(M) + \frac{(-1)^{m(m+1)/2}}{m!} \int [H(r) - r] dG^m \geq 0$$

with equality holding for one or several coincident spheres with coincident orientations or for ( $m = 1$ ) one or several coincident circles all transversed in the same direction each a number of times. Applied to the euclidean space ( $H(r) = r$ ), it yields the main theorem of Banchoff and Pohl [1].

#### REFERENCES

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