

IMAGINARY POWERS OF LAPLACE OPERATORS

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ABSTRACT. We show that if L is a second-order uniformly elliptic operator in divergence form on \mathbf{R}^d , then $C_1(1+|\alpha|)^{d/2} \leq \|L^{i\alpha}\|_{L^1 \rightarrow L^{1,\infty}} \leq C_2(1+|\alpha|)^{d/2}$. We also prove that the upper bounds remain true for any operator with the finite speed propagation property.

1. INTRODUCTION

Assume that $a_{ij} \in C^\infty(\mathbf{R}^d)$, $a_{ij} = a_{ji}$ for $1 \leq i, j \leq d$ and that $\kappa I \leq (a_{ij}) \leq \tau I$ for some positive constants κ and τ . We define a positive self-adjoint operator L on $L^2(\mathbf{R}^d)$ by the formula

$$(1) \quad L = - \sum \partial_i a_{ij} \partial_j.$$

We refer readers to [8] for the precise definition and basic properties of L . In particular, L admits a spectral resolution $E(t)$ and we can define the operator $L^{i\alpha}$ by the formula

$$L^{i\alpha} = \int_0^\infty t^{i\alpha} dE(t).$$

By spectral theory $\|L^{i\alpha}\|_{L^2 \rightarrow L^2} = 1$. It is well known that $L^{i\alpha}$ falls within the scope of classical Calderón-Zygmund theory (as described in [3] or [22]) and so it extends to a bounded operator on L^p , $1 < p < \infty$, and is also weak type (1,1). The main aim of this paper is to obtain the sharp estimate for the weak type (1,1) norm of $L^{i\alpha}$ in terms of α .

The study of imaginary powers of operators is an important part of the theory of operators of type ω with H^∞ functional calculus (see e.g., [6], [9] and [17]). What is perhaps more interesting and relevant from the point of view of this paper is that the weak type (1,1) norm of imaginary powers of self-adjoint operators can play a central role in the theory of spectral multipliers. See [5] and [15]. Imaginary powers of Laplace operators on compact Lie groups were also investigated in [20]. Theorem 2 below applied to Laplace operators on compact Lie groups gives the sharp endpoint result of Theorem 3 in [20], pp. 58. See also Corollary 4 of [20], pp. 121.

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However, the starting point for this paper is the following observation from [2]. If we denote the weak type (1,1) norm of an operator T on a measure space (X, μ) by $\|T\|_{L^1 \rightarrow L^1, \infty} = \sup \lambda \mu(\{x \in X : |Tf(x)| > \lambda\})$ where the supremum is taken over $\lambda > 0$ and functions f with $L^1(X)$ norm less than one, then for the standard Laplace operator on \mathbf{R}^d ,

$$(2) \quad C_1(1 + |\alpha|)^{d/2} \leq \|(-\Delta_d)^{i\alpha}\|_{L^1 \rightarrow L^1, \infty} \leq C_2(1 + |\alpha|)^{d/2} \log(1 + |\alpha|).$$

The classical Hörmander multiplier theorem (see [13]) states that a multiplier operator T_m on \mathbf{R}^d with multiplier m satisfies

$$(3) \quad \|T_m\|_{L^1 \rightarrow L^1, \infty} \leq C_s \sup_{t>0} \|\eta(\cdot)m(t)\|_{H_s} \leq A$$

for any $s > d/2$ and any $\eta \in C_c^\infty(\mathbf{R}_+)$ not identically zero. Here H_s is the Sobolev space of order s on \mathbf{R}^d . Since the Sobolev norm in (3) behaves like $(1 + |\alpha|)^s$ for the multiplier $m(x) = |x|^{i\alpha}$ of $(-\Delta)^{i\alpha}$, (2) shows that the exponent $d/2$ in Hörmander’s theorem is sharp. Furthermore, if (3) is satisfied with $A < \infty$, then the distribution $K = \hat{m}$ agrees with a locally integrable function away from the origin which satisfies

$$(4) \quad I(B) = \sup_{\substack{y \neq 0 \\ |x| \geq B|y|}} \int |K(x - y) - K(x)| dx \leq A$$

for $B \geq 2$ and Hörmander’s theorem actually shows that the weak type (1,1) norm of T_m is bounded by $I(B) + \|m\|_{L^\infty}^2 + B^d$. One can easily compute that for the convolution kernel K of $(-\Delta)^{i\alpha}$, the integral $I(B)$ is bounded above and below by $(1 + |\alpha|)^{d/2} \log(1 + |\alpha|/B)$. Hence Hörmander’s theorem gives the upper bound in (2). The lower bound is a simple consequence of the explicit formula for the kernel K of $(-\Delta)^{i\alpha}$. See for example, [21] pp. 51-52.

The main observation of this paper is to note that there is a slight improvement of the bound $I(B) + \|m\|_{L^\infty}^2 + B^d$ to $I(B) + (\|m\|_{L^\infty}^2 B^d)^{1/2}$. This can be achieved either by using C. Fefferman’s ideas in [11] of exploiting more information of L^2 bounds or by varying the level of the Calderón-Zygmund decomposition and optimising. Hence we will be able to remove the *log* term in (2). We will show that this more precise estimate holds for a general class of operators.

Theorem 1. *Suppose that L is defined by (1). Then*

$$(5) \quad C_1(1 + |\alpha|)^{d/2} \leq \|L^{i\alpha}\|_{L^1 \rightarrow L^1, \infty} \leq C_2(1 + |\alpha|)^{d/2}$$

for all $\alpha \in \mathbf{R}$.

Proof of the lower bound. We begin with some known estimates for the kernel $p_t(x, y)$ of the heat operator e^{-tL} associated to L . First, this kernel satisfies Gaussian bounds

$$(6) \quad C_1 \frac{1}{t^{d/2}} e^{-b_1 \rho^2(x,y)/t} \leq p_t(x, y) \leq C_2 \frac{1}{t^{d/2}} e^{-b_2 \rho^2(x,y)/t}$$

(see [8]) for some positive constants C_1, C_2, b_1 and b_2 and where $\rho(x, y)$ denotes the geodesic distance between x and y given by the Riemannian metric $(a_{i,j})$. In this setting of uniform ellipticity, $\kappa|x - y| \leq \rho(x, y) \leq \tau|x - y|$. Secondly, from the construction of a parametrix for the heat equation with respect to L (either via Hadamard’s construction, see §17.4 of [14], or using pseudodifferential operator

techniques, see chapter 7, §13 of [23]), we have for each $y \in \mathbf{R}^d$ a ball $B(y, r)$ such that for $x \in B(y, r)$ and $0 < t < 1$,

$$(7) \quad |p_t(x, y) - (\det a_{ij}(y))^{-1/2} (4\pi t)^{-d/2} e^{-\rho^2(x,y)/4t}| \leq C t^{1/2} t^{-d/2}.$$

Here we are using the fact that p_t is symmetric, $p_t(x, y) = p_t(y, x)$. From (6) and (7), we have for $x \in B(y, r)$ the bound

$$|p_t(x, y) - (\det a_{ij}(y))^{-1/2} (4\pi t)^{-d/2} e^{-\rho^2(x,y)/4t}| \leq C t^{1/4} t^{-d/2} \exp(-b' \rho(x, y)^2/t)$$

which translates into a bound for the kernel $K_{L^{i\alpha}}$ of $L^{i\alpha}$ since the functional calculus for L gives us the relationship

$$L^{i\alpha} = \Gamma(-i\alpha)^{-1} \int_0^\infty t^{-i\alpha-1} e^{-tL} dt$$

for $\alpha \neq 0$. Thus for $x \in B(y, r)$,

$$(8) \quad \begin{aligned} &|K_{L^{i\alpha}}(x, y) - (\det a_{ij}(y))^{-1/2} 4^{i\alpha} \pi^{-d/2} \gamma(\alpha) \rho(x, y)^{-d-i2\alpha}| \\ &\leq C |\Gamma(-i\alpha)|^{-1} \rho(x, y)^{-d+1/2} \end{aligned}$$

where $\gamma(\alpha) = \Gamma(i\alpha + d/2)/\Gamma(-i\alpha)$. Using (8) with $y = 0$ we obtain for λ large enough

$$\begin{aligned} &\mu(\{|K_{L^{i\alpha}}(x, 0)| \geq \lambda\}) \\ &\geq \mu(\{C_1 |\gamma(\alpha)| \rho^{-d}(x, 0) \geq 2\lambda\}) - \mu(\{C_2 |\Gamma(-i\alpha)| \rho^{-d+\frac{1}{2}}(x, 0) \geq \lambda\}) \\ &= \mu(B(0, (2C_1 |\gamma(\alpha)|/\lambda)^{1/d})) - \mu(B(0, (C_2 |\Gamma(-i\alpha)|/\lambda)^{1/(d-1/2)})) \\ &\geq C' |\gamma(\alpha)|/\lambda. \end{aligned}$$

Here μ is Lebesgue measure and the sets above have the further restriction that $x \in B(0, r)$. Since $K_{L^{i\alpha}}$ is smooth away from the diagonal, we see that $L^{i\alpha} \phi_\delta(x)$ tends to $K_{L^{i\alpha}}(x, 0)$ as $\delta \rightarrow 0$ for any $x \neq 0$ and any approximation of the identity $\{\phi_\delta\}$. Hence the above estimate shows that the weak type (1,1) norm of $L^{i\alpha}$ is bounded below by $|\gamma(\alpha)| = |\Gamma(i\alpha + d/2)/\Gamma(-i\alpha)| \sim (1 + |\alpha|)^{\frac{d}{2}}$ (see [10]).

The upper bound in Theorem 1 holds in a much more general setting which we describe now. Assume that (X, μ, ρ) is a space with measure μ and metric ρ . If $\|P\|_{L^2 \rightarrow L^\infty} < \infty$, then we can define the kernel K_P of the operator P by the formula

$$\langle P(\psi), \phi \rangle = \int P(\psi) \bar{\phi} d\mu = \int K_P(x, y) \psi(x) \overline{\phi(y)} d\mu(x) d\mu(y).$$

Note that $\sup_x \|K_P(x, \cdot)\|_{L^2} = \|P\|_{L^2 \rightarrow L^\infty}$. Next, we say that

$$(9) \quad \text{supp } K_P \subset \{(x, y) \in X^2 : \rho(x, y) \leq r\}$$

if $\langle P(\psi), \phi \rangle = 0$ for every $\phi, \psi \in L^2$ and every $r_1 + r_2 + r < \rho(x', y')$ such that $\psi(x) = 0$ for $\rho(x, x') > r_1$ and $\phi(x) = 0$ for $\rho(x, y') > r_2$. This definition (9) makes sense even if $\|P\|_{L^2 \rightarrow L^\infty} = \infty$. Now if L is a self-adjoint positive definite operator acting on $L^2(\mu)$, then we say that it satisfies the finite speed propagation property of the corresponding wave equation if

$$(10) \quad \text{supp } K_{C_t(\sqrt{L})} \subset \{(x, y) \in X^2 : \rho(x, y) \leq t\},$$

where $C_t(\sqrt{L}) = \int \cos(t\sqrt{\lambda}) dE(\lambda)$.

Theorem 2. *Suppose that L satisfies (10). Next assume that*

$$(11) \quad \|\exp(-tL)\|_{L^2 \rightarrow L^\infty}^2 \leq C_1 V_{d,D}(t^{1/2})^{-1} \leq C\mu(B(x, t^{1/2}))^{-1} \leq C_2 V_{d,D}(t^{1/2})^{-1}$$

for all $t > 0$ and $x \in X$, where $B(x, t)$ is a ball with radius t centred at x and

$$V_{d,D}(t) = \begin{cases} t^d & \text{for } t \leq 1, \\ t^D & \text{for } t > 1, \end{cases}$$

for $d, D \geq 0$. Then

$$\|L^{i\alpha}\|_{L^1 \rightarrow L^{1,\infty}} \leq C_2(1 + |\alpha|)^{\max(d,D)/2}$$

for all $\alpha \in \mathbf{R}$.

We remark that (10) and (11) are equivalent to having Gaussian upper bounds on the heat kernel and the associated volume growth on balls. See [18]. Furthermore, the upper bound in Theorem 1 follows from Theorem 2. Indeed, if $X = \mathbf{R}^d$, $\rho(x, y) = \tau|x - y|$ and μ is Lebesgue measure, then it is well known (see e.g. [8] and [19]) that (11) and (10) hold. We are going to prove Theorem 2 only in the case $d = D$. The argument for the other cases is similar.

2. PRELIMINARIES

The following lemma is a very simple but useful consequence of (10).

Lemma 1. *Assume that L satisfies (10) and that \hat{F} is a Fourier transform of an even bounded Borel function F with $\text{supp } \hat{F} \subset [-r, r]$. Then*

$$\text{supp } K_{F(\sqrt{L})} \subset \{(x, y) \in X^2 : \rho(x, y) \leq r\}.$$

Proof. If F is an even function, then by the Fourier inversion formula,

$$F(\sqrt{L}) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \hat{F}(t) C_t(\sqrt{L}) dt.$$

But since $\text{supp } \hat{F} \subset [-r, r]$,

$$F(\sqrt{L}) = \frac{1}{2\pi} \int_{-r}^r \hat{F}(t) C_t(\sqrt{L}) dt$$

and Lemma 1 follows from (10).

Lemma 2. *Let $\phi \in C_c^\infty(\mathbf{R})$ be even, $\phi \geq 0$, $\|\phi\|_{L^1} = 1$, $\text{supp}(\phi) \subset [-1, 1]$, and set $\phi_r(x) = 1/r \phi(x/r)$ for $r > 0$. Let Φ denote the Fourier transform of ϕ and Φ^r denote the Fourier transform of ϕ_r . If (11) and (10) hold, then the kernel $K_{\Phi^r(\sqrt{L})}$ of the self-adjoint operator $\Phi^r(\sqrt{L})$ satisfies*

$$(12) \quad \text{supp } K_{\Phi^r(\sqrt{L})} \subset \{(x, y) \in X^2; \rho(x, y) \leq r\}$$

and

$$(13) \quad |K_{\Phi^r(\sqrt{L})}(x, y)| \leq C r^{-d}$$

for all $r > 0$ and $x, y \in X$.

Proof. (12) follows from Lemma 1. For any $m \in \mathbf{N}$ and $r > 0$, we have the relationship

$$(I + rL)^{-m} = \frac{1}{m!} \int_0^\infty e^{-rtL} e^{-t} t^{m-1} dt$$

and so when $m > d/4$, (11) implies

$$(14) \quad \begin{aligned} \|(I + rL)^{-m}\|_{L^2 \rightarrow L^\infty} &\leq \frac{1}{m!} \int_0^\infty \|\exp(-rtL)\|_{L^2 \rightarrow L^\infty} e^{-t} t^{m-1} dt \\ &\leq C_1 r^{-d/4} \end{aligned}$$

for all $r > 0$. Now $\|(I + r^2L)^{-m}\|_{L^1 \rightarrow L^2} = \|(I + r^2L)^{-m}\|_{L^2 \rightarrow L^\infty}$ and so

$$\|\Phi^r(\sqrt{L})\|_{L^1 \rightarrow L^\infty} \leq \|(I + r^2L)^{2m} \Phi^r(\sqrt{L})\|_{L^2 \rightarrow L^2} \|(I + r^2L)^{-m}\|_{L^2 \rightarrow L^\infty}^2.$$

The L^2 operator norm of the first term is equal to the L^∞ norm of the function $(1 + r^2|t|)^{2m} \Phi(r\sqrt{|t|})$ which is uniformly bounded in $r > 0$ and so (13) follows by (14).

Next we recall the Calderón-Zygmund decomposition in the general setting of spaces of homogeneous type (see e.g. [3] or [22]).

Lemma 3. *There exists C such that, given $f \in L^1(X, \mu)$ and $\lambda > 0$, one can decompose f as*

$$f = g + b = g + \sum b_i$$

so that

1. $|g(x)| \leq C\lambda$, a.e. x and $\|g\|_{L^1} \leq C\|f\|_{L^1}$.
2. There exists a sequence of balls $B_i = B(x_i, r_i)$ such that the support of each b_i is contained in B_i and

$$\int |b_i(x)| d\mu(x) \leq C\lambda\mu(B_i).$$

3. $\sum \mu(B_i) \leq C \frac{1}{\lambda} \int |f(x)| d\mu(x)$.
4. There exists $k \in \mathbf{N}$ such that each point of X is contained in at most k of the balls $B(x_i, 2r_i)$.

We are now in a position to prove Theorem 2.

3. PROOF OF THEOREM 2

The proof follows closely the line of argument in [1] (which of course generalises to this setting). We are attempting to prove

$$\lambda \mu(\{x \in X : |L^{i\alpha} f(x)| \geq \lambda\}) \leq C(1 + |\alpha|)^{\frac{d}{2}} \|f\|_{L^1}.$$

As usual we start by decomposing f into $g + \sum b_i$ at the level of λ according to Lemma 3. We will follow the idea of C. Fefferman [11] of using more information of the L^2 operator norm (in our case, $\|L^{i\alpha}\|_{L^2 \rightarrow L^2} = 1$) by smoothing out the bad functions b_i at a scale smaller than the size of its support and considering this part of the good function where L^2 estimates can be used (see also [4]). In our case for each b_i , consider $\Phi^{s_i}(\sqrt{L})b_i$ where $s_i = \theta r_i$, $\theta = (1 + |\alpha|)^{-\frac{1}{2}}$, and let

$G = g + \sum \Phi^{s_i}(\sqrt{L})b_i$ be the modified good function. Hence $f = G + B$ where $B = \sum(I - \Phi^{s_i}(\sqrt{L}))b_i$ and we write

$$(15) \quad \begin{aligned} \lambda \mu(\{|L^{i\alpha} f(x)| \geq \lambda\}) &\leq \lambda \mu(\{|L^{i\alpha} G(x)| \geq \lambda/2\}) \\ &\quad + \lambda \mu(\{|L^{i\alpha} B(x)| \geq \lambda/2\}). \end{aligned}$$

The first term is less than $4/\lambda \|L^{i\alpha} G\|_{L^2}^2 \leq 4/\lambda \|G\|_{L^2}^2$. However, according to Lemma 2,

$$|\Phi^{s_i}(\sqrt{L})b_i(x)| \leq \int |K_{\Phi^{s_i}(\sqrt{L})}(x, y)b_i(y)| d\mu(y) \leq C(\theta r_i)^{-d} \|b_i\|_{L^1} \mathbb{1}_{B(x_i, 2r_i)}$$

and therefore by Lemma 3, $|G(x)| \leq C\theta^{-d}\lambda$ for a.e., x . Using Lemma 2 again which shows that the $L^p \rightarrow L^p$ operator norms of $\Phi^r(\sqrt{L})$ are uniformly bounded in $r > 0$, we also have that

$$\|G\|_{L^1} \leq \|g\|_{L^1} + C \sum \|\Phi^{s_i}(\sqrt{L})b_i\|_{L^1} \leq \|g\|_{L^1} + C \sum \|b_i\|_{L^1} \leq C\|f\|_{L^1}.$$

Therefore the first term in (15) is bounded by $(1 + |\alpha|)^{\frac{d}{2}} \|f\|_{L^1}$.

Since $\mu(\cup B(x_i, \theta^{-1}r_i)) \leq C\theta^{-d} \sum \mu(B_i) \leq C(1 + |\alpha|)^{\frac{d}{2}} \|f\|_{L^1}/\lambda$, then to bound the second term in (15), it suffices to show

$$(16) \quad \int_{x \notin \cup B_i^*} |L^{i\alpha} B(x)| d\mu(x) \leq C(1 + |\alpha|)^{\frac{d}{2}} \|f\|_{L^1},$$

where $B_i^* = B(x_i, \theta^{-1}r_i)$. Let $H^\alpha(t) = |t|^{2i\alpha}$ so that

$$L^{i\alpha} B(x) = \sum H^\alpha(1 - \Phi^{s_i})(\sqrt{L})b_i(x)$$

and therefore the left side of (16) is less than

$$\begin{aligned} \sum_i \int_{x \notin \cup_j B_j^*} \left| \int K_{H^\alpha(1 - \Phi^{s_i})(\sqrt{L})}(x, y)b_i(y) d\mu(y) \right| d\mu(x) \\ \leq \sum_i \int |b_i(y)| \int_{x \notin B_i^*} |K_{H^\alpha(1 - \Phi^{s_i})(\sqrt{L})}(x, y)| d\mu(x) d\mu(y). \end{aligned}$$

Since $F(L)^* = \overline{F}(L)$, we may interchange the roles of x and y , and so (16) will follow from Lemma 3 once we establish

$$(17) \quad \sup_{x, i} \int_{\rho(x, y) \geq \theta^{-1}r_i} |K_{H^\alpha(1 - \Phi^{s_i})(\sqrt{L})}(x, y)| d\mu(y) \leq C(1 + |\alpha|)^{\frac{d}{2}}.$$

We now fix $x \in X$ and i . Let $\eta \in C_c^\infty(\mathbf{R})$ be an even function supported in $\{t \in \mathbf{R} : 1 \leq |t| \leq 4\}$ such that

$$\sum_{n=-\infty}^\infty \eta(2^{-n}t) = 1 \text{ for all } t \neq 0.$$

We put $H_n^\alpha(t) = \eta(2^{-n}t)H^\alpha(t)$ so that

$$H^\alpha(1 - \Phi^{s_i})(\sqrt{L}) = \sum_n H_n^\alpha(1 - \Phi^{s_i})(\sqrt{L}).$$

Thus

$$(18) \quad \int_{y \notin B_i^*} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)| d\mu(y) \leq \sum_n \int_{y \notin B_i^*} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)| d\mu(y)$$

and we will estimate each term in the sum on the right side in terms of n and i , uniformly in $x \in X$.

Let $k_o = [d/2] + 1$ so that

$$\int_{y \notin B_i^*} (1 + 2^n \rho(x, y))^{-2k_o} d\mu(y) \leq C \int_{\theta^{-1} r_i}^\infty (1 + 2^n r)^{-2k_o} r^{d-1} dr \leq C 2^{-2nk_o} (\theta^{-1} r_i)^{d-2k_o}$$

and therefore by the Cauchy-Schwarz inequality,

$$(19) \quad \int_{y \notin B_i^*} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)| d\mu(y) \leq C 2^{-nk_o} (\theta^{-1} r_i)^{\frac{d}{2}-k_o} \cdot \left(\int_{\rho(x,y) \geq \theta^{-1} r_i} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)|^2 (1 + 2^n \rho(x, y))^{2k_o} d\mu(y) \right)^{1/2}.$$

We break up the integral on the right side of (19) where $2^n \rho(x, y)$ is roughly constant and consider

$$(20) \quad \sum_{2^j \geq 2^n r_i \theta^{-1}} 2^{2jk_o} \int_{2^{j-1-n} < \rho(x,y) \leq 2^{j-n}} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)|^2 d\mu(y).$$

Fix a nonnegative even $\varphi \in C_c^\infty(\mathbf{R})$ such that $\varphi = 1$ on $[-1/4, 1/4]$ and $\varphi = 0$ on $\mathbf{R} \setminus [-1/2, 1/2]$. Then the Fourier transforms of $H_n^\alpha(1 - \Phi^{s_i})$ and $H_n^\alpha(1 - \Phi^{s_i}) * (\delta - \hat{\varphi}_{2^{n-j}})$ agree on $\{\xi : |\xi| \geq 2^{j-1-n}\}$ and so by Lemma 1, the kernels of $H_n^\alpha(1 - \Phi^{s_i})(\sqrt{L})$ and $H_n^\alpha(1 - \Phi^{s_i}) * (\delta - \hat{\varphi}_{2^{n-j}})(\sqrt{L})$ agree on the set $\{(x, y) \in X^2 : \rho(x, y) \geq 2^{j-1-n}\}$. Here δ denotes the Dirac mass at 0. For each j , the integrals in (20) satisfy the bound

$$\int_{2^{j-1-n} < \rho(x,y) \leq 2^{j-n}} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y)|^2 d\mu(y) \leq \|K_{F_{n,j}^\alpha(\sqrt{L})}\|_{L^2 \rightarrow L^\infty}^2,$$

where we are defining $F_{n,j}^\alpha(t) = H_n^\alpha(1 - \Phi^{s_i}) * (\delta - \hat{\varphi}_{2^{n-j}})(t)$. So by (14), the right side of this inequality is bounded by $\|(I + 2^{-2n}L)^m F_{n,j}^\alpha(\sqrt{L})\|_{L^2 \rightarrow L^2}^2 2^{nd}$ as long as $m > d/4$. Everything then comes down to estimating the L^∞ norm of $(1 + 2^{-2n}t^2)^m F_{n,j}^\alpha(t)$. We make the following claim.

Claim. For each j, n and $m > d/4$,

$$(1 + 2^{-2n}t^2)^m |F_{n,j}^\alpha(t)| \leq C_m |\alpha|^{k_o} 2^{-jk_o} \min(1, (2^n r_i \theta)^2) \min(1, |\alpha| 2^{-j})$$

uniformly in $t \in \mathbf{R}$.

The claim shows that

$$\|K_{F_{n,j}^\alpha(\sqrt{L})}\|_{L^2 \rightarrow L^\infty} \leq C |\alpha|^{k_o} 2^{-jk_o} 2^{\frac{nd}{2}} \min(1, (2^n r_i \theta)^2) \min(1, |\alpha| 2^{-j})$$

and hence the sum in (20) is bounded by

$$\begin{aligned} & |\alpha|^{2k_o} 2^{nd} \min^2(1, (2^n r_i \theta)^2) \sum_{2^j \geq 2^n r_i \theta^{-1}} \min^2(1, |\alpha| 2^{-j}) \\ & \leq |\alpha|^{2k_o} 2^{nd} \min^2(1, (2^n r_i \theta)^2) \log(2 + \frac{|\alpha|}{2^n r_i \theta^{-1}}). \end{aligned}$$

Recall that θ and α are related so that $\theta|\alpha| = |\alpha|/(1 + |\alpha|)^{\frac{1}{2}} \leq \theta^{-1}$. Plugging this into (19) gives

$$\int_{y \notin B_i^*} |K_{H_n^\alpha(1-\Phi^{s_i})(\sqrt{L})}(x, y) d\mu(y) \leq \theta^{-d} (2^n r_i \theta)^{\frac{d}{2} - k_o} \min(1, (2^n r_i \theta)^2) \log(2 + \frac{1}{2^n r_i \theta})$$

and this makes the sum in (18) bounded by $\theta^{-d} = (1 + |\alpha|)^{\frac{d}{2}}$, proving (17) and hence Theorem 2.

Proof of the Claim. If $G_n(t) = H_n^\alpha(t)(1 - \Phi^{s_i}(t))$, then $F_{n,j}^\alpha(t) = 2^{(n-j)k_o} G_n^{(k_o)} * \hat{\psi}_{2^{n-j}}(t)$ where $\psi(\xi) = \xi^{-k_o}(1 - \varphi(\xi))$ (and so $\hat{\psi}$ is continuous, rapidly decreasing and has vanishing moments, $\int t^\ell \hat{\psi}(t) dt = 0$, $\ell = 0, 1, 2, \dots$). Hence

$$\begin{aligned} F_{n,j}^\alpha(t) &= 2^{(n-j)k_o} \int_{\mathbf{R}} [G_n^{(k_o)}(t-s) - G_n^{(k_o)}(t)] \hat{\psi}_{2^{n-j}}(s) ds \\ &= 2^{(n-j)k_o} \int_{\mathbf{R}} [G_n^{(k_o)}(t-2^{n-j}s) - G_n^{(k_o)}(t)] \hat{\psi}(s) ds. \end{aligned}$$

However $G_n(t) = \eta(2^{-n}t)|t|^{2i\alpha}(1 - \Phi(s_i t))$ and thereby each time we take a derivative, we gain a factor of 2^{-n} . $G_n^{(k_o)}(t)$ is thus a finite sum of terms of the form $\alpha^p 2^{-nk_o} \tilde{\eta}(2^{-n}t)|t|^{2i\alpha} \Psi(s_i t)$ where $\tilde{\eta} \in C_c^\infty(\mathbf{R})$, $\text{supp}(\tilde{\eta}) \subset \text{supp}(\eta)$ and Ψ is a Schwartz function which is $0(t^2)$ as $t \rightarrow 0$ (note that $\Phi'(0) = \int x\phi(x)dx = 0$ since ϕ is even). The worst power p is k_o which occurs when all derivatives land on the factor $|t|^{2i\alpha}$.

Without loss of generality, let us suppose that

$$G^{(k_o)}(t) = \alpha^{k_o} 2^{-nk_o} \eta(2^{-n}t)|t|^{2i\alpha} \Psi(s_i t).$$

From the above integral representation of $F_{n,j}^\alpha(t)$, we see that the main contribution to $(1 + 2^{-2n}t^2)^m |F_{n,j}^\alpha(t)|$ occurs when $|t| \sim 2^n$ and in this case,

$$|F_{n,j}^\alpha(t)| \leq C|\alpha|^{k_o} 2^{(n-j)k_o} 2^{-nk_o} \min(1, (s_i 2^n)^2) \leq C|\alpha|^{k_o} 2^{-jk_o} \min(1, (2^n r_i \theta)^2).$$

However we may write

$$F_{n,j}^\alpha(t) = -2^{(n-j)k_o} 2^{n-j} \int_0^1 \int_{\mathbf{R}} G_n^{(k_o+1)}(t - \sigma 2^{j-n} s) s \hat{\psi}(s) ds d\sigma$$

and therefore we also have

$$\begin{aligned} |F_{n,j}^\alpha(t)| &\leq C|\alpha|^{k_o+1} 2^{(n-j)k_o} 2^{n-j} 2^{-n(k_o+1)} \min(1, (s_i 2^n)^2) \\ &\leq C|\alpha|^{k_o} 2^{-jk_o} |\alpha| 2^{-j} \min(1, (2^n r_i \theta)^2), \end{aligned}$$

establishing the claim.

Remarks. Theorem 1 holds also for Laplace-Beltrami operators on compact manifolds of dimension d . The proof is essentially the same as the proof of Theorem 1.

The hypotheses of Theorem 2 are satisfied for Laplace operators on Lie groups of polynomial growth. However, if L is a sub-Laplacian on the three dimensional Heisenberg group, then $d = 4$ but

$$C_1(1 + |\alpha|)^{3/2} \leq \|L^{i\alpha}\|_{L^1 \rightarrow L^{1,\infty}} \leq C_\epsilon(1 + |\alpha|)^{3/2+\epsilon}.$$

(See [16]; see also [12].) The same estimates hold for a sub-Laplacian on $SU(2)$ for which $d = 4$ and $D = 0$ (see [7]). Thus there are situations where the upper bound is better than the one given by Theorem 2 and where the lower bound in Theorem 1 is false. For general groups of polynomial growth Theorem 2 gives the best known estimates; however as the above examples show, these bounds are not always best possible.

REFERENCES

- [1] G. Alexopoulos, Spectral multipliers on Lie groups of polynomial growth, *Proc. Amer. Math. Soc.* (3) **120** (1994), 973-979. MR **95j**:22016
- [2] M. Christ, L^p bounds for spectral multipliers on nilpotent groups, *Trans. Amer. Math. Soc.* (1) **328** (1991), 73-81. MR **92k**:42017
- [3] M. Christ, *Lectures on singular integral operators*, CBMS Regional Conference Series in Mathematics, 77, Amer. Math. Soc. Providence (1990). MR **92f**:42021
- [4] M. Christ, E.M. Stein, A remark on singular Calderón - Zygmund theory, *Proc. Amer. Math. Soc.* (3) **99** (1987), 71-75. MR **88c**:42030
- [5] M. Cowling, S. Meda, Harmonic analysis and ultracontractivity, *Trans. Amer. Math. Soc.* (2) **340** (1993), 733-752. MR **94b**:47050
- [6] M. Cowling, I. Doust, A. McIntosh, A. Yagi, Banach space operators with bounded H^∞ functional calculus, *J. Austral. Math. Soc. (Series A)* **60** (1996), 51-89. MR **97d**:47023
- [7] M. Cowling, A. Sikora, Spectral multipliers for a sub-Laplacian on $SU(2)$, *Math. Z.*, to appear.
- [8] E.B. Davies, *Heat kernels and spectral theory*, Cambridge Univ. Press (1989). MR **90e**:35123
- [9] G. Dore, A. Venni, On the closedness of the sum of two closed operators, *Math. Z.* (2) **196** (1987), 189-201. MR **88m**:47072
- [10] A. Erdélyi, W. Magnus, F. Oberhettinger, F.G. Tricomi, *Higher transcendental functions*. Vol. I. Mc Graw-Hill (1953). MR **15**:419i
- [11] C. Fefferman, Inequalities for strongly singular convolution operators, *Acta Math.* **124** (1970), 9-36. MR **41**:2468
- [12] W. Hebisch, Multiplier theorem on generalized Heisenberg groups, *Colloq. Math.* **65** (1993), 231-239. MR **94m**:43013
- [13] L. Hörmander, Estimates for translation invariant operators in L^p spaces, *Acta Math.* **104** (1960), 93-139. MR **22**:12389
- [14] L. Hörmander, *The analysis of linear partial differential operators III*, Springer-Verlag (1985). MR **87d**:35002a; corrected reprint MR **95h**:35255
- [15] S. Meda, A general multiplier theorem, *Proc. Amer. Math. Soc.* (3) **110** (1990), 639-647. MR **91f**:42010
- [16] D. Müller, E.M. Stein, On spectral multipliers for Heisenberg and related groups, *J. Math. Pures Appl.* **73** (1994), 413-440. MR **96m**:43004
- [17] J. Prüss, H. Sohr, On operators with bounded imaginary powers in Banach spaces, *Math. Z.* (3) **203** (1990), 429-452. MR **91b**:47030
- [18] A. Sikora, Sharp pointwise estimates on heat kernels, *Quart. J. Math. Oxford* (2) **47** (1996), 371-382. MR **97m**:58189
- [19] A. Sikora, On-diagonal estimates on Schrödinger semigroup kernels and reduced heat kernels, *Comm. Math. Phys.* **188** (1997), 233-249. MR **98k**:58213
- [20] E.M. Stein, *Topics in harmonic analysis related to the Littlewood-Paley Theory*, Princeton University Press (1970). MR **40**:6176

- [21] E.M. Stein, *Singular integrals and differentiability properties of functions*, Princeton University Press (1970). MR **44**:7280
- [22] E.M. Stein, *Harmonic analysis: real-variable methods, orthogonality, and oscillatory integrals*, Princeton University Press (1993). MR **95c**:42002
- [23] M. Taylor, *Partial differential equations, II*, Springer-Verlag, New York (1996). MR **98b**:35003

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