THE GEOMETRY OF FIELDS OF LINEAL ELEMENTS*

by

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1. Introduction. We shall begin by considering certain simple operations or transformations on the oriented lineal elements of the plane. A turn $T_\alpha$ converts each element into one having the same point and making a fixed angle $\alpha$ with the original direction. By a slide $S_k$, the line of the element remains the same and the point moves along the line a fixed distance $k$. These transformations together form a continuous group of three-parameters which we call the whirl group $W_3$. The group of whirls is isomorphic to the group of rigid motions $M_3$. These two three-parameter groups are commutative and together generate a continuous group of six-parameters which we term the whirl-motion group $G_6$. In preceding papers (see the bibliography at the end of this paper), Kasner and the author developed the geometry of this group $G_6$. In this paper, which is a continuation of the paper by the author The differential geometry of series of lineal elements, these Transactions, vol. 46 (1939), pp. 348–361, we shall give the differential geometry of fields of lineal elements with respect to the whirl-motion group $G_6$.

A set of $\infty^1$ elements is called a series; this includes a union (curve or point) as a special case. A collection of $\infty^2$ elements is termed a field, which of course corresponds to a differential equation of the first order, $F(x, y, y') = 0$. The totality of $\infty^3$ elements of the plane is called the opulence (as defined by Kasner).

In the earlier paper, we considered the tangent turbines, the osculating flat fields, and the osculating limaçon (circular) series of a given series $S$. We defined the curvature $\kappa$ and the torsion $\tau$ of any series. The curvature $\overline{\kappa}$ and the torsion $\overline{\tau}$ of a series $\overline{S}$ conjugate to a given series $S$ are given by the formulas $\overline{\kappa} = \kappa/\tau$, $\overline{\tau} = 1/\tau$. We proved the fundamental result that any two general (equi-parallel) series which have their curvatures and torsions the same functions of the angle $u$ (arc length $s$) are equivalent under the whirl-motion group $G_6$. This result establishes the intrinsic equations of any series in the geometry of the whirl-motion group $G_6$.

In the present paper, we shall derive the analogues of some of the classic theorems for a surface in a euclidean three-dimensional space. In particular,
we shall consider the Meusnier, the Euler, the Joachimsthal, and the Beltrami-Enneper theorems for a field in the geometry of the whirl-motion group $G_6$. The theory of geodesic series (minimum curvature) will be developed. We shall define the gaussian curvature $K$ of a field. Finally the theory of conjugate fields will be considered.

2. The tangent turbines of a series. A series which consists of $\infty^1$ non-parallel (parallel) elements is called a general (equiparallel) series. A general series is given by $v = v(u)$, $w = w(u)$, whereas an equiparallel series is given by $u = c$, $w = w(v)$, where $c$ is a constant. A general series possesses a point-union and a line-union, whereas an equiparallel series possesses only a point-union.

A turbine is the series which is obtained by applying a turn $T_\alpha$ to the elements of an oriented circle (the outer circle). It is nonlinear or linear according as this base circle is not or is a straight line.

A nonlinear turbine is a general series. Its point-union is a circle (the outer circle), and its line-union is also a circle (the inner circle). These two circles are concentric, and their common center is called the center of the turbine.

From the preceding remarks, we find that a nonlinear turbine may be constructed by applying a slide $S_s$ to the elements of an oriented circle (the inner circle). Thus the equations of a nonlinear turbine are

\begin{align}
v &= a \cos u + b \sin u + r, \\
w &= -a \sin u + b \cos u + s,
\end{align}

where $(a, b)$ are the cartesian coordinates of the center, $r$ is the radius of the inner circle, and $s$ is the constant distance of the slide $S_s$. We call $T(a, b, r, s)$ a set of nonlinear turbine coordinates.

A linear turbine is an equiparallel series whose base curve is a straight line. The equations of a linear turbine are

\begin{align}
u &= U - \omega, \\
v \cos \omega + w \sin \omega &= V,
\end{align}

where $(U, V)$ are the hessian coordinates of the base line and $\omega$ is the constant angle of the turn $T_\omega$. We call $T(U, V, \omega)$ a set of linear turbine coordinates. Obviously

$$T(U, V, \omega) = T(U + \pi, -V, \omega + \pi).$$

The angle $u_2 - u_1$ between any two elements is the angle between their lines. Two elements are parallel or supplementary (antiparallel) according as the angle between them is 0 or $\pi$. The distance $\sqrt{((v_2 - v_1)^2 + (w_2 - w_1)^2)}$ between two parallel elements is the distance between their points.
Two parallel elements are on a unique linear turbine. Two nonparallel elements are contained in a unique nonlinear turbine \( T \). The center of \( T \) is the intersection between the perpendicular bisector of the segment determined by the points of \( E_1 \) and \( E_2 \), and the angle bisector of the angle determined by the oriented lines of \( E_1 \) and \( E_2 \).

Two series \( S_1 \) and \( S_2 \) are said to be tangent (or to have contact of the first order) at a common element \( E \) if they have two (but not three) consecutive elements in common at \( E \). The two series \( S_1 \) and \( S_2 \) are said to be osculating (or to have contact of the second order) at \( E \) if they have three (but not four) consecutive elements in common at \( E \).

If a one-parameter family of series has the property that consecutive series have a common element, the family is called a set of enveloping series. The locus of intersection of consecutive series is termed the envelope. It is easy to prove that any series \( S_i \) of a set of enveloping series is tangent to the envelope \( S \) at any one of their common elements.

In the remainder of the paper, an accent will always mean total differentiation with respect to \( u \) unless otherwise specified.

**Theorem 1.** The tangent turbines of a general series are the \( \infty^1 \) nonlinear turbines whose parameter values are

\[
\begin{align*}
  a &= -v' \sin u - w' \cos u, \\
  b &= v' \cos u - w' \sin u, \\
  r &= v + w', \\
  s &= -v' + w.
\end{align*}
\]

On the other hand, the tangent turbines of an equiparallel series \( S \) are the \( \infty^1 \) linear turbines all of which possess the common direction of \( S \) and whose base lines are tangent to the base curve of \( S \).

It may be observed that two series \( S_1 \) and \( S_2 \) are tangent at a common element \( E \) if and only if they have the same tangent turbine at \( E \).

3. **Conjugate series.** Two turbines \( T \) and \( \bar{T} \) are said to be conjugate if they have the same circle as point-locus and the elements of the two turbines are symmetrically related to the elements of the circle. Two series \( S \) and \( \bar{S} \) are said to be conjugate if there exists a one-to-one correspondence between their elements in such a way that the tangent turbines of the two series at the corresponding elements are conjugate turbines.

The conjugate turbines \( \bar{T}_1 \) and \( \bar{T}_2 \) of two given turbines \( T_1 \) and \( T_2 \) (not both linear) do or do not possess a common element according as \( T_1 \) and \( T_2 \) do or do not possess a common element. The conjugate turbines of two intersecting linear turbines never possess a common element.

**Theorem 2.** For any general series \( S \), there always exists one and only one conjugate series \( \bar{S} \) which either consists of one element or is a general series. This
series \( \mathcal{S} \) is given by the equations

\[
\begin{align*}
\cos \hat{u} &= \frac{-a'r' + b's'}{a'^2 + b'^2}, \\
\sin \hat{u} &= \frac{-a's' - b'r'}{a'^2 + b'^2}, \\
\hat{v} &= a \cos \hat{u} + b \sin \hat{u} + r, \\
\hat{w} &= -a \sin \hat{u} + b \cos \hat{u} - s,
\end{align*}
\]

where \((a, b, r, s)\) are the parameter values of the tangent turbines of \( \mathcal{S} \).

If an equiparallel series \( \mathcal{S} \) is not a turbine, then there is no series conjugate to it.

4. The osculating flat fields of a series. A nonlinear flat field consists of the \( \infty^2 \) elements cocircular with a given element, called the central element. The equation of a nonlinear flat field \( \mathcal{II} \) is

\[
(5) \quad w = (v - \hat{v}) \cot \left( \frac{u - \hat{u}}{2} \right) - \hat{w},
\]

where \((\hat{u}, \hat{v}, \hat{w})\) are the hessian coordinates of the central element \( \mathcal{G} \) of \( \mathcal{II} \). We call \( \mathcal{II}(\hat{u}, \hat{v}, \hat{w}) \) a set of nonlinear flat field coordinates.

A linear flat field is the set of \( \infty^2 \) elements on \( \infty^1 \) parallel straight lines. Any linear field is given by \( u = \text{const} \).

The invariants between two nonlinear flat fields are identical with those between their central elements.

In a given flat field, there are \( \infty^2 \) turbines. The turbines which are contained in a nonlinear flat field \( \mathcal{II} \) are those whose conjugate turbines possess the central element \( \mathcal{G} \) of \( \mathcal{II} \). The turbines which are contained in a linear flat field \( \mathcal{II} \) are the linear turbines which have the common direction of \( \mathcal{II} \).

Three parallel elements determine a unique linear flat field. Three elements which are not all parallel and which do not lie on one turbine determine a unique nonlinear flat field \( \mathcal{II} \). The central element \( \mathcal{G} \) of \( \mathcal{II} \) is the single intersection of the conjugate turbines of the three turbines which pass through these elements.

Two elements of a flat field \( \mathcal{II} \) determine a turbine which lies entirely in \( \mathcal{II} \). Two flat fields (not both linear) intersect in a turbine. Two linear flat fields have no common elements.

The flat field which has three consecutive elements in common with a series \( \mathcal{S} \) at an element \( E \) of \( \mathcal{S} \) is called the osculating flat field of \( \mathcal{S} \) at \( E \).

**Theorem 3.** The osculating flat fields of a general series \( \mathcal{S} \) are the nonlinear flat fields whose central elements are the elements of the series \( \mathcal{S} \) conjugate to \( \mathcal{S} \).

If \( \mathcal{S} \) consists of only one element \( \mathcal{G} \), then \( \mathcal{S} \) is contained in the nonlinear flat field whose central element is \( \mathcal{G} \). In this case, we shall say that \( \mathcal{S} \) is a coflat series.
Any equiparallel series $S$ has one and only one osculating flat field, namely, the linear flat field in which it is contained.

5. **The osculating limaçon series of a general series.** Let $T$ be a nonlinear turbine, let $\mathcal{G}$ be a fixed element on the conjugate turbine $\bar{T}$ of $T$, and let $\gamma$ be a real number. Let $O$ be the point of $\mathcal{G}$, and let $P$ be the point of any element $E$ of $T$. On the line $(OP)$, let us select the points $P_i$ ($i=1, 2$) such that the distance $d(P, P_i) = 2\gamma$. Let $E_i$ be the element whose point is $P_i$ and whose direction is that of $E$. By this construction, to each element $E$ of $T$ there are associated two parallel elements $E_1$ and $E_2$. The totality of these elements $E_1, E_2$ is called a limaçon series with central turbine $T$ and radius $\gamma$.

Upon letting $C$ and $D$ denote

\begin{align}
C &= -2\gamma \sin \bar{u}/2, \\
D &= 2\gamma \cos \bar{u}/2,
\end{align}

we find that the equations of a limaçon series are

\begin{align}
v &= A \cos u + B \sin u + C \cos u/2 + D \sin u/2 + R, \\
w &= -A \sin u + B \cos u - C \sin u/2 + D \cos u/2 + S,
\end{align}

where $(A, B, R, S)$ are the parameters of the central turbine $\mathcal{G}$, and $\gamma$ is the radius of the limaçon series. We call $L(A, B, C, D, R, S)$ a set of limaçon series coordinates. Obviously


A limaçon series $L$ is contained in the nonlinear flat field $\Pi$ whose central element is $\mathcal{G}$. The centers of the tangent turbines of $L$, which of course are in $\Pi$, are on a circle with center $(A, B)$ and radius $\gamma$. We call this the associated circle of $L$. A limaçon series is uniquely determined by its flat field and its associated circle.

Three elements no two of which are parallel and which do not all lie on one turbine determine four limaçon series. Three elements only two of which are parallel determine two limaçon series. The flat field $\Pi$ of these limaçon series is the one determined by the given elements. Their associated circles are those which are tangent to all three of the angle bisectors of the angles formed by each of the oriented lines of these elements with the oriented line of the central element $\mathcal{G}$ of $\Pi$. In the first case, there are four circles, namely, the inscribed and escribed circles of the triangle formed by these lines. In the second case, there are only two circles since two of these three lines are parallel.

**Theorem 4.** The osculating limaçon series of a general series $S$ are those whose parameter values are
\[ A = a + 2r' \sin u + 2s' \cos u, \quad B = b - 2r' \cos u + 2s' \sin u, \]
\[ C = -4r' \sin u/2 - 4s' \cos u/2, \quad D = 4r' \cos u/2 - 4s' \sin u/2, \]
\[ R = r + 2s', \quad S = s - 2r', \]

where \((a, b, r, s)\) are the parameters of the tangent turbines of \(S\).

The envelope of the central turbines of the osculating limaçon series of a general series \(S\) is called the series of curvature of \(S\). This series is given by

\[ U = u + \pi, \quad V = v + 2w', \quad W = -2v' + w. \]

Theorem 5. The tangent turbines and the central turbines of any general series \(S\) have in common the series of curvature of \(S\).

6. The osculating circular series of an equiparallel series. An equiparallel series whose point-union is a circle with center \((A, B)\) and radius \(\gamma\) is called a circular series with center \((A, B)\) and radius \(\gamma\).

The osculating circular series of an equiparallel series \(S\) are those which possess the common direction of \(S\) and whose circles are the osculating circles of the point-union of \(S\).

7. The curvature and torsion of a general series. The curvature \(\kappa\) at an element \(E\) of a general series \(S\) is defined by the formula

\[ \kappa = (r'^2 + s'^2)^{1/2}, \]

where \((a, b, r, s)\) are the parameters of the tangent turbine of \(S\) at \(E\).

The quantity \(\kappa\) is one half of the radius of the osculating limaçon series \(L\) of \(S\) at \(E\), and also it is one half of the distance between the centers of the tangent and central turbines of \(S\) at \(E\).

The torsion \(\tau\) at an element \(E\) of a general series \(S\) is defined by the formula

\[ \tau = d\hat{u}/du, \]

where \(u\) and \(\hat{u}\) are the normal angles of the element \(E\) of \(S\) and the element \(\hat{E}\) which is the central element of the osculating flat field of \(S\) at \(E\).

The torsion \(\tau\) at an element \(E\) of a general series \(S\) is the rate of change of the angle of the osculating flat field per unit radian measure of the angle of the element \(E\).

8. The curvature of an equiparallel series. The curvature \(\kappa = 1/\gamma\) at an element \(E\) of an equiparallel series \(S\) is defined by the formula

\[ \kappa = \frac{1}{\gamma} = \frac{w''}{(1 + w'^2)^{3/2}}, \]

where the accent denotes differentiation with respect to \(v\).
The quantity \( \gamma = 1/\kappa \) is the radius of the osculating circular series of \( S \) at \( E \).

The torsion \( \tau \) of an equiparallel series is taken to be zero.

9. The osculating spherical fields of a general series. Let \( E \) denote an element, and let \( \Pi \) denote a nonlinear flat field with central element \( \mathcal{G} \). On the oriented line of \( E \), construct the element \( G \) which is in \( \Pi \). Let \( l \) be the line connecting the points of \( G \) and \( \mathcal{G} \). The perpendicular distance between the point of \( E \) and the line \( l \) is said to be the distance between \( E \) and \( \Pi \).

The set of \( \infty^2 \) elements \( E(u, v, w) \) which are at a constant distance \( \rho \) from a fixed nonlinear flat field \( \Pi(U, V, W) \) is called a spherical field \( \Sigma \). We term \( \Pi \) the central flat field and \( \rho \) the radius of \( \Sigma \). The equation of \( \Sigma \) is

\[
\begin{align*}
\omega &= (v - \bar{V}) \cot (u - \bar{U})/2 + \rho \csc (u - \bar{U})/2 - \bar{W}.
\end{align*}
\]

We call \( \Sigma(U, V, W, \rho) \) a set of spherical field coordinates. Obviously

\[
\Sigma(U, V, W, \rho) = \Sigma(U, V, W, -\rho).
\]

The integral curves of \( \Sigma \) are given in hessian line coordinates by the equation

\[
\begin{align*}
v &= -\rho \left[ \cos (u - \bar{U})/2 + \sin^2 (u - \bar{U})/2 \log \cot (u - \bar{U})/4 \right] \\
&\quad - C \cos (u - \bar{U}) + \bar{W} \sin (u - \bar{U}) + V + C,
\end{align*}
\]

where \( C \) is an arbitrary constant. If \( \rho = 0 \), then \( \Sigma \) becomes its central flat field \( \Pi \), and its integral curves are the \( \infty^1 \) circles which contain the central element \( \mathcal{G} \) of \( \Pi \). Otherwise if \( \rho \neq 0 \), the integral curves are transcendental.

Let \( C \) be the circular series whose center is the point of the central element \( \mathcal{G} \) of the central flat field \( \Pi \) of the spherical field \( \Sigma \), whose radius is the radius \( \rho \) of \( \Sigma \), and whose direction is that of \( \mathcal{G} \). We call \( C \) the associated circular series of \( \Sigma \). Obviously a spherical field \( \Sigma \) is uniquely determined by its associated circular series \( C \).

The only turbines in a spherical field \( \Sigma \) are the \( \infty^1 \) linear turbines whose conjugates are tangent to the associated circular series \( C \) of \( \Sigma \). These are the equiparallel series of \( \Sigma \). Thus a spherical field \( \Sigma \) contains no nonlinear turbines.

There are \( \infty^3 \) limaçon series in a spherical field \( \Sigma \). Their central turbines are contained in the central flat field \( \Pi \) of \( \Sigma \). If \( \rho \) is the radius of \( \Sigma \), \( \gamma \) is the radius of any one of these limaçon series \( L \), and \( \alpha \) is the angle between \( \Pi \) and the flat field of \( L \), then

\[
\rho = 2\gamma \sin \alpha/2.
\]

A least limaçon series of a spherical field \( \Sigma \) is any limaçon series of \( \Sigma \).
which has either one of the two equivalent properties: (1) its radius $\gamma$ is one half the radius $\rho$ of $\Sigma$, or (2) its flat field is supplementary or antiparallel to the central flat field $\Pi$ of $\Sigma$. There are $\infty^2$ least limaçon series in a spherical field $\Sigma$.

A spherical field $\Sigma$ intersects a nonlinear flat field $\Pi$, which is not parallel to the central flat field $\Pi$ of $\Sigma$ in a single limaçon series. If $\Pi_1$ and $\Pi$ are parallel but not identical, then $\Sigma$ intersects $\Pi_1$ in two linear turbines (which may be coincident or imaginary). If $\Pi_1$ and $\Pi$ are identical, then $\Sigma$ and $\Pi = \Pi_1$ have no common elements. A spherical field intersects a linear flat field in two linear turbines.

Two spherical fields $\Sigma_1$ and $\Sigma_2$ whose central flat fields $\Pi_1$ and $\Pi_2$ are not parallel intersect in two limaçon series. If $\Pi_1$ and $\Pi_2$ are parallel but not identical, then $\Sigma_1$ and $\Sigma_2$ intersect in four linear turbines (two of which may be coincident, or two or all four of which may be imaginary). If $\Pi_1$ and $\Pi_2$ are identical, then $\Sigma_1$ and $\Sigma_2$ have no common elements.

Four elements, at most two of which are parallel and which do not all lie in one flat field, determine eight spherical fields. Let us denote the four elements by $E_1, E_2, E_3, E_4$, where $E_3$ and $E_4$ are the two possible parallel ones. Now $E_1, E_2, E_j$ ($j = 3, 4$) determine four limaçon series. The associated circles of these are the inscribed and escribed circles of a triangle $T_j$ which has one vertex at the center $O$ of the turbine determined by $E_1$ and $E_2$. Let $L_{11}$ and $L_{12}$ denote the two limaçon series whose associated circles are the inscribed circle and the escribed circle opposite the vertex $O$ of $T_j$. Let $L_{13}$ and $L_{14}$ denote the remaining two limaçon series. The central turbines of $L_{21}, L_{22}, L_{23}, L_{24}$ will have a common element $F_1$, and hence will determine four nonlinear flat fields $\Pi_1$ ($i = 1, 2, 3, 4$). Similarly, the central turbines of $L_{33}, L_{34}, L_{43}, L_{44}$ will have a common element $F_2 \neq F_1$, and hence will determine four new nonlinear flat fields $\Pi_i$ ($i = 5, 6, 7, 8$). These eight flat fields $\Pi_i$ ($i = 1, \ldots, 8$) are the central flat fields of our eight spherical fields $\Sigma_i$. The radius $\rho_i$ of $\Sigma_i$ is the distance between $\Pi_i$ and any one of the four given elements.

The four elements $E_1, E_2, E_3, E_4$ such that only $E_1, E_2, E_3$ are parallel to each other but are not all on one turbine, determine six spherical fields. Construct the linear turbine $T_i$ determined by $E_i$ and $E_j$ ($i, j = 1, 2, 3$). Let $\Pi$ be the linear flat field which contains the conjugate turbine $\overline{T}_i$ of $T_i$. Construct the two linear turbines $\overline{T}_2$ and $\overline{T}_3$ contained in $\Pi$ such that their conjugate turbines $T_2$ and $T_3$ contain the remaining two elements. Our spherical fields are those whose associated circular series are tangent to $T_i$, $T_2$, and $T_3$.

The spherical field $\Sigma$ which has four consecutive elements in common with a general series $S$ at a given element $E$ of $S$ is called the osculating spherical field of $S$ at $E$. 
Theorem 6. The parameter values of the \( \infty \) osculating spherical fields of a general series \( S \) are

\[
\begin{align*}
\bar{U} &= u - 2 \arctan \frac{R'}{S'}, \\
\bar{V} &= A \cos \bar{U} + B \sin \bar{U} + R, \quad \bar{W} = -A \sin \bar{U} + B \cos \bar{U} - S, \\
\rho &= \frac{4(r'R' + s'S')}{(R'^2 + S'^2)^{1/2}},
\end{align*}
\]

where \((a, b, r, s)\) and \((A, B, R, S)\) are the parameters of the tangent and central turbines of \( S \).

Let \( v \) and \( w \) in (13) be replaced by functions of \( u \). Upon differentiating this result three times with respect to \( u \) and simplifying, we obtain

\[
\begin{align*}
\left( v - \bar{V} \right) \sin \left( u - \bar{U} \right)/2 + \left( w + \bar{W} \right) \cos \left( u - \bar{U} \right)/2 &= 2v' \cos \left( u - \bar{U} \right)/2 - 2w' \sin \left( u - \bar{U} \right)/2, \\
-\frac{\rho}{4} &= r' \sin \left( u - \bar{U} \right)/2 + s' \cos \left( u - \bar{U} \right)/2, \\
0 &= R' \cos \left( u - \bar{U} \right)/2 - S' \sin \left( u - \bar{U} \right)/2.
\end{align*}
\]

The second and last of these equations give the values of \( \bar{U} \) and \( \rho \) of (16). Upon replacing \( \rho \) in (13) by the value given in the second of these equations, and then solving this result and the first of the above equations for \( \bar{V} \) and \( \bar{W} \), we find their values to be those of (16). This completes the proof of Theorem 6.

An immediate consequence of Theorem 6 is

Theorem 7. The central flat fields of the osculating spherical fields of a general series \( S \) are the osculating flat fields of the series of curvature of \( S \) (the envelope of the central turbines of \( S \)).

If \( \alpha \) denotes the angle between the osculating flat fields of a general series \( S \) and its series of curvature, we find from (4) and (16)

\[
\sin \alpha/2 = \frac{r'R' + s'S'}{\left( r'^2 + s'^2 \right)^{1/2} \left( R'^2 + S'^2 \right)^{1/2}} = \frac{\rho}{2\gamma},
\]

where \( \rho \) and \( \gamma \) are the radii of the osculating spherical field \( \Sigma \) and the osculating limaçon series \( L \) at an element \( E \) of \( S \). Since the central turbine of \( L \) is also in the osculating flat field of the series of curvature of \( S \), we find that the following result holds.

Theorem 8. The osculating limaçon series \( L \) at an element \( E \) of a general series \( S \) is the intersection of the osculating flat field \( \Pi \) and the osculating spherical field \( \Sigma \) of \( S \) at \( E \).

From (4), (8), (10), and (11), we find
Substituting these into the last of equations (16), we obtain

**Theorem 9.** The radius of spherical curvature $\rho$ in terms of the curvature $\kappa$ and the torsion $\tau$ of the general series $S$ is

$$
\rho = \frac{4\kappa^2 \tau}{(\kappa^2 \tau^2 + 4\kappa'^2)^{1/2}}.
$$

Upon substituting the last of equations (19) into the formulas for the curvature and torsion of the series of curvature of a general series $S$, we obtain the theorem which follows.

**Theorem 10.** The curvature $\kappa_1$ and the torsion $\tau_1$ of the series of curvature of a general series $S$ in terms of the curvature $\kappa$ and the torsion $\tau$ of $S$ are

$$
\kappa_1 = \frac{\kappa^2 \tau + 4\kappa'^2}{(\kappa^2 \tau^2 + 4\kappa'^2)^{1/2}},
$$

$$
\tau_1 = \frac{\kappa^2 \tau^2 + 4\tau (2\kappa'^2 - \kappa \kappa'') + 4\kappa \kappa' \tau'}{\kappa^2 \tau^2 + 4\kappa'^2}.
$$

Differentiating (20) with respect to $u$, we find

**Theorem 11.** The derivative $\rho'$ of the radius of spherical curvature $\rho$ with respect to $u$ is

$$
\rho' = 4\kappa \kappa' \tau_1 / \kappa_1.
$$

It may be that the osculating spherical fields of a general series $S$ consist of only one spherical field, namely, the one in which it is contained. In that case, we shall say that $S$ is cospherical. From the preceding theorem, we deduce

**Theorem 12.** A general series $S$ is cospherical if and only if its series of curvature is coflat.

10. **The tangent flat fields of a field.** A set of $\infty^2$ elements of the plane is called a field. We shall omit from consideration the linear flat fields. That is, whenever we speak of a field, we shall understand it to be not a linear flat field. It is always possible to find a whirl-motion transformation such that any field $F$ is given by $w = w(u, v)$.

Let $v = v(u)$, $w = w(u, v)$ be a general series $S$ contained in the field $F$. Its tangent turbines are given by the parameter values

$$
a = - w_u \cos u - (\sin u + w_v \cos u)v',
$$

$$
b = - w_u \sin u + (\cos u - w_v \sin u)v',
$$

$$
r = v + w_u + v'w_v, \quad s = w - v'.
$$
From these equations, we conclude that the following proposition is true.

**Theorem 13.** The tangent turbines of all the series, contained in a field $F$ and passing through an element $E$ of $F$, constructed at $E$ are contained in a non-linear flat field.

The non-linear flat field of Theorem 13 is called the tangent flat field of $F$ at $E$. Its central element $E(\bar{u}, \bar{v}, \bar{w})$ is given by

$$\cos (\bar{u} - u) = -\frac{1 - w_v^2}{1 + w_v^2}, \quad \sin (\bar{u} - u) = -\frac{2w_v}{1 + w_v^2},$$

$$\bar{v} = v + \frac{2w_u}{1 + w_v^2}, \quad \bar{w} = -w - \frac{2w_w w_v}{1 + w_v^2}.$$  
(24)

We call $E$ the conjugate element of $E$ with respect to $F$.

Two fields $F_1$ and $F_2$ are said to be tangent at a common element $E$ if they have the same tangent flat field at $E$.

As an application of the above, we find that the tangent flat fields of a spherical field $\Sigma$ consist of the $\infty^1$ flat fields whose central elements are those of the associated circular series of $\Sigma$.

11. **One-parameter families of fields.** The equation

$$w = w(u, v, a)$$

defines a one-parameter family of fields. The series of intersection of any two consecutive fields of this family is called a characteristic. The locus of all the characteristics is a field, called the envelope of the family. The equations

$$w = w(u, v, a), \quad w_a(u, v, a) = 0$$

for each $a$ represents a characteristic of the family. When we eliminate $a$ from the above equations, the result is the equation of the envelope.

It may be easily proved by the preceding equations that the envelope is tangent to each member of the family at all elements of its characteristic.

The series of intersection of consecutive characteristics of a one-parameter family of fields is called the edge of regression. The eliminants with respect to $a$ of the equations

$$w = w(u, v, a), \quad w_a(u, v, a) = 0, \quad w_{aa}(u, v, a) = 0$$

(27)

give the equations of the edge of regression.

We may prove by (26) and (27) that the edge of regression is tangent to any characteristic at a common element.

12. **Developable fields.** The envelope of $\infty^1$ non-linear flat fields is called a developable field $F$. The series $S$ formed by the central elements of these $\infty^1$
tangent flat fields of $F$ is called the \textit{associated series} of $F$. The characteristics of $F$ are turbines. These are called the \textit{generators} of $F$.

Since each flat field is tangent to the envelope along its characteristic, it follows that \textit{the tangent flat field to a developable field $F$ is the same at all elements of a generator}.

An \textit{umbilical field} $F$ is a developable field whose associated series $S$ is an equiparallel series. Thus a spherical field $\Sigma$ is an umbilical field whose associated series is a circular series. The generators of an umbilical field $F$ are linear turbines. These are the conjugates of the tangent turbines of its associated equiparallel series $S$. The edge of regression of $F$ does \textit{not} exist. The equation of any umbilical field $F$ is

\[ w = v \cot \left( \frac{u - a}{2} + b(u) \right), \]

where $a$ is a constant.

A developable field $F$ is said to be \textit{general} if its associated series $S$ is a general series. The generators of a general developable field $F$ are nonlinear turbines. These are the tangent turbines of the edge of regression $R$. Since consecutive generators are the consecutive tangent turbines of $R$ at an element $E$ of $R$, the osculating flat field of $R$ at $E$ is that flat field of the family which contains these generators. But this flat field is tangent to the developable. \textit{Hence the osculating flat field at any element $E$ of the edge of regression $R$ of a general developable field $F$ is the tangent flat field of $F$ at $E$}. We find from this that \textit{the edge of regression $R$ and the associated series $S$ of a general developable field $F$ are conjugate series}.

The necessary and sufficient condition that a field $F: w = w(u, v)$ be developable is that its conjugate elements $E(\bar{u}, \bar{v}, \bar{w})$ of (24) consist of at most $\infty^1$ elements. These will then form the associated series $S$ of $F$. Hence upon setting the three jacobians of the three functions $\bar{u}$, $\bar{v}$, $\bar{w}$ of $(u, v)$ equal to zero, we obtain

\textbf{Theorem 14.} \textit{A field} $F: w = w(u, v)$ \textit{is a developable field if and only if}

\[ (1 + w_v + 2w_{uv})^2 - 4w_vw_{uv}(w_u + w_w) = 0. \]

A field $F: w = w(u, v)$ is a general developable field if and only if the function $w$ of $(u, v)$ satisfies the above equation and $w_w \neq 0$.

13. \textit{Conjugate fields}. The conjugate elements $E(\bar{u}, \bar{v}, \bar{w})$ of (24) of a field $F: w = w(u, v)$ form a field $\bar{F}: \bar{w} = \bar{w}(\bar{u}, \bar{v})$ if and only if $F$ is nondevelopable, or if and only if the function $w$ of $(u, v)$ does not satisfy (29). This field $\bar{F}$: $\bar{w} = \bar{w}(\bar{u}, \bar{v})$ is termed the \textit{conjugate field} of $F$. The equation of $\bar{F}$ is the eliminator with respect to $u$ and $v$ of the equations (24).
Since the conjugate field $\overline{F}$ consists of the central elements of the tangent flat fields of a non-developable field $F$ (and conversely), it follows that the tangent turbines of $\overline{F}$ are the conjugates of the tangent turbines of $F$ (and conversely). Hence for the special case when the tangent turbines are circles (the self-conjugate turbines), we deduce the following result.

**Theorem 15.** Two fields $F$ and $\overline{F}$ are conjugate if and only if their integral curves possess the same osculating circles.

14. **The gaussian curvature of a field.** The series of intersection between a nonlinear flat field and a given field $F$ is called a flat section of $F$. If $F$ is not a flat field, there are $\infty^3$ flat sections in $F$. There pass $\infty^2$ flat sections of $F$ through any element $E$. Finally there are $\infty^1$ flat sections of $F$ which contain a given element $E$ and which possess a fixed tangent turbine at $E$.

Let $S_1$ be any general series contained in a field $F$. There is a unique flat section $S$ which osculates $S_1$ at a given element $E$ of $S_1$. This flat section $S$ is the intersection between the field $F$ and the osculating flat field of $S_1$ at $E$. The two series $S_1$ and $S$ will have the same tangent turbine, the same osculating flat field, the same osculating limacon series, and the same curvature at $E$. Thus in order to study the curvatures and the osculating limacon series of any general series contained in a field $F$, it is necessary merely to study those of any flat section of $F$.

Next we shall seek to obtain the curvature $\kappa$ of any flat section $S$ of a field $F$: $w = w(u, v)$ at any element $E$ of $S$ in terms of the angle $\beta$ between the flat field of $S$ and the tangent flat field of $F$ at $E$. Upon eliminating $v'$ from the values of $r'$ and $s'$, the first derivatives with respect to $u$ of the last two parameters $r$ and $s$ of the tangent turbine of $S_1$ at $E$, we find

$$r' + w_s s' = (w_{uu} + w_u w_v) + (1 + w_v^2 + 2w_u w_v) v' + w_v v'^2. \tag{30}$$

We see from (4) and (24) that the angle $\beta$ satisfies the equation

$$\frac{r'}{s'} = \frac{\sin \beta/2 - w_v \cos \beta/2}{\cos \beta/2 + w_u \sin \beta/2}. \tag{31}$$

Solving the preceding two equations for $r'$ and $s'$, and then substituting these results into the curvature formula, we find that the value of the curvature $\kappa$ of a flat section $S$ of a field $F$ at any element $E$ of $S$ in terms of the angle $\beta$ between the flat field of $S$ and the tangent flat field of $F$ at $E$ is

$$\kappa = \frac{(w_{uu} + w_u w_v) + (1 + w_v^2 + 2w_u w_v) v' + w_v v'^2}{(1 + w_u^2)^{1/2} \sin \beta/2}. \tag{32}$$
When the angle $\beta = \pi$, we shall call the flat section a supplementary section and its curvature at $E$ the supplementary curvature $\kappa_s$. By the preceding equation, the value of the supplementary curvature $\kappa_s$ is

$$\kappa_s = \frac{(w_{uu} + w_u w_v) + (1 + w_v^2 + 2w_{uv})v' + w_{vv}v'^2}{(1 + w_v^2)^{1/2}}.$$  \hfill (33)

By the above two equations, we obtain the following analogue of Meusnier’s theorem in the geometry of the whirl-motion group $G_e$.

**Theorem 16.** Let $\kappa_s$ and $\kappa$ be the curvatures of a supplementary section and any other flat section which have the same tangent turbine at a common element $E$ of a field $F$. If $\alpha$ denotes the angle between these two flat sections, then

$$\kappa_s = \kappa \cos \alpha/2.$$  \hfill (34)

The above result shows that a supplementary section possesses the least curvature of all the flat sections of a field $F$ which pass through a given element $E$ in a given tangent turbine direction.

A field $F$ which is generated by a one-parameter family of turbines such that consecutive turbines of the family do not lie in a flat field is termed a ruled field. We shall say that a ruled field $F$ is general or special according as the turbines of the family are nonlinear or linear. A special ruled field $F$ is given by either $w_{vv} = 0$, or $w = v m(u) + b(u)$, which is not of the form (28).

A field $F: w = w(u, v)$ is called a general field if it is neither a special ruled field nor an umbilical field. Thus $F$ is a general field if and only if $w_{vv} \neq 0$. Of course, the general ruled and the general developable fields are all examples of general fields.

**Theorem 17.** At any element $E$ of a general field $F$, there is one and only one extremal (maximum or minimum) supplementary curvature $\kappa_0$. It is given by the formula

$$\kappa_0 = \frac{-(1 + w_v^2 + 2w_{uv})^2 + 4w_{vv}(w_{uu} + w_u w_v)}{4w_{vv}(1 + w_v^2)^{1/2}}.$$  \hfill (35)

For upon completing the square of the quadratic expression in $v'$ of (33), we find that (33) may be written in the form

$$\kappa_s = \frac{-(1 + w_v^2 + 2w_{uv})^2 + 4w_{vv}(w_{uu} + w_u w_v)}{4w_{vv}(1 + w_v^2)^{1/2}} + \frac{[(1 + w_v^2 + 2w_{uv}) + 2w_{vv}v']^2}{4w_{vv}(1 + w_v^2)^{1/2}}.$$  \hfill (36)
This will be a maximum or a minimum with respect to \( v' \) only when the squared bracket is zero. The remaining part of the above expression will give us the value of \( \kappa_0 \) of (35). Theorem 17 is completely proved.

Through any element \( E \) of a field \( F \), there passes a unique equiparallel series \( S \) contained in \( F \). The curvature of \( S \) at \( E \) is called the \textit{equiparallel curvature} \( \lambda \) of \( F \) at \( E \). It is given by

\[
\lambda = \frac{w_{vv}}{(1 + w_v^2)^{3/2}}.
\]

By means of (35) and (37), we find that (33) or (36) may be written in the form

\[
\kappa_s = \kappa_0 + \lambda \left( \frac{1}{v'} + \frac{w_v^2}{2w_{vv}} \right).
\]

If \( \delta \) denotes the distance between the centers of the tangent turbines of (23) of the extremal supplementary section and any supplementary section, we obtain the following analogue of Euler's theorem.

\textbf{Theorem 18.} Let \( \kappa_0 \) be the extremal supplementary curvature and \( \lambda \) the equiparallel curvature at an element \( E \) of a general field \( F \). If \( \delta \) is the distance between the centers of the tangent turbines of the extremal supplementary section and any supplementary section whose curvature at \( E \) is \( \kappa_s \), then

\[
\kappa_s = \kappa_0 + \lambda \delta^2.
\]

The \textit{gaussian curvature} \( K \) of any field \( F \) at any element \( E \) of \( F \) is given by the formula

\[
K = \frac{(1 + w_v^2 + 2w_{uv})^2 - 4w_{vv}(w_{uu} + w_u w_v)}{(1 + w_v^2)^3}.
\]

We note that \( K \) is zero if and only if \( F \) is developable.

\textbf{Theorem 19.} \textit{The gaussian curvature} \( K \) \textit{at any element} \( E \) \textit{of a general field} \( F \) \textit{is minus four times the product of the extremal supplementary curvature} \( \kappa_0 \) \textit{and the equiparallel curvature} \( \lambda \) \textit{at} \( E \). That is

\[
K = -4 \kappa_0 \lambda.
\]

By relations (23), (33), and (40), we now deduce the following proposition.

\textbf{Theorem 20.} \textit{If} \( \delta \) \textit{is the distance between the centers of the tangent turbines of any two supplementary sections whose curvatures are} \( \kappa_s \) \textit{and} \( \kappa'_s \) \textit{at a common element} \( E \) \textit{of a special ruled field} \( F \), \textit{then}
For an umbilical field $F$ of (28), all the supplementary curvatures at an element $E$ of $F$ are equal, and they have the common value

$$\kappa_\alpha = \kappa'_\alpha + K^{1/2}. $$

Through any element $E$ of a general field $F$, there is a single tangent turbine direction which gives the extremal supplementary curvature $\kappa_0$. This is called the principal direction of $F$ at $E$. Any general series $S$ of a general field $F$ such that the tangent turbine direction of any element $E$ of $S$ is a principal direction is called a principal series. The differential equation of all principal series of a general field $F$ is

$$2w_{vv'} + (1 + w^2 + 2w_{uu}) = 0.$$
\[
\frac{\partial \tilde{v}}{\partial \tilde{u}} = -w_x, \quad \frac{\partial \tilde{w}}{\partial \tilde{v}} = -w_y,
\]
\[
\frac{\partial^2 \tilde{w}}{\partial \tilde{u}^2} = \frac{-w_{xx}(1 + w_x^2) + 4w_xw_y(w_{xy}w_{yy} - w_{xx})}{(1 + w_x^2 + 2w_{xy})^2 - 4w_{xy}(w_{xx} + w_{yy})},
\]
\[
\frac{\partial^2 \tilde{w}}{\partial \tilde{v}^2} = \frac{-w_{yy}(1 + w_y^2) + 2(1 + w_y^2)(w_{xy}w_{yy} - w_{xx})}{(1 + w_y^2 + 2w_{xy})^2 - 4w_{xy}(w_{xx} + w_{yy})},
\]
\[
\frac{\partial^2 \tilde{w}}{\partial \tilde{u} \partial \tilde{v}} = \frac{-w_{xy}(1 + w_x^2)}{(1 + w_x^2 + 2w_{xy})^2 - 4w_{xy}(w_{xx} + w_{yy})}.
\]

By these equations, we deduce the following result.

**Theorem 22.** The product of the gaussian curvatures $K$ and $\tilde{K}$ at conjugate elements $E$ and $\tilde{E}$ of two conjugate fields $F$ and $\tilde{F}$ is unity. That is

\[
KK = 1.
\]

The series $S$ of a field $F$ and the series $\tilde{S}$ of the field $\tilde{F}$ conjugate to $F$ are said to be conjugate with respect to $F$ or $\tilde{F}$ if they are corresponding series under the transformation (24).

For any two series $S$ and $\tilde{S}$ conjugate with respect to two conjugate general fields $F$ or $\tilde{F}$, we obtain the following relation:

\[
1 + w_x^2 + 2w_{xy} + 2w_{yy} \frac{d\tilde{v}}{d\tilde{u}} = \frac{(1 + w_y^2)^2}{1 + w_x^2 + 2w_{xy} + 2w_{yy}w_{yy}'}.
\]

By substituting this into the supplementary curvature formula of $\tilde{F}$, we prove

**Theorem 23.** Let $E$ be any element of a general nondevelopable field $F$ and $\tilde{E}$ its conjugate element of the general field $\tilde{F}$ conjugate to $F$. The corresponding supplementary curvatures $\kappa_s$ and $\kappa_s'$, the extremal supplementary curvatures $\kappa_0$ and $\kappa_0'$, and the equiparallel curvatures $\lambda$ and $\tilde{\lambda}$ of $F$ and $\tilde{F}$ at $E$ and $\tilde{E}$ are related by the formulas

\[
\kappa_s = \frac{\kappa_s}{4\lambda(\kappa_s - \kappa_0)}, \quad \kappa_0 = \frac{1}{4\lambda}, \quad \tilde{\lambda} = \frac{1}{4\lambda_0}.
\]

For any two series $S$ and $\tilde{S}$ conjugate with respect to two conjugate special ruled fields $F$ or $\tilde{F}$, we find the following relation (since $w_{xy} = 0$)

\[
\frac{d\tilde{v}}{d\tilde{u}} = v' + \frac{2(1 + w_x^2)w_{uu} - 4w_xw_yw_{uv}}{(1 + w_x^2)(1 + w_y^2 + 2w_{xy})}.
\]
Upon substituting this into the supplementary curvature formula of \( \bar{F} \), we obtain the following result.

**Theorem 24.** Let \( E \) be any element of a special ruled field \( F \) and \( \bar{E} \) its conjugate element of the special ruled field \( \bar{F} \) conjugate to \( F \). The corresponding supplementary curvatures \( \kappa \) and \( \bar{\kappa} \) of \( F \) and \( \bar{F} \) at \( E \) and \( \bar{E} \) are related by the formula

\[
\bar{\kappa} = \frac{\kappa}{K}.
\]

The principal series of two conjugate fields \( F \) and \( \bar{F} \) are conjugate with respect to \( F \) or \( \bar{F} \). A principal general series of one of these two general fields corresponds by the transformation (24) to an equiparallel series of the other. Otherwise the equiparallel series of these two special ruled fields correspond to each other under the transformation (24).

16. **The osculating limaçon series of a field.** Upon solving the equations (30) and (31) for \( r' \) and \( s' \) and making use of the supplementary curvature formula (33), we find that their values are

\[
\begin{align*}
    r' &= \frac{\kappa_s(\sin \beta/2 - w_v \cos \beta/2)}{(1 + w_v^2)^{1/2} \sin \beta/2}, \\
    s' &= \frac{\kappa_s(\cos \beta/2 + w_v \sin \beta/2)}{(1 + w_v^2)^{1/2} \sin \beta/2}. 
\end{align*}
\]

Substituting these values into (8), we see that the parameters of the osculating limaçon series \( L \) of any flat section \( S \) of a field \( F : w = w(u, v) \) at an element \( E \) of \( F \) are

\[
\begin{align*}
    A &= a + \frac{2\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ \cos (\beta/2 - u) + w_v \sin (\beta/2 - u) \right], \\
    B &= b + \frac{2\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ - \sin (\beta/2 - u) + w_v \cos (\beta/2 - u) \right], \\
    C &= -\frac{4\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ \cos (\beta - u)/2 + w_v \sin (\beta - u)/2 \right], \\
    D &= \frac{4\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ \sin (\beta - u)/2 - w_v \cos (\beta - u)/2 \right], \\
    R &= r + \frac{2\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ \cos \beta/2 + w_v \sin \beta/2 \right], \\
    S &= s - \frac{2\kappa_s}{(1 + w_v^2)^{1/2} \sin \beta/2} \left[ \sin \beta/2 - w_v \cos \beta/2 \right],
\end{align*}
\]

where \((a, b, r, s)\) are the parameters of the tangent turbine \( T \) of (23) of \( S \) at \( E \),
\( k_s \) is the supplementary curvature at \( E \) of the supplementary section of \( F \) which is tangent to \( S \) at \( E \), and \( \beta \) is the angle between the flat field of \( S \) and the tangent flat field of \( F \) at \( E \).

Next let us consider all the \( \infty \) flat sections of the field \( F \) which pass through the element \( E \) of \( F \) and which possess the same tangent turbine \( T \) at \( E \). In the formulas (53) for the \( \infty \) osculating limacon series of these flat sections at \( E \), we observe that only the angle \( \beta \) is variable. The \( \infty \) central turbines of these osculating limacon series all contain the element \( E_i \) of the tangent turbine \( T \) which is supplementary (antiparallel) to \( E \). By (53), it may be proved after some calculation that these central turbines all are contained in the flat field II whose central element \( \bar{G}(U, V, W) \) is given by

\[
\begin{align*}
\cos(U - u) &= -\frac{1 - w^2}{1 + w^2}, \\
\sin(U - u) &= -\frac{2w}{1 + w^2}, \\
V &= v + \frac{2w_u}{1 + w^2} + \frac{4\kappa_s w_v}{(1 + w^2)^{1/2}}, \\
W &= -w - \frac{2w_u w_v}{1 + w^2} + \frac{4\kappa_s}{(1 + w^2)^{1/2}}.
\end{align*}
\]

From these equations and from Meusnier’s Theorem 16 (formula 34), we deduce

**Theorem 25.** Let us consider the \( \infty \) flat sections of a field \( F \) which pass through an element \( E \) of \( F \) and which possess the same tangent turbine \( T \) at \( E \). The \( \infty \) osculating limacon series of these flat sections at \( E \) generate a spherical field \( \Sigma \).

Let \( E_i \) be the element on the tangent turbine \( T \) which is supplementary (antiparallel) to the fixed element \( E \). Let \( T_s \) be the linear turbine whose direction is that of the central element \( \bar{E} \) of the tangent flat field of \( F \) at \( E \), and whose base line joins the points of \( \bar{E} \) and \( E_i \). The central element \( \bar{G} \) of the central flat field II of the spherical field \( \Sigma \) of Theorem 25 is on the linear turbine \( T_s \) and the distance of \( \bar{G} \) from \( \bar{E} \) is \( 4\kappa_s \). The radius \( \rho \) of \( \Sigma \) is also \( 4\kappa_s \).

If we vary the tangent turbine direction \( T \) of Theorem 25, there will result \( \infty \) spherical fields. The \( \infty \) central elements of their central flat fields will generate the linear turbine \( T_s \).

**17. The geodesic series of a field.** A series \( S \) of a field \( F \) is termed a geodesic series if its curvature at any element \( E \) of \( S \) does not exceed the curvature at \( E \) of any other series of \( F \) which is tangent to \( S \) at \( E \). By setting the partial derivative with respect to \( v'' \) of the curvature \( \kappa \) of any general series of a field \( F \) equal to zero, and solving the result for \( v'' \), we find that the differential equation of all the geodesic series of a field \( F \): \( w = w(u, v) \) is
There are \( \infty^4 \) geodesic series, all general, in a given field \( F \). The \( \infty^1 \) equi-
parallel series of a field \( F \) are also considered to be geodesic series. The geod-
esic series of a spherical field are its \( \infty^2 \) least limacon series, together with its
\( \infty^1 \) linear turbines. The geodesic series of a flat field are its \( \infty^2 \) turbines.

The curvature \( \kappa \) of any geodesic series \( S \) of a field \( F \) at any element \( E \) of \( S \)
is equal to the supplementary curvature \( \kappa_s \) of the supplementary section of \( F \)
which is tangent to \( S \) at \( E \).

To calculate the torsion \( \tau \) of a geodesic series \( S \), we proceed as follows. The
derivatives \( r' \) and \( s' \) with respect to \( u \) of the last two parameters \( r \) and \( s \)
of the tangent turbine \( T \) of \( S \) at any element \( E \) of \( S \) are given by (52) where
\( \beta = \pi \). Thus \( s'/r' = w_\nu \). By this and (4), we find that the normal angle \( \bar{u} \)
of the osculating flat field of \( S \) at \( E \) is

\[
\bar{u} = u + 2 \arctan w_\nu .
\]

Differentiating this with respect to \( u \), we see that the torsion \( \tau \) of any geodesic
series \( S \) at any element \( E \) of \( S \) is

\[
\tau = 1 + w_\nu^2 + 2w_\nu v + 2w_\nu v' .
\]

**Theorem 26.** The torsion \( \tau \) of a geodesic series \( S \) of a field \( F \) at an element \( E \)
of \( S \) is zero if and only if \( S \) is tangent to a principal series of \( F \) at \( E \). The neces-
sary and sufficient condition that a geodesic series be a principal series is that it
be coflat.

By equation (57), we obtain the following two results.

**Theorem 27.** Let \( \lambda \) be the equiparallel curvature of a general field \( F \) at an
element \( E \) of \( F \). Let \( \delta \) be the distance between the centers of the tangent turbines
of the extremal supplementary section and any geodesic series \( S \) through \( E \). The
torsion \( \tau \) of \( S \) is

\[
\tau = 2\lambda \delta .
\]

**Theorem 28.** The torsion \( \tau \) of any geodesic series \( S \) of a special ruled or an
umbilical field \( F \) at any element \( E \) of \( S \) is

\[
\tau = K^{1/2} .
\]

To define the geodesic curvature \( \kappa_0 \) of any series \( S \) of a field \( F \) at any ele-
ment \( E \) of \( S \), we proceed as follows. Let \( S_0 \) be the geodesic series tangent to \( S \).
at $E$. Let $E_s$ be the element on $S$ which makes an angle $\Delta u$ with $E$, and let $E_g$ be the element on $S_g$ which makes the same angle $\Delta u$ with $E$. At $E_s$ and $E_g$, construct the two tangent turbines $T_s$ and $T_g$ of $S$ and $S_g$. Let $\Delta \delta$ be the distance between the centers of $T_s$ and $T_g$. Then

$$\kappa_g = \frac{d\delta}{du} = \lim_{\Delta u \to 0} \frac{\Delta \delta}{\Delta u}$$

is defined to be the geodesic curvature of $S$ at $E$.

It is found that $\kappa_g = (1 + w^2)^{1/2}(v'' - v''')$, where $v''$ belongs to the geodesic series $S_v$. From this, it follows that the geodesic curvature $\kappa_g$ of any series $S$ of a field $F: w = w(u, v)$ at any element $E$ of $S$ is

$$\kappa_g = \frac{(1 + w^2)v'' - (w_u - w_v w_u) + 2w_v^2 w_u v' + w_v w_v v'^2}{(1 + w^2)^{1/2}}.$$

By means of (33), (52), and (61), we obtain the result which follows.

**Theorem 29.** Let $\beta$ be the angle between the osculating flat field of a general series $S$ of a field $F$ and the tangent flat field of $F$ at an element $E$ of $S$. Let $\kappa$ be the curvature and $\kappa_g$ the geodesic curvature of $S$ at $E$. Let $\kappa_s$ be the supplementary curvature of the supplementary section which is tangent to $S$ at $E$. Then

$$\kappa_s = \kappa \sin \beta / 2 = -\kappa g \tan \beta / 2, \quad \kappa_g = -\kappa \cos \beta / 2, \quad \kappa_s^2 + \kappa_g^2 = \kappa^2.$$

**18. The asymptotic series of a field.** Reciprocal directions at an element $E$ of a field $F$ may be defined as follows. Let $G$ be an element in $F$ adjacent to $E$. Let $ER$ be the turbine of intersection of the tangent flat fields of $F$ at $E$ and $G$. As $G$ tends to coincidence with $E$, the limiting tangent turbine directions of $EG$ and $ER$ are said to be reciprocal at $E$.

The necessary and sufficient condition that the tangent turbine directions $dv/du$ and $\delta u/\delta v$ be reciprocal are

$$2w_v \frac{dv}{du} + (1 + w^2 + 2w_v) \left( \frac{\delta v}{\delta u} + \frac{dv}{du} \right) + 2(w_{uu} + w_u w_v) = 0.$$

**Theorem 30.** Let $T_0$ be the principal tangent turbine (the principal direction), and let $T$ and $T_1$ be any other two tangent turbines at an element $E$ of a general field $F$. Let $\delta$ and $\delta_1$ be the distances of the centers of $T$ and $T_1$ from that of $T_0$. The directions of $T$ and $T_1$ are reciprocal if and only if

$$\delta \delta_1 = -\kappa_0 / \lambda,$$

where $\kappa_0$ is the extremal supplementary curvature and $\lambda$ is the equiparallel curvature of $F$ at $E$. 

...
Theorem 31. The two tangent turbine directions $T$ and $T'$ at an element $E$ of a special ruled field $F$ are reciprocal if and only if the sum of the supplementary curvatures $\kappa_s$ and $\kappa'_s$ of $T$ and $T'$ is zero. That is

$$\kappa_s + \kappa'_s = 0.$$  \hspace{1cm} (65)

Any two tangent turbine directions at an element $E$ of an umbilical field $F$ are reciprocal.

In general, given a one-parameter family of series $\phi(u, v) = \text{const.}$ of a general or special ruled field $F$, we can find another one-parameter family of series $\psi(u, v) = \text{const.}$ of $F$ such that the two tangent turbine directions of the two series of the two families passing through any element $E$ of $F$ are reciprocal at $E$.

The self-reciprocal directions of a field $F$ are called asymptotic directions.

Any series $S$ of a field $F$ whose tangent turbine direction at any element $E$ of $S$ is an asymptotic direction is termed an asymptotic series. The differential equation of all asymptotic series of a field $F$ is

$$w_{vv}v'^2 + (1 + w^2 + 2w_u)\nu' + (w_{uu} + w_u w_v) = 0.$$  \hspace{1cm} (66)

This means that a series $S$ is an asymptotic series if and only if the supplementary curvature $\kappa_s$ of the supplementary section tangent to $S$ at any element $E$ of $S$ is zero at $E$.

In a general field $F$, there are $2 \approx^1$ asymptotic series, all general series. A special ruled field $F$ possesses $2 \approx^1$ asymptotic series, namely, (1) the $\approx^1$ general series which satisfy (66), and (2) the $\approx^1$ equiparallel series of $F$. Every series of an umbilical field $F$ is an asymptotic series.

From Theorem 30 we pass to the following conclusion.

Theorem 32. Let $T_0$ be the principal tangent turbine (the principal direction), and let $T$ be any other tangent turbine at an element $E$ of a general field $F$. Let $\delta$ be the distance between the centers of $T_0$ and $T$. The tangent turbine direction $T$ is an asymptotic direction if and only if

$$\delta = (-\kappa_0/\lambda)^{1/2}. \hspace{1cm} (67)$$

The osculating flat field of a general series $S$ of a general or special ruled field $F$ at any element $E$ of $S$ will coincide with the tangent flat field of $F$ at $E$ if and only if the normal angles of the central elements of these two flat fields are identical. This means that the angle of equations (52) must be zero. Hence the supplementary curvature $\kappa_s$ of the supplementary section tangent to $S$ at $E$ is zero at $E$. Therefore $S$ is an asymptotic series. Thus we obtain

Theorem 33. A general series $S$ of a general or special ruled field $F$ is an
asymptotic series of $F$ if and only if its osculating flat fields coincide with the tangent flat fields of $F$ at the elements of $S$.

By this result and by (4) and (24), it follows that the torsion $\tau$ of an asymptotic series $S$ at any element $E$ of $S$ is the same as that of the geodesic series which is tangent to $S$ at $E$. Hence $\tau$ is given by (57). Upon squaring this value of $\tau$ and noting that $S$ satisfies (66), we obtain the following analogue of the Beltrami-Enneper theorem.

**Theorem 34.** The torsion $\tau$ of any asymptotic general series $S$ of a general or special ruled field $F$ at any element $E$ of $S$ is equal to the square root of the gaussian curvature $K$ of $F$ at $E$. That is,

$$\tau = K^{1/2}.$$  

Thus we have discussed in the geometry of the whirl-motion group $G_6$ the analogues of some of the classic theorems in the differential geometry of curves and surfaces embedded in a euclidean three-dimensional space. Of significant interest is the fact that our geometric configurations and invariants may be constructed by ordinary geometric means. Some of our results which seem to be completely analogous in content are nevertheless entirely distinct when we think of the meanings of the terms used.

**Bibliography**


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