

BOUNDARY BEHAVIOR OF HARMONIC FORMS ON A RANK ONE SYMMETRIC SPACE

BY

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ABSTRACT. We study the boundary behavior of 1-forms on a rank-one symmetric space M satisfying the equations $d\omega = 0 = \delta\omega$; the role of boundary is played by a nilpotent (Iwasawa) group \bar{N} of isometries of M . For forms satisfying certain H^p integrability conditions, we obtain the existence of boundary values in an appropriate sense, characterize these boundary values by means of fractional and singular integral operators on the group \bar{N} , and exhibit explicit isomorphisms between H^p spaces of forms on M and the ordinary L^p spaces of functions on the group \bar{N} .

Let M be a Riemannian manifold and let δ be the adjoint of the exterior differential d on M . The equation $d\omega = 0 = \delta\omega$ can be considered as a generalization of the classical Cauchy-Riemann equations. Their solutions were studied by Stein and Weiss in the case $M = \mathbf{R}^n \times \mathbf{R}^+$ with the euclidean metric (conjugate systems of harmonic functions) [9], by Korányi and Vági in the case when M is a euclidean ball in \mathbf{R}^n [7], and by Coifman and Weiss in the case $M = G \times \mathbf{R}^+$, G being a compact Lie group with the bi-invariant metric [1].

In this paper we consider the case when M is a noncompact symmetric space of rank one, define H^p spaces of 1-forms satisfying the above equations and study their boundary behavior. The role of boundary is played by a nilpotent group \bar{N} of isometries of M , which is the Cartan-conjugate of N in a fixed Iwasawa decomposition $G = KAN$ of the connected group of isometries of M [5].

A right action of the solvable group $\bar{S} = A\bar{N}$ induces a decomposition of the tangent bundle of M along the A and \bar{N} -directions ("vertical" and "horizontal" directions, by analogy with the upper half plane). It is shown that if a form ω is in certain H^p classes, its components along these directions have boundary values and that, moreover, the boundary values of the horizontal components can be obtained from the boundary values of the vertical

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component by means of fractional and singular integral operators on the group \bar{N} (the Riesz transforms). The kernels of these operators, introduced in §2, arise as left-invariant derivatives of a fundamental solution of a second-order hypoelliptic operator on the group \bar{N} . We are then able to recover a form ω in \mathbf{H}^p from the boundary value of its vertical component, and to establish in this manner canonical isomorphisms between the \mathbf{H}^p spaces of forms and the ordinary L^p spaces on the group \bar{N} .

1. **A class of vector fields on M .** The symmetric space M can be expressed as a homogeneous space $M = G/K$ where G is a semisimple group of isometries of M and K is a maximal compact subgroup of G . Let $\mathfrak{g}, \mathfrak{k}$ denote the Lie algebras of G and K , B the Killing form of \mathfrak{g} , and \mathfrak{p} the orthogonal complement of \mathfrak{k} in \mathfrak{g} relative to B . If $\pi: G \rightarrow G/K$ denotes the canonical projection, its differential at the identity, τ_* , identifies the subspace \mathfrak{p} of \mathfrak{g} with $T_0(M)$, the tangent space of M at the origin $o = \pi(e)$, and the invariant metric g on M can be chosen so that g_0 corresponds to the restriction of B to $\mathfrak{p} \times \mathfrak{p}$ under the above identification.

Let α be a maximal abelian subspace of \mathfrak{p} , $\Delta \subseteq \alpha^*$ the corresponding system of (restricted) roots, and for each $\alpha \in \Delta$ let \mathfrak{g}_α denote the corresponding root space. Let Δ_+ denote the system of positive roots relative to the choice of a fixed lexicographic ordering in α^* ; if $\alpha \in \Delta_+$, we write $\bar{\mathfrak{g}}_\alpha = \mathfrak{g}_{-\alpha}$. Then $\mathfrak{n} = \sum_{\alpha \in \Delta_+} \mathfrak{g}_\alpha$, $\bar{\mathfrak{n}} = \sum_{\alpha \in \Delta_+} \bar{\mathfrak{g}}_\alpha$ are nilpotent subalgebras of \mathfrak{g} , and if A, N, \bar{N} , are the connected subgroups of G with Lie algebras $\alpha, \mathfrak{n}, \bar{\mathfrak{n}}$, one has the Iwasawa decompositions $G = KAN$ and $G = \bar{N}AK$.

Now, $\bar{S} = \bar{N}A$ is a solvable subgroup of G and the above decomposition shows that every $p \in M$ can be uniquely written as $p = s \cdot o$ ($s \in \bar{S}$). We can therefore define a right action τ of \bar{S} on M by letting

$$\tau(s)(s' \cdot o) = s's \cdot o \quad (s, s' \in \bar{S}).$$

For each $X \in \bar{\mathfrak{s}} = \bar{\mathfrak{n}} + \alpha$ define a vector field \tilde{X} on M by letting

$$(1) \quad \tilde{X}_{na \cdot o} = \tau_*(\text{Ad}(a^{-1})X)_{na \cdot o} \quad (n \in \bar{N}, a \in A).$$

Now let X be any nonzero element of $\bar{\mathfrak{s}}$; for each fixed $a \in A$, the vector field $\tau_*(\text{Ad}(a^{-1})X)$ never vanishes on M , because $\text{Ad}(a^{-1})$ is an automorphism of $\bar{\mathfrak{s}}$ and τ is a free action; therefore, the same is true of the vector field \tilde{X} . It then follows that $X \rightarrow \tilde{X}$ maps a basis of $\bar{\mathfrak{s}}$ into a global frame of vector fields on M .

We shall now compute the Lie brackets and the inner products between vector fields induced in the above manner. Fix $a \in A$; the mapping $\Phi_a: \mathfrak{n} \rightarrow na \cdot o$ is a diffeomorphism from \bar{N} onto the submanifold $\bar{N}a \cdot o \subseteq M$, and if $X \in \bar{\mathfrak{n}}$, then $\tilde{X}_{na \cdot o} = (\Phi_a)_* X|_{na \cdot o}$. Therefore the vector fields \tilde{X} , ($X \in \bar{\mathfrak{n}}$) are tangent to the submanifolds $\bar{N}a \cdot o$, and the map $X \rightarrow \tilde{X}$ is a Lie algebra

homomorphism from \bar{n} into the Lie algebra of all smooth vector fields on M . On the other hand, if $n_0 \in \bar{N}$, $a_0 \in A$, the diffeomorphisms of \bar{S} given by $na \rightarrow nn_0a$ and $na \rightarrow na_0a$ commute with each other; it follows that if $H \in \alpha$, then $[\tilde{H}, \tilde{X}] = 0$ for every $X \in \bar{\mathfrak{g}}$.

Now, let $X \in \bar{\mathfrak{g}}$; the induced vector field \tilde{X} can be expressed as $\tilde{X}_{na \cdot o} = (na)_* (\pi_{*e}(\text{Ad}(a^{-1})X))$, and from the invariance of the Riemannian metric g , it follows that

$$(2) \quad g_{na \cdot o}(\tilde{X}, \tilde{Y}) = g_o(\pi_{*e}(\text{Ad}(a^{-1})X), \pi_{*e}(\text{Ad}(a^{-1})Y)) \quad (X, Y \in \bar{\mathfrak{g}}).$$

Let θ denote the Cartan involution associated to the decomposition $\mathfrak{g} = \mathfrak{k} + \mathfrak{p}$ and let (\cdot, \cdot) denote the inner product $-B(\cdot, \theta \cdot)$ on \mathfrak{g} . Then

$$(3) \quad \begin{aligned} g_o(\pi_{*e}(Z_1), \pi_{*e}(Z_2)) &= B(\frac{1}{2}(Z_1 - \theta Z_1), \frac{1}{2}(Z_2 - \theta Z_2)) \\ &= \frac{1}{2}(Z_1, Z_2) - \frac{1}{2}(Z_1, \theta Z_2). \end{aligned}$$

Now using (2) and (3), together with the fact that the subalgebras α , \mathfrak{n} , and $\bar{\mathfrak{n}} = \theta\mathfrak{n}$ are mutually orthogonal relative to the inner product (\cdot, \cdot) , one gets the following expression for the inner product of two arbitrary vector fields induced from $\bar{\mathfrak{g}}$:

$$(4) \quad g_{na \cdot o}((H_1 + X_1)^\sim, (H_2 + X_2)^\sim) = (H_1, H_2) + \frac{1}{2}(\text{Ad}(a^{-1})X_1, \text{Ad}(a^{-1})X_2),$$

where $H_1, H_2 \in \alpha$, $X_1, X_2 \in \bar{\mathfrak{n}}$.

In particular, if X is a root vector corresponding to the root $-\alpha$ ($\alpha \in \Delta_+$) and H is in α , then the induced vector fields \tilde{H} and \tilde{X} are orthogonal, \tilde{H} has constant length equal to $(H, H)^{1/2}$, and the length of \tilde{X} is given by

$$(5) \quad g_{na \cdot o}(\tilde{X}, \tilde{X}) = \frac{1}{2}e^{2\alpha(\log a)}(X, X).$$

In §3 it will be useful to have at our disposal a formula for the codifferential (or "divergence") $\delta\omega$ of a 1-form ω in terms of the vector fields introduced above. First of all, if Y_1, \dots, Y_n is an orthonormal frame defined on an open subset of a Riemannian manifold (V, g) , and ω is a 1-form on V , then $\delta\omega = \sum_i (Y_i \omega(Y_i) - t(Y_i)\omega(Y_i))$, where $t(Y_i) = \sum_j g([Y_i, Y_j], Y_j)$; this expression is easily derived from any of the standard definitions of the operator δ .

Now choose a basis $\{H_i, X_{\alpha,j}\}$ of the Lie algebra $\bar{\mathfrak{g}}$, such that $H_i \in \alpha$, $X_{\alpha,j} \in \bar{\mathfrak{g}}_\alpha$, and orthonormal with respect to the inner product $(\cdot, \cdot) = -B(\cdot, \theta \cdot)$. Letting e^α be the function on M whose value at the point $na \cdot o$ ($n \in \bar{N}$, $a \in A$) is $e^{\alpha(\log a)}$, from definition (1) it follows that $\tilde{H}_i = \tau_*(H_i)$ and $\tilde{X}_{\alpha,j} = e^\alpha \tau_*(X_{\alpha,j})$. Therefore (5) implies that $\{\tau_*(H_i), \sqrt{2}\tau_*(X_{\alpha,j})\}$ is an orthonormal frame on M . Since τ_* is a Lie algebra homomorphism, the relations $[H_i, X_{\alpha,j}] = -\alpha(H_i)X_{\alpha,j}$, $[X_{\alpha,j}, X_{\beta,k}] \in \bar{\mathfrak{g}}_{\alpha+\beta}$ imply that, relative to this ortho-

normal frame, $t(\tau_*(X_{\alpha,j})) = 0$ for all α, j , and $t(\tau_*(H_i)) = -\sum_{\alpha \in \Delta_+} m_\alpha \alpha(H_i) = -2\rho(H_i)$ (where as usual we put $m_\alpha = \dim \mathfrak{g}_\alpha$, $\rho = \frac{1}{2} \sum_{\alpha \in \Delta_+} m_\alpha \alpha$). Hence

$$\delta\omega = \sum_i \tau_*(H_i)\omega(\tau_*(H_i)) + 2 \sum_\alpha \sum_j \tau_*(X_{\alpha,j})\omega(\tau_*(X_{\alpha,j})) + 2 \sum_i \rho(H_i)\omega(\tilde{H}_i).$$

Now $\sum_i \rho(H_i)H_i = H_\rho$, the vector dual to ρ ; since $\tau_*(X_{\alpha,j}) = e^{-\alpha} \tilde{X}_{\alpha,j}$ and the vector field $\tilde{X}_{\alpha,j}$ annihilates the function $e^{-\alpha}$, one obtains

$$(6) \quad \delta\omega = \sum_i \tilde{H}_i \omega(\tilde{H}_i) + 2\omega(\tilde{H}_\rho) + 2 \sum_{\alpha \in \Delta_+} \sum_j e^{-2\alpha} \tilde{X}_{\alpha,j} \omega(\tilde{X}_{\alpha,j}).$$

Now assume that the rank of M is one. One can then choose a positive root α such that either $\Delta_+ = \{\alpha\}$ or $\Delta_+ = \{\alpha, 2\alpha\}$. Let $H_0 \in \mathfrak{a}$ be the vector such that $\alpha(H_0) = 1$; then

$$(H_0, H_0) = B(H_0, H_0) = \text{trace}(\text{ad } H_0)^2 = 2(m_\alpha + 4m_{2\alpha}).$$

Now set $\eta = (16(m_\alpha + 4m_{2\alpha}))^{-1}$ and $m = m_\alpha + 2m_{2\alpha}$; then the vector $(8\eta)^{1/2}H_0$ has unit length, and since $(H_\rho, H_0) = \rho(H_0) = \frac{1}{2}(m_\alpha + 2m_{2\alpha}) = m/2$, one obtains $H_\rho = 4\eta m H_0$. Let $\{X_i\}_{i=1}^{m_\alpha}$ and $\{Y_j\}_{j=1}^{m_{2\alpha}}$ be orthonormal bases of $\bar{\mathfrak{g}}_\alpha$ and $\bar{\mathfrak{g}}_{2\alpha}$, respectively; then (6) now implies

$$(7) \quad \begin{aligned} \delta\omega &= 8\eta \tilde{H}_0 \omega(\tilde{H}_0) + 8\eta m \omega(\tilde{H}_0) + 2e^{-2\alpha} \sum_i \tilde{X}_i \omega(\tilde{X}_i) \\ &+ 2e^{-4\alpha} \sum_j \tilde{Y}_j \omega(\tilde{Y}_j). \end{aligned}$$

For notational convenience we replace the vector field \tilde{H}_0 (which has constant length) by the vector field $\tilde{W} = e^{2\alpha} \tilde{H}_0$, which grows on the order of the vector fields induced by the elements of $\bar{\mathfrak{g}}_{2\alpha}$. Since $\tilde{H}_0(e^{-2\alpha}) = -2e^{-2\alpha}$ and $\tilde{H}_0 \omega(\tilde{H}_0) = e^{-4\alpha} \tilde{W} \omega(\tilde{W}) - 2e^{-2\alpha} \omega(\tilde{W})$, we obtain from (7):

LEMMA 1. *In the notation above,*

$$\begin{aligned} \frac{1}{8\eta} e^{2\alpha} \delta\omega &= e^{-2\alpha} \tilde{W} \omega(\tilde{W}) + (m - 2)\omega(\tilde{W}) + \frac{1}{4\eta} \sum_i \tilde{X}_i \omega(\tilde{X}_i) \\ &+ \frac{1}{4\eta} e^{-2\alpha} \sum_j \tilde{Y}_j \omega(\tilde{Y}_j). \end{aligned}$$

2. Riesz transforms on \bar{N} . This section is concerned with some analysis on the nilpotent group \bar{N} . In particular, we will construct some integral operators on this group which will play a central role in the characterization of the boundary values of harmonic forms on the space M .

From now on we assume that the rank of M is one; thus, $\bar{\mathfrak{n}} = \bar{\mathfrak{g}}_\alpha \oplus \bar{\mathfrak{g}}_{2\alpha}$. Let $\{X_i\}$, $1 \leq i \leq m_\alpha$, be an orthonormal basis of $\bar{\mathfrak{g}}_\alpha$, and define

$$(8) \quad L = \sum_i X_i^2;$$

then L is a differential operator on \bar{N} which is independent of the choice of orthonormal basis of \mathfrak{g}_α . Note that L is not elliptic, unless $\bar{g}_{2\alpha} = (0)$ (in this case, which occurs when M is a real hyperbolic space, \bar{N} is isomorphic to \mathbf{R}^n and L becomes the standard Laplacian); however, since any basis of \bar{g}_α constitutes a system of generators for the Lie algebra $\bar{\mathfrak{n}}$, a result of Hörmander [4] implies that L is always hypoelliptic. Our first objective will be to obtain an explicit fundamental solution for this operator.

For $X \in \mathfrak{g}$, set $|X| = (X, X)^{1/2}$. The following lemma holds without restriction on the rank of M .

LEMMA 2. *Let α be a restricted root, $X \in \bar{g}_\alpha$, and $Y \in \bar{g}_{2\alpha}$. Then $[[Y, \theta X]] = \sqrt{2}|\alpha||X||Y|$. If $\{X_i\}$, $1 \leq i \leq m_\alpha$, is an orthonormal basis of \bar{g}_α , then*

$$\sum_{i=1}^{m_\alpha} [[X, X_i]]^2 = 2m_{2\alpha}|\alpha|^2|X|^2.$$

PROOF. We have

$$[[Y, \theta X]]^2 = -B([Y, \theta X], [\theta Y, X]) = B([[Y, \theta X], X], \theta Y).$$

Since $[\bar{g}_\alpha, \bar{g}_{2\alpha}] = (0)$, $[[Y, \theta X], X] = -[Y, [X, \theta X]]$ by Jacobi's identity. Now $[X, \theta X] = -B(X, \theta X)H_\alpha = |X|^2H_\alpha$ and $[Y, H_\alpha] = 2\alpha(H_\alpha)Y = 2|\alpha|^2Y$. Therefore

$$[[Y, \theta X]]^2 = -2|\alpha|^2|X|^2B(Y, \theta Y) = 2|\alpha|^2|X|^2|Y|^2,$$

proving the first identity.

Now let $\{X_i\}$, $1 \leq i \leq m_\alpha$, be as above, and choose an orthonormal basis $\{Y_j\}$, $1 \leq j \leq m_{2\alpha}$, of $\bar{g}_{2\alpha}$. Then

$$[[X, X_i]]^2 = \sum_{j=1}^{m_{2\alpha}} (Y_j, [X, X_i])^2 = \sum_{j=1}^{m_{2\alpha}} ([Y_j, \theta X], X_i)^2;$$

adding over $i = 1, \dots, m_\alpha$, one gets

$$\sum_{i=1}^{m_\alpha} [[X, X_i]]^2 = \sum_{j=1}^{m_{2\alpha}} [[Y_j, \theta X]]^2.$$

But the first part of the lemma shows that every term in the last sum is equal to $2|\alpha|^2|X|^2$, finishing the proof.

Any element $n \in \bar{N}$ can be written uniquely as $n = \exp(X + Y)$, with $X \in \bar{g}_\alpha$, $Y \in \bar{g}_{2\alpha}$. A function F on \bar{N} will be called *biradial* if there is a function $f(u, v)$ of two real variables such that

$$F(\exp(X + Y)) = f(|X|^2, |Y|^2).$$

LEMMA 3. Let F be a smooth biradial function on \bar{N} , $F(\exp(X + Y)) = f(|X|^2, |Y|^2) = f(u, v)$; let $X' \in \bar{g}_\alpha$, $Y' \in \bar{g}_{2\alpha}$ and let L be the operator defined by (8). Then, for $n = \exp(X + Y)$,

- (i) $(X'F)(n) = 2(X', X)\partial f/\partial u + ([X, X'], Y)\partial f/\partial v,$
- (ii) $(Y'F)(n) = 2(Y', Y)\partial f/\partial v,$
- (iii) $(LF)(n) = 4|X|^2\partial^2 f/\partial u^2 + 2|\alpha|^2|X|^2|Y|^2\partial^2 f/\partial v^2$
 $+ 2m_\alpha\partial f/\partial u + m_{2\alpha}|\alpha|^2|X|^2\partial f/\partial v.$

PROOF. For $t \in \mathbf{R}$, one has

$$\exp(X + Y)\exp tX' = \exp(X + tX' + Y + \frac{1}{2}t[X, X']).$$

Therefore

$$F(n \exp tX') = f(|X + tX'|^2, |Y + \frac{1}{2}t[X, X']|^2),$$

and

$$\frac{d}{dt}F(n \exp tX') = 2((X, X') + t|X'|^2)\frac{\partial f}{\partial u} + ((Y, [X, X']) + \frac{1}{2}t|[X, X']|^2)\frac{\partial f}{\partial v}.$$

Therefore

$$(X'F)(n) = \left. \frac{d}{dt}F(n \exp tX') \right|_{t=0} = 2(X', X)\frac{\partial f}{\partial u} + ([X, X'], Y)\frac{\partial f}{\partial v},$$

showing (i); also

$$\begin{aligned} (X'^2F)(n) &= \left. \frac{d^2}{dt^2}F(n \exp tX') \right|_{t=0} \\ (9) \quad &= 4(X, X')^2\frac{\partial^2 f}{\partial u^2} + (Y, [X, X'])^2\frac{\partial^2 f}{\partial v^2} + 4(X, X')(Y, [X, X'])\frac{\partial^2 f}{\partial u\partial v} \\ &\quad + 2|X'|^2(\partial f/\partial u) + \frac{1}{2}[X, X']^2(\partial f/\partial v). \end{aligned}$$

Now, if $\{X_i\}$, $1 \leq i \leq m_\alpha$, is an orthonormal basis of \bar{g}_α , one has $\sum_i (X, X_i)^2 = |X|^2$, $\sum_i |X_i|^2 = m_\alpha$, $\sum_i (X, X_i)(Y, [X, X_i]) = (Y, [X, X]) = 0$; also, Lemma 2 implies

$$\begin{aligned} \sum_i (Y, [X, X_i])^2 &= \sum_i ([Y, \theta X], X_i)^2 = \|[Y, \theta X]\|^2 \\ &= 2|\alpha|^2 |X|^2 |Y|^2 \end{aligned}$$

and

$$\sum_i \|[X, X_i]\|^2 = 2m_{2\alpha} |\alpha|^2 |X|^2.$$

Letting $X' = X_i$ in (9), it now follows that $(LF)(n) = \sum_i (X_i^2 F)(n)$ is given by (iii). Since (ii) is clear, this proves the lemma.

For each $\epsilon \geq 0$, define a function $\|n\|_\epsilon$ on \bar{N} by $\|n\|_\epsilon = \eta(|X|^2 + \epsilon^2)^2 + |Y|^2$, where $n = \exp(X + Y)$, $X \in \bar{g}_\alpha$, $Y \in \bar{g}_{2\alpha}$, and η denotes the constant $|\alpha|^2/8 = (16(m_\alpha + 4m_{2\alpha}))^{-1}$. One verifies that under the "dilations" of \bar{N} induced by the group A ,

$$(10) \quad \|ana^{-1}\|_\epsilon = e^{-4\alpha(\log a)} \|n\|_{\epsilon'}, \quad \epsilon' = e^{\alpha(\log a)} \epsilon, \quad a \in A.$$

In particular, $\|n\|_0 = \eta|X|^4 + |Y|^2$ is a "gauge" on \bar{N} , in the sense of Korányi and Vági [6], which satisfies the homogeneity condition

$$(11) \quad \|ana^{-1}\|_0 = e^{-4\alpha(\log a)} \|n\|_0, \quad a \in A.$$

From now on we shall assume that $m = m_\alpha + 2m_{2\alpha} > 2$; this excludes the cases when M is the real hyperbolic space of dimension 2 or 3. Now, set $k = \frac{1}{4}(m_\alpha + 2m_{2\alpha} - 2) = \frac{1}{4}(m - 2)$; since $\dim \bar{N} = m_\alpha + m_{2\alpha}$, the function $\|n\|_\epsilon^{-k-1}$ is integrable on \bar{N} for every $\epsilon > 0$; one can then introduce a constant β by $\beta^{-1} = -4k \eta m_\alpha \int_{\bar{N}} \|n\|_1^{-k-1} dn$.

THEOREM 1. *The function $G(n) = \beta \|n\|_0^{-k}$ is a fundamental solution for the operator L .*

PROOF. A straightforward application of Lemma 3 to the biradial function $\|n\|_\epsilon^{-k}$ shows that

$$(12) \quad L(\|n\|_\epsilon^{-k}) = -4k \eta m_\alpha \epsilon^2 \|n\|_\epsilon^{-k-1}.$$

In particular, $LG(n) = 0$ for all $n \neq e$.

We claim now that $L(\beta \|n\|_\epsilon^{-k})$ is an approximate identity as $\epsilon \rightarrow 0$. In fact, by (12) one has

$$\int_{\bar{N}} L(\beta \|n\|_\epsilon^{-k}) dn = \left[\int_{\bar{N}} \|n\|_1^{-k-1} dn \right]^{-1} \int_{\bar{N}} \epsilon^2 \|n\|_\epsilon^{-k-1} dn.$$

But, via the change of variables $n \rightarrow a^{-1}na$, with $a = \exp \epsilon H_0$, $\alpha(H_0) = 1$, formula (10) shows that

$$\int_{\bar{N}} \varepsilon^2 \|n\|_{\varepsilon}^{-k-1} dn = \int_{\bar{N}} \|n\|_1^{-k-1} dn \quad (\varepsilon > 0).$$

Therefore $L(\beta\|n\|_{\varepsilon}^{-k})$ is a positive L^1 -function on \bar{N} with L^1 -norm one. Moreover, the same change of variables shows that for any $\delta > 0$,

$$\begin{aligned} \int_{\|n\|>\delta} L(\|n\|_{\varepsilon}^{-k}) dn &= -4km_{\alpha}\eta \int_{\|n\|>\delta} \varepsilon^2 \|n\|_{\varepsilon}^{-k-1} dn \\ &= -4km_{\alpha}\eta \int_{\|n\|>\delta/\varepsilon^4} \|n\|_1^{-k-1} dn, \end{aligned}$$

so that $\lim_{\varepsilon \rightarrow 0} \int_{\|n\|>\delta} L(\|n\|_{\varepsilon}^{-k}) dn = 0$.

Now, if f is a continuous function on \bar{N} with compact support, and g is smooth, then $D(f * g) = f * Dg$ for any left-invariant differential operator D on \bar{N} , where the convolution is given by $f * g(n) = \int_{\bar{N}} f(n_1)g(n^{-1}, n) dn_1$. Hence

$$L(f * G) = \lim_{\varepsilon \rightarrow 0} L(f * \beta\|n\|_{\varepsilon}^{-k}) = \lim_{\varepsilon \rightarrow 0} f * L(\beta\|n\|_{\varepsilon}^{-k}) = f.$$

This finishes the proof of the theorem.

REMARK. In the case when \bar{N} is the Heisenberg group ($M =$ complex hyperbolic space) this fundamental solution was obtained by Folland [2].

DEFINITION. For $Z \in \bar{n}$, the Z -Riesz kernel is the function $r_Z(n) = ZG(n)$, $n \neq e$.

A straightforward application of Lemma 3 gives the following expressions of the above kernels:

$$r_Z(n) = -k\beta\|n\|_0^{-k-1} [4\eta|X|^2(Z, X) + ([X, Z], Y)] \quad \text{for } Z \in \bar{g}_{\alpha}, \text{ and}$$

$$r_Z(n) = -2k\beta\|n\|_0^{-k-1} (Y, Z) \quad \text{for } Z \in \bar{g}_{2\alpha}.$$

Under the adjoint action of A on \bar{N} , these kernels satisfy the homogeneity properties

$$r_Z(a^{-1}na) = e^{-(m-1)\alpha(\log a)} r_Z(n) \quad \text{for } Z \in \bar{g}_{\alpha}, \text{ and}$$

$$r_Z(a^{-1}na) = e^{-m\alpha(\log a)} r_Z(n) \quad \text{for } Z \in \bar{g}_{2\alpha};$$

in the latter case, one also has $r_Z(n^{-1}) = -r_Z(n)$. Thus, the functions r_Z can be considered as fractional integral kernels (for $Z \in \bar{g}_{\alpha}$) and singular integral kernels (for $Z \in \bar{g}_{2\alpha}$), and a standard argument gives

THEOREM 2. *The convolution operators defined by $R_Z f = f * r_Z$ are:*

- (a) *bounded operators from $L^p(\bar{N})$ into $L^p(\bar{N})$, for $1 < p < \infty$ and $Z \in \bar{g}_{2\alpha}$,*
- (b) *bounded operators from $L^p(\bar{N})$ into $L^q(\bar{N})$, for $1 < p < m$, $1/q = 1/p - 1/m$, and $Z \in \bar{g}_{\alpha}$.*

For more general results on fractional and singular integral operators on nilpotent groups see Stein [8] and Korányi and Vági [6].

We also note that these Riesz transforms satisfy the formal properties

$$(13) \quad \begin{aligned} X \circ R_Y - Y \circ R_X &= R_{[X,Y]} & (X, Y \in \bar{n}), \\ \sum_i X_i \circ R_{X_i} &= \text{identity} & (\{X_i\} \text{ orthonormal basis of } \bar{u}_\alpha). \end{aligned}$$

We finish this section with a technical result that will be needed later.

LEMMA 4. *If F_1, F_2 are biradial functions on \bar{N} , then $F_1 * F_2 = F_2 * F_1$.*

PROOF. With $n = \exp(X + Y)$, $n_1 = \exp(X_1 + Y_1)$ and $F_i(\exp(X + Y)) = f_i(|X|^2, |Y|^2)$, one has

$$\begin{aligned} F_1 * F_2(n_1) &= \int_{\bar{N}} F_1(n) F_2(n^{-1}n_1) dn \\ &= \int f_1(|X|^2, |Y|^2) f_2(|X_1 - X|^2, |Y_1 - Y - \frac{1}{2}[X, X_1]|^2) dX dY. \end{aligned}$$

Under the orthogonal change of variables $X \rightarrow X' = 2((X, X_1)/|X_1|^2)X_1 - X$, one has, $|X_1 - X'| = |X_1 - X|$ and $[X, X_1] = -[X', X_1]$. Thus,

$$\begin{aligned} F_1 * F_2(n_1) &= \int f_1(|X|^2, |Y|^2) f_2(|X_1 - X|^2, |Y_1 - Y + \frac{1}{2}[X, X_1]|^2) dX dY \\ &= \int_{\bar{N}} F_1(n) F_2(nn^{-1}) dn = F_2 * F_1(n_1), \end{aligned}$$

proving the lemma.

3. Boundary values of harmonic forms. We regard the nilpotent group \bar{N} as a boundary for the symmetric space M , as, for example, in [5]. A function Φ on M is said to be *uniformly in L^p* if there exists a constant K such that $\int_{\bar{N}} |\Phi(na \cdot o)|^p dn < K$ for all $a \in A$. If the family of functions $\{\Phi_a, a \in A\}$ on \bar{N} given by $\Phi_a(n) = \Phi(na \cdot o)$ converges to a function φ on \bar{N} as $a \rightarrow \infty$ (in L^p , uniformly, etc.), one simply says that Φ *converges to φ* or that φ is *the boundary value of Φ* . The following facts are well known: Let $1 < p < \infty$ and let Φ be a harmonic function on M which is uniformly in L^p . Then Φ has a boundary value $\varphi \in L^p(\bar{N})$ to which it converges in L^p and almost everywhere. The Poisson kernel is the function P on $\bar{N} \times A$ defined by

$$P(n, a) = P_a(n) = \exp\{-2\rho(H(a^{-1}na) - \log a)\}$$

where for each $g \in G$, $H(g)$ is the unique element in \mathfrak{a} such that $g = k \exp H(g)n$ ($k \in K, n \in N$). If $\varphi \in L^p(\bar{N})$ one defines the Poisson integral of φ as the function Φ on M given by

$$\Phi(na \cdot o) = \varphi * P_a(n) = \int_N \varphi(n_1) P_a(n^{-1}n) dn_1;$$

then Φ is harmonic on M , it is uniformly in L^p , and its boundary value is the original function φ .

We shall now define H^p spaces of differential 1-forms. The need for different integrability conditions for different components will be apparent in Theorems 3 and 4.

DEFINITION. For $1 < p < m$ ($m = m_\alpha + 2m_{2\alpha}$), let q be given by $1/q = 1/p - 1/m$, and let H^p be the space of all 1-forms on M such that

- (i) $d\omega = 0 = \delta\omega$.
- (ii) The functions $\omega(\tilde{X})$ ($X \in \bar{g}_\alpha$), are uniformly in L^q .
- (iii) The functions $\omega(\tilde{W}), \omega(\tilde{Y})$ ($Y \in \bar{g}_{2\alpha}$) and $\sum_i \tilde{X}_i \omega(\tilde{X}_i)$ ($\{\tilde{X}_i\}$ orthonormal basis of \bar{g}_α) are uniformly in L^p .

REMARK. It should be pointed out that the components $\omega(\tilde{W})$ and $\omega(\tilde{X})$ are *not* harmonic (except when X belongs to the center of \bar{n} , a fact that is used below). Although replacing these vector fields by the infinitesimal isometries induced by \bar{n} would give harmonic components, such vector fields are inappropriate in the present context.

Our next objective is to associate a form $\omega_f \in H^p$ to each function $f \in L^p(\bar{N})$. For this we will need a kernel Q on M defined by

$$Q(na \cdot o) = Q_a(n) = c\tilde{W}(P_a * G(n)) = c\tilde{W} \int_N P_a(g)G(g^{-1}n) dg$$

where G is the fundamental solution for the operator L introduced in §2, and c is a constant to be specified later.

Under conjugation by A , the kernel Q satisfies the homogeneity property

$$(14) \quad Q_a(a^{-1}na) = e^{-m\alpha(\log a)} Q_{a'}(n).$$

In fact, for the Poisson kernel one has $P_a(a^{-1}na) = e^{-m\alpha(\log a)} P_{a'}(n)$ and therefore, $\tilde{W}P_a(a^{-1}na) = e^{-(m+2)\alpha(\log a)} \tilde{W}P_{a'}(n)$; on the other hand, $G(a^{-1}na) = e^{-(m-2)\alpha(\log a)} G(n)$ so that (14) follows by the definition of Q as the convolution of $\tilde{W}P_a$ with G .

LEMMA 5. For an appropriate choice of the constant c , the kernel Q_a is an approximate identity on \bar{N} , that is:

- (i) $Q_a(n) > 0$ and $Q_a \in L^1(\bar{N})$,
- (ii) $\int_{\bar{N}} Q_a(n) dn = 1$ for all $a \in A$,
- (iii) $\lim_{t \rightarrow \infty} \int_{\bar{N}(\epsilon)} Q_{a_t}(n) dn = 1$ for all $\epsilon > 0$, where $a_t = \exp tH_0$ and $\bar{N}(\epsilon) = \{n \in \bar{N}: \|n\| \leq \epsilon\}$.

PROOF. The proof of the integrability of Q_a is more involved than one would perhaps expect.

First of all, since $P_{a_t} * G(n)$ is a monotone function of t , the kernel $\tilde{W}P_{a_t} * G(n) = \tilde{W}P_{a_t} * G(n) = e^{2t} (d/dt) P_{a_t} * G(n)$ has constant sign; therefore, by Fubini's Theorem, the integrability of Q_a will follow from the existence of the iterated integral

$$(15) \quad \int_{\bar{g}_\alpha} \int_{\bar{g}_{2\alpha}} (\tilde{W}P_a * G)(\exp(X + Y)) dX dY$$

(in this proof we systematically use X, X' and Y, Y' to denote elements of \bar{g}_α and $\bar{g}_{2\alpha}$ respectively). Secondly, because of the homogeneity property (14) it is enough to consider the case $a = e = \text{identity}$.

Now, if $n = \exp(X + Y)$ and $n' = \exp(X' + Y')$, then

$$n'^{-1}n = \exp(X - X' + Y - Y' + \frac{1}{2}[X, X']),$$

where $X - X' \in \bar{g}_\alpha$ and $Y - Y' + \frac{1}{2}[X, X'] \in \bar{g}_{2\alpha}$. We then have

$$(16) \quad \begin{aligned} & \tilde{W}P_e * G(|X|, |Y|) \\ &= \int_{\bar{n}} \tilde{W}P_e(|X'|, |Y'|) G(|X - X'|, |Y - Y' + \frac{1}{2}[X, X']|) dX dY \end{aligned}$$

where we write $F(|X|, |Y|) = F(\exp(X + Y))$ whenever F is a biradial function of \bar{N} . Integrating over $\bar{g}_{2\alpha}$, and since the integrand in (16) has constant sign as a function of Y , we can exchange the order of the integrations over $\bar{g}_{2\alpha}$ and \bar{n} and get

$$(17) \quad \begin{aligned} & \int_{\bar{g}_{2\alpha}} \tilde{W}P_e * G(|X|, |Y|) dY \\ &= \int_{\bar{n}} \tilde{W}P_e(|X'|, |Y'|) \left(\int_{\bar{g}_{2\alpha}} G(|X - X'|, |Y - Y' + \frac{1}{2}[X, X']|) dY \right) dX' dY' \\ &= \int_{\bar{n}} \tilde{W}P_e(|X'|, |Y'|) \left(\int_{\bar{g}_{2\alpha}} G(|X - X'|, |Y|) dY \right) dX' dY' \end{aligned}$$

because of invariance under translations. Since

$$G(|X - X'|, |Y|) = \beta(\eta|X - X'|^4 + |Y|^2)^{-k},$$

the change of variable $Y \rightarrow \eta^{1/2}|X - X'|^2 Y$ shows that

$$\int_{\bar{g}_{2\alpha}} G(|X - X'|, |Y|) dY = \beta\eta^{1/2m_{2\alpha}-4k} |X - X'|^{2m_{2\alpha}-4k} \int_{\bar{g}_{2\alpha}} (1 + |Y|^2)^{-k} dY.$$

Since under the assumption $m_\alpha + 2m_{2\alpha} > 2$ one has

$$2k = \frac{1}{2}(m_\alpha + 2m_{2\alpha} - 2) > m_{2\alpha} = \dim \bar{g}_{2\alpha},$$

the last integral exists, and we conclude that

$$(18) \quad \int_{\bar{g}_{2\alpha}} \tilde{W}P_2 * G(|X|, |Y|) dY \\ = C_1 \int_{\bar{\pi}} \tilde{W}P_e(|X'|, |Y'|) |X - X'|^{2m_{2\alpha} - 4k} dX' dY'$$

where C_1 is a constant. Now, let $\Sigma = \{X \in \bar{g}_\alpha : |X| = 1\}$ and introduce polar coordinates in \bar{g}_α : $X' = \rho\xi$, $0 \leq \rho$, $\xi \in \Sigma$; since $2m_{2\alpha} - 4k = -m_\alpha + 2$ the last integral becomes

$$(19) \quad \int_0^\infty I(\rho) \left[\int_\Sigma |X - \rho\xi|^{-m_\alpha + 2} d\xi \right] \rho^{m_\alpha - 1} d\rho$$

where $I(\rho) = \int_{\bar{g}_{2\alpha}} \tilde{W}P_e(\rho, |Y|) dY$. Now, consider the function

$$X \rightarrow \int_\Sigma |\rho^{-1}X - \xi|^{-m_\alpha + 2} d\xi;$$

it is continuous on \bar{g}_α , smooth and harmonic for $|X| \neq \rho$ (with respect to the Laplacian associated with the inner product (\cdot, \cdot)), invariant under the corresponding rotation group, equal to $\sigma =$ measure of Σ for $X = 0$, and not identically constant. One can therefore see that

$$\int_\Sigma |\rho^{-1}X - \xi|^{-m_\alpha + 2} d\xi = \begin{cases} \sigma & \text{for } |X| \leq \rho, \\ \sigma \rho^{m_\alpha - 2} |X|^{-m_\alpha + 2} & \text{for } |X| \geq \rho, \end{cases}$$

which implies that (19) can be rewritten as

$$(20) \quad \sigma |X|^{-m_\alpha + 2} \int_0^{|X|} I(\rho) \rho^{m_\alpha - 1} d\rho + \sigma \int_{|X|}^\infty I(\rho) \rho d\rho.$$

In order to evaluate $I(\rho)$ we make use of the formula [3, p. 65]

$$P_\alpha(|X|, |Y|) = e^{-mt} [(e^{-2t} + c|X|^2)^2 + 4c|Y|^2]^{-m/2}$$

where $c = (4(m_\alpha + 4m_{2\alpha}))^{-1}$, $m = m_\alpha + 2m_{2\alpha}$, which yields

$$I(\rho) = \int_{\bar{g}_{2\alpha}} \tilde{W}P_e(\rho, |Y|) dY \\ = -m \int_{\bar{g}_{2\alpha}} [(1 + c\rho^2)^2 + 4c|Y|^2]^{-m/2 - 1} \\ \times [(1 + c\rho^2)^2 + 4c|Y|^2 - 2(1 + c\rho^2)] dY.$$

After the change of variable $Y \rightarrow (1 + c\rho^2)Y$, this becomes

$$I(\rho) = \text{const} \cdot [2b(1 + c\rho^2)^{-(m_\alpha+m_{2\alpha}+1)} - (1 + c\rho^2)^{-(m_\alpha+m_{2\alpha})}]$$

where

$$b = \left[\int_{\bar{g}_{2\alpha}} (1 + 4c|Y|^2)^{-m/2+1} dY \right]^{-1} \int_{\bar{g}_{2\alpha}} (1 + 4c|Y|^2)^{-m/2} dY.$$

Evaluating these two integrals gives $b = m^{-1}(m_\alpha + m_{2\alpha})$. Now, substituting this expression for $I(\rho)$ into (20), one checks by differentiation that (20) equals $C(1 + c|X|^2)^{-(m_\alpha+m_{2\alpha}-1)}$, where the constant C is given by $C = \sigma(m_\alpha - 2) \cdot [2cm(m_\alpha + m_{2\alpha} - 1)]^{-1}$. Therefore,

$$(21) \quad \int_{\bar{g}_{2\alpha}} \tilde{W}P_e * G(|X|, |Y|) dY = \text{const} \cdot (1 + c|X|^2)^{-(m_\alpha+m_{2\alpha}-1)}.$$

Since $2(m_\alpha + m_{2\alpha} - 1) > m_\alpha = \dim \bar{g}_\alpha$, the above expression is integrable over \bar{g}_α ; we have therefore shown that (15) exists, so that $\tilde{W}P_a * G$ is integrable. Now choose c so that $Q_e = c\tilde{W}P_e * G$ has integral equal to one; then $Q_e > 0$ and by the homogeneity property (14) one has

$$(22) \quad Q_a(n) = e^{m\alpha(\log a)} Q_e(a^{-1}na)$$

so that $Q_a > 0$ for all $a \in A$. This finishes the proof of (i).

Integrating (22) and recalling that the Jacobian of the change of variable $n \rightarrow ana^{-1}$ is $e^{-m\alpha(\log a)}$, one gets

$$\int_N Q_a(n) dn = e^{m\alpha(\log a)} \int_N Q_e(a^{-1}na) dn = \int_N Q_e(n) dn = 1,$$

proving (ii). On the other hand, letting $a = a_t$ in (22) and integrating over $\bar{N}(\epsilon)$ one gets

$$\int_{\bar{N}(\epsilon)} Q_{a_t}(n) dn = e^{mt} \int_{\bar{N}(\epsilon)} Q_e(a_t^{-1}na_t) dn = \int_{a_t^{-1}\bar{N}(\epsilon)a_t} Q_e(n) dn.$$

Since $\|a_t^{-1}na_t\| = e^{4t}\|n\|$, it follows that $a_t^{-1}\bar{N}(\epsilon)a_t = \bar{N}(e^{4t}\epsilon)$. Therefore, as $t \rightarrow +\infty$, the last integral converges to $\int_N Q_e(n) dn = 1$. This shows (iii) and finishes the proof of the lemma.

REMARK. Although the kernel Q_a is not an elementary function in general, formula (21) gives its integral over $\bar{g}_{2\alpha}$. In particular, when $\bar{g}_{2\alpha} = (0)$ (case of real hyperbolic space) one gets the explicit expression

$$Q_{a_t}(X) = C e^{(m_\alpha-2)t} / (e^{2t} + |X|^2)^{m_\alpha-1}.$$

Now, for each function f on \bar{N} for which the following convolutions make sense, define a 1-form ω_f on M by

$$(23) \quad \begin{aligned} \omega_f(\tilde{W})_{na \cdot o} &= f * Q_a(n), \\ \omega_f(\tilde{X})_{na \cdot o} &= f * P_a * r_X(n) \quad (X \in \bar{\pi}). \end{aligned}$$

As before, all convolutions are taken on the group \bar{N} . Note that, if f is sufficiently nice, then $f * P * G$ exists and ω_f is just the exterior derivative of this function; but this convolution is not in general defined for $f \in L^p(\bar{N})$ and p in the full range $1 < p < m$.

THEOREM 3. *Let $1 < p < m$ and let $f \in L^p(\bar{N})$. Then*

- (i) ω_f is well defined and belongs to \mathbf{H}^p ,
- (ii) $\omega_f(\tilde{W})$ converges in L^p and a.e. to f ,
- (iii) for each $X \in \bar{g}_\alpha$ (resp. $X \in \bar{g}_{2\alpha}$), $\omega_f(\tilde{X})$ converges in L^q (resp. L^p) and a.e. to the Riesz transform $R_X f = f * r_X$.

PROOF. For $f \in L^p(\bar{N})$, the convolutions $f * P$ and $f * Q$ are smooth functions on M which are uniformly in $L^p(\bar{N})$. Since the Riesz transforms are defined in $L^p(\bar{N})$, it follows that ω_f is a well-defined differential form on M .

Now notice that if $f(n)$ and $F(na \cdot o) = F_a(n)$ are functions on \bar{N} and M , respectively, then derivatives of the convolution $f * F_a$ with respect to the vector fields \tilde{X} satisfy $\tilde{X}(f * F_a) = f * XF_a$. Therefore, for $X, Y \in \bar{\pi}$ one has

$$\tilde{X}\omega_f(\tilde{Y}) - \tilde{Y}\omega_f(\tilde{X}) = f * P * (Xr_Y - Yr_X) = f * P * r_{[X,Y]} = \omega_f([\tilde{X}, \tilde{Y}]).$$

Also, since \tilde{X} and \tilde{W} commute, $\tilde{X}\omega_f(\tilde{W}) - \tilde{W}\omega_f(\tilde{X}) = 0 = \omega_f([\tilde{X}, \tilde{W}])$. Since the vector fields \tilde{X}, \tilde{W} ($X \in \bar{\pi}$) span the tangent space to M at every point, it follows that ω_f is closed. In order to show that it is coclosed, we may assume that f is, say, continuous of compact support, the result for a general $f \in L^p(\bar{N})$ following by a standard approximation argument. Then $\delta\omega_f = \delta d(f * P * G) = \Delta(f * P * G)$, where Δ is the Laplace-Beltrami operator on M ; since P and G are both biradial, the function $f * P * G = f * G * P$ is harmonic, and so is annihilated by Δ .

The Poisson integral of f is uniformly in $L^p(\bar{N})$ and it converges in L^p to its boundary value f . Since convolution against the Riesz kernels r_X ($X \in \bar{g}_\alpha$) is a bounded operator from $L^p(\bar{N})$ into $L^q(\bar{N})$, one concludes that the function $\omega_f(\tilde{X})$ (which, we recall, is not always harmonic) is uniformly in $L^q(\bar{N})$ and that it converges to $R_X f = f * r_X$ in L^q -norm. The same argument applies to the components $\omega(\tilde{X})$ when $X \in \bar{g}_{2\alpha}$, giving a bounded operator from L^p into L^p . If X_1, \dots, X_n is an orthonormal basis of \bar{g}_α , $\sum X_i r_{X_i}$ is the Dirac-distribution on \bar{N} ; therefore $\sum_i \tilde{X}_i \omega(\tilde{X}_i) = \sum f * X_i (P * r_{X_i}) = f * P$ and so this function is also uniformly in $L^p(\bar{N})$. Finally, Q being an approximate identity, the component $\omega_f(\tilde{W}) = f * Q$ is clearly uniformly in $L^p(\bar{N})$ and it converges in L^p -norm to f . Q.E.D.

In the next theorem we show that the mapping $f \rightarrow \omega_f$ from $L^p(\bar{N})$ into \mathbf{H}^p is actually an isomorphism onto. Together with Theorem 3 this implies that every form ω in \mathbf{H}^p has boundary values in the appropriate sense, and that the boundary values of the components of ω along the vector fields \tilde{X} ($X \in \bar{n}$) are precisely the Riesz transforms of the boundary value of the component of ω along the vector field \tilde{W} .

THEOREM 4. *Let $\omega \in \mathbf{H}^p$. Then there exists a unique $f \in L^p(\bar{N})$ such that $\omega = \omega_f$.*

PROOF. Let X_1, \dots, X_{m_α} be an orthonormal basis of \bar{g}_α , $\tilde{L} = \sum \tilde{X}_i^2$, and let Y be in the center of \bar{n} . For any closed form ξ on M , a straightforward computation using Lemma 1 shows that $\Delta(\sum_i \tilde{X}_i \xi(\tilde{X}_i)) = \tilde{L}\delta\xi$ and $\Delta\xi(\tilde{Y}) = \tilde{Y}\delta\xi$. Therefore $\omega \in \mathbf{H}^p$ implies that the functions $\sum \tilde{X}_i \omega(\tilde{X}_i)$ and $\omega(\tilde{Y})$ are harmonic; they can then be written as Poisson integrals $\sum_i \tilde{X}_i \omega(\tilde{X}_i) = f * P$ and $\omega(\tilde{Y}) = \varphi * P$, with $f, \varphi \in L^p(\bar{N})$. Let ω_f be the form associated to f by equation (23); then $\omega_f(\tilde{Y}) = f * P * r_Y = \sum \tilde{X}_i \omega(\tilde{X}_i) * YG$, and since Y is in the center of \bar{n} , this last expression is equal to $\tilde{Y}(\sum_i \tilde{X}_i \omega(\tilde{X}_i)) * G$. But $\tilde{Y}(\sum \tilde{X}_i \omega(\tilde{X}_i)) = \tilde{L}\omega(\tilde{Y})$, so we can write

$$(24) \quad \omega_f(\tilde{Y}) = \tilde{L}\omega(\tilde{Y}) * G = \tilde{L}(\varphi * P) * G = \varphi * \tilde{L}P * G.$$

Now the functions G, P and $\tilde{L}P$ are biradial and since G is a fundamental solution for the operator L , one has

$$\tilde{L}P * G = G * \tilde{L}P = \tilde{L}(G * P) = \tilde{L}(P * G) = P.$$

Substituting in (24) we get $\omega_f(\tilde{Y}) = \varphi * P = \omega(\tilde{Y})$.

Now set $\omega' = \omega - \omega_f$; then $\omega' \in \mathbf{H}^p$ and if Y is in the center of \bar{n} , we have shown above that $\omega'(\tilde{Y}) = 0$; let X be an arbitrary element of \bar{n} ; then

$$\tilde{Y}\omega'(\tilde{X}) = \tilde{X}\omega'(\tilde{Y}) + \omega'([\tilde{Y}, \tilde{X}]) = 0.$$

Consequently, the functions $\omega'(\tilde{X})$ are constant on the orbits of the center of \bar{N} under the action τ ; on the other hand, $\omega' \in \mathbf{H}^p$ implies that $\omega'(\tilde{X})$ is in L^p of the orbits of $\tau(\bar{N})$, so that these functions cannot be constant on the submanifolds $\tau(\exp \bar{g}_{2\alpha})$ unless they are identically zero. We then conclude that $\omega'(\tilde{X}) = 0$ for all $X \in \bar{n}$, and exactly the same argument shows that $\omega'(\tilde{W}) = 0$. Thus $\omega' \equiv 0$, proving that $\omega = \omega_f$, as claimed.

Theorems 3 and 4 can be summarized by

COROLLARY. *Let $\omega \in \mathbf{H}^p$; then $\omega(\tilde{W}), \omega(\tilde{X}), \omega(\tilde{Y})$ have boundary values in L^p, L^q , and L^p respectively, and if $f =$ boundary value of $\omega(\tilde{W})$, then $\omega = \omega_f$.*

REMARK. As observed before, the results of §§2-3 hold under the assumption $m = m_\alpha + 2m_{2\alpha} \geq 3$. In the two remaining cases $m = 1$ (usual upper-

half plane) and $m = 2$ (three-dimensional real hyperbolic space), appropriate modifications in the definitions of the Riesz transforms and in the length of the vector field \tilde{W} , yield similar results. In particular, for $m = 1$, one gets the classical relation between conjugate harmonic functions in the upper-half plane and the Hilbert transform on the line.

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