

**DETERMINATION OF THE LEADING COEFFICIENT  $a(x)$   
IN THE HEAT EQUATION  $u_t = a(x)\Delta u$**

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**Abstract.** This note deals with the parabolic inverse problem of determination of the leading coefficient in the heat equation with an extra condition at the terminal. After introducing a new variable, we reformulate the problem as a nonclassical parabolic equation along with the initial and boundary conditions. The existence of a solution is established by means of the Schauder fixed-point theorem.

**1. Introduction.** Recently, considerable effort was made in dealing with inverse problems in partial differential equations. These inverse problems not only have the intrinsic mathematical interests, but also have a variety of applications in industry and engineering sciences. It is known that an inverse problem is not well-posed in general. An important task is to formulate the problem properly and to find the conditions that ensure its well-posedness. In the present work, we study the inverse problem of finding  $a(x) > 0$  and  $u(x, t)$ , which satisfy:

$$u_t = a(x)\Delta u, \quad (x, t) \in Q_T; \quad (1.1)$$

$$u(x, t) = g(x, t), \quad (x, t) \in S_T = \partial\Omega \times [0, T]; \quad (1.2)$$

$$u(x, 0) = u_0(x), \quad x \in \Omega, \quad (1.3)$$

along with an extra condition

$$u(x, T) = u_1(x), \quad x \in \Omega, \quad (1.4)$$

where  $T > 0$  is fixed and  $Q_T = \Omega \times (0, T]$ , where  $\Omega$  is a bounded domain in  $\mathbf{R}^n$ .

When an unknown coefficient appears in the lower-order terms, various results are obtained in [1, 2, 4] (also see [7] and the references therein). The uniqueness of solution of the problem (1.1)–(1.4) was studied in [8]. In the present work we shall follow the idea of [4] to establish the existence for the problem (1.1)–(1.4). After

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Received October 28, 1991.

1991 *Mathematics Subject Classification.* Primary 35R25, 35R30.

The first author is partially supported by National Science Foundation DMS-90-24986, USA. The second author is partially supported by NSERC Canada.

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introducing a new variable, we have a nonlinear parabolic equation with the involvement of a trace-type functional as the leading coefficient. To avoid the degeneracy of the equation, we construct appropriate auxiliary functions and deduce some a priori estimates. The Schauder fixed point is used to prove the existence.

**2. The assumptions and the main result.** Let  $v(x, t) = u_t(x, t)$ . Then the extra condition (1.4) implies that

$$a(x) = \frac{v(x, T)}{\Delta u_1(x)},$$

provided  $\Delta u_1(x) \neq 0$ .

Now we differentiate Eq. (1.1) with respect to  $t$ ; then  $v(x, t)$  satisfies:

$$v_t = \frac{v(x, T)}{\Delta u_1(x)} \Delta v, \quad (x, t) \in Q_T; \tag{2.1}$$

$$v(x, t) = g_t(x, t), \quad (x, t) \in S_T; \tag{2.2}$$

$$v(x, 0) = k(x)v(x, T), \quad x \in \Omega, \tag{2.3}$$

where

$$k(x) = \frac{\Delta u_0(x)}{\Delta u_1(x)}.$$

Because of the nonlocal term  $v(x, T)$ , Eq. (2.1) is nonclassical. Moreover, the initial condition (2.3) is not known. In the sequel a solution of problem (1.1)–(1.4) or (2.1)–(2.3) is always understood in the classical sense.

**PROPERTY.** The problems (1.1)–(1.4) and (2.1)–(2.3) are equivalent if  $\Delta u_1 \neq 0$  in  $\bar{\Omega}$ .

Indeed, we have seen that if  $u(x, t)$ ,  $a(x)$  is a solution of problem (1.1)–(1.4), then  $v(x, t)$  is a solution of problem (2.1)–(2.3). Conversely, assuming that  $v(x, t)$  is a solution of problem (2.1)–(2.3), we easily verify that

$$u(x, t) = u_0(x) + \int_0^t v(x, \tau) d\tau, \quad a(x) = \frac{v(x, T)}{\Delta u_1}$$

is a solution of the inverse problem (1.1)–(1.4). Therefore, we shall investigate the problem (2.1)–(2.3).

Throughout this paper the following conditions are assumed:

H(1) The functions  $u_0(x), u_1(x) \in C^{4+\alpha}(\bar{\Omega})$ ,

$$\Delta u_0(x) \geq 0, \quad 0 < \Delta u_1(x) \leq M_0 \quad \text{in } \bar{\Omega}.$$

H(2) The function  $g(x, t) \in C^{4+\alpha, 2+\alpha/2}(S_T)$ , and

$$0 < g_0 \leq g_t(x, t) \leq G_0, \quad g_t(x, 0) = k(x)g_t(x, T) \quad \text{for } x \in \partial\Omega$$

$$\text{and } \frac{d^2}{T} \leq \frac{2g_0}{M_0 e^{3/2}},$$

where  $d = \text{MD}(\Omega)$  is the *minimum diameter* of  $\Omega$ , i.e., the infimum of distances between pairs  $\Pi_1, \Pi_2$  of parallel planes such that  $\Omega$  is contained in the strip determined by  $\Pi_1$  and  $\Pi_2$ .

H(3) The function  $k(x)$  satisfies

$$0 \leq k(x) \leq \exp \left( \frac{g_0 T}{e^{3/2} M_0} e^{-2d} \right).$$

The essential difficulty lies in that Eq. (2.1) may be degenerate, i.e.,  $v(x, T)$  may become zero at some points in  $\Omega$ . This would easily be avoided by using the maximum principle if the initial and the boundary data were uniformly positive; however, our initial condition is given by a relation between the initial and final states. The condition H(2) is physically reasonable since we require that  $a(x)$  is positive, which is equivalent to saying that  $v(x, T) > 0$  on  $\Omega$ . This is the case if the surrounding temperature is high enough.

The main result is

**THEOREM.** Under the conditions H(1)–H(3), the problem (1.1)–(1.4) admits a solution.

**3. Proof.** We shall use the Schauder fixed-point theorem to prove the result.

*Proof of Theorem.* Without loss of generality, we may assume that  $0 \in \partial\Omega$  and that  $\Omega$  lies in the strip  $0 \leq x_1 \leq d$ . Let

$$K = \left\{ w(x) \in C^\alpha(\bar{\Omega}) : k_0 \leq w(x) \leq G_0(e^{2d} - e^{x_1}) \text{ for } x \in \Omega, \right. \\ \left. w(x) = \frac{g_t(x, 0)}{k(x)} \text{ for } x \in \partial\Omega \text{ and } \|w\|_{C^\alpha(\bar{\Omega})} \leq k_1 \right\},$$

where the positive constants  $k_0$  and  $k_1$  will be specified later.

Obviously, each  $w(x)$  in  $K$  is bounded from above. For each  $w(x) \in K$ , we consider the problem:

$$v_t = \frac{w(x)}{\Delta u_1(x)} \Delta v, \quad (x, t) \in Q_T; \tag{3.1}$$

$$v(x, t) = g_t(x, t), \quad (x, t) \in S_T, \tag{3.2}$$

$$v(x, 0) = k(x)w(x), \quad x \in \Omega. \tag{3.3}$$

The standard theory of parabolic equations (cf. [6]) implies the problem admits a unique classical solution

$$v(x, t; w) \in C(\bar{Q}_T) \cap C^{2+\alpha, 1+\alpha/2}(Q_T).$$

Moreover, that for any  $t_0 > 0$  ( $t_0 < T$ ),

$$v(x, t; w) \in C^{2+\alpha, 1+\alpha/2}(\bar{Q}_{Tt_0}),$$

where  $Q_{Tt_0} = Q_T \cap \{(x, t) : x \in \bar{\Omega}, t_0 \leq t \leq T\}$  and

$$\|v\|_{C^{2+\alpha, 1+\alpha/2}(\bar{Q}_{Tt_0})} \leq C_0, \tag{3.4}$$

where  $C_0$  depends only on  $k_0, k_1$  and known data. Also, by Krylov-Safanov's  $C^\alpha$ -estimate,  $\|v\|_{C^{\alpha, \alpha/2}(\bar{Q}_{Tt_0})} \leq C_1$ , where  $C_1$  depends only on  $k_0$  and known data and is independent of  $k_1$ . Hence we can take  $k_1 = C_1$ .

Now we define a mapping  $M$  from  $K$  into  $C^{2+\alpha}(\overline{\Omega})$  as follows:

$$M: w \in K \rightarrow v(x, T; w) \in C^{2+\alpha}(\overline{\Omega}) \subset C^\alpha(\overline{\Omega}),$$

where  $v(x, t; w)$  is the solution of problem (3.1)–(3.3).

We first show that  $M$  is a continuous mapping from  $C^\alpha(\overline{\Omega})$  to  $C^\alpha(\overline{\Omega})$ . Let  $\{w_n(x)\} \subset K$  with  $w_n(x) \rightarrow w(x)$  in  $C^\alpha(\overline{\Omega})$  as  $n$  tends to infinity. Let  $v_n(x, t)$  and  $v(x, t)$  be the corresponding solutions of Eqs. (3.1)–(3.3), respectively. Then the function  $U(x, t) = v(x, t) - v_n(x, t)$  satisfies

$$LU = U_t - \frac{w(x)}{\Delta u_1(x)} \Delta U = \frac{\Delta v_n}{\Delta u_1} (w - w_n), \quad (x, t) \in Q_T;$$

$$U(x, t) = 0 \quad \text{on } S_T; \quad U(x, 0) = k(x)[w(x) - w_n(x)] \quad \text{on } \Omega.$$

By Green's representation, we have

$$U(x, t) = \int_{\Omega} G(x, y; t, 0) k(y) [w(y) - w_n(y)] dy + \int_0^t \int_{\Omega} G \frac{\Delta v_n}{\Delta u_1} [w - w_n] dy d\tau,$$

where  $G(x, y; t, \tau)$  is the Green's function corresponding to the operator  $L$ . It follows by Eq. (3.4) that

$$\max_{\overline{Q}_T} |U(x, t)| \leq C_1 \|w - w_n\|_0 + C_1 \|w - w_n\|_0 \rightarrow 0$$

as  $n \rightarrow \infty$ . Using the Schauder estimate on  $Q_{Tt_0}$ , we have

$$\|U\|_{C^{2+\alpha, 1+\alpha/2}(\overline{Q}_{Tt_0})} \leq C[\|U\|_0 + \|w - w_n\|_{C^\alpha(\overline{\Omega})}] \rightarrow 0$$

as  $n \rightarrow \infty$ . In particular, as  $n \rightarrow \infty$ ,

$$\|v(x, T) - v_n(x, T)\|_{C^\alpha(\Omega)} \rightarrow 0.$$

The compactness of  $M$  is clear since the embedding operator from  $C^{2+\alpha}(\overline{\Omega})$  into  $C^\alpha(\overline{\Omega})$  is compact. In order to apply the Schauder fixed-point theorem, it remains to prove that the mapping  $M$  maps  $K$  into itself.

Note that for  $x \in \partial\Omega$ ,  $v(x, 0) = k(x)w(x) = g_t(x, 0)$ , it follows by H(2) that for  $x \in \partial\Omega$

$$v(x, T) = g_t(x, T) = \frac{g_t(x, 0)}{k(x)}.$$

We shall next construct a subsolution for  $v(x, t) \equiv v(x, t; w)$ . Let  $\lambda = M_0/k_0$ . For  $x = (x_1, x_2, \dots, x_n) \in \overline{\Omega}$ , we introduce an auxiliary function

$$\psi(x, t) = \frac{C^*}{\sqrt{t}} \exp\left(-\frac{\lambda(x_1 - \xi_1)^2}{4t}\right),$$

where  $\xi = (\xi_1, 0, \dots, 0)$  is a fixed point, which lies in the outside of  $R^n \setminus \overline{\Omega}$  and  $C^*$  is a positive constant to be determined later. Then

$$\psi_t(x, t) = C^* \left[ -\frac{1}{2} + \frac{\lambda(x_1 - \xi_1)^2}{4t} \right] \frac{1}{\sqrt{t^3}} \exp\left(-\frac{\lambda(x_1 - \xi_1)^2}{4t}\right).$$

We choose  $\xi$  such that for all  $x = (x_1, \dots, x_n) \in \Omega$ ,

$$D \leq |x_1 - \xi_1| \leq 2D;$$

this is possible if we choose  $D$  such that  $D \geq d$ . With the above choice of  $\xi$ , we have  $\psi_t(x, t) \geq 0$  on  $Q_T$  if we choose  $D$  such that  $D^2\lambda = 2T$ . By a direct calculation, we see

$$\Delta\psi = \lambda C^* \left[ -\frac{1}{2} + \frac{\lambda(x_1 - \xi_1)^2}{4t} \right] \frac{1}{\sqrt{t^3}} \exp\left(-\frac{\lambda(x_1 - \xi_1)^2}{4t}\right).$$

It follows that

$$\Delta\psi = \frac{1}{\lambda}\psi_t \geq 0.$$

Hence

$$\begin{aligned} \psi_t - \frac{w(x)}{\Delta u_1(x)}\Delta\psi &\leq \psi_t - \frac{k_0}{M_0}\Delta\psi \\ &= \psi_t - \frac{1}{\lambda}\Delta\psi = 0. \end{aligned}$$

Moreover, since  $\xi \notin \bar{\Omega}$ , we have, for all  $x \in \Omega$ ,

$$\psi(x, 0) = \lim_{t \rightarrow 0} \psi(x, t) = 0.$$

Furthermore, on  $S_T$ , as  $|x_1 - \xi_1| \geq D$ , we have

$$\begin{aligned} \psi(x, t) &\leq \frac{C^*}{\sqrt{t}} \exp\left(-\frac{\lambda D^2}{4t}\right) \\ &\leq \frac{C^*}{\sqrt{T}} \sqrt{\frac{T}{t}} \exp\left(-\frac{\lambda D^2 T}{4T t}\right) \\ &\leq \frac{C^*}{\sqrt{T}} \sup_{0 < s < \infty} \left[ \sqrt{s} \exp\left(-\frac{s}{2}\right) \right] \\ &= \frac{C^*}{\sqrt{T}} \frac{1}{\sqrt{e}} \\ &\leq g_0, \end{aligned}$$

if we choose  $C^* = \sqrt{e}\sqrt{T}g_0$ . It follows that

$$\psi(x, t) \leq g_t(x, t) \quad \text{on } S_T.$$

By the comparison principle, one obtains

$$v(x, t) \geq \psi(x, t), \quad (x, t) \in \bar{Q}_T.$$

In particular, on  $\bar{\Omega}$ ,

$$\begin{aligned} v(x, T) &\geq \psi(x, T) \\ &\geq \sqrt{e}g_0 \exp\left(-\frac{\lambda D^2}{T}\right) \\ &\geq \frac{1}{\sqrt{e^3}}g_0. \end{aligned}$$

Therefore, if we take  $k_0 = e^{-3/2}g_0$ , then we have, on  $\bar{\Omega}$ ,

$$Mw = v(x, T; w) \geq k_0 \quad \text{for all } x \in \bar{\Omega}.$$

With our choice of the constants  $D$ ,  $\lambda$ , and  $k_0$  and also using assumption H(2), we get

$$D = \sqrt{\frac{2T}{\lambda}} = \sqrt{\frac{2Tk_0}{M_0}} = \sqrt{\frac{2Tg_0}{M_0e^{3/2}}} \geq d,$$

which is exactly what we assumed in the proof.

To show  $v(x, T) \leq G_0(e^{2d} - e^{x_1})$ , we introduce another auxiliary function:

$$\varphi(x, t) = G_0e^{\gamma(T-t)}[e^{2d} - e^{x_1}].$$

Then

$$\begin{aligned} \varphi_t - \frac{w(x)}{\Delta u_1(x)}\Delta\varphi &= G_0e^{\gamma(T-t)} \left[ -\gamma(e^{2d} - e^{x_1}) + \frac{w(x)}{\Delta u_1(x)}e^{x_1} \right] \\ &\geq G_0e^{\gamma(T-t)} \left( -\gamma e^{2d} + \frac{k_0}{M_0}e^{x_1} \right) \\ &\geq G_0e^{\gamma(T-t)} \left( -\gamma e^{2d} + \frac{k_0}{M_0} \right) \\ &= 0, \end{aligned}$$

if we choose  $\gamma = (k_0/M_0)e^{-2d}$ . Recalling the definition for  $k_0$ , we conclude that  $e^{\gamma T} = \exp((g_0T/e^{3/2}M_0)e^{-2d})$ . Thus, for  $x \in \Omega$ ,

$$\begin{aligned} \varphi(x, 0) &= e^{\gamma T}G_0(e^{2d} - e^{x_1}) \\ &\geq e^{\gamma T}w(x) \\ &\geq k(x)w(x) \quad [\text{by assumption H(3)}] \\ &= v(x, 0). \end{aligned}$$

On the boundary  $S_T$ ,

$$\begin{aligned} \varphi(x, t) &\geq G_0(e^{2D} - e^{x_1}) \\ &\geq G_0 = \max |g_t| \\ &\geq v(x, t). \end{aligned}$$

Again by the comparison principle, we have

$$v(x, t) \leq \varphi(x, t) \quad \text{on } \bar{Q}_T.$$

It follows that

$$v(x, T) \leq G_0(e^{2D} - e^{x_1}).$$

Thus, the mapping  $M$  is from  $K$  into itself. By the Schauder fixed-point theorem, the mapping  $M$  admits a fixed point, which is a solution of the problem (2.1)–(2.3). This completes our proof.

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