

## A LOWER BOUND FOR SUMS OF EIGENVALUES OF THE LAPLACIAN

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ABSTRACT. Let  $\lambda_k(\Omega)$  be the  $k$ th Dirichlet eigenvalue of a bounded domain  $\Omega$  in  $\mathbb{R}^n$ . According to Weyl's asymptotic formula we have

$$\lambda_k(\Omega) \sim C_n(k/V(\Omega))^{2/n}.$$

The optimal in view of this asymptotic relation lower estimate for the sums  $\sum_{j=1}^k \lambda_j(\Omega)$  has been proven by P.Li and S.T.Yau (*Comm. Math. Phys.* **88** (1983), 309-318). Here we will improve this estimate by adding to its right-hand side a term of the order of  $k$  that depends on the ratio of the volume to the moment of inertia of  $\Omega$ .

### 1. INTRODUCTION

Let  $\Omega \subseteq \mathbb{R}^n$  be a bounded open set and let  $0 < \lambda_1(\Omega) \leq \lambda_2(\Omega) \leq \dots$  denote the eigenvalues (repeated with multiplicity) of the Dirichlet Laplacian on  $\Omega$ , that is, of the eigenvalue problem

$$(1.1) \quad \begin{aligned} \Delta u + \lambda u &= 0 \text{ in } \Omega, \\ u &= 0 \text{ on } \partial\Omega. \end{aligned}$$

The asymptotic behavior of  $\lambda_k(\Omega)$  as  $k \rightarrow \infty$  relates to geometric properties of the open set  $\Omega$ . In fact Weyl's asymptotic formula asserts that

$$(1.2) \quad \lambda_k(\Omega) \sim C_n \left( \frac{k}{V(\Omega)} \right)^{2/n} \text{ as } k \rightarrow \infty$$

where  $V(\Omega)$  is the volume of  $\Omega$  and  $C_n = (2\pi)^2 \omega_n^{-2/n}$  with  $\omega_n$  being the volume of the unit ball in  $\mathbb{R}^n$ . Pólya proved in [4] that the above asymptotic relation is in fact a one-sided inequality if  $\Omega$  is a plane domain that tiles  $\mathbb{R}^2$  (and his proof also works in  $\mathbb{R}^n$ ) and he conjectured, for any domain  $\Omega$  in  $\mathbb{R}^n$ , the inequality

$$(1.3) \quad \lambda_k(\Omega) \geq C_n \left( \frac{k}{V(\Omega)} \right)^{2/n}$$

for all  $k \geq 1$ .

In this direction Lieb [3] proved an inequality like (1.3) for any domain  $\Omega$  in  $\mathbb{R}^n$  but with a constant  $\tilde{C}_n$  that differs from the constant  $C_n$  by a factor.

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Then P.Li and S.T.Yau [2] proved that on the average (1.3) is true for any domain  $\Omega$  in  $\mathbb{R}^n$ , that is,

$$(1.4) \quad \sum_{j=1}^k \lambda_j(\Omega) \geq \frac{nC_n}{n+2} k^{\frac{n+2}{n}} V(\Omega)^{-\frac{2}{n}}$$

which is sharp in view of Weyl's asymptotic formula. This also gives a lower bound for each individual eigenvalue, better than previously known and tending to be optimal as the dimension  $n \rightarrow \infty$ . This inequality was complemented by P.Kröger in [1] who gave an upper bound for the sums of the eigenvalues depending on geometric properties of  $\Omega$  that have to do with the behavior of the volume of the  $\varepsilon$ -neighbourhoods of the boundary  $\partial\Omega$ . Using this he obtained close to optimal estimates for individual eigenvalues that however depend on these geometric assumptions.

Here we will obtain an improvement of the estimate (1.4). Let  $I(\Omega)$  denote the "moment of inertia" of  $\Omega$ , that is,

$$(1.5) \quad I(\Omega) = \min_{a \in \mathbb{R}^n} \int_{\Omega} |x - a|^2 dx.$$

Then we have the following.

**Theorem 1.** *For any bounded open set  $\Omega \subseteq \mathbb{R}^n$  and any  $k \geq 1$  we have*

$$(1.6) \quad \sum_{j=1}^k \lambda_j(\Omega) \geq \frac{nC_n}{n+2} k^{\frac{n+2}{n}} V(\Omega)^{-\frac{2}{n}} + M_n k \frac{V(\Omega)}{I(\Omega)}$$

where the constant  $M_n$  depends only on the dimension.

In fact one may take  $M_n = \frac{c}{n+2}$ ,  $c$  being independent of  $n$ , as the proof will show. The proof will follow in part the argument of Li and Yau in [2].

## 2. LOWER ESTIMATE FOR SUMS OF EIGENVALUES

In this section we will prove Theorem 1.

By translating the open set  $\Omega$  we may assume that

$$(2.1) \quad I(\Omega) = \int_{\Omega} |x|^2 dx.$$

We now fix a  $k \geq 1$  and let  $u_1, \dots, u_k$  denote an orthonormal set of eigenfunctions of (1.1) corresponding to the set of eigenvalues  $\lambda_1(\Omega), \lambda_2(\Omega), \dots, \lambda_k(\Omega)$ . We consider the Fourier transform of each eigenfunction

$$(2.2) \quad f_j(\xi) = \hat{u}_j(\xi) = (2\pi)^{-n/2} \int_{\Omega} u_j(x) e^{ix \cdot \xi} dx.$$

Plancherel's Theorem implies that  $f_1, \dots, f_k$  is an orthonormal set in  $\mathbb{R}^n$ . Since  $u_1, \dots, u_k$  are also orthonormal in  $L^2(\Omega)$ , Bessel's inequality implies that for every  $\xi \in \mathbb{R}^n$

$$(2.3) \quad \sum_{j=1}^k |f_j(\xi)|^2 \leq (2\pi)^{-n} \int_{\Omega} |e^{ix \cdot \xi}|^2 dx = (2\pi)^{-n} V(\Omega)$$

and since

$$(2.4) \quad \nabla f_j(\xi) = (2\pi)^{-n/2} \int_{\Omega} i x u_j(x) e^{ix \cdot \xi} dx$$

that

$$(2.5) \quad \sum_{j=1}^k |\nabla f_j(\xi)|^2 \leq (2\pi)^{-n} \int_{\Omega} |i x e^{ix \cdot \xi}|^2 dx = (2\pi)^{-n} I(\Omega).$$

Since each  $u_j$  vanishes on the boundary of  $\Omega$  it is easy to see that (see [2])

$$(2.6) \quad \int_{\mathbb{R}^n} |\xi|^2 |f_j(\xi)|^2 d\xi = \int_{\Omega} |\nabla u_j(x)|^2 dx = \lambda_j(\Omega)$$

for each  $j$ . Hence setting

$$(2.7) \quad F(\xi) = \sum_{j=1}^k |f_j(\xi)|^2$$

we have  $0 \leq F(\xi) \leq (2\pi)^{-n} V(\Omega)$ ,

$$(2.8) \quad |\nabla F(\xi)| \leq 2 \left( \sum_{j=1}^k |f_j(\xi)|^2 \right)^{1/2} \left( \sum_{j=1}^k |\nabla f_j(\xi)|^2 \right)^{1/2} \leq 2(2\pi)^{-n} \sqrt{V(\Omega)I(\Omega)}$$

for every  $\xi \in \mathbb{R}^n$  and moreover

$$(2.9) \quad \int_{\mathbb{R}^n} F(\xi) d\xi = k \text{ and } \int_{\mathbb{R}^n} |\xi|^2 F(\xi) d\xi = \sum_{j=1}^k \lambda_j(\Omega).$$

Now let  $F^*(\xi) = \phi(|\xi|)$  denote the decreasing radial rearrangement of  $F$  where we may assume (by approximating  $F$ ) that the decreasing function  $\phi : [0, +\infty) \rightarrow [0, (2\pi)^{-n} V(\Omega)]$  is absolutely continuous. Setting  $\mu(t) = |\{F^* > t\}| = |\{F > t\}|$  the coarea formula implies that

$$(2.10) \quad \mu(t) = \int_t^{(2\pi)^{-n} V(\Omega)} \int_{\{F=s\}} |\nabla F|^{-1} d\sigma_s ds.$$

Since  $F^*$  is radial we have  $\mu(\phi(s)) = |\{F^* > \phi(s)\}| = \omega_n s^n$  and so differentiating we get  $n\omega_n s^{n-1} = \mu'(\phi(s))\phi'(s)$  for almost every  $s$ . But (2.8), (2.10) and the isoperimetric inequality imply that, with  $\rho = 2(2\pi)^{-n} \sqrt{V(\Omega)I(\Omega)}$ ,

$$(2.11) \quad \begin{aligned} -\mu'(\phi(s)) &= \int_{\{F=\phi(s)\}} |\nabla F|^{-1} d\sigma_{\phi(s)} \geq \rho^{-1} \text{Vol}_{n-1}(\{F = \phi(s)\}) \\ &\geq \rho^{-1} n\omega_n s^{n-1} \end{aligned}$$

and so

$$(2.12) \quad -\rho \leq \phi'(s) \leq 0$$

for almost every  $s$ .

Now (2.9) implies that

$$(2.13) \quad k = \int_{\mathbb{R}^n} F(\xi) d\xi = \int_{\mathbb{R}^n} F^*(\xi) d\xi = n\omega_n \int_0^\infty s^{n-1} \phi(s) ds$$

and

$$(2.14) \quad \sum_{j=1}^k \lambda_j(\Omega) = \int_{\mathbb{R}^n} |\xi|^2 F(\xi) d\xi \geq \int_{\mathbb{R}^n} |\xi|^2 F^*(\xi) d\xi = n\omega_n \int_0^\infty s^{n+1} \phi(s) ds$$

since  $\xi \rightarrow |\xi|^2$  is radial and increasing.

To prove Theorem 1 we will also need the following lemma.

**Lemma 1.** *Let  $n \geq 1, \rho, A > 0$  and  $\psi : [0, +\infty) \rightarrow [0, +\infty)$  be decreasing (and absolutely continuous) such that*

$$(2.15) \quad -\rho \leq \psi'(s) \leq 0$$

and

$$(2.16) \quad \int_0^\infty s^{n-1} \psi(s) ds = A.$$

Then

$$(2.17) \quad \int_0^\infty s^{n+1} \psi(s) ds \geq \frac{1}{n+2} (nA)^{\frac{n+2}{n}} \psi(0)^{-\frac{2}{n}} + \frac{A\psi(0)^2}{6(n+2)\rho^2}.$$

*Proof.* By considering the function  $\alpha\psi(\beta t)$  for appropriate  $\alpha, \beta > 0$  we may assume that  $\rho = 1$  and  $\psi(0) = 1$ . We also assume that  $B = \int_0^\infty s^{n+1} \psi(s) ds < \infty$ , otherwise there is nothing to prove, and so  $\lim_{j \rightarrow \infty} T_j^{n+1} \psi(T_j) = 0$  for some sequence  $(T_j)$  with  $T_j \rightarrow \infty$  as  $j \rightarrow \infty$ .

Let  $h(s) = -\psi'(s)$  for  $s \geq 0$ . Then  $0 \leq h(s) \leq 1$  and  $\int_0^\infty h(s) ds = \psi(0) = 1$ . Moreover integration by parts shows that

$$(2.18) \quad \int_0^\infty s^n h(s) ds = n \int_0^\infty s^{n-1} \psi(s) ds = nA$$

since  $\lim_{j \rightarrow \infty} T_j^{n+1} \psi(T_j) = 0$ , and

$$(2.19) \quad \int_0^\infty s^{n+2} h(s) ds = \lim_{T \rightarrow \infty} (-T^{n+2} \psi(T) + (n+2) \int_0^T s^{n+1} \psi(s) ds) \leq (n+2)B.$$

Next let  $a \geq 0$  be such that

$$(2.20) \quad \int_a^{a+1} s^n ds = \int_0^\infty s^n h(s) ds = nA.$$

Such an  $a$  exists since by the same argument as in Lemma 1 of [2] one can easily show that the assumptions on  $h$  imply  $\int_0^\infty s^n h(s) ds \geq \int_0^1 s^n ds$ . Indeed this follows by integrating the inequality  $(s^n - 1)(h(s) - \xi(s)) \geq 0$  over  $[0, +\infty)$  where  $\xi$  is the characteristic function of the interval  $[0, 1]$ . We also choose  $\lambda, \mu \in \mathbb{R}$  such that the function

$$(2.21) \quad q(s) = s^{n+2} - \lambda s^n + \mu$$

satisfies  $q(a) = q(a+1) = 0$ . Since the derivative  $q'(s)$  has at most one zero in  $[0, +\infty)$  we conclude that  $q(s) < 0$  in  $(a, a+1)$  and  $q(s) > 0$  in  $[0, +\infty) \setminus (a, a+1)$  (and also  $\lambda, \mu \geq 0$ ). Thus letting  $\chi(s)$  denote the characteristic function of the interval  $[a, a+1]$ , the assumptions on  $h$  imply that

$$(2.22) \quad q(s)(\chi(s) - h(s)) \leq 0 \text{ on } [0, +\infty).$$

Integrating the inequality (2.22), taking into account the choice of  $a$  and using (2.19) we have

$$(2.23) \quad (n + 2)B \geq \int_0^\infty s^{n+2}h(s)ds \geq \int_a^{a+1} s^{n+2}ds.$$

To estimate the last integral we take  $\tau > 0$  to be chosen later and integrate the inequality

$$(2.24) \quad ns^{n+2} - (n + 2)\tau^2s^n + 2\tau^{n+2} \geq 2\tau^n(s - \tau)^2$$

(that can be proved by dividing the left-hand side by  $(s - \tau)^2$  over  $[a, a + 1]$  to get, also using (2.23), that

$$(2.25) \quad \begin{aligned} n(n + 2)B - (n + 2)\tau^2nA + 2\tau^{n+2} &\geq 2\tau^n \int_a^{a+1} (s - \tau)^2ds \geq \\ &\geq 2\tau^n \int_{-1/2}^{1/2} t^2dt = \frac{\tau^n}{6}. \end{aligned}$$

Now choosing  $\tau = (nA)^{1/n}$  we get

$$(2.26) \quad B \geq \frac{1}{n + 2}(nA)^{\frac{n+2}{n}} + \frac{A}{6(n + 2)}$$

and this completes the proof of the Lemma. □

To complete the proof of Theorem 1 we apply Lemma 1 to the function  $\phi$  with  $A = (n\omega_n)^{-1}k$ ,  $\rho = 2(2\pi)^{-n}\sqrt{V(\Omega)I(\Omega)}$  and get in view of (2.14) that

$$(2.27) \quad \sum_{j=1}^k \lambda_j(\Omega) \geq \frac{n}{n + 2}\omega_n^{-\frac{2}{n}}k^{\frac{n+2}{n}}\phi(0)^{-\frac{2}{n}} + \frac{ck\phi(0)^2}{(n + 2)\rho^2}$$

where  $c$  is any constant such that  $0 < c < \frac{1}{6}$ .

Now observe that  $0 < \phi(0) \leq (2\pi)^{-n}V(\Omega)$  and that if  $R$  is such that  $\omega_nR^n = V(\Omega)$ , then  $I(\Omega) \geq \int_{B(R)} |x|^2 dx = \frac{n\omega_nR^{n+2}}{n + 2}$  and so

$$(2.28) \quad \rho \geq 2(2\pi)^{-n}\sqrt{\frac{n}{n + 2}\omega_n^{-\frac{2}{n}}V(\Omega)^{\frac{n+2}{n}+1}} \geq (2\pi)^{-n}\omega_n^{-\frac{1}{n}}V(\Omega)^{\frac{n+1}{n}}.$$

On the other hand the function  $g(t) = \frac{n}{n + 2}\omega_n^{-\frac{2}{n}}k^{\frac{n+2}{n}}t^{-\frac{2}{n}} + \frac{ckt^2}{(n + 2)\rho^2}$  would be decreasing on  $(0, (2\pi)^{-n}V(\Omega)]$  if  $g'((2\pi)^{-n}V(\Omega)) \leq 0$  which in view of (2.28) and since  $k \geq 1$  will be satisfied if

$$(2.29) \quad c < (2\pi)^2\omega_n^{-\frac{4}{n}}.$$

It is easy to see that we can choose  $c$  independent of  $n$  that satisfies (2.29). Then we can replace  $\phi(0)$  by  $(2\pi)^{-n}V(\Omega)$  in (2.27) which gives inequality (1.6) and so completes the proof of Theorem 1.

## REFERENCES

- [1] P.Kröger: Estimates for sums of Eigenvalues of the Laplacian, *Jour. Funct. Anal.* **126** (1994), 217-227. MR **95j**:58173
- [2] P.Li, S.T.Yau: On the Schrödinger equation and the eigenvalue problem, *Comm. Math. Phys.* **88** (1983), 309-318. MR **84k**:58225
- [3] E.Lieb: The number of bound states of one-body Schrödinger operators and the Weyl problem, *Proc. Sym. Pure Math.* **36** (1980), 241-252. MR **82i**:35134
- [4] G.Pólya: On the eigenvalues of vibrating membranes, *Proc. London Math. Soc.* (3) **11** (1961), 419-433. MR **23**:B2256
- [5] B.Simon: Weak trace ideals and the number of bound states of Schrödinger operators, *Trans. Amer. Math. Soc.* **224** (1976), 367-380. MR **54**:11109
- [6] R.S.Strichartz: Estimates for sums of eigenvalues for domains in homogeneous spaces, *Jour. Funct. Anal.* **137** (1996), 152-190. MR **97g**:58172

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