

ON A SOBOLEV INEQUALITY WITH REMAINDER TERMS

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ABSTRACT. In this note we consider the Sobolev inequality

$$\|\Delta \phi\|_2 \geq S_2 \|\phi\|_{\frac{2N}{N-4}}, \quad N > 4, \quad \phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where S_2 is the best Sobolev constant and $\mathcal{D}_0^{2,2}(\mathbb{R}^N)$ is the space obtained by taking the completion of $C_0^\infty(\mathbb{R}^N)$ with the norm $\|\Delta \phi\|_2$. We prove here a refined version of this inequality,

$$\|\Delta \phi\|_2^2 - S_2^2 \|\phi\|_{\frac{2N}{N-4}}^2 \geq \alpha d^2(\phi, M_2), \quad N > 4, \quad \phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where α is a positive constant, the distance is taken in the Sobolev space $\mathcal{D}_0^{2,2}(\mathbb{R}^N)$, and M_2 is the set of solutions which attain the Sobolev equality. This generalizes a result of Bianchi and Egnell (*A note on the Sobolev inequality*, J. Funct. Anal. 100 (1991), 18-24), which was posed by Brezis and Lieb (*Sobolev inequalities with remainder terms*, J. Funct. Anal. 62 (1985), 73-86). regarding the classical Sobolev inequality

$$\|\nabla \phi\|_2 \geq S_1 \|\phi\|_{\frac{2N}{N-2}}, \quad \phi \in \mathcal{D}_0^{1,2}(\mathbb{R}^N).$$

A key ingredient in our proof is the analysis of eigenvalues of the fourth order equation

$$\Delta^2 v - \mu S_2^{p+1} U^{\frac{8}{N-4}} v = 0, \quad v \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where $p = \frac{N+4}{N-4}$ and U is the unique radial function in M_2 with $\|\Delta U\|_2 = 1$. We will show that the eigenvalues μ of the above equation are discrete:

$$\mu_1 = 1, \mu_2 = \mu_3 = \cdots = \mu_{N+2} = p < \mu_{N+3} \leq \cdots$$

and the corresponding eigenfunction spaces are

$$V_1 = \{U\}, V_p = \left\{ \frac{\partial U}{\partial y_j}, j = 1, \dots, N, x \cdot \nabla U + \frac{N-4}{2} U \right\}, \dots$$

1. INTRODUCTION

In this note we consider the Sobolev inequality

$$(1.1) \quad \|\Delta \phi\|_2 - S_2 \|\phi\|_{\frac{2N}{N-4}} \geq 0, \quad N > 4, \quad \phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where $\mathcal{D}_0^{2,2}(\mathbb{R}^N)$ is the completion of the space of smooth functions with compact support under the norm $\|\Delta \phi\|_2$, and S_2 is the best Sobolev constant. We assume $N > 4$ throughout this paper, and we denote by $\|\cdot\|_p$ the L^p norm in \mathbb{R}^N .

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In [2], Brezis and Lieb asked if the following refined classical Sobolev inequality holds:

$$(1.2) \quad \|\nabla \phi\|_2^2 - S_1^2 \|\phi\|_{\frac{2N}{N-2}}^2 \geq \alpha d^2(\phi, M_1), \quad N > 2, \quad \phi \in \mathcal{D}_0^{1,2}(\mathbb{R}^N),$$

where $\mathcal{D}_0^{1,2}(\mathbb{R}^N)$ is the completion of $C_0^\infty(\mathbb{R}^N)$ under the norm of $\|\nabla \phi\|_2$ when $N \geq 3$, $\alpha > 0$ and M_1 is the $(N+2)$ -dimensional manifold which consists of all solutions attaining the classical Sobolev inequality. In [1], Bianchi and Egnell gave an affirmative answer to this question.

Some related results were proved by Brezis and Nirenberg in [3], among many other things. They showed in [3] that if $q < \frac{N}{N-2}$, then there is a positive number A such that

$$(1.3) \quad \|\nabla \phi\|_2^2 - S_1^2 \|\phi\|_{\frac{2N}{N-2}}^2 \geq A \|\phi\|_q^2, \quad N > 2, \quad \phi \in \mathcal{D}_0^{1,2}(\Omega),$$

where A only depends on N, q and the bounded domain Ω . Moreover, the result (1.3) is proved to be sharp in the sense that it fails for $q = \frac{N}{N-2}$.

The following refinement of (1.3) was also proved in [2]:

$$(1.4) \quad \|\nabla \phi\|_2^2 - S_1^2 \|\phi\|_{\frac{2N}{N-2}}^2 \geq A \|\phi\|_{\frac{N}{N-2}, w}^2, \quad N > 2, \quad \phi \in \mathcal{D}_0^{1,2}(\Omega),$$

where $\|\cdot\|_{q,w}$ is the weak L^q norm. An alternate proof of (1.4) was given in [1].

The result of Brezis and Nirenberg [3] was generalized to $1 < p < N$ and $q < \frac{N(p-1)}{N-p}$ by Egnell, Pacella and Tricarico in [5]. They proved in [5] that there is a positive number A depending only N, q, p and the bounded domain Ω such that

$$(1.5) \quad \|\nabla \phi\|_p^p - S_p^p \|\phi\|_{p^*}^p \geq A \|\phi\|_q^p, \quad N > 2, \quad \phi \in \mathcal{D}_0^{1,2}(\Omega),$$

where $p^* = \frac{Np}{N-p}$ and S_p is the best Sobolev constant in the Sobolev embedding theorem with the critical exponent p^* . Furthermore, (1.5) fails for $q = \frac{N(p-1)}{N-p}$.

The purpose of this paper is to generalize the result of [1] and prove the following refined inequality of (1.1):

$$(1.6) \quad \|\Delta \phi\|_2^2 - S_2^2 \|\phi\|_{\frac{2N}{N-4}}^2 \geq \alpha d^2(\phi, M_2),$$

where $M_2 = \{\phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N) : \|\Delta \phi\|_2 = S_2 \|\phi\|_{\frac{2N}{N-4}}\}$.

By Theorem 2.1 of [4], M_2 is an $(N+2)$ -dimensional manifold and consists of functions of the form

$$(1.7) \quad \phi(x) = cU_{\lambda,y}(x) = c\lambda U(\lambda^{\frac{2}{N-4}}(x-y)),$$

where $c \in \mathbb{R}, \lambda \in \mathbb{R}_+, U(x) = k_0(1+|x|^2)^{-\frac{N-4}{2}}$ and k_0 is chosen so that $\|\Delta U\|_2 = 1$. The best constant is

$$S_2 = \pi^2(N+2)N(N-2)(N-4) (\Gamma(N/2)/\Gamma(N))^{4/N}.$$

Hence

$$M_2 = \{cU_{\lambda,y} : c \in \mathbb{R}, \lambda \in \mathbb{R}_+, y \in \mathbb{R}^N\}.$$

(In fact, C.-S. Lin recently showed (see Theorem 1.4 in [7]) that any solution of the equation

$$(1.8) \quad (-\Delta)^2 u = u^{\frac{N+4}{N-4}}, \quad u > 0 \text{ in } \mathbb{R}^N \quad (N \geq 5)$$

is symmetric around some point and is of the form (1.7).)

We define the distance between this manifold and a function $\phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N)$ as

$$d(\phi, M_2) = \inf_{u \in M_2} \|\Delta(\phi - u)\|_2 = \inf_{c, \lambda, y} \|\Delta(\phi - cU_{\lambda, y})\|_2.$$

Note that $d(c\lambda\phi(\lambda^{\frac{2}{N-4}}(\cdot - y)), M_2) = cd(\phi, M_2)$.

We shall prove

Theorem 1.9. *There is a positive constant α depending only on the dimension N such that*

$$\|\Delta\phi\|_2^2 - S_2^2\|\phi\|_{\frac{2N}{N-4}}^2 \geq \alpha d^2(\phi, M_2).$$

Furthermore, this result is sharp in the sense that it is false if the left hand side is replaced with $d(\phi, M_2)^\beta \|\Delta\phi\|_2^{2-\beta}$, where $\beta < 2$.

We refer the reader to [1], [2], [5], [6], [8] and [9] for more information on Sobolev inequalities of such type.

A key ingredient in our proof is the analysis of eigenvalues of the following fourth order equation:

$$(1.10) \quad \Delta^2 v - \mu S_2^{p+1} U^{\frac{8}{N-4}} v = 0, v \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where $p = \frac{N+4}{N-4}$. We will show that the eigenvalues μ of (1.10) are discrete and

$$(1.11) \quad \mu_1 = 1, \mu_2 = \mu_3 = \cdots = \mu_{N+2} = p < \mu_{N+3} \leq \cdots.$$

The corresponding eigenfunction spaces are

$$(1.12) \quad V_1 = \{U\}, V_p = \left\{ \frac{\partial U}{\partial y_j}, j = 1, \dots, N, x \cdot \nabla U + \frac{N-4}{2} U \right\}.$$

This result is interesting and may have applications in analyzing blow up problems involving the biharmonic operator. We prove this in Section 2. In Section 3, we prove Theorem 1.9.

2. AN EIGENVALUE PROBLEM

In this section, we solve the eigenvalue problem

$$(2.1) \quad \Delta^2 v - \mu S_2^{p+1} U^{\frac{8}{N-4}} v = 0, v \in \mathcal{D}_0^{2,2}(\mathbb{R}^N),$$

where $U = k_0(1 + |x|^2)^{-\frac{N-4}{2}}$. Note that U satisfies

$$\Delta^2 U - S_2^{p+1} U^p = 0, U(0) = \max_{z \in \mathbb{R}^N} U(z)$$

and $p = \frac{N+4}{N-4}$.

Note also that (2.1) is a fourth order eigenvalue problem. We will show the following

Theorem 2.2. *The eigenvalues μ of (2.1) are discrete and*

$$(2.3) \quad \mu_1 = 1, \mu_2 = \mu_3 = \cdots = \mu_{N+2} = p < \mu_{N+3} \leq \cdots.$$

The corresponding eigenfunction spaces are

$$(2.4) \quad V_1 = \{U\}, V_p = \left\{ \frac{\partial U}{\partial y_j}, j = 1, \dots, N, x \cdot \nabla U + \frac{N-4}{2} U \right\}.$$

Proof. We decompose the fourth order equation (2.1) into a system of two second order equations.

We first note that the eigenvalues of $\Delta_{S^{N-1}}$ are given by

$$(2.5) \quad \lambda_1 = 0, \lambda_2 = \cdots = \lambda_{N+1} = N - 1, \lambda_{N+1} < \lambda_{N+2}.$$

Let $\Psi_i(\theta)$ be the corresponding eigenfunctions, i.e., $\Delta_{S^{N-1}}\Phi_i = -\lambda_i\Phi_i$ for each i . Set

$$\phi_i(r) = \int_{S^{N-1}} v(x)\Psi_i(\theta)d\theta$$

and

$$w_i(r) = - \int_{S^{N-1}} \Delta v(x)\Psi_i(\theta)d\theta.$$

Then we obtain a system of two equations

$$(2.6) \quad \Delta \phi_i - \frac{\lambda_i}{r^2}\phi_i + w_i = 0,$$

$$(2.7) \quad \Delta w_i - \frac{\lambda_i}{r^2}w_i + \mu S_2^{p+1}U^{p-1}\phi_i = 0.$$

We now show that

$$(2.8) \quad \phi_i = 0, \text{ if } i > N + 2, \mu \leq p.$$

We will prove it by contradiction. Suppose now that $\mu \leq p$ and $\phi_i \neq 0, i > N + 2$. Without loss of generality, we assume that

$$(2.9) \quad \phi_i < 0 \text{ in } (0, r_1), \phi_i(r_1) = 0.$$

Since $i > N + 2, \phi_i(0) = w_i(0) = \phi_i'(0) = w_i'(0) = 0$.

We now compare (ϕ_i, w_i) with a new pair of functions.

Observe that U satisfies $\Delta^2 U - S_2^{p+1}U^p = 0$, and let $U_1 = -\Delta U$. Then both U and U_1 are radial and monotone decreasing functions. We observe that for any radial function $f(r)$ we have

$$(\Delta f)' = (f''(r) + \frac{N-1}{r}f'(r))' = f'''(r) + \frac{N-1}{r}f''(r) - \frac{N-1}{r^2}f'(r).$$

Then (U', U_1') satisfies

$$(2.10) \quad \Delta U' - \frac{N-1}{r^2}U' + U_1' = 0.$$

We note

$$\Delta^2 U = -\Delta U_1 = -\left(U_1'' + \frac{N-1}{r}U_1'\right);$$

thus

$$U_1'' + \frac{N-1}{r}U_1' + S_2^{p+1}U^p = 0.$$

Differentiating both sides of the above identity with respect to r , we get

$$U_1''' + \frac{N-1}{r}U_1'' - \frac{N-1}{r^2}U_1' + pS_2^{p+1}U^{p-1}U' = 0.$$

This is equivalent to

$$(2.11) \quad \Delta U_1' - \frac{N-1}{r^2}U_1' + pS_2^{p+1}U^{p-1}U' = 0.$$

Multiplying (2.7) by U' and integrating over B_r we get

$$\begin{aligned}
 (2.12) \quad 0 &= \int_{B_r} \mu S_2^{p+1} U^{p-1} \phi_i U' + \int_{B_r} \left(\Delta w_i - \frac{\lambda_i}{r^2} \phi_i \right) U' \\
 &= (\mu - p) \int_{B_r} S_2^{p+1} U^{p-1} \phi_i U' + p \int_{B_r} S_2^{p+1} U^{p-1} \phi_i U' + \int_{B_r} \left(\Delta w_i - \frac{\lambda_i}{r^2} \phi_i \right) U' \\
 &= (\mu - p) \int_{B_r} S_2^{p+1} U^{p-1} \phi_i U' - \int_{B_r} \left[\Delta U'_1 - \frac{N-1}{r^2} U'_1 \right] \phi_i \\
 &\quad + \int_{B_r} \left[\Delta w_i - \frac{\lambda_i}{r^2} w_i \right] U'.
 \end{aligned}$$

The second term in (2.12) can be seen as follows:

$$\begin{aligned}
 & - \int_{B_r} \left[\Delta U'_1 - \frac{N-1}{r^2} U'_1 \right] \phi_i \\
 &= - \int_{B_r} \operatorname{div}(\nabla U'_1) \phi_i + \int_{B_r} \frac{N-1}{r^2} U'_1 \phi_i \\
 &= - \int_{\partial B_r} (\nabla U'_1 \cdot \nu) \phi_i + \int_{B_r} (\nabla U'_1) \cdot \phi_i + \int_{B_r} \frac{N-1}{r^2} U'_1 \phi_i \\
 &= - \int_{\partial B_r} \frac{\partial U'_1}{\partial \nu} \phi_i + \int_{\partial B_r} U'_1 \frac{\partial \phi_i}{\partial \nu} - \int_{B_r} U'_1 \Delta \phi_i + \int_{B_r} \frac{N-1}{r^2} U'_1 \phi_i.
 \end{aligned}$$

By using equation (2.6), the second term of (2.12) becomes

$$\int_{B_r} \frac{N-1-\lambda_i}{r^2} U'_1 \phi_i + \int_{\partial B_r} \left(U'_1 \frac{\partial \phi_i}{\partial \nu} - \phi_i \frac{\partial U'_1}{\partial \nu} \right) + \int_{B_r} U'_1 w_i.$$

The third term of (2.12) is

$$\int_{B_r} \left[\Delta w_i - \frac{\lambda_i}{r^2} w_i \right] U' = \int_{\partial B_r} \frac{\partial w_i}{\partial \nu} U' - \int_{\partial B_r} w_i \frac{\partial U'}{\partial \nu} + \int_{B_r} w_i \Delta U' - \int_{B_r} \frac{\lambda_i}{r^2} w_i U'.$$

Note that by (2.10) we have $\Delta U' = \frac{N-1}{r^2} U' - U'_1$; thus the third term of (2.12) is equal to

$$\int_{B_r} \frac{N-1-\lambda_i}{r^2} U' w_i + \int_{\partial B_r} \left(U' \frac{\partial w_i}{\partial \nu} - w_i \frac{\partial U'}{\partial \nu} \right) - \int_{B_r} w_i U'_1.$$

Therefore

$$\begin{aligned}
 0 &= (\mu - p) \int_{B_r} S_2^{p+1} U^{p-1} \phi_i U' + \int_{B_r} \frac{N-1-\lambda_i}{r^2} U'_1 \phi_i + \int_{\partial B_r} \left(U'_1 \frac{\partial \phi_i}{\partial \nu} - \phi_i \frac{\partial U'_1}{\partial \nu} \right) \\
 &\quad + \int_{B_r} \frac{N-1-\lambda_i}{r^2} U' w_i + \int_{\partial B_r} \left(U' \frac{\partial w_i}{\partial \nu} - w_i \frac{\partial U'}{\partial \nu} \right) \\
 &= (\mu - p) \int_{B_r} S_2^{p+1} U^{p-1} \phi_i U' + \int_{B_r} \frac{N-1-\lambda_i}{r^2} (U'_1 \phi_i + U' w_i) \\
 &\quad + \int_{\partial B_r} \left(U'_1 \frac{\partial \phi_i}{\partial \nu} - \phi_i \frac{\partial U'_1}{\partial \nu} \right) + \int_{\partial B_r} \left(U' \frac{\partial w_i}{\partial \nu} - w_i \frac{\partial U'}{\partial \nu} \right) \\
 &= I_1(r) + I_2(r) + I_3(r) + I_4(r),
 \end{aligned}$$

where $I_i(r)$, $i = 1, 2, 3, 4$, are defined at the last equality.

We now choose appropriate r and estimate each term. Recall that $\phi_i < 0$ in $(0, r_1)$, $\phi_i(r_1) = 0$ (we can take $r_1 = \infty$ if $\phi_i < 0$ in $(0, \infty)$).

We first deduce from equations (2.6) and (2.7) some useful facts:

1) $U'(r) < 0$, $U_1'(r) < 0$ if $r \neq 0$.
 2) $\Delta w_i - \frac{\lambda_i}{r^2} w_i = -\mu S_2^{p+1} U^{p-1} \phi_i$. Thus w_i cannot have a positive local maximum in $(0, r_1)$.

3) There exists some $r^* \in (0, r_1)$ such that $w_i(r^*) < 0$. This is because ϕ_i attains at least one local minimum in $(0, r_1)$ and by using equation (2.7).

4) By 2) and 3), $w_i(r) < 0$ for $r \in (0, r_2)$ for some $r_2 > 0$ and $w_i(r_2) = 0$.

If $r_1 = r_2$, then we take $r = r_1 = r_2$. In this case, $I_1(r) \leq 0$ since $\mu \leq p$, $\phi_i, U' < 0$; $I_2(r) \leq 0$ since $\lambda_i \geq N - 1$; $I_3(r) \leq 0$ since $\frac{\partial \phi_i}{\partial \nu} \geq 0$; $I_4(r) \leq 0$ since $\frac{\partial w_i}{\partial \nu} \geq 0$.

Therefore, we have $I_1(r) = I_2(r) = I_3(r) = I_4(r) = 0$. This is a contradiction since $I_3(r) < 0$ and $I_4(r) < 0$.

The rest of the cases can be discussed in the following:

Case 1: $r_2 < r_1$.

In this case, we take $r = r_2$. Then $I_1(r_2) \leq 0$, $I_2(r_2) \leq 0$, $I_4(r_2) \leq 0$. We only need to know if $I_3(r_2) \leq 0$. To this end, we consider the function

$$w = r^{N-1} \phi_i' U_1' - r^{N-1} U_1'' \phi_i.$$

By an easy calculation we get

$$w' = r^{N-1} \left[\phi_i'' + \frac{N-1}{r} \phi_i' \right] U_1' - r^{N-1} \left[U_1''' + \frac{N-1}{r} U_1'' \right] \phi_i.$$

By using equations (2.6) and (2.11) we can easily see that in (r_2, r_1)

$$w' = (\lambda_i - (N-1)) r^{N-3} \phi_i U_1' - r^{N-1} w_i U_1' + p S_2^{p+1} U^{p-1} U' r^{N-1} \phi_i > 0,$$

since $w_i > 0$ in (r_2, r_1) by fact 2), $U_1' < 0$, $\phi_i < 0$ and $\lambda_i > N - 1$.

Thus

$$w(r_2) = r_2^{N-1} \phi_i'(r_2) U_1'(r_2) - r_2^{N-1} U_1''(r_2) \phi_i(r_2) < w(r_1) < 0$$

and

$$I_3(r_2) = \int_{\partial B_{r_2}} r^{1-N} w(r) < 0.$$

This is again a contradiction.

Case 2: $r_1 < r_2$.

In this case,

$$\Delta \phi_i - \frac{\lambda_i}{r^2} \phi_i = -w_i > 0, \text{ for } r \in (0, r_2).$$

Hence $\phi_i > 0$ for $r \in (r_1, r_2)$.

Now we take $r = r_1$. Then $I_1(r_1) \leq 0$, $I_2(r_1) \leq 0$, $I_3(r_1) \leq 0$. Thus we only need to show $I_4(r_1) \leq 0$. Consider

$$h = r^{N-1} w_i' U' - r^{N-1} U'' w_i.$$

We note that

$$\begin{aligned} h' &= r^{N-1} \left[w_i'' + \frac{N-1}{r} w_i' \right] U' - r^{N-1} \left[U''' + \frac{N-1}{r} U'' \right] w_i \\ &= r^{N-1} (\Delta w_i) U' - r^{N-1} (\Delta U') w_i. \end{aligned}$$

By using equations (2.7) and (2.10) we thus have

$$h' = (\lambda_i - (N - 1)) w_i U' r^{N-3} - p S_2^{p+1} U^{p-1} \phi_i U' r^{N-1} + r^{N-3} (N - 1) U_1' w_i > 0$$

for $r \in (r_1, r_2)$ in $\lambda_i > N - 1$. Hence $h(r_1) < h(r_2) = r_2^{N-1} w_i'(r_2) U'(r_2) < 0$. Therefore we have

$$I_4(r_1) = \int_{B_{r_1}} r^{1-N} h(r) < 0.$$

In conclusion, we have that if $\mu \leq p$, then $\phi_i \equiv 0$ for $i > N + 2$. Hence, when $\mu \leq p$, the eigenfunction space has at most $(N + 2)$ dimensions. On the other hand, when $\mu = 1$, $v = U$ is a solution; when $\mu = p$, $\frac{\partial U}{\partial y_j}$ ($1 \leq j \leq N$) and $\frac{\partial U_{\lambda,0}}{\partial \lambda}|_{\lambda=1} = x \cdot \nabla U + \frac{N-4}{2} U$ are eigenfunctions.

Note that $U, \frac{\partial U}{\partial y_j}$ ($1 \leq j \leq N$), and $\frac{\partial U_{\lambda,0}}{\partial \lambda} = x \cdot \nabla U + \frac{N-4}{2} U$ are linearly independent. Thus (2.2) has the following solutions:

$$\begin{aligned} \mu = 1, v &= \alpha U, \\ \mu = p, v &\in \text{span} \left\{ \frac{\partial U}{\partial y_j}, \frac{\partial U_{\lambda}}{\partial \lambda} \Big|_{\lambda=1} \right\}. \end{aligned}$$

□

Consider the operator

$$L_{\lambda,y} = \frac{1}{S_2^{p+1}} U_{\lambda,y}^{1-p} \Delta^2, \text{ on } L^2(U_{\lambda,y}^{p-1}).$$

Since the imbedding

$$\mathcal{D}_0^{2,2}(\mathbb{R}^N) \rightarrow L^2(U_{\lambda,y}^{p-1})$$

is compact, the spectrum is discrete.

Consider

$$(2.13) \quad L_{\lambda,y} v - \mu v = 0.$$

Then we have

$$(2.14) \quad \Delta^2 v - \mu S_2^{p+1} U_{\lambda,y}^{p-1} v = 0.$$

By a simple scaling argument, we have that μ does not depend on λ, y . Moreover, by Theorem 2.1 we have

Lemma 2.15. *Let $\mu_i, i = 1, \dots$, denote the eigenvalues of (2.13) in increasing order. Then $\mu_1 = 1$ is simple with eigenfunction $U_{\lambda,y}$ and $\mu_2 = \dots = \mu_{N+2} = p$ with the corresponding $(N + 1)$ -dimensional eigenfunction space spanned by $(\partial_\lambda U_{\lambda,y}, \nabla_y U_{\lambda,y})$. Furthermore, eigenvalues do not depend on λ, y , and $\mu_{N+3} > \mu_2$.*

3. PROOF OF THE THEOREM

The main ingredient in the proof of the theorem is contained in the lemma below, where the behavior near M_2 is studied.

Lemma 3.1. *There is a constant α , depending only on the dimension, such that*

$$\|\Delta \phi\|_2^2 - S_2^2 \|\phi\|_{\frac{2N}{N-4}}^2 \geq \alpha d^2(\phi, M_2) + o(d(\phi, M_2)^2),$$

for all $\phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N)$ with $d(\phi, M_2) < \|\Delta \phi\|_2$.

Proof. Since M_2 is an $(N+2)$ -dimensional manifold embedded in $\mathcal{D}_0^{2,2}(\mathbb{R}^N)$,

$$(c, \lambda, y) \in \mathbb{R} \times \mathbb{R}_+ \times \mathbb{R}^N \rightarrow cU_{\lambda,y} \in \mathcal{D}_0^{2,2}(\mathbb{R}^N).$$

Take $\phi \in \mathcal{D}_0^{2,2}(\mathbb{R}^N)$ with

$$\begin{aligned} d &= d(\phi, M_2) \\ &= \inf_{c,\lambda,y} \|\Delta(\phi - cU_{\lambda,y})\|_2^2 \\ &= \inf_{c,\lambda,y} \left(\|\Delta\phi\|_2^2 + c^2 - 2c \int_{\mathbb{R}^N} \Delta\phi \cdot \Delta U_{\lambda,y} \right) \\ &< \|\Delta\phi\|_2^2. \end{aligned}$$

It is easy to see that the infimum above is attained at a point $(c_0, \lambda_0, y_0) \in \mathbb{R} \times \mathbb{R}_+ \times \mathbb{R}^N$ with $c_0 \neq 0$. Since $M_2 \setminus \{0\}$ is a smooth manifold and the tangential space at (c_0, λ_0, y_0) is given by

$$TM_{c_0U_{\lambda_0,y_0}} = \text{span} \{U_{\lambda_0,y_0}, \partial_\lambda U_{\lambda_0,y_0}, \nabla_y U_{\lambda_0,y_0}\},$$

we must have that $(\phi - c_0U_{\lambda_0,y_0})$ is perpendicular to $TM_{c_0U_{\lambda_0,y_0}}$. By Lemma 2.15, we necessarily have

$$\int_{\mathbb{R}^N} |\Delta(\phi - c_0U_{\lambda_0,y_0})|^2 \geq \mu_{N+3} S_2^{p+1} \int_{\mathbb{R}^N} U_{\lambda_0,y_0}^{p-1} (\phi - c_0U_{\lambda_0,y_0})^2.$$

Let $\phi = c_0U_{\lambda_0,y_0} + dv$, $\|\Delta v\|_2 = 1$. Then

$$\|\Delta\phi\|_2^2 = d^2 + c_0^2 \|\Delta U_{\lambda_0,y_0}\|_2^2 = d^2 + c_0^2,$$

$$\begin{aligned} \int_{\mathbb{R}^N} \phi^{p+1} &= \int_{\mathbb{R}^N} (c_0U_{\lambda_0,y_0} + dv)^{p+1} \\ &= \int_{\mathbb{R}^N} c_0^{p+1} U_{\lambda_0,y_0}^{p+1} + d(p+1)c_0^p \int_{\mathbb{R}^N} U_{\lambda_0,y_0}^p v \\ &\quad + d^2 \frac{p(p+1)}{2} |c_0|^{p-1} \int_{\mathbb{R}^N} U_{\lambda_0,y_0}^{p-1} v^2 + o(d^2). \end{aligned}$$

Since v is perpendicular to U_{λ_0,y_0} , then $\int_{\mathbb{R}^N} U_{\lambda_0,y_0}^p v dx = 0$, i.e.,

$$\int_{\mathbb{R}^N} \Delta v \cdot \Delta U_{\lambda_0,y_0} = 0 = \int_{\mathbb{R}^N} v \Delta^2 U_{\lambda_0,y_0} = S_2^{p+1} \int_{\mathbb{R}^N} v U_{\lambda_0,y_0}^p.$$

Thus

$$\int_{\mathbb{R}^N} \phi^{p+1} \leq c_0^{p+1} S_2^{-(p+1)} + d^2 |c_0|^{p-1} \frac{p(p+1)}{2\mu_{N+3} S_2^{p+1}} + o(d^2).$$

Therefore

$$\begin{aligned} \left(\int_{\mathbb{R}^N} \phi^{p+1} \right)^{\frac{2}{p+1}} &\leq \left(c_0^{p+1} S_2^{-(p+1)} + d^2 |c_0|^{p-1} \frac{p(p+1)}{2\mu_{N+3} S_2^{p+1}} + o(d^2) \right)^{\frac{2}{p+1}} \\ &= c_0^2 S_2^{-2} \left(1 + d^2 c_0^{-2} \frac{p(p+1)}{2\mu_{N+3}} + o(d^2) \right)^{\frac{2}{p+1}} \\ &= c_0^2 S_2^{-2} \left(1 + \frac{2}{p+1} d^2 c_0^{-2} \frac{p(p+1)}{2\mu_{N+3}} + o(d^2) \right) \\ &= c_0^2 S_2^{-2} + \frac{d^2 p}{\mu_{N+3}} S_2^{-2} + o(d^2). \end{aligned}$$

Thus

$$\begin{aligned} & \|\Delta\phi\|_2^2 - S_2^2\|\phi\|_{\frac{2N}{N-4}}^2 \\ & \geq d^2 + c_0^2 - S_2^2\left(c_0^2S_2^{-2} + \frac{d^2pS_2^{-2}}{\mu_{N+3}} + o(d^2)\right) \\ & \geq d^2\left(1 - \frac{p}{\mu_{N+3}}\right) + o(d^2). \end{aligned}$$

This is true when d is small. Thus the lemma holds with $\alpha = (1 - \frac{p}{\mu_{N+3}})$. To see that this is the best possible result we can argue as follows:

Take $\phi = U + dv$, where v is the $(N+3)$ th eigenfunction of $L_{1,0}$ and d is a small positive number. We then have $d(\phi, M_2) = d$ if d is small and the closest point on M is U . \square

Proof of Theorem 1.9. The fact that the result is sharp follows from the last part of the proof of the lemma above. Assume that the theorem were not true, then we could find a sequence $\{\phi_m\}$ such that

$$\frac{\|\Delta\phi_m\|_2^2 - S_2^2\|\phi_m\|_{\frac{2N}{N-4}}^2}{d(\phi_m, M_2)^2} \rightarrow 0, \text{ as } m \rightarrow \infty.$$

By homogeneity we can assume that $\|\Delta\phi_m\|_2 = 1$, and after selecting a subsequence we can assume that $d(\phi_m, M_2) \rightarrow L \in [0, 1]$. Note that $d(\phi_m, M_2) \leq \|\Delta\phi_m\|_2 = 1$. If $L = 0$, then we have a contradiction by Lemma 3.1 above. The other possibility is that $L > 0$. In this case we must have

$$\|\Delta\phi_m\|_2^2 - S_2^2\|\phi_m\|_{\frac{2N}{N-4}}^2 \rightarrow 0, \|\Delta\phi_m\|_2^2 = 1.$$

By P. L. Lions' concentration and compactness principle (see Corollary 1.2 of Section I.4 in Part I of P. L. Lions [6]) we have two sequences of numbers λ_m, y_m so that

$$\lambda_m\phi_m(\lambda_m^{\frac{4}{N-4}}(x - y_m)) \rightarrow +U \text{ (or } -U) \in \mathcal{D}_0^{2,2}(\mathbb{R}^N) \text{ as } m \rightarrow \infty.$$

This implies

$$d(\phi_m, M_2) = d\left(\lambda_m\phi_m\left(\lambda_m^{\frac{4}{N-4}}(x - y_m)\right), M_2\right) \rightarrow 0 \text{ as } m \rightarrow \infty.$$

This is a contradiction.

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