

DISCRETE APPROXIMATION OF THE POISSON-VLASOV SYSTEM *

BY

JACK SCHAEFFER

Carnegie-Mellon University

Abstract. Solutions of the Poisson–Vlasov system of equations that have spherical symmetry are considered. A numerical scheme which exploits the symmetry is proposed and is shown to converge in a pointwise sense. Prior convergence results due to Neunzert conclude convergence in only a weak topology although they require no symmetry.

The motion of a continuous mass (charge) distribution acting under the influence of its own gravitational (electric) field without collisions is described by the Poisson–Vlasov system of equations:

$$\psi_t + v \cdot \nabla_x \psi + F(x, t) \cdot \nabla_v \psi = 0, \quad (1)$$

$$F(x, t) = -\nabla_x U(x, t), \quad (2)$$

$$U(x, t) = \gamma \int \frac{\sigma(y, t)}{|x - y|} dy, \quad (3)$$

$$\sigma(x, t) = \int \psi(x, v, t) dv. \quad (4)$$

Here $\psi = \psi(x, v, t)$ ($x \in \mathbf{R}^3$, $v \in \mathbf{R}^3$, $t \geq 0$) is a scalar function describing the density in phase space, σ is the mass (charge) density, and F is the gravitational (electric) field. The constant γ is -1 in the gravitational case and $+1$ in the electrical case.

The Cauchy problem is to solve (1), (2), (3), and (4) with the initial condition

$$\psi(x, v, 0) = \dot{\psi}(x, v) \quad (5)$$

for all $t > 0$. Local existence of classical solutions was established by Kurth [4]. Batt [2] showed global existence in the gravitational case for spherically symmetric data (i.e., $\dot{\psi}$ depends only on $|x|$, $|v|$, and $x \cdot v$). Wollman [6] extended this result to the electrical case, $\gamma = +1$. Horst [3] further extended these results to include data with cylindrical symmetry (i.e., $\dot{\psi}$ depends only on $\sqrt{x_1^2 + x_2^2}$, $\sqrt{v_1^2 + v_2^2}$, $x_1 v_1 + v_2 v_2$, x_3 , and v_3). Recently, Bardos

*Received March 11, 1985.

and Degond [1] showed global existence for data that is sufficiently small in a certain norm but need not have any symmetry. The general existence problem is still open.

A natural way to generate a discrete approximation to a solution of the above Cauchy problem is to approximate the data function $\dot{\psi}$ by a large number of point masses and then follow the time evolution of the resulting many-body problem; this is called the particle-in-cell method. There is no hope for obtaining pointwise convergence for such a method since the initial error of approximating the data by point masses immediately introduces a large pointwise error. However, Neunzert [5] has established the convergence of the particle-in-cell method in a weak topology under quite general conditions. Here we consider smooth data with spherical symmetry and propose a discrete model of the continuous system, (1)–(5). In order to preserve the symmetry, the basic unit of approximation will be a spherical shell of mass (charge) rather than a point mass. Preserving the symmetry permits better estimates and this way pointwise convergence is established.

Preliminaries. We will consider only data $\dot{\psi}$ of the form $\dot{\psi}(x, v) = \dot{\phi}(|x|, |v|, \cos^{-1}(\frac{x \cdot v}{|x||v|}))$ where $\dot{\phi} \in C_0^1((0, \infty) \times (0, \infty) \times (0, \pi))$. Hence it is natural to use the variables $r = |x|$, $u = |v|$, and $\alpha = \cos^{-1}(x \cdot v/|x||v|)$. Batt showed [2] that the resulting solution ψ is continuously differentiable in x , v , and t for all $t > 0$ and that ψ depends only on r , u , α , and t (i.e., the time evolution preserves the symmetry). We will write $\psi(x, v, t) = \phi(r, u, \alpha, t)$ and note that

$$U(x, t) = \gamma \int_0^\infty \frac{\rho(\eta, t)}{\max(r, \eta)} 4\pi \eta^2 d\eta \quad (6)$$

and

$$F(x, t) = \gamma M(r, t) r^{-3} x \quad (7)$$

where

$$M(r, t) = \int_0^r \rho(\eta, t) 4\pi \eta^2 d\eta \quad \text{and} \quad \rho(r, t) = \sigma(x, t). \quad (8)$$

Given $(r, u, \alpha) \in (0, \infty) \times (0, \infty) \times (0, \pi)$ the characteristic associated with this point is given by $z = z(r, u, \alpha, t)$, $q = q(r, u, \alpha, t)$, and $\theta = \theta(r, u, \alpha, t)$, where

$$z_t = q \cos \theta \quad \text{and} \quad z(r, u, \alpha, 0) = r, \quad (9)$$

$$q_t = \gamma M(z, t) z^{-2} \cos \theta \quad \text{and} \quad q(r, u, \alpha, 0) = u, \quad (10)$$

$$\theta_t = (-\gamma M(z, t) z^{-2} q^{-1} - qz^{-1}) \sin \theta \quad \text{and} \quad \theta(r, u, \alpha, 0) = \alpha. \quad (11)$$

Note that $zq \sin \theta$ is independent of t , so we also have

$$z_{tt} - (ru \sin \alpha)^2 z^{-3} - \gamma M(z, t) z^{-2} = 0, \quad (12)$$

$z(r, u, \alpha, 0) = r$, and $z_t(r, u, \alpha, 0) = u \cos \alpha$.

Let \mathcal{S} be the support of $\dot{\phi}$ and recall that \mathcal{S} is assumed to be a compact subset of $(0, \infty) \times (0, \infty) \times (0, \pi)$. Hence there exist positive constants R_0 , l_0 , and L_0 such that if $(r, u, \alpha) \in \mathcal{S}$ then $r \leq R_0$ and $l_0 \leq ru \sin \alpha \leq L_0$. A consequence of work of Batt [2] and Wollman [6] is that there exists a positive constant U_0 (depending only on $\dot{\phi}$) such that

$q(r, u, \alpha, t) \leq U_0$ for all $(r, u, \alpha) \in \mathcal{S}$ and $t \geq 0$. Hence $z(r, u, \alpha, t) \leq R_0 + U_0 t$ if $(r, u, \alpha) \in \mathcal{S}$. Also since $zq \sin \theta$ is independent of t , $l_0 \leq ru \sin \alpha = zq \sin \theta \leq zq \leq zU_0$ so $l_0 U_0^{-1} \leq z(r, u, \alpha, t)$ for $(r, u, \alpha) \in \mathcal{S}$ and $t \geq 0$. ϕ is constant on characteristics so $\|\phi(\cdot, \cdot, \cdot, t)\|_\infty = \|\dot{\phi}\|_\infty$; hence $\rho(r, t) = \int_0^{U_0} \int_0^\pi \phi(r, u, \alpha, t) 2\pi \sin \alpha d\alpha u^2 du \leq 2\pi^2 U_0^3 \|\dot{\phi}\|_\infty$. To summarize, there exist positive constants r_0 , R_0 , l_0 , L_0 , U_0 , and ρ_0 (depending only on $\dot{\phi}$) such that

$$r_0 \leq z(r, u, \alpha, t) \leq R_0 + U_0 t, \quad (13)$$

$$l_0 \leq ru \sin \alpha = zq \sin \theta \leq L_0, \quad (14)$$

$$q(r, u, \alpha, t) \leq U_0 \quad (15)$$

for all $(r, u, \alpha, t) \in \mathcal{S} \times [0, \infty)$ and

$$\rho(r, t) \leq \rho_0 \quad \text{for all } r \geq 0 \text{ and } t \geq 0. \quad (16)$$

The letter C will denote a positive generic constant depending only on $\dot{\phi}$, and $C(t)$ will denote an increasing positive function of t depending only on $\dot{\phi}$.

Another fact that will be used is that the time evolution preserves measure in \mathbf{R}^6 . If the characteristics of the original system are given by $X(x_0, v_0, t)$ and $V(x_0, v_0, t)$, where $X_t = V$, $X(x_0, v_0, 0) = x_0$ and $V_t = F(X, t)$, $V(x_0, v_0, 0) = v_0$, then this says

$$\begin{aligned} & \int \int f(r, u, \alpha) \phi(r, u, \alpha, t) dv dx \\ &= \int \int f\left(|x|, |v|, \cos^{-1}\left(\frac{x \cdot v}{|x||v|}\right)\right) \psi(x, v, t) dv dx \\ &= \int \int f\left(|X|, |V|, \cos^{-1}\left(\frac{X \cdot V}{|X||V|}\right)\right) \psi(x_0, v_0, 0) dv_0 dx_0 \\ &= \int \int f(z, q, \theta) \dot{\phi}(r, u, \alpha) dv_0 dx_0 \end{aligned} \quad (17)$$

for all $f \in L^\infty((0, \infty) \times (0, \infty) \times (0, \pi))$. If f is taken to be identically equal to one, we get

$$\int \int \phi(r, u, \alpha, t) dv dx = \int \int \dot{\phi}(r, u, \alpha) dv_0 dx_0,$$

which is mass conservation. Take $M = \int \int \dot{\phi}(r, u, \alpha) dv_0 dx_0$ so that

$$M = \int \int \phi(r, u, \alpha, t) dv dx = \int_0^\infty \rho(r, t) 4\pi r^2 dr. \quad (18)$$

Another conserved quantity is the energy

$$\begin{aligned} E &= \int \int \psi(x, v, t) v^2 dv dx + \int \sigma(x, t) U(x, t) dx \\ &= \int \int \phi(r, u, \alpha, t) u^2 dv dx + \gamma \int \rho(r, t) \int_0^\infty \frac{\rho(\eta, t) 4\pi \eta^2}{\max(r, \eta)} d\eta dx. \end{aligned} \quad (19)$$

Before introducing the discrete model let us consider the following continuous dependence lemma which involves only the continuous solution.

LEMMA I. There exists an increasing function $C_1(t)$ with $C_1(0) \geq 1$ (depending only on ϕ) such that for all $t \geq 0$, $(r, u, \alpha) \in \mathcal{S}$, and $(\eta, w, \beta) \in \mathcal{S}$

$$\begin{aligned} & |z(r, u, \alpha, t) - z(\eta, w, \beta, t)| + |z_t(r, u, \alpha, t) - z_t(\eta, w, \beta, t)| \\ & \leq C_1(t)|(r, u, \alpha) - (\eta, w, \beta)|. \end{aligned} \quad (20)$$

This holds for $\gamma = 1$ and $\gamma = -1$.

Proof. For brevity we will temporarily use the notation $z_1(t) = z(r, u, \alpha, t)$, $z_2(t) = z(\eta, w, \beta, t)$, and $Q(t) = \sup\{|z_1(\tau) - z_2(\tau)| + |z'_1(\tau) - z'_2(\tau)|: 0 \leq \tau \leq t\}$.

From equation (12) for $t > 0$

$$\begin{aligned} |z''_1(t) - z''_2(t)| &= |(ru \sin \alpha)^2 z_1^{-3}(t) + \gamma M(z_1(t), t) z_1^{-2}(t)| \\ & \quad - |(\eta w \sin \beta)^2 z_2^{-3}(t) - \gamma M(z_2(t), t) z_2^{-2}(t)| \\ & \leq |(ru \sin \alpha)^2 - (\eta w \sin \beta)^2| z_1^{-3}(t) + (\eta w \sin \beta)^2 |z_1^{-3}(t) - z_2^{-3}(t)| \\ & \quad + |M(z_1(t), t) - M(z_2(t), t)| z_1^{-2}(t) + M(z_2(t), t) |z_1^{-2}(t) - z_2^{-2}(t)| \\ & \leq (ru \sin \alpha + \eta w \sin \beta) |ru \sin \alpha - \eta w \sin \beta| r_0^{-3} \\ & \quad + L_0^2 z_1^{-3}(t) z_2^{-3}(t) |z_2^3(t) - z_1^3(t)| + |M(z_1(t), t) - M(z_2(t), t)| r_0^{-2} \\ & \quad + M z_1^{-2}(t) z_2^{-2}(t) |z_2^2(t) - z_1^2(t)| \\ & \leq 2L_0 r_0^{-3} |ru \sin \alpha - \eta w \sin \beta| \\ & \quad + L_0 r_0^{-6} (z_1^2(t) + z_1(t) z_2(t) + z_2^2(t)) |z_1(t) - z_2(t)| \\ & \quad + r_0^{-2} |M(z_1(t), t) - M(z_2(t), t)| + M r_0^{-4} (z_1(t) + z_2(t)) |z_1(t) - z_2(t)|. \end{aligned} \quad (21)$$

Now note that

$$\begin{aligned} |ru \sin \alpha - \eta w \sin \beta| &\leq ru |\sin \alpha - \sin \beta| + r|u - w| \sin \beta + |r - \eta| w \sin \beta \\ &\leq L_0 |\alpha - \beta| + R_0 |u - w| + U_0 |r - \eta| \leq C|(r, u, \alpha) - (\eta, w, \beta)| \end{aligned}$$

and

$$\begin{aligned} |M(z_1(t), t) - M(z_2(t), t)| &= \left| \int_{z_1(t)}^{z_2(t)} 4\pi s^2 \rho(s, t) ds \right| \\ &\leq 4\pi (R_0 + U_0 t)^2 \rho_0 |z_1(t) - z_2(t)|. \end{aligned}$$

Using these in (21) we have

$$\begin{aligned} |z''_1(t) - z''_2(t)| &\leq 2L_0 r_0^{-3} C |(r, u, \alpha) - (\eta, w, \beta)| + L_0 r_0^{-6} 3(R_0 + U_0 t)^2 |z_1(t) - z_2(t)| \\ & \quad + r_0^{-2} 4\pi (R_0 + U_0 t)^2 \rho_0 |z_1(t) - z_2(t)| + M r_0^{-4} 2(R_0 + U_0 t) |z_1(t) - z_2(t)| \\ & = C |(r, u, \alpha) - (\eta, w, \beta)| + C(t) |z_1(t) - z_2(t)|. \end{aligned} \quad (22)$$

By (22) we have for $t \geq 0$

$$\begin{aligned}
 |z_1(t) - z_2(t)| + |z_1'(t) - z_2'(t)| & \\
 & \leq Q(0) + \int_0^t (|z_1'(\tau) - z_2'(\tau)| + |z_1''(\tau) - z_2''(\tau)|) d\tau \\
 & \leq Q(0) + \int_0^t (Q(\tau) + C|(r, u, \alpha) - (\eta, w, \beta)| + C(\tau)Q(\tau)) d\tau \\
 & \leq |r - u| + |u \cos \alpha - w \cos \beta| + Ct|(r, u, \alpha) - (\eta, w, \beta)| \\
 & \quad + (1 + C(t)) \int_0^t Q(\tau) d\tau \\
 & \leq C(t)|(r, u, \alpha) - (\eta, w, \beta)| + C(t) \int_0^t Q(\tau) d\tau.
 \end{aligned}$$

Now for $0 \leq \tau \leq t$,

$$|z_1(\tau) - z_2(\tau)| + |z_1'(\tau) - z_2'(\tau)| \leq C(t)|(r, u, \alpha) - (\eta, w, \beta)| + C(t) \int_0^t Q(s) ds$$

so

$$Q(t) \leq C(t)|(r, u, \alpha) - (\eta, w, \beta)| + C(t) \int_0^t Q(s) ds.$$

Since $C(t)$ is increasing we have

$$Q(t) \leq C(T)|(r, u, \alpha) - (\eta, w, \beta)| + C(T) \int_0^t Q(s) ds$$

for all $t \in [0, T]$ and $T > 0$. By Gronwall's inequality

$$Q(t) \leq C(T)|(r, u, \alpha) - (\eta, w, \beta)| \exp(C(T)t)$$

for $t \in [0, T]$ so $Q(T) \leq C(T)|(r, u, \alpha) - (\eta, w, \beta)|$. This completes the proof of Lemma I.

The discrete model. To discretize the continuous mass distribution $\dot{\phi}$ assume that \mathcal{S} is the union of disjoint connected sets S_1, \dots, S_N . Let

$$\delta = \max_i \sup \left\{ \sqrt{(r - \bar{r})^2 + (u - \bar{u})^2 + (\alpha - \bar{\alpha})^2} : (r, u, \alpha), (\bar{r}, \bar{u}, \bar{\alpha}) \in S_i \right\}$$

and assume that $\delta < \min\{1, r_0\}$. For each $i \in \{1, 2, \dots, N\}$, choose $(r_i, u_i, \alpha_i) \in S_i$ and define $L_i = r_i u_i \sin \alpha_i$, $M_i = \iint_{S_i} \dot{\phi}(r, u, \alpha) dv dx$, and $z_i(t) = z(r_i, u_i, \alpha_i, t)$.

From Eq. (12) $z_i''(t) - L_i^2 z_i^{-3}(t) - \gamma \bar{M}(z_i(t), t) z_i^{-2}(t) = 0$ with $z_i(0) = r_i$ and $z_i'(0) = u_i \cos \alpha_i$, so for the discrete model we define the quantities $\bar{z}_1(t), \bar{z}_2(t), \dots, \bar{z}_N(t)$ (which approximate $z_1(t), z_2(t), \dots, z_N(t)$) and $\bar{M}(r, t)$ by the system of differential equations:

$$\bar{z}_i'' - L_i^2 (\bar{z}_i)^{-3} - \gamma \bar{M}(\bar{z}_i, t) (\bar{z}_i)^{-2} = 0 \quad (23)$$

with $\bar{z}_i(0) = r_i$ and $\bar{z}_i'(0) = u_i \cos \alpha_i$ for $i = 1, 2, \dots, N$ and

$$\bar{M}(r, t) = \sum_{i=1}^N M_i \xi(r - \bar{z}_i(t)), \quad (24)$$

where we define $\xi: \mathbf{R} \rightarrow \mathbf{R}$ by

$$\xi(r) = \begin{cases} 0 & \text{if } r \leq 0, \\ r\delta^{-1} & \text{if } 0 < r < \delta, \\ 1 & \text{if } \delta \leq r. \end{cases} \quad (25)$$

$\bar{M}(r, t)$ is Lipschitz continuous in r for $r \geq 0$ and $r_0 \leq \bar{z}_i(0) \leq R_0$ for each i so there exists $T > 0$ such that a unique solution to the system (23) and (24) exists for $0 \leq t < T$ and satisfies

$$\frac{1}{2}r_0 \leq \bar{z}_i(t) \leq 2(R_0 + U_0t) \quad (26)$$

for each i and for $0 \leq t < T$. We will take T to be the maximal such time and note that T is heavily dependent on the discretization; however it will be shown later that T tends to infinity as δ tends to zero.

THEOREM I. Let $\dot{\phi} \in C_0^1((0, \infty) \times (0, \infty) \times (0, \pi))$, $z(r, u, \alpha, t)$ as in (9), (10), and (11), $M(r, t)$ as in (8), and $\bar{z}_i(t)$, $\bar{M}(r, t)$ as above. There exists a positive increasing function $C_2(t)$ (which is dependent only on $\dot{\phi}$ and not on δ) such that for all $i \in \{1, 2, \dots, N\}$, $(r, u, \alpha) \in S_i$, and $t \in [0, T)$:

$$|z(r, u, \alpha, t) - \bar{z}_i(t)| + |z_i(r, u, \alpha, t) - \bar{z}_i'(t)| \leq C_2(t)\delta \quad (27)$$

and

$$|M(\eta, t) - \bar{M}(\eta, t)| \leq C_2(t)\delta \quad \text{for all } \eta \geq 0. \quad (28)$$

This holds for $\gamma = 1$ and $\gamma = -1$.

Proof. Define

$$\bar{\bar{M}}(r, t) = \sum_{i=1}^N M_i \xi(r - z_i(t)) \quad (29)$$

for $r \geq 0$ and $t \geq 0$. The proof is divided into parts A, B, and C. Part A will estimate $M - \bar{M}$ and part B will estimate $\bar{M} - \bar{\bar{M}}$.

Part A. Fix $t > 0$ and $(\eta, w, \beta) \in \mathcal{S}$. Choose i so that $(\eta, w, \beta) \in S_i$ and define $I = \{j: z_j(t) \leq z(\eta, w, \beta, t)\}$. Note that if $(r, u, \alpha) \in S_j$, then by Lemma I $|z(r, u, \alpha, t) - z_j(t)| \leq C_1(t)\delta$. Hence, if $(r, u, \alpha) \in S_j$ and $j \in I$,

$$z(r, u, \alpha, t) \leq z_j(t) + C_1(t)\delta \leq z(\eta, w, \beta, t) + C_1(t)\delta,$$

that is, if $j \in I$, then

$$S_j \subset \{(r, u, \alpha): z(r, u, \alpha, t) \leq z(\eta, w, \beta, t) + C_1(t)\delta\}$$

and hence

$$\bigcup_{j \in I} S_j \subset \{(r, u, \alpha): z(r, u, \alpha, t) \leq z(\eta, w, \beta, t) + C_1(t)\delta\}. \quad (30)$$

Now denoting $z = z(\eta, w, \beta, t)$,

$$\bar{\bar{M}}(z, t) = \sum_{j=1}^N M_j \xi(z - z_j(t)) \leq \sum_{j \in I} M_j$$

since $j \notin I$ implies $z - z_j(t) \leq 0$ which implies $\xi(z - z_j(t)) = 0$ by (25). Also by (30)

$$\begin{aligned} \sum_{j \in I} M_j &= \iint_{\cup_{j \in I} S_j} \dot{\phi}(r, u, \alpha) \, dv \, dx \\ &\leq \iint_{\{(r, u, \alpha): z(r, u, \alpha, t) \leq z + C_1(t)\delta\}} \dot{\phi}(r, u, \alpha) \, dv \, dx \\ &= \iint_{\{r \leq z + C_1(t)\delta\}} \phi(r, u, \alpha, t) \, dv \, dx = M(z + C_1(t)\delta, t) \\ &= M(z, t) + \int_z^{z + C_1(t)\delta} 4\pi s^2 \rho(s, t) \, ds \leq M(z, t) + C(t)\delta. \end{aligned}$$

Therefore

$$\overline{\overline{M}}(z, t) \leq M(z, t) + C(t)\delta. \tag{31}$$

To bound $\overline{\overline{M}}$ from below define $J = \{j: z_j(t) \leq z - C_1(t)\delta\}$, where $z = z(\eta, w, \beta, t)$ as before. If $(r, u, \alpha) \in S_j$ and $z(r, u, \alpha, t) \leq z - 2C_1(t)\delta$, then by Lemma I, $z_j(t) \leq z(r, u, \alpha, t) + C_1(t)\delta \leq z - C_1(t)\delta$, that is $j \in J$. Therefore,

$$\mathcal{S} \cap \{(r, u, \alpha): z(r, u, \alpha, t) \leq z - 2C_1(t)\delta\} \subset \bigcup_{j \in J} S_j.$$

If $j \in J$ then $z_j(t) \leq z - C_1(t)\delta \leq z - \delta$, so $z - z_j(t) \geq \delta$, and hence $\xi(z - z_j(t)) = 1$. Therefore

$$\begin{aligned} \overline{\overline{M}}(z, t) &= \sum_{j=1}^N M_j \xi(z - z_j(t)) \geq \sum_{j \in J} M_j = \iint_{\cup_{j \in J} S_j} \dot{\phi}(r, u, \alpha) \, dv \, dx \\ &\geq \iint_{\mathcal{S} \cap \{(r, u, \alpha): z(r, u, \alpha, t) \leq z - 2C_1(t)\delta\}} \dot{\phi}(r, u, \alpha) \, dv \, dx \\ &= \iint_{\{(r, u, \alpha): z(r, u, \alpha, t) \leq z - 2C_1(t)\delta\}} \dot{\phi}(r, u, \alpha) \, dv \, dx \\ &= \iint_{\{r \leq z - 2C_1(t)\delta\}} \phi(r, u, \alpha, t) \, dv \, dx \\ &= M(\max\{0, z - 2C_1(t)\delta\}, t) \\ &= M(z, t) - \int_{\max\{0, z - 2C_1(t)\delta\}}^z 4\pi s^2 \rho(s, t) \, ds \\ &\geq M(z, t) - C(t)\delta. \end{aligned}$$

Combining this and (31) we have

$$|M(z, t) - \overline{\overline{M}}(z, t)| \leq C(t)\delta \tag{32}$$

for $z = z(\eta, w, \beta, t)$ and $(\eta, w, \beta) \in \mathcal{S}$.

We will need (32) for all z , not just for $z \in z(\mathcal{S}, t)$. Note that if $r_1 \geq R_0 + U_0 t + \delta$ then $M(r_1, t) = \overline{\overline{M}}(r_1, t) = M$, and if $0 \leq r_1 \leq r_0$ then $M(r_1, t) = \overline{\overline{M}}(r_1, t) = 0$. Hence we consider $t \geq 0$ and $r_1 \in [r_0, R_0 + U_0 t + 1]$, such that $r_1 \notin z(\mathcal{S}, t)$. Define $r_2 = \inf\{\eta: \eta \geq r_0 \text{ and } (\eta, r_1) \cap z(\mathcal{S}, t) = \emptyset\}$ and note that $r_0 \leq r_2 \leq r_1$ and $(r_2, r_1) \cap z(\mathcal{S}, t) = \emptyset$.

We have

$$\begin{aligned}
M(r_1, t) - M(r_2, t) &= \int_{r_2}^{r_1} 4\pi s^2 \rho(s, t) ds \\
&= \iint_{\{(x, v): |x| \in (r_2, r_1)\}} \phi(r, u, \alpha, t) dv dx \\
&= \iint_{\{(\eta, w, \beta): z(\eta, w, \beta, t) \in (r_2, r_1)\}} \mathring{\phi}(\eta, w, \beta) dv dx \\
&= 0 \text{ since } \{(\eta, w, \beta): z(\eta, w, \beta, t) \in (r_2, r_1)\} \cap \mathcal{S} \quad (33)
\end{aligned}$$

is empty.

$z_j(t) \notin (r_2, r_1)$ for all j so by (25)

$$\begin{aligned}
\bar{M}(r_1, t) - \bar{M}(r_2, t) &= \sum_{j=1}^N M_j(\xi(r_1 - z_j(t)) - \xi(r_2 - z_j(t))) \\
&\leq \sum_{\{j: z_j(t) \in (r_2 - \delta, r_2]\}} M_j.
\end{aligned}$$

If $z_j(t) \in (r_2 - \delta, r_2]$ then by Lemma I

$$\begin{aligned}
z(S_j, t) &\subset [z_j(t) - C_1(t)\delta, z_j(t) + C_1(t)\delta] \\
&\subset [r_2 - \delta - C_1(t)\delta, r_2 + C_1(t)\delta]
\end{aligned}$$

so

$$\begin{aligned}
0 &\leq \bar{M}(r_1, t) - \bar{M}(r_2, t) \\
&\leq \iint_{\substack{\cup S_j \\ j: z_j(t) \in (r_2 - \delta, r_2]}} \mathring{\phi}(r, u, \alpha) dv dx \\
&\leq \iint_{\{(r, u, \alpha): |z(r, u, \alpha, t) - r_2| \leq \delta + C_1(t)\delta\}} \mathring{\phi}(r, u, \alpha) dv dx \\
&= \iint_{\{(x, v): ||x| - r_2| \leq \delta + C_1(t)\delta\}} \phi(r, u, \alpha, t) dv dx \\
&= \int_{\max\{0, r_2 - \delta - C_1(t)\delta\}}^{r_2 + \delta + C_1(t)\delta} 4\pi s^2 \rho(s, t) ds \\
&\leq 4\pi (r_2 + 1 + C_1(t))^2 \rho_0 2(1 + C_1(t))\delta \\
&\leq C(t)\delta, \text{ since } r_2 \leq R_0 + U_0 t + 1. \quad (34)
\end{aligned}$$

Recall that by definition $r_2 = \inf\{\eta: \eta \geq r_0 \text{ and } (\eta, r_1) \cap z(\mathcal{S}, t) = \emptyset\}$. The map $(r, u, \alpha) \rightarrow z(r, u, \alpha, t)$ is continuous by Lemma I and \mathcal{S} is compact; hence $z(\mathcal{S}, t)$ is compact. If $r_2 > r_0$ then $r_2 \in z(\mathcal{S}, t)$ and so by (32)

$$|M(r_2, t) - \bar{M}(r_2, t)| \leq C(t)\delta. \quad (35)$$

If $r_2 = r_0$ then $M(r_2, t) = \bar{\bar{M}}(r_2, t) = 0$ so (35) still holds. Now combining (33), (34), and (35) we have

$$\begin{aligned} & |M(r_1, t) - \bar{\bar{M}}(r_1, t)| \\ & \leq |M(r_1, t) - M(r_2, t)| + |M(r_2, t) - \bar{\bar{M}}(r_2, t)| + |\bar{\bar{M}}(r_2, t) - \bar{\bar{M}}(r_1, t)| \\ & \leq 0 + C(t)\delta + C(t)\delta. \end{aligned}$$

Therefore

$$|M(r, t) - \bar{\bar{M}}(r, t)| \leq C(t)\delta \quad \text{for all } r \geq 0 \text{ and } t \geq 0. \quad (36)$$

Part B. In order to estimate $\bar{M} - \bar{\bar{M}}$ let us define for $0 \leq t < T$

$$\|z - \bar{z}\|(t) = \sup_{0 \leq \tau \leq t} \max_{1 \leq i \leq N} |z_i(\tau) - \bar{z}_i(\tau)|. \quad (37)$$

Let $r \geq 0$ and $t \in [0, T)$. If $z_j(t) \geq r + \|z - \bar{z}\|(t)$,

$$r - z_j(t) \leq -\|z - \bar{z}\|(t) \leq 0,$$

and

$$r - \bar{z}_j(t) \leq r - (z_j(t) - \|z - \bar{z}\|(t)) \leq 0,$$

so

$$\xi(r - z_j(t)) = \xi(r - \bar{z}_j(t)) = 0.$$

Similarly, if $z_j(t) \leq r - \|z - \bar{z}\|(t) - \delta$, then

$$r - z_j(t) \geq \|z - \bar{z}\|(t) + \delta \geq \delta$$

and

$$r - \bar{z}_j(t) \geq r - (z_j(t) + \|z - \bar{z}\|(t)) \geq \delta,$$

so

$$\xi(r - z_j(t)) = \xi(r - \bar{z}_j(t)) = 1.$$

Hence

$$\begin{aligned} |\bar{M}(r, t) - \bar{\bar{M}}(r, t)| &= \left| \sum_{j=1}^N M_j(\xi(r - \bar{z}_j(t)) - \xi(r - z_j(t))) \right| \\ &\leq \sum_{j \in K} M_j \end{aligned} \quad (38)$$

where

$$K = \left\{ j: |r - z_j(t)| < \|z - \bar{z}\|(t) + \delta \right\}.$$

Now if $j \in K$ and $(\eta, w, \beta) \in S_j$ then

$$|z_j(t) - r| < \|z - \bar{z}\|(t) + \delta$$

and by Lemma I

$$|z(\eta, w, \beta, t) - z_j(t)| \leq C_1(t)\delta$$

so

$$|z(\eta, w, \beta, t) - r| < \|z - \bar{z}\|(t) + \delta + C_1(t)\delta.$$

Therefore

$$\bigcup_{j \in K} S_j \subset \{(\eta, w, \beta) : |z(\eta, w, \beta, t) - r| < \varepsilon\}$$

where $\varepsilon = \|z - \bar{z}\|(t) + \delta + C_1(t)\delta$. Now by (38) we have

$$\begin{aligned} |\bar{M}(r, t) - \overline{\bar{M}}(r, t)| &\leq \sum_{j \in K} M_j \\ &= \iint_{\bigcup_{j \in K} S_j} \dot{\phi}(\eta, w, \beta) \, dv \, dx \\ &\leq \iint_{\{(\eta, w, \beta) : |z(\eta, w, \beta, t) - r| < \varepsilon\}} \dot{\phi}(\eta, w, \beta) \, dv \, dx \\ &= \iint_{\{(x, v) : \|x\| - r < \varepsilon\}} \psi(x, v, t) \, dv \, dx \\ &= \int_{\max\{0, r-\varepsilon\}}^{r+\varepsilon} 4\pi s^2 \rho(s, t) \, ds \leq 4\pi(r + \varepsilon)^2 \rho_0 2\varepsilon. \end{aligned} \quad (39)$$

Recall that $\varepsilon = \|z - \bar{z}\|(t) + \delta + C_1(t)\delta$. Now $|z_j(\tau) - \bar{z}_j(\tau)| \leq |z_j(\tau)| + |\bar{z}_j(\tau)| \leq 3(R_0 + U_0 t)$ for all $\tau \in [0, t]$ so $\|z - \bar{z}\|(t) \leq 3(R_0 + U_0 t)$ and $\varepsilon \leq 3(R_0 + U_0 t) + 1 + C_1(t)$. If $r \leq 2R_0 + 2U_0 t + 1$ we have by (39),

$$\begin{aligned} |\bar{M}(r, t) - \overline{\bar{M}}(r, t)| &\leq 4\pi(5(R_0 + U_0 t) + 2 + C_1(t))^2 \rho_0 2\varepsilon \\ &= C(t)\varepsilon = C(t)(\|z - \bar{z}\|(t) + \delta + C_1(t)\delta) \\ &= C(t)\|z - \bar{z}\|(t) + C(t)\delta. \end{aligned}$$

If $r > 2R_0 + 2U_0 t + 1$, then $\bar{M}(r, t) = \overline{\bar{M}}(r, t) = M$ so for all $r \geq 0$ and $t \in [0, T)$

$$|\bar{M}(r, t) - \overline{\bar{M}}(r, t)| \leq C(t)\|z - \bar{z}\|(t) + C(t)\delta. \quad (40)$$

Combining (40) and (36) we get for $r \geq 0$ and $t \in [0, T)$

$$\begin{aligned} |M(r, t) - \bar{M}(r, t)| &\leq |M(r, t) - \overline{\bar{M}}(r, t)| + |\overline{\bar{M}}(r, t) - \bar{M}(r, t)| \\ &\leq C(t)\delta + C(t)\|z - \bar{z}\|(t) + C(t)\delta \\ &= C(t)\delta + C(t)\|z - \bar{z}\|(t). \end{aligned} \quad (41)$$

Part C. In this section we use (41) and the differential equations for z_i and \bar{z}_i to estimate $\|z - \bar{z}\|(t)$. From Eqs. (12) and (23)

$$\begin{aligned}
|z_i''(t) - \bar{z}_i''(t)| &= |L_i^2 z_i^{-3} + \gamma M(z_i, t) z_i^{-2} - L_i^2 (\bar{z}_i)^{-3} - \gamma \bar{M}(\bar{z}_i, t) (\bar{z}_i)^{-2}| \\
&\leq L_i^2 |z_i^{-3} - (\bar{z}_i)^{-3}| + M(z_i, t) |z_i^{-2} - (\bar{z}_i)^{-2}| \\
&\quad + (\bar{z}_i)^{-2} |M(z_i, t) - \bar{M}(\bar{z}_i, t)| \\
&\leq L_0^2 z_i^{-3} (\bar{z}_i)^{-3} |z_i^3 - (\bar{z}_i)^3| + M z_i^{-2} (\bar{z}_i)^{-2} |z_i^2 - (\bar{z}_i)^2| \\
&\quad + (\bar{z}_i)^{-2} (|M(z_i, t) - M(\bar{z}_i, t)| + |M(\bar{z}_i, t) - \bar{M}(\bar{z}_i, t)|) \quad (42)
\end{aligned}$$

for $0 \leq t < T$. Note that

$$\begin{aligned}
|M(z_i, t) - M(\bar{z}_i, t)| &= \left| \int_{\bar{z}_i}^{z_i} 4\pi s^2 \rho(s, t) ds \right| \\
&\leq 4\pi (\max\{z_i, \bar{z}_i\})^2 \rho_0 |z_i - \bar{z}_i| \leq C(t) \|z - \bar{z}\|(t)
\end{aligned}$$

and by (41)

$$|M(\bar{z}_i, t) - \bar{M}(\bar{z}_i, t)| \leq C(t)\delta + C(t)\|z - \bar{z}\|(t)$$

so

$$|M(z_i, t) - M(\bar{z}_i, t)| + |M(\bar{z}_i, t) - \bar{M}(\bar{z}_i, t)| \leq C(t)\delta + C(t)\|z - \bar{z}\|(t).$$

Using this in (42) yields for $0 \leq t < T$

$$\begin{aligned}
&|z_i''(t) - \bar{z}_i''(t)| \\
&\leq L_0^2 z_i^{-3} (\bar{z}_i)^{-3} (z_i^2 + z_i(\bar{z}_i) + (\bar{z}_i)^2) |z_i - \bar{z}_i| \\
&\quad + M z_i^{-2} (\bar{z}_i)^{-2} (z_i + \bar{z}_i) |z_i - \bar{z}_i| + (\bar{z}_i)^{-2} (C(t)\delta + C(t)\|z - \bar{z}\|(t)) \\
&\leq L_0^2 r_0^{-3} (\frac{1}{2}r_0)^{-3} 7(R_0 + U_0 t)^2 \|z - \bar{z}\|(t) \\
&\quad + M r_0^{-2} (\frac{1}{2}r_0)^{-2} 3(R_0 + U_0 t) \|z - \bar{z}\|(t) + (\frac{1}{2}r_0)^{-2} (C(t)\delta + C(t)\|z - \bar{z}\|(t)) \\
&= C(t)\delta + C(t)\|z - \bar{z}\|(t). \quad (43)
\end{aligned}$$

Define for $0 \leq t < T$

$$Q(t) = \sup_{0 \leq \tau \leq t} \max_{1 \leq i \leq N} (|z_i(\tau) - \bar{z}_i(\tau)| + |z_i'(\tau) - \bar{z}_i'(\tau)|).$$

Note that $z_i(0) = \bar{z}_i(0) = r_i$ and $z_i'(0) = \bar{z}_i'(0) = u_i \cos \alpha_i$, so by (43)

$$\begin{aligned} |z_i(t) - \bar{z}_i(t)| + |z_i'(t) - \bar{z}_i'(t)| &\leq \int_0^t (|z_i'(\tau) - \bar{z}_i'(\tau)| + |z_i''(\tau) - \bar{z}_i''(\tau)|) d\tau \\ &\leq \int_0^t (Q(\tau) + C(\tau)\delta + C(\tau)Q(\tau)) d\tau \\ &\leq C(t)\delta + C(t) \int_0^t Q(\tau) d\tau. \end{aligned}$$

Hence for all $\tau \in [0, t]$ and $i \in \{1, 2, \dots, N\}$,

$$\begin{aligned} |z_i(\tau) - \bar{z}_i(\tau)| + |z_i'(\tau) - \bar{z}_i'(\tau)| &\leq C(\tau)\delta + C(\tau) \int_0^\tau Q(s) ds \\ &\leq C(t)\delta + C(t) \int_0^t Q(s) ds \end{aligned}$$

and therefore $Q(t) \leq C(t)\delta + C(t) \int_0^t Q(s) ds$ for $0 \leq t < T$. For $\tau \in [0, t]$ and $t < T$, $Q(\tau) \leq C(t)\delta + C(t) \int_0^\tau Q(s) ds$ since $C(t)$ is increasing. By Gronwall's inequality we get for $0 \leq \tau \leq t$, $Q(\tau) \leq C(t)\delta \exp(C(t)\tau)$ so

$$Q(t) \leq C(t)\delta \quad \text{for } 0 \leq t < T. \quad (44)$$

Finally by (44) and Lemma I if $(r, u, \alpha) \in S_i$ and $0 \leq t < T$ then

$$\begin{aligned} |z(r, u, \alpha, t) - \bar{z}_i(t)| + |z_i(r, u, \alpha, t) - \bar{z}_i'(t)| \\ \leq |z(r, u, \alpha, t) - z_i(t)| + |z_i(t) - \bar{z}_i(t)| \\ + |z_i(r, u, \alpha, t) - z_i'(t)| + |z_i'(t) - \bar{z}_i'(t)| \\ \leq C_1(t)\delta + C(t)\delta = C(t)\delta. \end{aligned}$$

Also by (41) and (44) for $0 \leq r$ and $0 \leq t < T$

$$\begin{aligned} |M(r, t) - \bar{M}(r, t)| &\leq C(t)\delta + C(t)\|z - \bar{z}\|(t) \\ &\leq C(t)\delta + C(t)C(t)\delta = C(t)\delta. \end{aligned}$$

This completes the proof of Theorem I.

COMMENT. Recall that the existence time T depends on δ . We can now use Theorem I to show that T tends to infinity as δ tends to zero: Let $T_0 > 0$ and $\delta \in (0, r_0/(3C_2(T_0)))$. Then for all j and for $0 \leq t < \min(T, T_0)$ we have by (27)

$$\begin{aligned} \frac{2}{3}r_0 \leq r_0 - C_2(T_0)\delta \leq r_0 - C_2(t)\delta \leq z_j(t) - |\bar{z}_j(t) - z_j(t)| \\ \leq \bar{z}_j(t) \leq z_j(t) + |\bar{z}_j(t) - z_j(t)| \\ \leq R_0 + U_0t + C_2(t)\delta \leq R_0 + U_0t + \frac{1}{3}r_0 < \frac{3}{2}(R_0 - U_0t). \end{aligned}$$

Now if $T \leq T_0$ the solution $\bar{z}_1, \bar{z}_2, \dots, \bar{z}_N$ can be continued beyond $t = T$, which is a contradiction. Therefore $T > T_0$.

Recall from (19) that the energy is conserved; we may write

$$\begin{aligned}
E &= \int \int \phi(r, u, \alpha, t) u^2 dv dx + \gamma \int \rho(r, t) \int_0^\infty \frac{\rho(\eta, t)}{\max(r, \eta)} 4\pi\eta^2 d\eta dx \\
&= \int \int \phi(r, u, \alpha, t) \left(u^2 + \gamma \int_0^\infty \frac{\rho(\eta, t)}{\max(r, \eta)} 4\pi\eta^2 d\eta \right) dv dx \\
&= \int \int \phi(r, u, \alpha, t) \left(u^2 + \gamma \int \int \frac{\phi(r_0, u_0, \alpha_0, t)}{\max(r, r_0)} dv_0 dx_0 \right) dv dx \\
&= \int \int \dot{\phi}(r, u, \alpha) (q^2(r, u, \alpha, t)) \\
&\quad + \gamma \int \int \frac{\dot{\phi}(r_0, u_0, \alpha_0) dv_0 dx_0}{\max(z(r, u, \alpha, t), z(r_0, u_0, \alpha_0, t))} dv dx. \tag{45}
\end{aligned}$$

Also since

$$z_t(r, u, \alpha, t) = q(r, u, \alpha, t) \cos(\theta(r, u, \alpha, t))$$

and

$$\begin{aligned}
z(r, u, \alpha, t) q(r, u, \alpha, t) \sin(\theta(r, u, \alpha, t)) &= ru \sin \alpha, \\
q^2(r, u, \alpha, t) &= q^2(r, u, \alpha, t) \cos^2(\theta(r, u, \alpha, t)) + q^2(r, u, \alpha, t) \sin^2(\theta(r, u, \alpha, t)) \\
&= z_t^2(r, u, \alpha, t) + (ru \sin \alpha)^2 z^{-2}(r, u, \alpha, t). \tag{46}
\end{aligned}$$

Substituting (46) into (45) yields

$$\begin{aligned}
E &= \int \int \dot{\phi}(r, u, \alpha) \left(z_t^2(r, u, \alpha, t) \right. \\
&\quad \left. + \frac{(ru \sin \alpha)^2}{z^2(r, u, \alpha, t)} + \gamma \int \int \frac{\dot{\phi}(r_0, u_0, \alpha_0) dv_0 dx_0}{\max(z(r, u, \alpha, t), z(r_0, u_0, \alpha_0, t))} \right) dv dx. \tag{47}
\end{aligned}$$

(47) suggests a quantity which will be nearly conserved by the discrete system.

COROLLARY I. There exists a positive increasing function $C_3(t)$ (depending only on $\dot{\phi}$) such that for all $t \in [0, T)$

$$\left| E - \sum_{i=1}^N M_i \left((\bar{z}'_i(t))^2 + L_i^2 (\bar{z}_i(t))^{-2} + \sum_{j=1}^N \frac{\gamma M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right) \right| \leq C_3(t) \delta. \tag{48}$$

Proof. For brevity denote $z = z(r, u, \alpha, t)$ and $z_0 = z(r_0, u_0, \alpha_0, t)$. Note that

$$\begin{aligned}
\left| \frac{1}{\max(z, z_0)} - \frac{1}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right| &\leq \left| \frac{1}{\max(z, z_0)} - \frac{1}{\max(\bar{z}_i(t), z_0)} \right| \\
&\quad + \left| \frac{1}{\max(\bar{z}_i(t), z_0)} - \frac{1}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right| \\
&\leq \frac{|z - \bar{z}_i(t)|}{(\frac{1}{2}r_0)^2} + \frac{|z_0 - \bar{z}_j(t)|}{(\frac{1}{2}r_0)^2}. \tag{49}
\end{aligned}$$

If $(r, u, \alpha) \in S_i$ and $(r_0, u_0, \alpha_0) \in S_j$, then by Theorem I and (49)

$$\left| \frac{1}{\max(z, z_0)} - \frac{1}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right| \leq \frac{C_2(t)\delta}{(r_0/2)^2} + \frac{C_2(t)\delta}{(r_0/2)^2} = \frac{8C_2(t)\delta}{r_0^2}.$$

Therefore if $(r, u, \alpha) \in S_i$, then

$$\begin{aligned} & \left| \int \int \frac{\dot{\phi}(r_0, u_0, \alpha_0)}{\max(z, z_0)} dv_0 dx_0 - \sum_{j=1}^N \frac{M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right| \\ &= \left| \sum_{j=1}^N \left(\iint_{S_j} \frac{\dot{\phi}(r_0, u_0, \alpha_0)}{\max(z, z_0)} dv_0 dx_0 - \frac{M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right) \right| \\ &= \left| \sum_{j=1}^N \iint_{S_j} \dot{\phi}(r_0, u_0, \alpha_0) \left(\frac{1}{\max(z, z_0)} - \frac{1}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right) dv_0 dx_0 \right| \\ &\leq \sum_{j=1}^N \iint_{S_j} \dot{\phi}(r_0, u_0, \alpha_0) \frac{8C_2(t)\delta}{r_0^2} dv_0 dx_0 = 8Mr_0^{-2}C_2(t)\delta = C(t)\delta. \end{aligned} \tag{50}$$

Again for $(r, u, \alpha) \in S_i$ we have

$$\begin{aligned} & \left| \left(z_i^2 + \frac{(ru \sin \alpha)^2}{z^2} \right) - \left((\bar{z}'_i(t))^2 + \frac{L_i^2}{(\bar{z}_i(t))^2} \right) \right| \\ &\leq \left| z_i^2 - (\bar{z}'_i(t))^2 \right| + z^{-2} \left| (ru \sin \alpha)^2 - L_i^2 \right| + L_i^2 \left| z^{-2} - (\bar{z}_i(t))^{-2} \right| \\ &\leq |z_i + \bar{z}'_i(t)| |z_i - \bar{z}'_i(t)| + r_0^{-2} (ru \sin \alpha + L_i) |ru \sin \alpha - L_i| \\ &\quad + L_0^2 z^{-2} (\bar{z}_i(t))^{-2} (\bar{z}_i(t) + z) |\bar{z}_i(t) - z|. \end{aligned}$$

By Theorem I, $|z_i - \bar{z}'_i(t)| \leq C_2(t)\delta$ so

$$|z_i + \bar{z}'_i(t)| \leq |2z_i| + |\bar{z}'_i(t) - z_i| \leq 2U_0 + C_2(t)\delta$$

and

$$\begin{aligned} & \left| z_i^2 + \frac{(ru \sin \alpha)^2}{z^2} - \left((\bar{z}'_i(t))^2 + \frac{L_i^2}{(\bar{z}_i(t))^2} \right) \right| \\ &\leq (2U_0 + C_2(t)\delta)C_2(t)\delta + r_0^{-2}2L_0C\delta + L_0^2r_0^{-2}(\frac{1}{2}r_0)^{-2}3(R_0 + U_0t)C_2(t)\delta \\ &\leq C(t)\delta. \end{aligned} \tag{51}$$

Now to finish the proof we rewrite E as in (47) and use the bounds given by (50) and (51):

$$\begin{aligned}
& \left| E - \sum_{i=1}^N M_i \left((\bar{z}'_i(t))^2 + L_i^2 (\bar{z}_i(t))^{-2} + \sum_{j=1}^N \frac{\gamma M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right) \right| \\
&= \left| \sum_{i=1}^N \iint_{S_i} \dot{\phi}(r, u, \alpha) \left(z_i^2 + \frac{(ru \sin \alpha)^2}{z^2} + \gamma \iint \frac{\dot{\phi}(r_0, u_0, \alpha_0)}{\max(z, z_0)} dv_0 dx_0 \right. \right. \\
&\quad \left. \left. - (\bar{z}'_i(t))^2 - L_i^2 (\bar{z}_i(t))^{-2} - \gamma \sum_{j=1}^N \frac{M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} dv dx \right) \right| \\
&\leq \sum_{i=1}^N \iint_{S_i} \dot{\phi}(r, u, \alpha) \left(\left| z_i^2 + \frac{(ru \sin \alpha)^2}{z^2} - (\bar{z}'_i(t))^2 - \frac{L_i^2}{(\bar{z}_i(t))^2} \right| \right. \\
&\quad \left. + \left| \iint \frac{\dot{\phi}(r_0, u_0, \alpha_0)}{\max(z, z_0)} dv_0 dx_0 - \sum_{j=1}^N \frac{M_j}{\max(\bar{z}_i(t), \bar{z}_j(t))} \right| \right) dv dx \\
&\leq \sum_{i=1}^N \iint_{S_i} \dot{\phi}(r, u, \alpha) (C(t)\delta + C(t)\delta) dv dx \\
&= 2MC(t)\delta.
\end{aligned}$$

This completes the proof of Corollary I.

REFERENCES

- [1] C. Bardos and P. Degond, *Global existence for the Vlasov–Poisson equation in 3 space variables with small initial data*. Preprint.
- [2] J. Batt, *Global symmetric solutions of the initial value problem of stellar dynamics*, J. Differential Equations **25**, 342–364 (1977)
- [3] E. Horst, *On the classical solutions of the initial value problem for the unmodified nonlinear Vlasov equation*, Parts 1 and 2, Math. Methods Appl. Sci. **3**, 229–248 (1981); **4**, 19–32 (1982)
- [4] R. Kurth, *Das Anfangswertproblem der Stelardynamik*, Z. Astrophys. **30**, 213–229 (1952)
- [5] H. Neunzert, *Mathematical investigations on particle-in-cell methods*, Fluid Dynamics Trans. **9**, 229–254 (1978)
- [6] S. Wollman, *The spherically symmetric Vlasov–Poisson system*, J. Differential Equations **35**, 30–35 (1980)